

Finite difference schemes for time-dependent problems, part I

In the previous lecture we have seen how to approximate the solution to $-p \partial_x^2 u + q \partial_x u + ru = f$.

Let us now discuss how to solve the time-dependent problem

$$\partial_t u - p \partial_x^2 u + q \partial_x u + ru = f.$$

For simplicity we will set $q=0$ and $r=0$ and only discuss the "diffusion" equation $\partial_t u - p \partial_x^2 u = f$. (You can simply add back the missing terms from the previous lecture...)

Recall that for the spatial discretization we chose

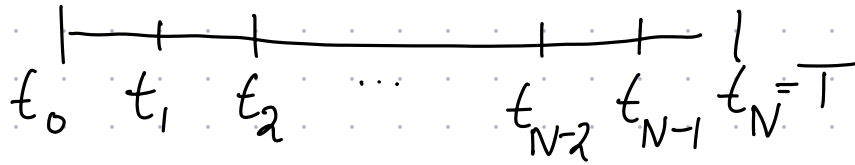
Spatial discretization: Subdivide interval $\Omega = [a, b]$ into $M+1$ subintervals of equal size: $x_i = a + ih$, $h = \frac{b-a}{M+1}$

$$\begin{array}{ccccccc} | & | & | & & | & | & | \\ x_0 = a & x_1 & x_2 & \dots & x_{M-1} & x_M & x_{M+1} = b \end{array}$$

index i

We will now augment this by discretizing in time:

Temporal discretization: Subdivide interval $I = [t_0, T]$ into N subintervals of equal size: $t_n = t_0 + \tau n$, $\tau = \frac{T - t_0}{N}$



index n

And we approximate now: $U_i^n \approx u(t_n, x_i)$

$$\frac{\partial_t u}{\tau} - \frac{\partial_x^2 u}{h^2} = f$$

$$\frac{U_i^{n+1} - U_i^n}{\tau} - \frac{1}{h^2} (U_{i+1}^n - 2U_i^n + U_{i-1}^n) = f(t_n, x_i)$$

at time t_n at time t_n

forward Euler:

Combining this with boundary and initial conditions:

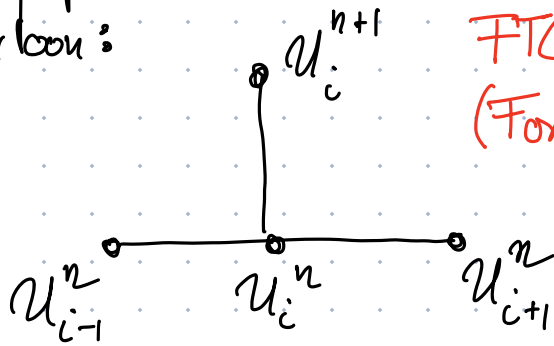
Task (Forward Euler)

Given $\{U_i^n\}_{i=0}^{M+1}$ at time t_n and boundary values g_a, g_b construct $\{U_i^{n+1}\}_{i=0}^{M+1}$ at time t_{n+1} by setting

$$(*) \quad U_i^{n+1} = U_i^n + \tau \frac{1}{h^2} (U_{i+1}^n - 2U_i^n + U_{i-1}^n) + \tau f(t_n, x_i)$$

$$\text{and } U_0^{n+1} = g_a \text{ and } U_{M+1}^{n+1} = g_b$$

This is again an explicit time-marching scheme. We can simply compute u_i^{n+1} without the need to solve a linear system! (Carson: FTCS
(Forward time centered space))



But, it is only first-order in time (in contrast to second order in space). Worse, there is another problem: For this method to work we need a time step size restriction:

$$\tau \text{ has to be chosen such that } \tau \leq \frac{1}{2\rho} h^2.$$

"parabolic CFL condition"

named after Richard Courant, Kurt Friedrichs, Hans Lewy

Remark: The CFL condition can be derived by analyzing carefully how (*) propagates and "amplifies" error. Say, the approximation $\{u_i^n\}$ at time t_n has the error $e_i^n = u(t_n, x_i) - u_i^n$, then how large is the error of $\{u_i^{n+1}\}$ at time t_{n+1} , $e_i^{n+1} = u(t_{n+1}, x_i) - u_i^{n+1}$?

Such an analysis is known as "**von Neumann stability analysis**".
 Cutting a lot of corners, let's assume that the error is given
 by a combination of "error modes"

$$\hat{e}_m^n = a^n e^{i k m}$$

\hat{e} writing m instead of \hat{e}

a = amplification factor
 k = wave number

(i = imaginary unit)

that each solve (*) : $\hat{e}_m^{n+1} = (1 - 2\rho \tau/h^2) \hat{e}_m^n + \rho \tau/h^2 (\hat{e}_{m-1}^n + \hat{e}_{m+1}^n)$.

Substituting: $a^{n+1} e^{i k m} = a^n e^{i k m} + \rho \tau/h^2 (\underbrace{e^{-i k} - 2 + e^{i k}}_{= -4 \sin^2(k/2)}) a^n e^{i k m}$

and dividing by $a^n e^{i k m}$: $a = 1 - 4\rho \tau/h^2 \sin^2(k/2)$

≥ 0 (which we have not discussed)

We now want that the amplification factor a
 has an absolute value not greater than one: $|a| \leq 1$.

This implies that $1 \geq a \geq 1 - 4\rho \tau/h^2 \geq -1$
 always true \uparrow \sin is at most one

Rearranging: $4\rho \tau/h^2 \leq 2$.

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If we want to avoid this time-step size restriction we
 need to use an **implicit time stepping** method:

$$\partial_t u - \rho \partial_x^2 u = f$$

$$\frac{u_i^{n+1} - u_i^n}{\tau} - \star = \star$$

trapezoidal rule: half at time t_n and half at time t_{n+1} !

$$\frac{u_i^{n+1} - u_i^n}{\tau} - \frac{\tau}{2h^2} \left(\underbrace{u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1}}_{\text{at time } t_{n+1}} + \underbrace{u_{i+1}^n - 2u_i^n + u_{i-1}^n}_{\text{at time } t_n} \right)$$

$$= \frac{\tau}{2} \left(f(t_n, x_i) + f(t_{n+1}, x_i) \right)$$

Writing this out again:

Task (Crank-Nicolson scheme)

Given $\{u_i^n\}_{i=0}^{M+1}$ at time t_n and boundary values g_a, g_b
construct $\{u_i^{n+1}\}_{i=0}^{M+1}$ at time t_{n+1} by solving

$$u_i^{n+1} - \rho \frac{\tau}{2h^2} (u_{i+1}^{n+1} - 2u_i^{n+1} + u_{i-1}^{n+1})$$

$$= u_i^n + \rho \frac{\tau}{2h^2} (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$$

$$+ \frac{\tau}{2} \left\{ f(t_n, x_i) + f(t_{n+1}, x_i) \right\}$$

for $i=1, \dots, M$

and $u_0^{n+1} = g_a$ and $u_{M+1}^{n+1} = g_b$

Remark: In this case the stability analysis results in

$$a (1 + 2\rho \tau/h^2 \sin(k/2)) = (1 - 2\rho \tau/h^2 \sin(k/2))$$

$$\Rightarrow |a| = \left| \frac{1 - 2\rho \tau/h^2 \sin(k/2)}{1 + 2\rho \tau/h^2 \sin(k/2)} \right| \leq 1 \text{ unconditionally.}$$

Cartoon for Crank-Nicolson scheme:

