

1) $N=1$

→ $\left(\frac{\partial U}{\partial V}\right)_T = 0$; $dU = T ds - p dv = C_v dT + \left[T \left(\frac{\partial p}{\partial T}\right)_V - p\right] dv$
 Enumerando

Significa que $T \left(\frac{\partial p}{\partial T}\right)_V = p$ (i) y $\left(\frac{\partial U}{\partial T}\right)_V = C_v$

U Independiente de V

Otra: $\left(\frac{\partial H}{\partial p}\right)_T = 0$; $dH = T ds + v dp = C_p dT + \left[v - T \left(\frac{\partial v}{\partial T}\right)_p\right] dp$

→ $v = T \left(\frac{\partial v}{\partial T}\right)_p$ (ii)

de (i) y (ii) de integrar:

(i) $p = T a(v)$

(ii) $v = b(p)T$

Si dividimos:

$p b(p) = v a(v) =$

Por que se cumple la igualdad:

$p b(p) = v a(v) = Cte = R$

Pq hay variable diferente a cada lado.

Reemplazando en (i) o (ii) queda:

$p v = R T$, Independiente de $C_v(T)$

Entropía

Cualquier fórmula. TdS

$$TdS = C_V(T) dT + T \left(\frac{\partial P}{\partial T} \right)_V dV = C_V dT + \frac{RT}{V} dV$$

En principio no podemos suponer $C_V = \text{cte}$, pero
estamos demostrando que es un gas ideal, necesitamos
ser lo más general posible.

Entonces

$$S = S_0 + \int_{T_0}^T \frac{C_V dT}{T} + R \ln \left(\frac{V}{V_0} \right)$$

$$U = \int_{T_0}^T C_V dT$$

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Homogénea en N

$$S(T, \lambda V, \lambda N) = \lambda S(T, V, N)$$

Si usamos $\lambda = \frac{1}{N}$

$$\Rightarrow NS(T, \frac{V}{N}, 1)$$

Entonces por la ecuación de estado

$$\boxed{\frac{PV}{N} = RT}$$

Gas ideal $\Rightarrow C_V = \text{cte}$ experimentalmente.

$$S = S_0 N + N C_V \ln \left(\frac{T}{T_0} \right) + N R \ln \left(\frac{V}{V_0 N} \right)$$

$$U = U_0 N + N C_V T$$

dependiendo T de S

$$U = N c_v T_0 e^{\left[\frac{S}{N c_v} - \frac{S_0}{c_v} - \frac{R}{c_v} \ln \left(\frac{V}{N V_0} \right) \right]}$$

$$\rightarrow S(U, V, N) = N S_0 + N c_v \ln \left[\frac{U}{N c_v T_0} \right] + R N \ln \left(\frac{V}{N V_0} \right)$$

con $C_v = mR$ y $U = N_0 c_v T_0$ y $v_0 = \frac{V_0}{N_0}$

$$S = N S_0 + N R \ln \left[\frac{V}{V_0} \left(\frac{U}{U_0} \right)^m \left(\frac{N_0}{N} \right)^{m+1} \right]$$

2) $du = p_1 dv_1 + p_2 dv_2$

Proceso Isotermico

$$dW = - (dA_1 + dA_2)$$

$$A = U - TS = N c_v T + U_0 \cdot N + N T \left[S_0 + c_v \ln \left(\frac{T}{T_0} \right) + R \ln \left(\frac{V}{N V_0} \right) \right]$$

a N y T cte:

$$\Delta A = - N T R \ln \left(\frac{V_f}{V_i} \right)$$

\therefore

$$W = N T R \ln \left(\frac{V_{f1} V_{f2}}{V_{i1} V_{i2}} \right)$$

3) $F(T, L) = -K(T)[L - L_0(T)]$
 vemos la expansión térmica.

$$dL = \left(\frac{\partial L}{\partial T} \right)_F dT + \left(\frac{\partial L}{\partial F} \right)_T dF$$

Expansión
Térmica.

$$dL = \left(\frac{1}{K} F + L_0' \right) dT + \left(\frac{\partial L}{\partial F} \right)_T dF$$

Condición para que sea 0

$F \approx 0$ y L_0 no depende de T

$\Rightarrow L \approx L_0$

es decir está orientado verticalmente.

$$F(T, L) = -K(T)[L - L_0]$$

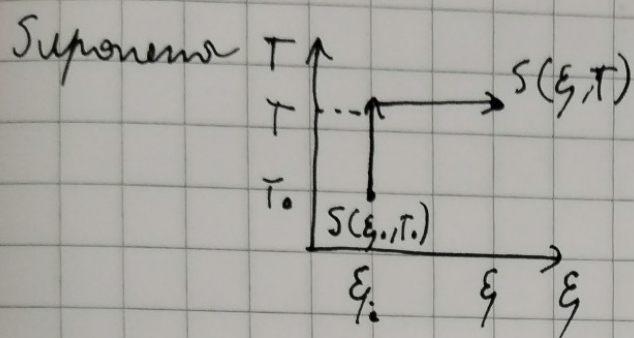
$$dU = T dS - F d\epsilon$$

$$T dS = C_\epsilon dT + T \left(\frac{\partial F}{\partial T} \right)_\epsilon d\epsilon = C_\epsilon dT - T K' \epsilon d\epsilon$$

$$\text{Sabemos } \left(\frac{\partial C_\epsilon}{\partial \epsilon} \right)_T = -T \epsilon K''$$

$$C_\epsilon = -\frac{1}{2} \epsilon^2 K''(T) + g(T)$$

para calcular S ~~temos~~ podemos inventar um processo que
 Não convendo d, p, q, N e podemos integrar C_p , p que
 depende de ξ .



$$\Delta S = \underbrace{\int_{T_0}^T C_{\xi} \frac{dT}{T}}_{\text{se anula } \xi = \text{cte}} - \underbrace{\int_{\xi_0}^{\xi} T \xi' K' d\xi'}_{\text{se anula Isotermico}} + \int_T^T C_{\xi} \frac{dT}{T} + \int_{\xi_0}^{\xi} T \xi' K' d\xi'$$

$$S = S_0 + \int_{T_0}^T C_{\xi} \frac{dT}{T} + \int_{\xi_0}^{\xi} T \xi' K' d\xi' ; \text{ podemos usar } \xi_0 = 0$$

$$S = S_0 + \int_{T_0}^T C_{\xi} \frac{dT}{T} - \frac{1}{2} \xi^2 K'(T)$$

$$U = U_0 + \int_{T_0}^T C_{\xi} \frac{dT}{T} + \frac{1}{2} [K(T) - T K'(T)] \xi^2$$

$$\text{definimos } U(T, 0) = U_0 + \int_{T_0}^T C_{\xi} \frac{dT}{T} \quad \text{e} \quad S(T, 0) = S_0 + \int_{T_0}^T C_{\xi} \frac{dT}{T}$$

$$A = U - TS = A(T, 0) + \frac{1}{2} K(T) \xi^2$$