

Problema 1:

$$(a) \quad S(T) = \alpha \frac{T}{T_0} + \beta \cos\left(\frac{\pi T}{2 T_0}\right)$$

$$\rightarrow dS = \left[\frac{\alpha}{T_0} - \frac{\beta\pi}{2T_0} \sin\left(\frac{\pi T}{2 T_0}\right) \right] dT$$

$$\therefore T dS = \left[\alpha \frac{T}{T_0} - \frac{\beta\pi}{2} \frac{T}{T_0} \sin\left(\frac{\pi T}{2 T_0}\right) \right] dT$$

Ahora, sabemos que $dQ = C dT$; C : calor específico molar
y junto con $dQ = T dS$, tenemos que

$$\left\{ C = \alpha \frac{T}{T_0} - \frac{\beta\pi}{2} \frac{T}{T_0} \sin\left(\frac{\pi T}{2 T_0}\right) \right\}$$

$$\text{Ahora, } C < 0 \Rightarrow \alpha \frac{T}{T_0} < \frac{\beta\pi}{2} \frac{T}{T_0} \sin\left(\frac{\pi T}{2 T_0}\right)$$

$$\frac{2\alpha}{\beta\pi} < \sin\left(\frac{\pi T}{2 T_0}\right) \Rightarrow \text{Arcsin}\left(\frac{2\alpha}{\beta\pi}\right) < \frac{\pi T}{2 T_0}$$

$$\therefore \left\{ T > T_0 \cdot \frac{2}{\pi} \text{Arcsin}\left(\frac{2\alpha}{\beta\pi}\right) \right\}$$

$$(b) \quad dQ = C dT \Rightarrow Q = \int_{T_1}^{T_2} C dT$$

$$Q = \frac{\alpha}{T_0} \int_{T_1}^{T_2} T dT - \frac{\beta\pi}{2T_0} \int_{T_1}^{T_2} T \sin\left(\frac{\pi T}{2 T_0}\right) dT$$

$$I_1 = \int_{T_1}^{T_2} T dT = \frac{T^2}{2} \Big|_{T_1}^{T_2} = \frac{1}{2} (T_2^2 - T_1^2) = \frac{1}{2} \left(\frac{4T_0^2}{9} - \frac{T_0^2}{9} \right) = \frac{3T_0^2}{2 \cdot 9}$$

$$\boxed{I_1 = \frac{T_0^2}{6}}$$

$$I_2 = \int_{T_1}^{T_2} T \sin\left(\frac{\pi T}{2 T_0}\right) dT$$

$$u = T \quad du = dT$$

$$dv = \sin\left(\frac{\pi T}{2 T_0}\right) dT$$

$$v = -\frac{2T_0}{\pi} \cos\left(\frac{\pi T}{2 T_0}\right)$$

$$I_2 = -\frac{2T_0 T}{\pi} \cos\left(\frac{\pi T}{2T_0}\right) \Big|_{T_1}^{T_2} + \int_{T_1}^{T_2} \frac{2T_0}{\pi} \cos\left(\frac{\pi T}{2T_0}\right) dT$$

$$I_2 = \frac{2T_0}{\pi} \left\{ -T_2 \cos\left(\frac{\pi T_2}{2T_0}\right) + T_1 \cos\left(\frac{\pi T_1}{2T_0}\right) + \frac{2T_0}{\pi} \sin\left(\frac{\pi T_2}{2T_0}\right) - \frac{2T_0}{\pi} \sin\left(\frac{\pi T_1}{2T_0}\right) \right\}$$

$$I_2 = \frac{2T_0}{\pi} \left\{ -\frac{2T_0}{3} \cos\left(\frac{\pi}{3}\right) + \frac{T_0}{3} \cos\left(\frac{\pi}{6}\right) + \frac{2T_0}{\pi} \left[\sin\left(\frac{\pi}{3}\right) - \sin\left(\frac{\pi}{6}\right) \right] \right\}$$

$$I_2 = \frac{2T_0}{\pi} \left\{ -\frac{T_0}{3} + \frac{T_0 \sqrt{3}}{6} + \frac{2T_0}{\pi} \left(\frac{\sqrt{3}-1}{2} \right) \right\} = \frac{2T_0^2}{\pi} \left\{ \frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right\}$$

$$\therefore Q = \frac{\alpha}{T_0} \cdot I_1 - \frac{\beta \pi}{2T_0} \cdot I_2 = \frac{\alpha}{T_0} \cdot \frac{T_0^2}{6} - \frac{\beta \pi}{2T_0} \cdot \frac{2T_0^2}{\pi} \left\{ \frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right\}$$

$$Q = T_0 \left\{ \frac{\alpha}{6} - \beta \left(\frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right) \right\}$$

