

# Tensor Asymmetry $A_{zz}$ in the $x > 1$ Region

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A Letter of Intent to Jefferson Lab PAC 42

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## Abstract

We propose the first measurement of the tensor asymmetry  $A_{zz}$  in the quasi-elastic region through the tensor polarized  $D(e, e')X$  channel; an asymmetry that is sensitive to the nucleon-nucleon potential. Previous measurements of  $A_{zz}$  have been used to extract  $b_1$  in the DIS region and  $T_{20}$  in the elastic region. In the quasi-elastic region,  $A_{zz}$  data will be used to compare light cone calculations with variation nucleon-nucleon calculations, and is an important quantity to determine for understanding tensor effects, such as the dominance of  $pn$  correlations in nuclei.

In the quasi-elastic region,  $A_{zz}$  was first calculated in 1988 by Frankfurt and Strikman, using the Hamada-Johnstone and Reid soft-core wave functions [1]. Recent calculations by M. Sargsian revisit  $A_{zz}$  in the  $x > 1$  range using virtual-nucleon and light-cone methods, which differ by up to a factor of two [2] and can be discriminated experimentally at the  $3 - 6\sigma$  level. This potential measurement has stirred the interest of a number of theorists, and will be proposed in full as calculations solidify.

We propose an experimental determination of  $A_{zz}$  utilizing the same equipment as the E13-12-011  $b_1$  experiment. Three different  $Q^2$  values will be measured over the course of 30 days, with 8.6 additional days of overhead. The measurements are less sensitive to systematic uncertainties than E13-12-011, so this experiment would be utilized in parallel to better understand the in-beam conditions and time-dependent systematic effects of a tensor polarized target for the  $b_1$  experiment.

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# 1 Background

The deuteron is the simplest composite nuclear system, and in many ways it is as important to understanding bound states in QCD as the hydrogen atom was to understanding bound systems in QED. Our experimental and theoretical understanding of the deuteron remains unsatisfying.

Due to their small size and simple structure, tensor polarized deuterons are ideal for studying nucleon-nucleon interactions. Tensor polarization enhances the D-state contribution, which compresses the deuteron [3], making the system more sensitive to short-range QCD effects. Understanding the nucleon-nucleon potential of the deuteron is essential for understanding short-range correlations as they are largely dependent on the tensor force [4]. We can resolve the short-range structure of nuclei on the level of nucleon and hadronic constituents by utilizing processes that transfer to the nucleon constituents both energy and momentum larger than the scale of the NN short-range correlations, particularly at  $Q^2 > 1 \text{ (GeV/c)}^2$ .

By taking a ratio of cross sections from electron scattering from tensor-polarized and unpolarized deuterons, the S and D-wave states can be disentangled, leading to a fuller understanding of the repulsive nucleon core. A measurement of  $A_{zz}$  is sensitive to the  $\frac{D^2-SD}{S^2+D^2}$  ratio and its evolution with increasing minimal momentum of the struck nucleon. Originally calculated by L. Frankfurt and M. Strikman [1], this has recently been revisited by M. Sargsian, who calculated  $A_{zz}$  in this region using a light cone approach and a virtual nucleon approach. The calculations vary by up to a factor of 2, and can be experimentally determined at the  $3 - 6\sigma$  level as discussed in this letter.

For the lower  $Q^2$  region, W. Van Orden has calculations in progress using different nucleon-nucleon potentials, as well as different prescriptions for handling the reactions mechanisms [5]. Similar calculations are currently being finalized for the approved  $D(e, e'p)n$  at high  $Q^2$ , high  $p_m$  experiment. Once completed, he will turn his attention to tensor polarization observables in the low  $Q^2$  region.

Additionally, measuring  $A_{zz}$  in the quasi-elastic region will fill a gap in measurements performed on tensor polarized deuterium. It is directly proportional to the observable used in the elastic region to measure  $T_{20}$ , by  $A_{zz} \propto T_{20}$ . Due to the large acceptance of the SHMS spectrometer, we will be taking data in the  $x = 2$  elastic range as well that, at low  $Q^2$ , may be able to resolve the discrepancy of the deuteron structure functions that were measured at JLab in Halls A and C [6]. This will be investigated further as a potential parasitic measurement when the full proposal is submitted. In the deep inelastic region,  $A_{zz}$  will soon be measured to extract the tensor structure function  $b_1$  by the relation  $A_{zz} \propto \frac{b_1}{F_1^D}$ . Not only will measuring  $A_{zz}$  in the quasi-elastic region provide information necessary for understanding the properties of the deuteron and contribution from the tensor force, but it will be the first experiment to bridge a gap in measurements of electron scattering from tensor-polarized deuterons.

## 1.1 Probing the Deuteron Wavefunction

It was suggested for some time [7] that to resolve the microscopic structure of nuclei one needs to study scattering at sufficiently large momentum transfer and large relative momenta of the produced nucleons. This logic was confirmed [4] by a series of experiments at SLAC [8] and



Figure 1: The AV18 [11] deuteron wave-function, showing the dominance of the D-state (dashed) in comparison to the S-state (dotted) in the full wavefunction (solid) at high momentum ( $k > 300$  MeV/c).

JLab [9, 10] that directly observed short-range correlations (SRC) in a series of nuclei, and established a similar effect of SRC in the deuteron and in heavier nuclei with  $pn$  correlations giving the dominant contribution. Hence, the deuteron serves as a “hydrogen atom” for the studies of the microscopic short-range structure of the nuclei since it is the simplest nuclei that follows SRC scaling.

To achieve further progress, it is necessary to improve our knowledge of the deuteron wave function at high momenta, and to separate the S and D contributions to the high momentum component of the deuteron. The dominance of the D-wave at a large range of the nucleon momenta is expected in a range of the theoretical models, as demonstrated in Fig. 1, but experimentally it was probed in a rather indirect way via measurement of  $T_{20}$  for the deuteron form factor [6]. Still, the knowledge of S/D ratio for large momenta is rather poor. Indeed, all wavefunctions are constrained by low energy data to reproduce the S/D ratio at small momenta while the overall probability of the D-wave in the deuteron differs by a factor up to 1.5, leading to a large difference of the S/D ratio at large momenta.

The S and D-states are related to the tensor asymmetry  $A_{zz}$  by [1]

$$A_{zz} \propto \frac{\frac{1}{2}w^2(k) - u(k)w(k)\sqrt{2}}{u^2(k) + w^2(k)}, \quad (1)$$

where  $u(k)$  is the S-state wave function and  $w(k)$  is the D-state wave function. Additionally, measuring  $A_{zz}$  at lower  $Q^2$  will map out the transition from hadronic to partonic degrees of freedom.

Ratios of inclusive cross sections at  $x > 1$  has demonstrated an early onset of the scaling of the ratios when plotted as a function of the light-cone fraction of the struck nucleon momentum. As a result, the ratios provide a direct measurement of the ratio of the high momentum components in nuclei. Similarly, one can expect that in the case of scattering from the polarized deuteron we expect the early scaling for the asymmetry when plotted as a function of the minimal struck

nucleon momentum or the light cone fraction in the  $A(e, e')$  case. It was observed at JLab that the scaling of the ratios set in starting at  $Q^2 \sim 1 \text{ GeV}^2$  [9] so covering the range of  $Q^2$  up to  $2 \text{ GeV}^2$  will be sufficient to measure the S/D ratios in an interesting momentum range.

It is worth noting here that in addition to comparing predictions for the different wave functions, one expects to be able to distinguish between non-relativistic and light cone quantum mechanic models. The principal difference between the models is the relation between the spectator momentum and momentum in the wave function. In the nonrelativistic model they coincide, while in the light cone model the relation is non-linear starting at  $k \sim 250 \text{ MeV}/c$ . This difference is most clearly manifested in the scattering from the polarized deuteron due to a strong dependence of the S/D ratio on the nucleon momentum.

## 1.2 Study of the Relativistic NN Bound System

One of the important issues in studying of nuclear structure at short distances is the relativistic description of the bound system. This is an important issue also in understanding the QCD medium effect with recent studies indicating that parton distribution modifications in nuclei are proportional to the high momentum component of nuclear wave function.

The deuteron is the simplest bound system and naturally any self-consistent attempt to understand the relativistic effects in the bound nuclear systems should start with the deuteron. The issue of the relativistic description of the deuteron has long history with extensive research that started in late 1970's [12, 13, 14, 7].

The experimental studies of the relativistic effects in the deuteron up to now include the large  $Q^2$  elastic  $ed$  scattering [15], however due to complexities in the reaction mechanism [16] the relativistic effects were difficult to isolate.

The inclusive  $D(e, e')X$  experiments from tensor-polarized deuterons at  $Q^2 > 1 \text{ GeV}^2$  and  $x > 1$  region gives a new possibility to probe the relativistic structure of the deuteron. In this case the use of the tensor polarized deuteron allows us to prepare the nucleus in the most compact state in which, due to the absence of the pure S-wave<sup>2</sup> contribution, the system in average is sensitive to the higher moment of the nucleon in the deuteron. At large  $Q^2 > 1 \text{ GeV}^2$  kinematics, the probed longitudinal momenta of the bound nucleon  $p_z \approx m_N(1 - x)$ , or the light cone momentum fraction  $\alpha \geq x$ . Because of these kinematic conditions and the absence of the large S-wave<sup>2</sup> contribution, one expects a measurable relativistic effects already at  $x \leq 1.2$ .

The biggest advantage is that one expects less uncertainty due to the choice of the NN potential and reaction dynamics due to relatively small values of the bound nucleon momenta involved ( $\geq 200 \text{ MeV}/c$ ).

The sensitivity to relativistic effects is estimated using the theoretical calculations based on two very different approaches. The first approach treats the virtuality of the bound nucleon within a description of the deuteron in the lab. frame with treating the interacting nucleon as being virtual (virtual nucleon, or VN, approximation) by taking the residue over the energy of the spectator nucleon. In this case, the deuteron wave function satisfies the covariant equation of two-nucleon bound system with spectator being on energy shell [17, 18].

Another approach is based on the observation that high energy processes evolve along the light-cone (LC). Therefore, it is natural to describe the reaction within the light-cone non-covariant

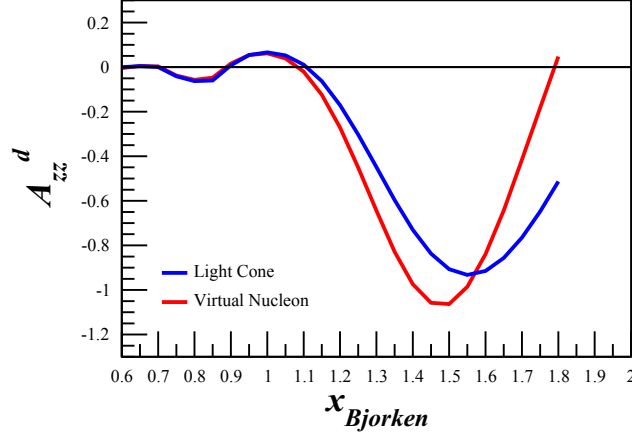


Figure 2: The  $A_{zz}$  observable calculated at  $Q^2 = 1.5 \text{ (GeV/c)}^2$  using the light-cone and virtual nucleon models. Calculations provided by M. Sargsian.

framework [7]. Negative energy states do not enter in this case, though one has to take into account so called instantaneous interactions. In the approximation when non-nucleonic degrees of freedom in the deuteron wave function can be neglected, one can unambiguously relate the light-cone wave functions to those calculated in the lab. frame by introducing the LC  $pn$  relative three momentum,

$$k = \sqrt{\frac{m^2 + p_t^2}{\alpha(2 - \alpha)}} - m^2. \quad (2)$$

In Fig. 2, the prediction for VN [17] and LC [8] approximations are given for the highest  $Q^2$  kinematics proposed. As was previously mentioned, a measurable difference is predicted to be observable already at  $x \geq 1.2$ . to the choice of the wave function.

### 1.3 Interest from Theorists

The measurement proposed in this letter has stirred interest in a number of theorists who are working on calculations. Many of these are on-going and are expected to be completed in the coming months to further the physics motivation of a measurement of  $A_{zz}$  in the quasi-elastic region at various momentum-transfer.

The light cone and virtual nucleon calculations of M. Sargsian [19] and M. Strikman [20] are already available for  $A_{zz}$  and are presented in this document. Early calculations of the M. Strikman light-cone model indicate a potential measurable discrepancy based on the input of different NN potentials, but the results are very preliminary and require further investigation. S. Liuti has agreed to join in this theoretical effort, stating “This is an important measurement to know, and should be calculated more thoroughly.” [21]

Models involving 6-quark calculations of quasi-elastic  $A_{zz}$  can be calculated by G. Miller [22] and have been motivated by the collaboration. In his own words, he states “This measurement was

a highlighted need early at JLab. A new measurement at higher  $Q^2$  would be very interesting. In principle such could test my model. I could calculate the influence of my 6-quark configurations on elastic scattering.”

Continuing his interest from DIS  $b_1$  calculations, W. Cosyn is developing calculations of the quasi-elastic contribution to inclusive deuteron scattering, which will be the dominant contribution in the  $x > 1$  regime [23]. His calculations, which include final-state interactions, can be modified to include  $A_{zz}$  and expects to have them within the coming months. In his words, “I hope to do some calculations soon and could easily do them for the kinematics in your proposal.”

In addition, W. Van Orden has calculations in progress using different nucleon-nucleon potentials, as well as different prescriptions for handling the reactions mechanisms [5]. Similar calculations are currently being finalized for the approved  $D(e, e'p)n$  at high  $Q^2$ , high  $p_m$  experiment. Once completed, he will turn his attention to tensor polarization observables in the low  $Q^2$  region and investigate the effects from differing NN potentials.

In summary, we are encouraged that several theorists have begun serious efforts to calculate  $A_{zz}$  in the  $x > 1$  region using a variety of models. We anticipate that the PAC’s guidance will further encourage and focus these efforts.

## 2 The Proposed Experiment

We propose to measure the tensor asymmetry  $A_{zz}$  from inclusive electron scattering from polarized deuterons in the quasi-elastic region of  $0.80 < x < 1.75$ ,  $1.0 \text{ (GeV/c)}^2 < Q^2 < 1.9 \text{ (GeV/c)}^2$ , and  $0.59 < W < 1.09 \text{ GeV}$  using the Hall C HMS and SHMS spectrometers at forward angle using a solid polarized  $\text{ND}_3$  target.

### 2.1 Experimental Method

The measured double differential cross section for a spin-1 target is characterized by a vector polarization  $P_z$  and tensor polarization  $P_{zz}$  is expressed as,

$$\frac{d^2\sigma_p}{d\Omega dE'} = \frac{d^2\sigma_u}{d\Omega dE'} \left( 1 - P_z P_B A_1 + \frac{1}{2} P_{zz} A_{zz} \right), \quad (3)$$

where,  $\sigma_p$  ( $\sigma_u$ ) is the polarized (unpolarized) cross section,  $P_B$  is the incident electron beam polarization, and  $A_1$  ( $A_{zz}$ ) is the vector (tensor) asymmetry of the virtual-photon deuteron cross section. This allows us to write the polarized tensor asymmetry with positive tensor polarization using an unpolarized electron beam as

$$A_{zz} = \frac{2}{P_{zz}} \left( \frac{\sigma_p - \sigma_u}{\sigma_u} \right). \quad (4)$$

The tensor polarization is given by

$$P_{zz} = \frac{n_+ - 2n_0 + n_-}{n_+ + n_- + n_0}, \quad (5)$$



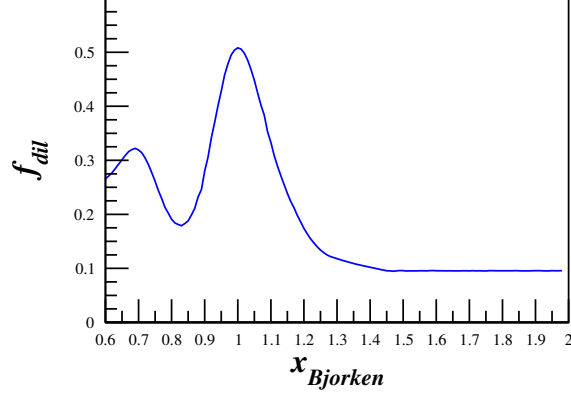


Figure 3: The estimated dilution factor, in this case at  $Q^2 = 1.5 \text{ (GeV/c)}^2$ , is expected to drop off at high  $x$  until it reaches the SRC plateau region. This effect will be counteracted by using a high-luminosity solid target.

where  $n_m$  represents the population in the  $m_z = +1, -1$ , or 0 state.

Eq. 4 reveals that the asymmetry  $A_{zz}$  compares two different cross sections measured under different polarization conditions of the target: positively tensor polarized and unpolarized. To obtain the relative cross section measurement in the same configuration, the same target cup and material will be used at alternating polarization states (polarized vs. unpolarized), and the magnetic field providing the quantization axis will be oriented along the beamline at all times. This field will always be held at the same value, regardless of the target material polarization state. This process, identical to that used for the E12-13-011  $b_1$  measurement, ensures that the acceptance remains consistent within the stability ( $10^{-4}$ ) of the super conducting magnet.

Since many of the factors involved in the cross sections cancel in the ratio, Eq. 4 can be expressed in terms of the charge normalized, efficiency corrected numbers of tensor polarized ( $N_p$ ) and unpolarized ( $N_u$ ) counts,

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_p - N_u}{N_u} \right). \quad (6)$$

The dilution factor  $f$  corrects for the presence of unpolarized nuclei in the target and is defined by

$$f = \frac{N_D \sigma_D}{N_N \sigma_N + N_D \sigma_D + \sum_A N_A \sigma_A}, \quad (7)$$

where  $N_D$  is the number of deuterium nuclei in the target and  $\sigma_D$  is the corresponding inclusive double differential scattering cross section,  $N_N$  is the nitrogen number of scattered nuclei with cross section  $\sigma_N$ , and  $N_A$  is the number of other scattering nuclei of mass number  $A$  with cross section  $\sigma_A$ . As has been noted in previous work [1], the dilution factor at high  $x$  drops off considerably until the SRC plateau region, as shown in Fig. 3. By using a high-luminosity solid target and a low scattering angle  $\theta_e$ , this effect will be counteracted.

	$E_0$ (GeV)	$Q^2$ (GeV <sup>2</sup> )	$E'$ (GeV)	$\theta_{e'}$ (deg.)	Rates (kHz)	PAC Time (hours)
SHMS	8.8	1.5	8.36	8.2	0.43	600
SHMS	6.6	0.7	6.50	8.2	3.19	90
SHMS	2.2	0.3	2.11	14.4	3.73	30
HMS	2.2	0.3	2.11	14.9	2.92	30

Table 1: Summary of the central kinematics and physics rates using the Hall C spectrometers.

The dilution factor can be written in terms of the relative volume ratio of ND<sub>3</sub> to LHe in the target cell, otherwise known as the packing fraction  $p_f$ . In our case of a cylindrical target cell oriented along the magnetic field, the packing fraction is exactly equivalent to the percentage of the cell length filled with ND<sub>3</sub>.

