

# Introduction to Quantum Computation, UPB

## Winter 2022, Assignment 2

To be completed by: Thursday, October 27

Eli Kogan-Wang

### 1 Exercises

1. Use the spectral decompositions of  $X$  and  $Z$  to prove that  $HXH^\dagger = Z$ . (Do not simply write out the matrices and multiply!) Why does this immediately also yield that  $HZH^\dagger = X$ ?

$$X = (|+\rangle\langle+| - |-\rangle\langle-|)$$

$$Z = (|0\rangle\langle 0| - |1\rangle\langle 1|)$$

Since for  $H$  we have  $H = H^\dagger = H^{-1}$ , we can write  $HXH^\dagger = HXH$ .

For the Hadamard gate, we have  $H|0\rangle = |+\rangle$  and  $H|1\rangle = |-\rangle$ , so we can write:

$$H = |+\rangle\langle 0| + |-\rangle\langle 1| \quad (\text{trivial, since Orthonormal-Change-of-Basis})$$

Additionally, we have  $H|+\rangle = |0\rangle$  and  $H|-\rangle = |1\rangle$ , so we can write:

$$H = |0\rangle\langle+| + |1\rangle\langle-| \quad (\text{trivial, since Orthonormal-Change-of-Basis})$$

Now, we can write:

$$\begin{aligned} HXH^\dagger &= H(|+\rangle\langle+| - |-\rangle\langle-|)H \\ &= (|0\rangle\langle+| + |1\rangle\langle-|)(|+\rangle\langle+| - |-\rangle\langle-|)(|+\rangle\langle 0| + |-\rangle\langle 1|) \end{aligned}$$

We will not write out the  $\langle-|+\rangle$  and  $\langle+|-\rangle$  terms, since they are zero.

$$\begin{aligned} &= (|0\rangle\langle+| + |1\rangle\langle-|)(|+\rangle\langle+| - |-\rangle\langle-|)(|+\rangle\langle 0| + |-\rangle\langle 1|) \\ &= (|0\rangle\langle+| + |1\rangle\langle-|)(|+\rangle\langle+| - |-\rangle\langle-|)(|+\rangle\langle 0| + |-\rangle\langle 1|) \\ &= (|0\rangle\langle 0| - |1\rangle\langle 1|) \\ &= Z \end{aligned}$$

From this it immediately follows, that  $HZH^\dagger = X$ :

$$\begin{aligned} HXH^\dagger &= Z & | \cdot H & \quad (\text{since } H = H^\dagger = H^{-1}) \\ \iff HX &= ZH & | H^\dagger. \\ \iff X &= HZH^\dagger \end{aligned}$$

2. (a) Write out the 4-dimensional vector for  $(\alpha|0\rangle + \beta|1\rangle) \otimes (\gamma|0\rangle + \delta|1\rangle)$ .

$$\begin{bmatrix} \alpha \cdot \begin{bmatrix} \gamma \\ \delta \end{bmatrix} \\ \beta \cdot \begin{bmatrix} \gamma \\ \delta \end{bmatrix} \end{bmatrix} = \begin{bmatrix} \alpha\gamma \\ \alpha\delta \\ \beta\gamma \\ \beta\delta \end{bmatrix} = \alpha\gamma|00\rangle + \alpha\delta|01\rangle + \beta\gamma|10\rangle + \beta\delta|11\rangle$$

(b) Let  $\mathcal{B}_1 = \{|\psi_1\rangle, |\psi_2\rangle\}$  and  $\mathcal{B}_2 = \{|\phi_1\rangle, |\phi_2\rangle\}$  be two orthonormal bases for  $\mathbb{C}^2$ . Prove that

$$\mathcal{B}_3 = \{|\psi_1\rangle \otimes |\phi_1\rangle, |\psi_1\rangle \otimes |\phi_2\rangle, |\psi_2\rangle \otimes |\phi_1\rangle, |\psi_2\rangle \otimes |\phi_2\rangle\}$$

is an orthonormal basis for  $\mathbb{C}^4$ . In other words, show that for each  $|v\rangle \in \mathcal{B}_3$ ,  $\| |v\rangle \|_2 = 1$ , and for all pairs of distinct  $|v\rangle, |w\rangle \in \mathcal{B}_3$ ,  $\langle v|w\rangle = 0$ .

Every  $|v\rangle \in \mathcal{B}_3$  is of the form  $|\psi_i\rangle \otimes |\phi_j\rangle$ , where  $i, j \in \{1, 2\}$ .

Let  $i, j \in \{1, 2\}$ . Now let:

$$|\psi_i\rangle = \begin{bmatrix} \alpha \\ \beta \end{bmatrix}$$

$$|\phi_j\rangle = \begin{bmatrix} \gamma \\ \delta \end{bmatrix}$$

Since  $\mathcal{B}_1, \mathcal{B}_2$  are orthonormal bases, we have  $\alpha^2 + \beta^2 = 1$  and  $\gamma^2 + \delta^2 = 1$ .