(c) Ya que $Q = \Delta U - W \rightarrow W = \Delta U - Q$

$$W = C_V (T_2 - T_1) - T_0 \left[\frac{\alpha}{6} - \beta \left(\frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right) \right]$$

$$W = C_V \left(\frac{2T_0}{3} - \frac{T_0}{3} \right) - T_0 \left[\frac{\alpha}{6} - \beta \left(\frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right) \right]$$

$$W = T_0 \left[\frac{C_V}{3} - \frac{\alpha}{6} + \beta \left(\frac{\sqrt{3}-2}{6} + \frac{\sqrt{3}-1}{\pi} \right) \right]$$

Problema 2 : (a) Partiendo combinando el 1º y 2º principio de la termodinámica: $TdS = dU + PdV$,
o bien,

$$dS = \frac{1}{T} dU + \frac{P}{T} dV$$

si $U = U(T, V) \rightarrow dU = \left(\frac{\partial U}{\partial T}\right)_V dT + \left(\frac{\partial U}{\partial V}\right)_T dV$

$$\therefore dS = \frac{1}{T} \left(\frac{\partial U}{\partial T}\right)_V dT + \left[\frac{1}{T} \left(\frac{\partial U}{\partial V}\right)_T + \frac{P}{T}\right] dV$$

Entonces,

$$\left(\frac{\partial S}{\partial T}\right)_V = \frac{1}{T} \left(\frac{\partial U}{\partial T}\right)_V$$

$$\left(\frac{\partial S}{\partial V}\right)_T = \frac{1}{T} \left(\frac{\partial U}{\partial V}\right)_T + \frac{P}{T}$$

Ahora, ya que dS es diferencial exacta,

$$\frac{\partial^2 S}{\partial T \partial V} = \frac{\partial^2 S}{\partial V \partial T}$$

$$\frac{\partial}{\partial T} \left[\frac{1}{T} \left(\frac{\partial U}{\partial V}\right)_T + \frac{P}{T} \right] = \frac{\partial}{\partial V} \left[\frac{1}{T} \left(\frac{\partial U}{\partial T}\right)_V \right]$$

$$-\frac{1}{T^2} \left(\frac{\partial U}{\partial V}\right)_T + \frac{1}{T} \frac{\partial}{\partial T} \left(\frac{\partial U}{\partial V}\right)_T + \frac{\partial}{\partial T} \left(\frac{P}{T}\right) = \frac{1}{T} \frac{\partial}{\partial V} \left(\frac{\partial U}{\partial T}\right)_V$$

$$\Rightarrow \left(\frac{\partial U}{\partial V}\right)_T = T^2 \frac{\partial}{\partial T} \left(\frac{P}{T}\right)$$

$$\left(\frac{\partial U}{\partial V}\right)_T = T^2 \left\{ \frac{1}{T} \left(\frac{\partial P}{\partial T}\right)_V - \frac{P}{T^2} \right\}$$

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial P}{\partial T}\right)_V - P$$

Ya que $P = f(V)T$, entonces

$$\left(\frac{\partial U}{\partial V}\right)_T = T f(V) - f(V)T = 0$$

De esta manera, U no depende de V .

$$(b) \quad P = \frac{1}{3} \mu(T) \quad ; \quad \mu = \frac{U}{V} \Rightarrow P = \frac{U}{3V} \wedge U = 3PV$$

usamos $\underbrace{\left(\frac{\partial U}{\partial V}\right)_T}_{\mu(T)} = T \underbrace{\left(\frac{\partial P}{\partial T}\right)_V}_{\frac{1}{3} \frac{d\mu}{dT}} - \underbrace{P}_{\frac{1}{3} \mu(T)}$

$$\mu(T) = T \frac{d\mu}{dT} - \frac{1}{3} \mu(T)$$

$$\frac{4}{3} \mu(T) = T \frac{d\mu}{dT} \Rightarrow \frac{d\mu}{\mu} = 4 \frac{dT}{T}$$

$$\Rightarrow \ln \mu = \ln T^4 + \underbrace{\ln C}_{\text{cte.}}$$

$$\Rightarrow \boxed{\mu(T) = C T^4}$$

← Gas de fotones.

