

# AMSC 661      HW #1

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$$(1) \quad y' = f(t, y), \quad y(t_0) = y_0,$$

**Theorem 1.** Consider IVP (1). Suppose  $f$  is a continuous function of  $t$  and  $y$  defined on the cylinder

$$Q := \{t_0 \leq t \leq T, \|y - y_0\| \leq r\}$$

and  $\|f\| \leq M$  on  $Q$ .

- (1) Then IVP (1) has a solution which exists for  $0 \leq t - t_0 \leq \min(\frac{r}{M}, T - t_0)$ .
- (2) If, in addition,  $f$  is Lipschitz in  $y$ , i.e., for some constant  $L$ ,

$$\|f(t, y_1) - f(t, y_2)\| \leq L\|y_1 - y_2\|,$$

then the solution is unique.

**Example 1.**

$$(2) \quad y' = y^2, \quad y(t_0) = y_0.$$

The solution is obtained by the following sequence of steps:

$$\frac{dy}{y^2} = dt, \quad -\frac{1}{y} = t - C, \quad -\frac{1}{y_0} = t_0 - C, \quad y = \frac{y_0}{1 + y_0(t_0 - t)}.$$

The solution blows up at time  $t^* = t_0 + 1/y_0$ . Thus, in this example, the solution to IVP (2) exists, it is unique, but it blows up at a finite time. Hence, we must choose  $T < t^*$ .

**Exercise 1.** Apply Theorem 1, part 1, to Example 1: fix  $r \in \mathbb{R}$ , define the cylinder  $Q$ , express  $M$  via  $r$ , and find the interval of time where the existence of the solution is guaranteed.

Fix  $r \in \mathbb{R}$ . Define cylinder  $Q = \{(t, y) : t_0 \leq t \leq T, |y - y_0| \leq r\}$ .

which means

$$y_0 - r \leq y \leq y_0 + r$$

since  $f = y^2$ , to bound  $f$  inside  $Q$

$$|f(t, y)| = |y^2| \leq \max_{y \in [y_0 - r, y_0 + r]} y^2 = M$$

$$\Rightarrow M = \max \{ (y_0 - r)^2, (y_0 + r)^2 \}$$

$$\text{set } M = (y_0 + r)^2$$

Now, by Theorem 1, the solution exist for

$$0 \leq t-t_0 \leq \min\left(\frac{r}{\mu}, T-t_0\right)$$

$$\Rightarrow t-t_0 \leq \min\left(\frac{r}{(Y_{t_0}+r)^2}, T-t_0\right)$$

which is consistent with the example 1.

**Example 2.**

$$(3) \quad y' = 2\sqrt{y}, \quad y(0) = 0.$$

One obvious solution is  $y(t) \equiv 0$ . The other solution,  $y(t) = t^2$ , is obtained by the following sequence of steps:

$$\frac{dy}{2\sqrt{y}} = dt, \quad \sqrt{y} = t + C, \quad 0 = C, \quad y = t^2.$$

Furthermore, there is a one-parameter family of solutions

$$y(t) = \begin{cases} 0, & 0 \leq t \leq t_0, \\ (t - t_0)^2, & t \geq t_0, \end{cases}$$

for any  $0 \leq t_0 < \infty$ . Thus, in this example, the solution to IVP (3) exists for  $0 \leq T < \infty$  but it is not unique.

**Exercise 2.** Apply Theorem 1, parts 1 and 2, to Example 2. As in the previous exercise, find the interval where the existence of the solution is guaranteed. Then check that the function  $f(y) = 2\sqrt{y}$  is not Lipschitz at  $y = 0$ .

Define cylinder  $\Omega = \{(t, y) \mid t_0 \leq t \leq T, |y - y_0| \leq r\}$

since  $f(y) = 2\sqrt{y}$ , we bound  $L$  in  $\Omega \Rightarrow$

$$M = \max_{y \in [y_0 - r, y_0 + r]} 2\sqrt{y}$$

$$\Rightarrow M = 2\sqrt{y_0 + r}$$

Then, by Theorem 1 part 1

$$t - t_0 \leq \frac{r}{M} = \frac{r}{2\sqrt{y_0 + r}}$$

which agrees with ex. 2. Furthermore, to check Lipschitz cond.

$$|f(y_1) - f(y_2)| = |2\sqrt{y_1} - 2\sqrt{y_2}| \leq L|y_1 - y_2|$$

$$\Rightarrow \frac{2|\sqrt{y_1} - \sqrt{y_2}|}{|y_1 - y_2|} \leq L$$

where we assumed  $y_1 \neq y_2$ . Using the mean value theorem,  $\exists c$

between  $y_1$  and  $y_2$  such that

$$\frac{d}{dy} (2\sqrt{y}) \Big|_{y=c} = \frac{1}{\sqrt{c}}$$

as  $y \rightarrow 0$ ,  $\frac{1}{\sqrt{c}} \rightarrow \infty$ , hence,  $f(y)$  is not Lipschitz.