Since:

$$|\psi_i\rangle \otimes |\phi_j\rangle = \begin{bmatrix} \alpha\gamma \\ \alpha\delta \\ \beta\gamma \\ \beta\delta \end{bmatrix}$$

We need to show that  $\alpha^2\gamma^2 + \alpha^2\delta^2 + \beta^2\gamma^2 + \beta^2\delta^2 = 1$ .

$$\begin{aligned} \alpha^2\gamma^2 + \alpha^2\delta^2 + \beta^2\gamma^2 + \beta^2\delta^2 &= \alpha^2 \cdot (\gamma^2 + \delta^2) + \beta^2 \cdot (\gamma^2 + \delta^2) \\ &= \alpha^2 \cdot 1 + \beta^2 \cdot 1 \\ &= 1 \end{aligned}$$

We have shown that, for each  $|v\rangle \in \mathcal{B}_3$ ,  $\| |v\rangle \|_2 = 1$ .

Now, let  $|v\rangle, |w\rangle \in \mathcal{B}_3$  be distinct. Then,  $|v\rangle, |w\rangle$  are of the form  $|\psi_i\rangle \otimes |\phi_j\rangle$  and  $|\psi_k\rangle \otimes |\phi_l\rangle$ , respectively, where  $i, j, k, l \in \{1, 2\}$ . And  $(i, j) \neq (k, l)$  or in other words,  $i \neq k$  or  $j \neq l$ .

Now:

$$\begin{aligned} \langle v|w\rangle &= (\langle\psi_i| \otimes \langle\phi_j|)(|\psi_k\rangle \otimes |\phi_l\rangle) \\ &= \langle\psi_i|\psi_k\rangle \cdot \langle\phi_j|\phi_l\rangle \\ &\text{Since } i \neq k \text{ or } j \neq l \text{ and the } \psi\text{'s and } \phi\text{'s are orthogonal,} \\ &= 0 \end{aligned}$$

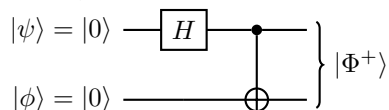
We have successfully shown, that  $\mathcal{B}_3$  is an orthonormal basis for  $\mathbb{C}^4$ .

3. (a) Prove that  $(Z \otimes Y)^\dagger = Z \otimes Y$ . Do not write out any matrices explicitly; rather, you must use the properties of the tensor product, dagger, and  $Y$ .

$$\begin{aligned} (Z \otimes Y)^\dagger &= Z^\dagger \otimes Y^\dagger && |\dagger \text{ distributes over } \otimes \\ &= Z \otimes Y && |Z, Y \text{ are self-adjoint} \end{aligned}$$

- (b) In class, we saw a quantum circuit which, given starting state  $|0\rangle \otimes |0\rangle$ , prepared the Bell state  $|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$ . In fact, that circuit is a change of basis matrix, mapping the standard basis  $|00\rangle, |01\rangle, |10\rangle, |11\rangle$  to the Bell basis  $|\Phi^+\rangle, |\Psi^+\rangle, |\Phi^-\rangle, |\Psi^-\rangle$ .

In fact, in class we saw:



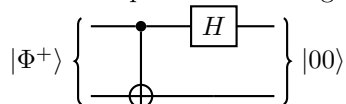
Which is equivalent to:

$$\text{CNOT}(H \otimes I) |0\rangle |0\rangle = |\Phi^+\rangle$$

- i. Write down a quantum circuit which maps  $|\Phi^+\rangle$  to  $|0\rangle \otimes |0\rangle$ . (Hint: Write out the circuit from class as a sequence of matrix operations. Then, recalling that the inverse of any unitary operation  $U$  is given by  $U^\dagger$ , take the inverse of this entire sequence by taking the dagger.)

$$\begin{aligned} (\text{CNOT}(H \otimes I))^{-1} &= ((H \otimes I)^{-1} \text{CNOT}^{-1}) \\ &= ((H \otimes I)^\dagger \text{CNOT}^\dagger) \\ &= ((H^\dagger \otimes I^\dagger) \text{CNOT}) \\ &= (H \otimes I) \text{CNOT} \end{aligned}$$

This corresponds to the diagram:



- ii. Your circuit from 3(b)(i) is actually a change of basis which maps the Bell basis *back* to the standard basis. To verify this, run your circuit from 3(b)(i) on input  $|\Psi^-\rangle$  and check that the output is  $|1\rangle \otimes |1\rangle$ .

$$\begin{aligned} (H \otimes I) \text{CNOT} |\Psi^-\rangle &= (H \otimes I) \text{CNOT} \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle) \\ &= (H \otimes I) \frac{1}{\sqrt{2}} \text{CNOT}(|01\rangle - |10\rangle) \\ &= (H \otimes I) \frac{1}{\sqrt{2}} (\text{CNOT} |01\rangle - \text{CNOT} |10\rangle) \\ &= (H \otimes I) \frac{1}{\sqrt{2}} (|01\rangle - |11\rangle) \\ &= (H \otimes I) \frac{1}{\sqrt{2}} ((|0\rangle - |1\rangle) \otimes |1\rangle) \\ &= (H \otimes I) \left( \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \right) \otimes |1\rangle \\ &= (H \otimes I) (|-\rangle \otimes |1\rangle) \\ &= ((H |-\rangle) \otimes (I |1\rangle)) \\ &= |1\rangle \otimes |1\rangle \\ &= |11\rangle \end{aligned}$$