Problema 3 : Ec. de estado: $\boxed{\bar{E} = \bar{E}_0 - \alpha (T - T_0)}$
 $\Rightarrow \frac{d\bar{E}}{dT} = -\alpha$

La energía interna:

$$U = U_0 + \left(\bar{E} - T \frac{d\bar{E}}{dT} \right) z + C_z (T - T_0)$$

$$\therefore U = U_0 + (\bar{E}_0 - \alpha T + \alpha T_0 + \alpha T) z + C_z (T - T_0)$$

$$U = U_0 + (\bar{E}_0 + \alpha T_0) z + C_z (T - T_0)$$

$$U = U_0 + (\bar{E}_0 + \alpha T_0) z + \frac{C_z}{\alpha} (\bar{E}_0 - \bar{E})$$

$$\rightarrow dU = (\bar{E}_0 + \alpha T_0) dz - \frac{C_z}{\alpha} d\bar{E}$$

Weg, $dQ = dU - \bar{E} dz$

$$\rightarrow \boxed{dQ = (\bar{E}_0 - \bar{E} + \alpha T_0) dz - \frac{C_z}{\alpha} d\bar{E}} \quad (*)$$

(a) Para representar un ciclo de Carnot necesitamos determinar las curvas adiabáticas e isotérmicas. De la $\bar{E} d\bar{E}$,

$$\bar{E} = \bar{E}_0 + \alpha T_0 - \alpha T$$

Si $T_f < T_c \Rightarrow \bar{E}_f > \bar{E}_c$, donde $\bar{E}_f = (\bar{E}_0 + \alpha T_0) - \alpha T_f$
 $\bar{E}_c = (\bar{E}_0 + \alpha T_0) - \alpha T_c$

las adiabáticas las obtenemos de (*), $dQ = 0$,

$$(\bar{E}_0 - \bar{E} + \alpha T_0) dz - \frac{C_z}{\alpha} d\bar{E} = 0$$

$$\rightarrow \frac{\alpha}{C_z} dz = \frac{d\bar{E}}{(\bar{E}_0 - \bar{E} + \alpha T_0)}$$

$$\frac{\alpha}{C_z} \int_{z_i}^z dz = \int_{\bar{E}_i}^{\bar{E}} \frac{d\bar{E}}{\bar{E}_0 - \bar{E} + \alpha T_0}$$

$$\frac{\alpha}{C_z} (z - z_i) = \ln \left(\frac{\bar{E}_0 - \bar{E}_i + \alpha T_0}{\bar{E}_0 - \bar{E} + \alpha T_0} \right)$$

$$\frac{\bar{E}_0 - \bar{E} + \alpha T_0}{\bar{E}_0 - \bar{E}_i + \alpha T_0} = e^{-\frac{\alpha}{C_z} (z - z_i)} = \frac{\bar{E}_0 - \bar{E}_i + \alpha T_0 + \bar{E}_i - \bar{E}}{\bar{E}_0 - \bar{E}_i + \alpha T_0} =$$

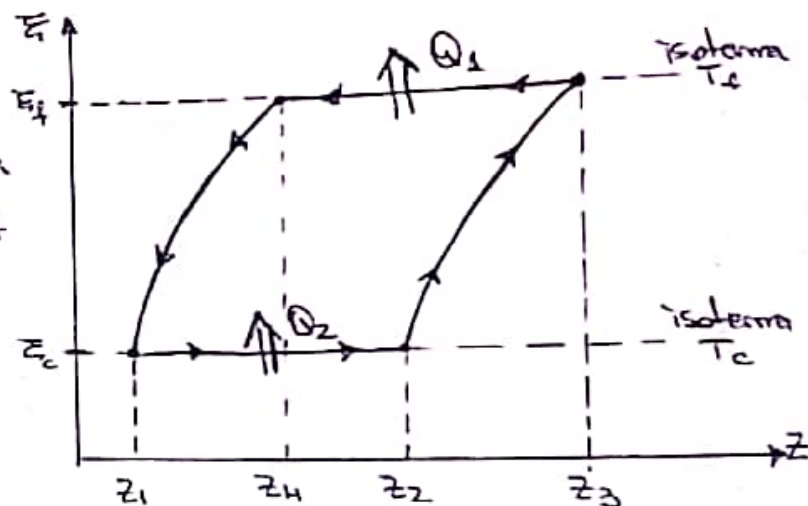
$$= 1 + \frac{(\bar{E}_i - \bar{E})}{\bar{E}_0 - \bar{E}_i + \alpha T_0}$$