$y=0$ .

**Exercise 5.** Show that the three-step Adams-Basforth method

$$(34) \quad u_{n+1} = u_n + h \left( \frac{23}{12} f(t_n, u_n) - \frac{4}{3} f(t_{n-1}, u_{n-1}) + \frac{5}{12} f(t_{n-2}, u_{n-2}) \right)$$

is consistent of order 3.

We will check the local truncation error  $\tau$  given as

$$\tau_{n+1} = y_{n+1} - y_n - h \left( \frac{23}{12} f_n - \frac{4}{3} f_{n-1} + \frac{5}{12} f_{n-2} \right)$$

Taylor expanding, we have

$$y_{n+1} = y_n + h y'_n + \frac{h^2}{2} y''_n + \frac{h^3}{6} y'''_n + O(h^4)$$

$$= y_n + h f_n + \frac{h^2}{2} f'_n + \frac{h^3}{6} f''_n + O(h^4)$$

$$f_{n-1} = f_n - h f'_n + \frac{h^2}{2} f''_n + O(h^3)$$

$$f_{n-2} = f_n - 2h f'_n + 2h^2 f''_n + O(h^3)$$

Substituting:

$$\begin{aligned} \tau_{n+1} &= h f_n + \frac{h^2}{2} f'_n + \frac{h^3}{6} f''_n - h \left[ \frac{23}{12} f_n - \frac{4}{3} \left( f_n - h f'_n + \frac{h^2}{2} f''_n \right) \right. \\ &\quad \left. + \frac{5}{12} \left( f_n - 2h f'_n + 2h^2 f''_n \right) + O(h^3) \right] \end{aligned}$$

$$\begin{aligned} &= h \left\{ f_n + \frac{h}{2} f'_n + \frac{h^2}{6} f''_n \right\} - h \left\{ f_n + \frac{h}{2} f'_n - \frac{h^2}{6} f''_n \right\} + O(h^4) \\ &= O(h^4) \quad \Rightarrow \text{consistent of order 3.} \end{aligned}$$

4. (3 pts) Show that the implicit midpoint rule (a symplectic method)

$$k = f \left( t_n + \frac{h}{2}, u_n + \frac{h}{2}k \right), \quad u_{n+1} = u_n + hk. \quad (1)$$

is consistent (at least) of order 2.

Taylor expanding  $k$  around  $t_n$ :

$$k = f \left( t_n + \frac{h}{2}, u(t_n) + \frac{hk}{2} \right)$$

$$= f(t_n, u(t_n)) + \frac{h}{2} \frac{\partial f}{\partial t}(t_n, u(t_n)) + \frac{h}{2} \frac{\partial^2 f}{\partial t^2}(t_n, u(t_n)) k + O(h^2)$$

Solve for  $k$

$$k = f(t_n, u(t_n)) + \frac{h}{2} \left[ \frac{\partial f}{\partial t}(t_n, v_1) + \frac{\partial^2 f}{\partial t \partial u}(t_n, v_1) f(t_n, v_1) \right] + O(h^2)$$

$$= f(t_n, v_1) + \frac{h}{2} f'(t_n, v_1) + O(h^2)$$

Substituting back to the scheme

$$u_{n+1} = v_n + h \left[ f + \frac{h}{2} f' + O(h^2) \right]$$

$$\begin{aligned} \Rightarrow v_{n+1} &= u_{n+1} - v_n + h \left[ f + \frac{h}{2} f' + O(h^2) \right] \\ &= \frac{h^3}{6} f''(t_n, u(t_n)) + O(h^4) \end{aligned}$$

$\Rightarrow$  2nd order consistent!