$$e^{-\frac{\alpha}{T_0}(z-z_1)} - 1 = \frac{E_i - E_f}{E_0 - E_i + \alpha T_0} = -\frac{E_f - E_i}{E_0 - E_i + \alpha T_0}$$

$$\therefore \boxed{E_f(z) = E_i + (E_0 - E_i + \alpha T_0) (1 - e^{-\frac{\alpha}{T_0}(z-z_1)})}$$

ciclo de Carnot:

- i) Carga isotérmica a T_c
- ii) Carga adiabática; la temperatura disminuye desde T_c hasta T_f
- iii) Descarga isotérmica a T_f
- iv) Descarga adiabática; la temperatura aumenta desde T_f hasta T_c .



1) En la carga isotérmica, $E = E_c \rightarrow dE = 0$

$$\therefore \boxed{Q_2 = (E_0 - E_c + \alpha T_0)(z_2 - z_1) > 0}$$

$$\rightarrow \boxed{Q_2 = (E_0 - E_c + \alpha T_0) \Delta z_c} \quad \boxed{\Delta z_c = z_2 - z_1} \text{ : Carga}$$

2) En la descarga isotérmica, $E = E_f \rightarrow dE = 0$

$$\therefore \boxed{Q_1 = (E_0 - E_f + \alpha T_0)(z_4 - z_3) < 0}$$

$$\rightarrow \boxed{|Q_1| = (E_0 - E_f + \alpha T_0) \Delta z_d} \quad \boxed{\Delta z_d = z_3 - z_4} \text{ : Descarga}$$

b) - La eficiencia $\eta = 1 - \frac{|Q_1|}{Q_2} = 1 - \frac{(E_0 - E_f + \alpha T_0) \Delta z_d}{(E_0 - E_c + \alpha T_0) \Delta z_c}$

Pero, de la E d E

$$\alpha T = E_0 - E + \alpha T_0 \Rightarrow \alpha T_f = E_0 - E_f + \alpha T_0$$

$$\alpha T_c = E_0 - E_c + \alpha T_0$$

Luego, $\eta = 1 - \frac{\alpha T_f}{\alpha T_c} \cdot \frac{\Delta z_d}{\Delta z_c} = 1 - \frac{T_f}{T_c} \cdot \frac{\Delta z_d}{\Delta z_c}$

Ya que es un ciclo de Carnot, $\eta = \eta_c = 1 - T_f/T_c$, y así:

$$\boxed{\Delta z_d = \Delta z_c}$$

Problema 4 $C_v = 3R \cdot \frac{4\pi^4}{5} \left(\frac{T}{\Theta}\right)^3$

$TdS = m C_v dT + \underbrace{PdV}_{=0} \text{ (proceso termodinámico)}$

$\therefore dS = 3mR \cdot \frac{4\pi^4}{5} \frac{T^2}{\Theta^3} dT$

$\int_{S_1}^{S_2} dS = \frac{3mR \cdot 4\pi^4}{5 \Theta^3} \int_{T_1}^{T_2} 3T^2 dT$

$\frac{\Delta S}{R} = \frac{4\pi^4 \cdot m}{5} \left(\frac{T_2^3 - T_1^3}{\Theta^3} \right)$

Datos: $T_1 = 10 \text{ K}$, $T_2 = 350 \text{ K}$; $\Theta = 2230 \text{ K}$; $m = 1.2 \text{ g}$.

Propiedades: $12 \left(\frac{\text{g}}{\text{mol}} \right) \Rightarrow \left. \begin{array}{l} 12 \text{ g} = 1 \text{ mol} \\ 1.2 \text{ g} = m \end{array} \right\} \Rightarrow m = 0.1 \text{ mol}.$

$\therefore \left\{ \frac{\Delta S}{R} \approx 0.03 \text{ mol} \right\}$