5. (8pts)

$$y_n = f_n$$

- (a) Use the method of undetermined coefficients to determine  $a_0, a_1, b_0$ , and  $b_1$  that make the linear two-step explicit method consistent of as high order as possible:

$$u_{n+1} - a_0 u_n - a_1 u_{n-1} = h(b_0 f_n + b_1 f_{n-1}). \quad (2)$$

To ensure consistency, we will assume the method is exact for  $y(t)=1$ ,  $y(t)=t$ ,  $y(t)=t^2$ , polynomials of degree as high as possible.

$$\underline{y(t)=1 \text{ gives: } 1 - a_0 - a_1 = 0 \Rightarrow a_0 + a_1 = 1} \quad \color{red}{\times}$$

$$\underline{y(t)=t \text{ gives: } (1 - a_0 - a_1)t_n + h(1 + a_1) = h(b_0 + b_1) \Rightarrow 1 + a_1 = b_0 + b_1}$$

$$\underline{y(t)=t^2 \text{ gives: } (1 - a_0 - a_1)t_n^2 + (2h + 2a_1h - 2h^2b_0 - 2h^2b_1)t_n + (h^2 - a_1h^2 + 2h^2b_1) = 0}$$

$$\Rightarrow 2 + 2a_1 - 2b_0 - 2b_1 = 0 \Rightarrow 1 + a_1 = b_0 + b_1 \quad \color{red}{\times}$$

$$1 - a_1 + 2b_1 = 0 \Rightarrow a_1 = 1 + 2b_1 \quad \color{red}{\times}$$

Solving the system of eqns  $\color{red}{(\times)}$ , one gets

$$a_0 = 0, \quad a_1 = 0, \quad b_0 = 2, \quad b_1 = 0$$

$$\Rightarrow \underline{\underline{u_{n+1} - u_{n-1} = 2h f_n}}$$

Leap-frog method: 2<sup>nd</sup> order consistent.

5. (8pts)

- (a) Use the method of undetermined coefficients to determine  $a_0$ ,  $a_1$ ,  $b_0$ , and  $b_1$  that make the linear two-step explicit method consistent of as high order as possible:

$$u_{n+1} - a_0 u_n - a_1 u_{n-1} = h(b_0 f_n + b_1 f_{n-1}). \quad (2)$$

Let us start by Taylor expanding:

$$y_{n+1} = y_n + hy_n' + \frac{h^2 y_n''}{2} + \frac{h^3}{6} y_n''' + O(h^4)$$

$$y_{n-1} = y_n - hy_n' + \frac{h^2}{2} y_n'' - \frac{h^3}{6} y_n''' + O(h^4)$$

$$\underline{f_{n+1} = f_n - hf_n' + \frac{h^2}{2} f_n'' - \frac{h^3}{6} f_n''' + O(h^4)}$$

$$\stackrel{\text{LHS}}{\Rightarrow} \underbrace{(1-a_0-a_1)y_n}_{\text{red}} + \underbrace{(1+a_1)y_n'}_{\text{blue}} h + \underbrace{(1-a_1)\frac{y_n''}{2} h^2}_{\text{green}} + \underbrace{(1+a_1)\frac{h}{6} y_n'''}_{\text{purple}} + \dots$$

$$\stackrel{\text{RHS}}{\Rightarrow} h \left[ \underbrace{(b_0+b_1)f_n}_{\text{blue}} - \underbrace{hb_1 f_n'}_{\text{green}} + \underbrace{\frac{h^2}{2} f_n'' b_1}_{\text{purple}} + O(h^3) \right]$$

Since we demand  $y_n' = f_n$ , we must have

$$(1-a_0-a_1) = 0 \quad \text{•}$$

$$(1+a_1) = (b_0+b_1) \quad \text{•}$$

$$(1-a_1)\frac{1}{2} = -b_1 \quad \text{•}$$

$$\frac{1+a_1}{6} = \frac{b_1}{2} \quad \text{•}$$

Solving the system of eqns, we get!

$$\begin{aligned} \frac{(1-a_1)}{2} + \frac{1+a_1}{3} &= 0 \Rightarrow 3-3a_1+2+2a_1 = 0 \\ &\Rightarrow 5-a_1=0 \\ &\Rightarrow a_1=5 \end{aligned}$$

$$\Rightarrow b_1 = 2$$

$$\Rightarrow b_0 = 4$$

$$\Rightarrow c_0 = -4$$

Hence the scheme needs:

$$u_{n+1} + 4u_n - 5u_{n-1} = h(4f_n + 2f_{n-1})$$

This ensures  $\tau \sim O(h^3)$ , hence 2<sup>nd</sup> order

consistent

(b) Show that this method is unstable by applying it to the ODE  $y' = 0$  and mimicking the technique shown in Section 3.5.

for  $y' = 0$ , we have

$$u_{n+1} + 4u_n - 5u_{n-1} = 0$$

with the characteristic equation

$$r^2 + 4r - 5 = 0$$

with roots  $r_1 = 1, r_2 = -5$ , the general solution will be of the form of  $A r_1^n + B r_2^n$ , thus, any slight perturbation that will make  $B \neq 0$ , will grow infinitely and diverge.

(c) Apply this method to the 2D gravity problem with a unit-circle solution:

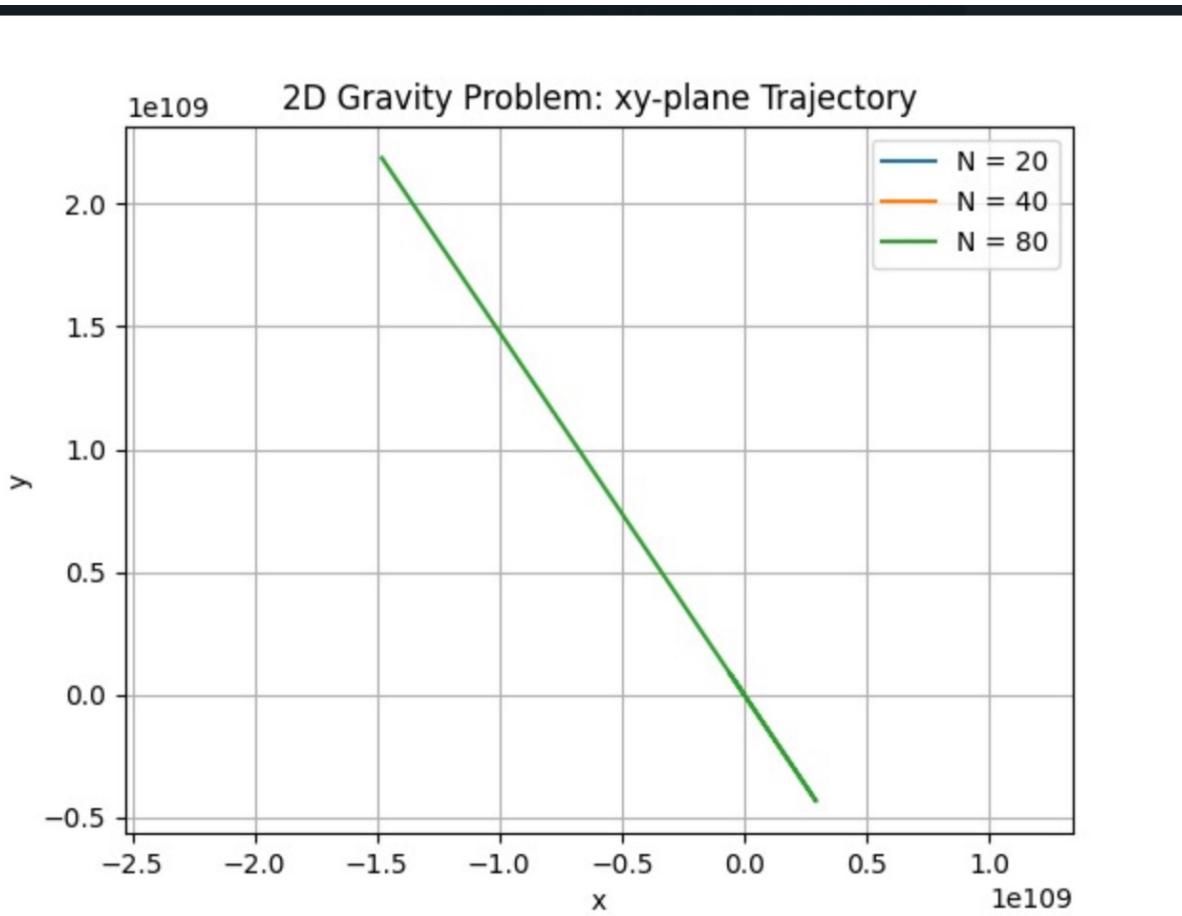
$$\frac{d}{dt} \begin{bmatrix} x \\ y \\ u \\ v \end{bmatrix} = \begin{bmatrix} u \\ v \\ -\frac{x}{x^2+y^2} \\ -\frac{y}{x^2+y^2} \end{bmatrix}, \quad \begin{bmatrix} x \\ y \\ u \\ v \end{bmatrix}(0) = \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix} \quad (3)$$

Integrate the numerical solution for two periods. Show that the solution blows up by reporting the norm of the solution at time  $4\pi$  for  $h = 2\pi/N$  with  $N = 20, 40, 80$ . Also, plot the  $x$  and  $y$  components of the solution in the  $xy$  plane.

We implemented the code and observed that the norm blows up as expected:

```
For N = 20, h = 0.3142, norm at t = 4π: 474578335703168451700850688.0000
For N = 40, h = 0.1571, norm at t = 4π: 1210668141203986655650221859108781593405687928153178112.0000
C:\Users\ekrem\Desktop\AMSC661\hw1.py:10: RuntimeWarning: overflow encountered in scalar power
  dudt = -x / r**3
C:\Users\ekrem\Desktop\AMSC661\hw1.py:11: RuntimeWarning: overflow encountered in scalar power
  dvdt = -y / r**3
For N = 80, h = 0.0785, norm at t = 4π: 26610149085859516499407786099251966827105915487970803063150720277212110786845480361712331906012909597131538432.0000
```

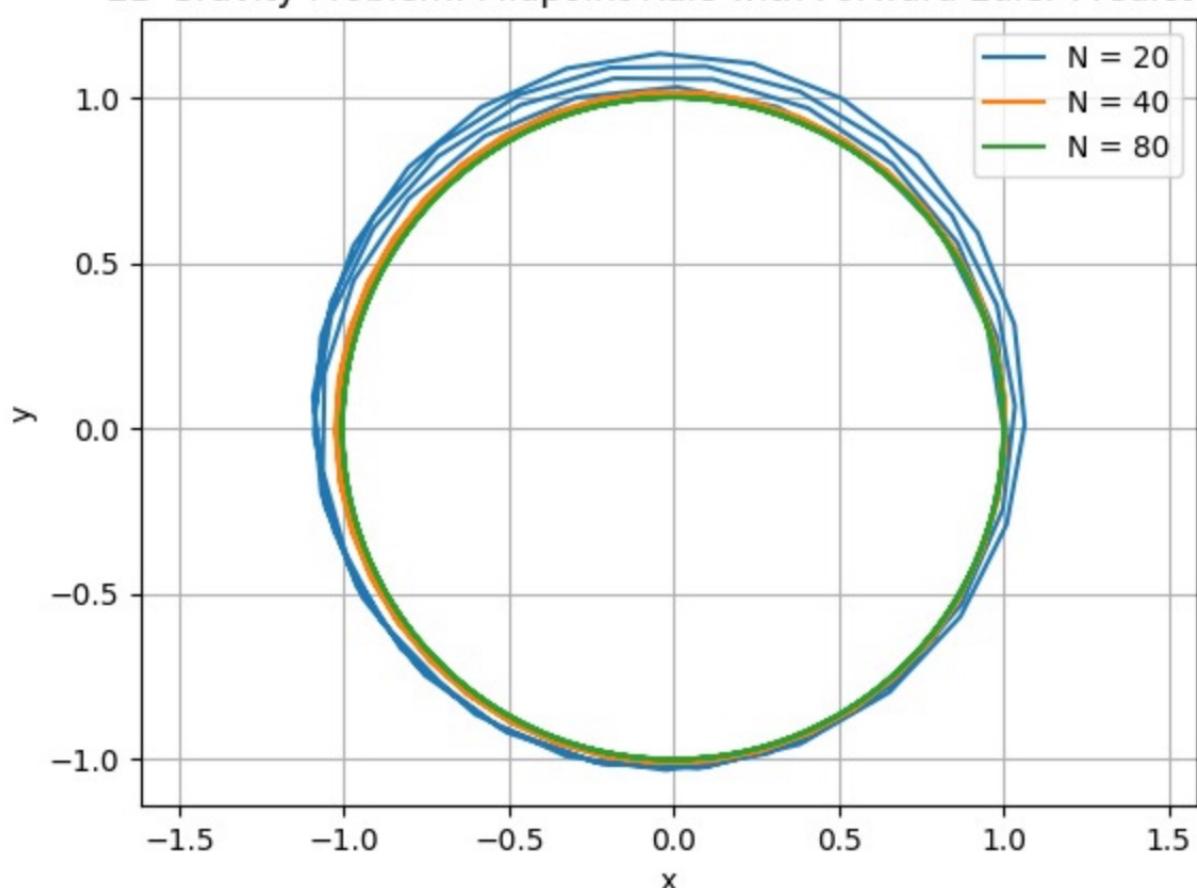
Hence, orbits are not meaningful when plotted:



- (d) Apply the midpoint rule with the Forward Euler predictor to the 2D gravity problem with a unit-circle solution. Use the time interval  $8\pi$ , set  $h = 2\pi/N$  with  $N = 20, 40, 80$  and plot the numerical solutions.

We repeated part c , this time with midpoint rule with forward Euler:

2D Gravity Problem: Midpoint Rule with Forward Euler Predictor



The solution is stable for the time interval  $T = 8\pi$  as seen.

Codes can be found:

<https://github.com/ekremdemirboga/AMSC661/tree/main/HW1>