If the time is evenly split between scattering off of polarized and unpolarized ND<sub>3</sub>, the time necessary to achieve the desired precision  $\delta A$  is:

$$T = \frac{N_p}{R_p} + \frac{N_u}{R_u} = \frac{8}{f^2 P_{zz}^2} \left( \frac{R_p(R_u + R_p)}{R_u^3} \right) \frac{1}{\delta A_{zz}^2} \quad (8)$$

where  $R_{p(u)}$  is the polarized (unpolarized) rate and  $N_{p(u)}$  is the total estimated number of polarized (unpolarized) counts to achieve the uncertainty  $\delta A_{zz}$ .

## 2.2 Kinematics

We propose to measure the tensor asymmetry  $A_{zz}$  for  $0.80 < x < 1.75$ ,  $1.0 \text{ (GeV/c)}^2 < Q^2 < 1.9 \text{ (GeV/c)}^2$ , and  $0.59 < W < 1.09 \text{ GeV}$  after applying kinematics cuts. Fig. 4 shows the planned kinematic coverage utilizing the Hall C HMS and SHMS spectrometers at forward angle.

The polarized ND<sub>3</sub> target is discussed in section 2.4. The magnetic field of the target will be held constant along the beamline at all times, while the target state is alternated between a polarized and unpolarized state. The tensor polarization and packing fraction used in the rates estimate are 30% and 0.65, respectively. The dilution fraction in the range of this measurement is shown in Fig. 5. With an incident electron beam current of 90 nA, the expected deuteron luminosity is  $1.3 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ .

The momentum bite and the acceptance were assumed to be  $\Delta P = \pm 8\%$  and  $\Delta\Omega = 5.6 \text{ msr}$  for the HMS, and  $\Delta P = {}^{+20\%}_{-8\%}$  and  $\Delta\Omega = 4.4 \text{ msr}$  for the SHMS. For the choice of the kinematics, special attention was taken onto the angular and momentum limits of the spectrometers: for the HMS,  $10.5^\circ \leq \theta \leq 85^\circ$  and  $1 \leq P_0 \leq 7.3 \text{ GeV/c}$ , and for the SHMS,  $5.5^\circ \leq \theta \leq 40^\circ$  and  $2 \leq P_0 \leq 11 \text{ GeV/c}$ . In addition, the opening angle between the spectrometers is physically constrained to be larger than  $17.5^\circ$ . The dilution factors and projected uncertainties in  $A_{zz}$  are summarized in Table 2 and displayed in Fig. 6.

A total of 30 days of beam time is requested for production data, with an additional 8.6 days of expected overhead.

$x$	$Q^2 = 1.5 \text{ (GeV/c)}^2$			$Q^2 = 0.7 \text{ (GeV/c)}^2$			$Q^2 = 0.3 \text{ (GeV/c)}^2$		
	$f_{dil}$	$\delta A_{zz}^{stat}$ $\times 10^{-2}$	$\delta A_{zz}^{sys}$ $\times 10^{-2}$	$f_{dil}$	$\delta A_{zz}^{stat}$ $\times 10^{-2}$	$\delta A_{zz}^{sys}$ $\times 10^{-2}$	$f_{dil}$	$\delta A_{zz}^{stat}$ $\times 10^{-2}$	$\delta A_{zz}^{sys}$ $\times 10^{-2}$
0.80	0.205	0.58	2.03	0.175	0.71	0.74	0.298	0.46	0.74
0.90	0.274	0.44	0.58	0.375	0.31	1.68	0.462	0.29	1.68
1.00	0.507	0.23	1.01	0.518	0.22	0.02	0.521	0.27	0.02
1.10	0.333	0.47	0.21	0.409	0.35	1.99	0.431	0.40	1.63
1.20	0.174	1.28	2.36	0.264	0.74	3.98	0.301	0.74	3.25
1.30	0.120	2.63	6.28	0.174	1.42	5.98	0.193	1.34	4.88
1.40	0.108	4.08	10.2	0.156	2.46	7.97	0.144	2.16	6.50
1.50	0.096	6.56	12.7	0.133	3.31	9.96	0.100	3.61	8.13
1.60	0.096	8.96	12.8	0.110	5.49	12.0	0.086	4.62	9.75
1.70	0.095	12.1	10.7	0.096	7.88	13.9	0.063	7.13	11.4
1.80	0.096	15.5	7.18	0.096	10.3	14.0	0.056	8.87	13.0

Table 2: Summary of the expected statistical uncertainty after combining overlapping x-bins. Values represent the statistics weighted average of all events that satisfy our kinematic cuts.

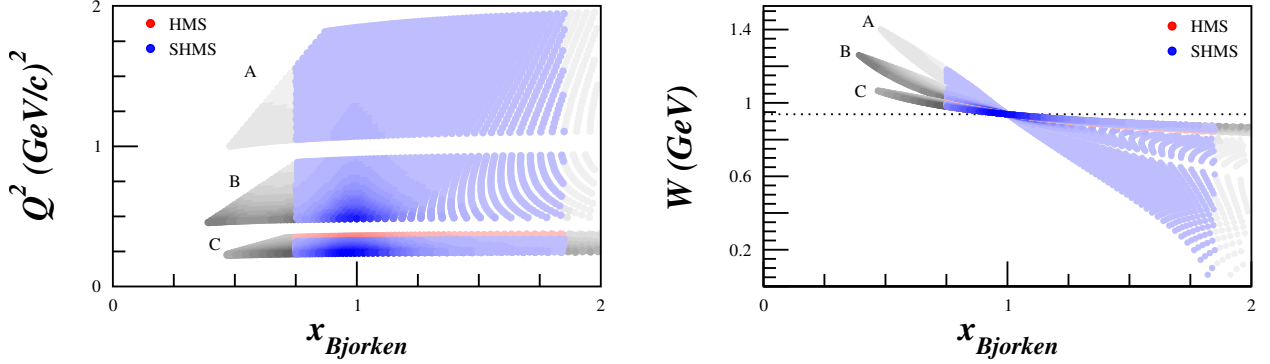


Figure 4: Kinematic coverage for central spectrometer settings at  $Q^2 = 1.5 \text{ (GeV/c)}^2$  (A),  $0.7 \text{ (GeV/c)}^2$  (B), and  $0.3 \text{ (GeV/c)}^2$  (C). The HMS is only used for setting C, and its coverage largely falls under the SHMS coverage. The grey regions are not included in our statistics estimates since they fall outside of  $0.80 < x < 1.75$ . Darker shading represents areas with higher statistics, and the dotted line in the  $W$  plot indicates nucleon mass.

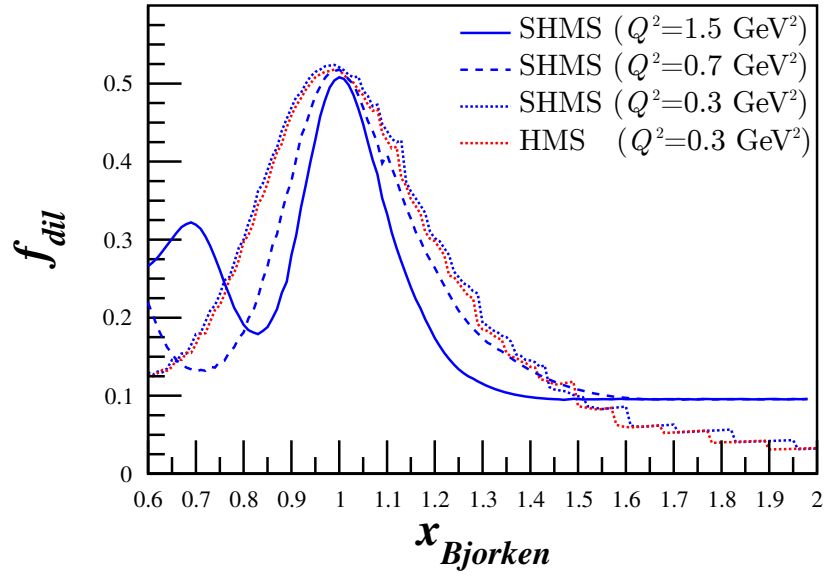


Figure 5: Projected dilution factor covering the entire  $x$  range to be measured using a combination of P. Bosted's [24] and M. Sargsian's [19] code for the SHMS and HMS.

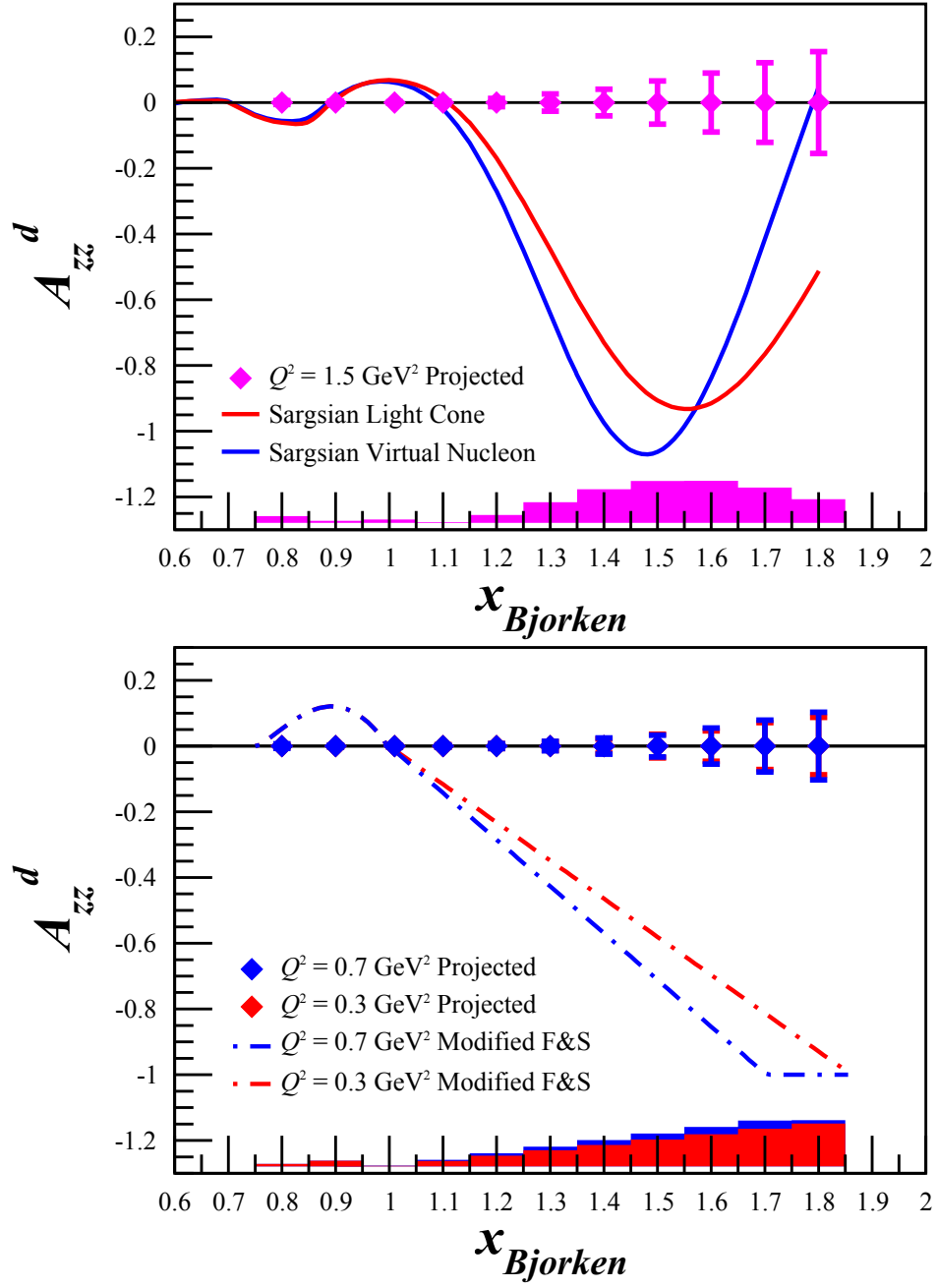


Figure 6: Projected statistical errors for the tensor asymmetry  $A_{zz}$  with 30 days of beam time. The band represents the systematic uncertainty. Also shown for  $Q^2 = 1.5 \text{ (GeV/c)}^2$  are calculations provided by M. Sargsian for using a light cone and virtual nucleon model, and for  $Q^2 = 0.3$  and  $0.7 \text{ (GeV/c)}^2$  a modified Frankfurt and Strikman model [1] that estimates the peak shifts in  $x$  expected due to the SRC scaling changing with  $Q^2$  [25].

## 2.3 Uncertainty Estimates

We discuss here the expected experimental and systematic uncertainties that we expect to contribute to the measurement.

### 2.3.1 Statistical Uncertainty

To investigate the statistical uncertainty we start with the equation for  $A_{zz}$  using measured counts for polarized data ( $N_p$ ) and unpolarized data ( $N_u$ ),

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_p}{N_u} - 1 \right). \quad (9)$$

The statistical error with respect to counts is then

$$\delta A_{zz} = \frac{2}{fP_{zz}} \sqrt{\left( \frac{\delta N_p}{N_u} \right)^2 + \left( \frac{N_p \delta N_u}{N_u^2} \right)^2}. \quad (10)$$

For  $\delta N_{p(u)} = \sqrt{N_{p(u)}}$ , the uncertainty becomes

$$\delta A_{zz} = \frac{2}{fP_{zz}} \sqrt{\frac{N_p(N_u + N_p)}{N_u^3}}, \quad (11)$$

which can't be simplified further due to the large expected asymmetry.

The number of counts was calculated using a combination of P. Bosted's [24] and M. Sargsian's [19] code. The Bosted code was used for the lowest  $Q^2$  setting, where effects of SRC scaling are expected to be negligible, and for  $x < 1.1$  to accurately determine the quasi-elastic peak. The Sargsian code was used for the higher  $Q^2$  settings at  $x > 1.1$  due to its inclusion of SRC scaling effects.

### 2.3.2 Systematic Uncertainty

Table 3 shows a list of the scale dependent uncertainties contributing to the systematic error in  $A_{zz}$ . With careful uncertainty minimization in polarization the relative error in vector polarization,  $P_z$ , can be less than or equal to 3.9%, as was demonstrated for the proton in the recent E08-027/E08-007 experiment [26] and nearly as good for the deuteron using multiple techniques to measure the NMR signal as discussed in [27]. With the use of a positive tensor enhanced target it has been projected to be able to achieve a relative error in  $P_{zz}$  better than 12% [27]. The uncertainty from the dilution in the polarized target is estimated to be about 6% over the range of kinematics points of interest. We consider separately the uncertainty in the packing fraction of the ammonia target contributes at a level of less than 3%. Charge calibration and detector efficiencies are expected to be known better to 1%.

Source	Systematic
$P_{zz}$ Polarimetry	12%
Dilution Factor	6.0%
Packing Fraction	3.0%
Trigger/Tracking Efficiency	1.0%
Acceptance	0.5%
Charge Determination	1.0%
Detector Resolution and Efficiency	1.0%
Total	14%

Table 3: Estimates of the scale dependent contributions to the systematic error of  $A_{zz}$ .

### Time Dependent Systematic Effects

Eq. 6 involves the ratio of counts, which leads to cancellation of several first order systematic effects. However, the fact that the two data sets will not be taken simultaneously leads to a sensitivity to time dependent variations which will be monitored and suppressed. The typical size of these types of effects are small compared to the large asymmetries predicted for most of the proposed kinematic regions, but must be carefully monitored whenever the expected asymmetry is small, such as at a zero crossing.

To investigate the systematic differences in the time dependent components of the integrated counts, we need to consider the effects from calibration, efficiency, acceptance, and luminosity between the two polarization states.

Fluctuations in luminosity due to target density variation can easily be kept to a minimum by keeping the material beads at the same temperature for both polarization states by control of the microwave and the LHe evaporation. The He vapor pressure reading gives accuracy of material temperature changes at the level of  $\pm 0.1\%$ .

The beam charge asymmetries between two helicity states using the luminosity monitors for experiment E06-010 has been shown to be at the level of  $4 \times 10^{-5}$  with a width of  $2 \times 10^{-4}$ . An additional estimate on the change in the BCM calibration constant is seen in experiment E08-027 resulting in a absolute deviation of  $2 \times 10^{-4}$  over the course of six days. We expect to be able to minimize long term drifts by careful thermal isolation of the BCMs.

The acceptance of each cup can only change as a function of time if the magnetic field changes. The capacity to set, reset, and hold the target superconducting magnet to a desired holding field causes a field uncertainty of  $\delta B/B = 0.01\%$ . This implies that, like the cup length  $l$ , the acceptance  $\mathcal{A}$  for each polarization state is the same.

In order to look at the effect on  $A_{zz}$  due to drifts in beam current monitor calibration and detector efficiency, we rewrite Eq. 6 explicitly in terms of the raw measured counts  $N_p^c$  and  $N_u^c$ ,

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_p^c}{N_u^c} - 1 \right)$$

$$= \frac{2}{fP_{zz}} \left( \frac{Q\varepsilon l\mathcal{A}}{Q_1\varepsilon_1 l\mathcal{A}} \frac{N_p}{N_u} - 1 \right) \quad (12)$$

where  $Q$  represents the accumulated charge, and  $\varepsilon$  is the detector efficiency. The target length  $l$  and acceptance  $\mathcal{A}$  are identical in both states to first order.

We can then express  $Q_1$  as the change in beam current measurement calibration that occurs in the time it takes to collect data in one polarization state before switching to another, such that  $Q_1 = Q(1 - dQ)$ . In this notation,  $dQ$  is a dimensionless ratio of changes in different polarization states and would ideally be equal to zero. A similar representation is used for drifts in detector efficiency leading to,

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_p Q(1 - dQ)\varepsilon(1 - d\varepsilon)}{N_u Q\varepsilon} - 1 \right). \quad (13)$$

which simplifies to,

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_p}{N_u} (1 - dQ - d\varepsilon + dQd\varepsilon) - 1 \right). \quad (14)$$

We obtain estimates of  $dQ$  and  $d\varepsilon$  from previous experimental studies. During the JLab transversity experiment E06-010, the detector drift was measured such that the normalized yield over a three month period indicated little change ( $< 1\%$ ). These measurement were then used to show that for short time (20 minutes periods between target spin flip), the detector drift was estimated to be less than 1% times the ratio of the time period between target spin flip and three months. For the present experiment we use the same estimate except for the period between target polarization states used is  $\approx 36$  hours leading to an overall drift  $d\varepsilon \approx 0.01\%$ . A similar approach is used to establish an estimate for  $dQ$  using studies from the data from the E08-027 experiment resulting in  $d\varepsilon \approx 0.01\%$ .

To express  $A_{zz}$  in terms of the estimated experimental drifts in efficiency and current measurement we can write,

$$A_{zz} = \frac{2}{fP_{zz}} \left( \frac{N_1}{N} - 1 \right) \pm \frac{2}{fP_{zz}} d\xi. \quad (15)$$

This leads to a contribution to  $A_{zz}$  on the order of  $1 \times 10^{-3}$ ,

$$dA_{zz}^{drift} = \pm \frac{2}{fP_{zz}} d\xi = \pm 3.7 \times 10^{-3}. \quad (16)$$

Naturally detector efficiency can drift for a variety of reasons, for example including fluctuations in gas quality, HV drift or drifts in the spectrometers magnetic field. All of these types of variation as can be realized both during the experiment though monitoring as well as systematic studies of the data collected. Checks on the consistency of the cross section data that can be use ensuring the quality of each run will be used in the asymmetry analysis. Regression can be use to correct for any long term drifts that are of a non-stochastic nature. Each of these systematic effects can mitigate the systematic uncertainty to  $\sim 0.001$ . In the kinematic region proposed here,  $A_{zz}$  is expected to be large, on the order of 0.1 to 1.0, making any absolute errors on this scale only critical as the data and models pass through the x-axis.



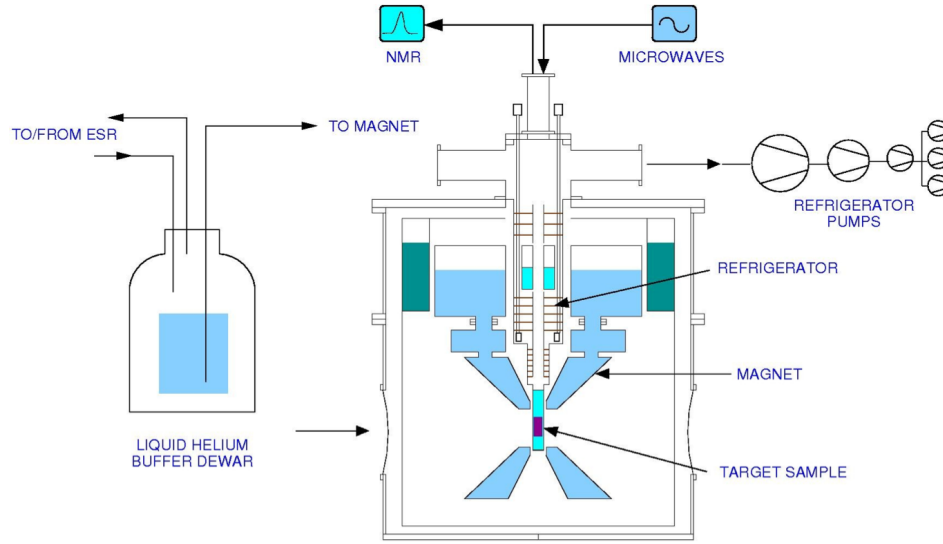


Figure 7: Cross section view of the JLab/UVa polarized target. The proposed experiment will use the modified Hall B magnet, where the backwards-scattering cone is blocked with quench protection circuitry. Figure courtesy of C. Keith.

## 2.4 Polarized Target

This experiment will use the JLab/UVa dynamically polarized solid  $\text{ND}_3$  target operated in longitudinal mode. The target is typically operated with a specialized slow raster and beamline instrumentation capable of characterizing the low current 50-100 nA beam. All of these requirements have been met previously in Hall C. The polarized target (see Fig. 7), has been successfully used in experiments E143, E155, and E155x at SLAC, and E93-026, E01-006 and E07-003, E08-027 and E08-007 at JLab. A similar target was used in Hall B for the EG1, EG4, and DVCS experiments.

The JLab/UVa target underwent significant renovation and improvement [28] during the recent g2p run. The magnet was replaced early in the run, and the target then performed consistently. A new 1 K refrigerator and target insert were designed and constructed by the JLab target group. The cryogenic pumping system has been overhauled. In particular, the older Alcatel 2060H rotary vane pumps have been replaced with new Pfeiffer DU065 magnetically coupled rotary vane pumps, and the pump controls are being refurbished. The target motion system has been rebuilt from scratch.

The target operates on the principle of Dynamic Nuclear Polarization, to enhance the low temperature (1 K), high magnetic field (5 T) polarization of solid materials by microwave pumping. The polarized target assembly contains several target cells of 3.0 cm length that can be selected individually by remote control to be located in the uniform field region of a superconducting Helmholtz pair. The permeable target cells are immersed in a vessel filled with liquid Helium and maintained at 1 K by use of a high power evaporation refrigerator. The coils have a  $50^\circ$  conical shaped aperture along the beam axis which allow for unobstructed forward scattering.

The target material is exposed to microwaves to drive the hyperfine transition which aligns the nucleon spins. The heating of the target by the beam causes a drop of a few percent in the

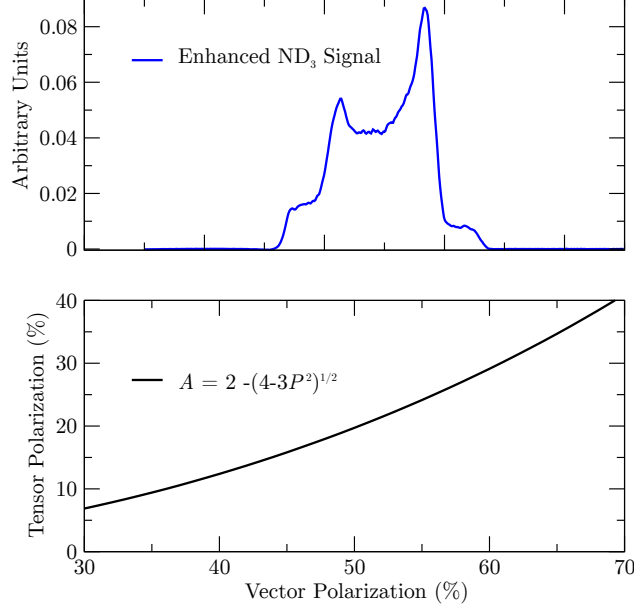


Figure 8: **Top:** NMR signal for  $\text{ND}_3$  with a vector polarization of approximately 50% from the GEN experiment. **Bottom:** Relationship between vector and tensor polarization in equilibrium, and neglecting the small quadrupole interaction.

polarization, and the polarization slowly decreases with time due to radiation damage. Most of the radiation damage can be repaired by periodically annealing the target, until the accumulated dose reached is greater than about  $0.5 \times 10^{17} \text{ e}^-/\text{cm}^2$ , at which time the target material needs to be replaced.

### 2.4.1 Polarization Analysis

The three Zeeman sublevels of the deuteron system ( $m = -1, 0, 1$ ) are shifted unevenly due to the quadrupole interaction [29]. This shift depends on the angle between the magnetic field and the electrical field gradient, and gives rise to two separate transition energies. Hence, the unique double peaked response displayed in Fig. 8. When the system is at thermal equilibrium with the solid lattice, the deuteron polarization is known from:

$$P_z = \frac{4 + \tanh \frac{\mu B}{2kT}}{3 + \tanh^2 \frac{\mu B}{2kT}} \quad (17)$$

where  $\mu$  is the magnetic moment, and  $k$  is Boltzmann's constant. The vector polarization can be determined by comparing the enhanced signal with that of the TE signal (which has known polarization). This polarimetry method is typically reliable to about 3.9% relative.

Similarly, the tensor polarization is given by:

$$P_{zz} = \frac{4 + \tanh^2 \frac{\mu B}{2kT}}{3 + \tanh^2 \frac{\mu B}{2kT}} \quad (18)$$

From Eqs. 17 and 18, we find:

$$P_{zz} = 2 - \sqrt{4 - 3P_z^2}$$

In addition to the TE method, polarizations can be determined by analyzing NMR lineshapes as described in [30] with a typical 7% relative uncertainty. At high polarizations, the intensities of the two transitions differ, and the NMR signal shows an asymmetry  $R$  in the value of the two peaks, as shown in Fig. 8. The vector polarization is then given by:

$$P_z = \frac{R^2 - 1}{R^2 + R + 1} \quad (19)$$

and the tensor polarization is given by:

$$P_{zz} = \frac{R^2 - 2R + 1}{R^2 + R + 1} \quad (20)$$

This measuring technique can be used as a compliment to the TE method resulting in reduced uncertainty in polarization.

#### 2.4.2 Tensor Polarization Enhancement

It is possible to enhance tensor polarization using RF irradiation on the oriented deuterium nuclei to manipulate the alignment. Applying a saturating RF field on the pedestal of the smaller transition equalizes the substate  $m = +1$  and  $m = 0$  populations over 2/3 of the NMR signal. This equalization over the range of a single pedestal leads to enhancement in tensor polarization with only a small loss to the overall area ( $\sim 2\%$ ). Very recent studies at UVA using deuterated butanol have indicated that the tensor polarization can be increased by using a modified hole burning technique. The result will be investigated in the near future, and the method applied to  $\text{ND}_3$ . The studies also indicate that microwaves used during DNP does not interfere with the saturation from the RF irradiation when sufficient power is used. This implies that RF over the pedestal can be done the same time DNP is performed to enhance the area while taking beam in an experiment. Research and development is ongoing to study various techniques to optimize tensor enhancement for nuclear experiments targets.

#### 2.4.3 Depolarizing the Target

To move from polarized to unpolarized measurements, the target polarization will be annihilated using destructive NMR loop field changes and destructive DNP microwave pumping. During unpolarized data taking the incident electron beam heating is enough to remove the thermal equilibrium polarization.

We are able to verify that the target is in the unpolarized state via NMR measurements. The target material will be kept at 1 K for polarized and unpolarized data collection, and the target field will be held constant for both states as well. These consistencies are used to minimize the systematic differences in the polarized and unpolarized data collection. To minimize systematic effects over time, the polarization condition will be switched twice in a 72 hour period, as shown in Fig. 10. This will be sufficient to account for drift in integrated charge accumulation.

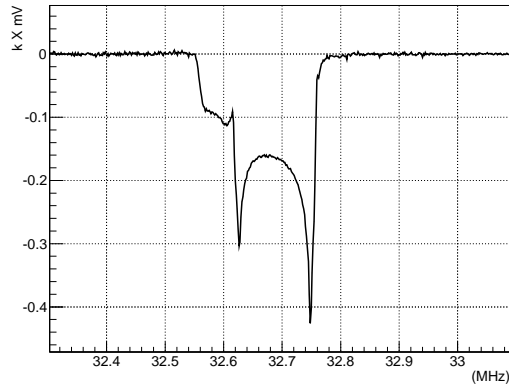


Figure 9: The deuterium magnetic resonance line shape showing the recent achievement of high tensor polarization of deuterated butanol after RF saturation of a pedestal at the UVA polarized target lab accomplished during their April 2014 cool-down.

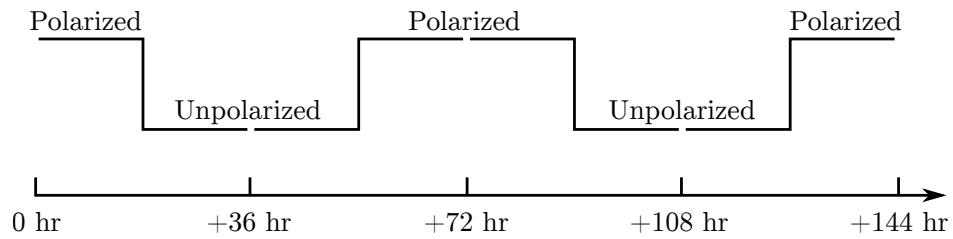


Figure 10: A visual demonstration of how the polarization cycle will happen over a 72 hour period to reduce time-dependent systematic effects. For the two lower  $Q^2$  measurements, the cycle will happen over a 18 hour period.

#### 2.4.4 Dilution Factor

To derive the dilution factor, we first start with the ratio of polarized to unpolarized counts. In each case, the number of counts that are actually measured, neglecting the small contributions of the thin aluminium cup window materials, NMR coils, etc., are

$$N_1 = Q_1 \varepsilon_1 \mathcal{A} l_1 [(\sigma_N + 3\sigma_1)p_f + \sigma_{He}(1 - p_f)], \quad (21)$$

and

$$N = Q \varepsilon \mathcal{A} l [(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)]. \quad (22)$$

where  $Q$  represents accumulated charge,  $\varepsilon$  is the dectector efficiency,  $\mathcal{A}$  the cup acceptance, and  $l$  the cup length.

For this calculation we assume similar charge accumulation such that  $Q \simeq Q_1$ , and that the efficiencies stay constant, in which case all factors drop out of the ratio leading to

$$\begin{aligned} \frac{N_1}{N} &= \frac{(\sigma_N + 3\sigma_1)p_f + \sigma_{He}(1 - p_f)}{(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)} \\ &= \frac{(\sigma_N + 3\sigma(1 + A_{zz}P_{zz}/2))p_f + \sigma_{He}(1 - p_f)}{(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)} \\ &= \frac{[(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)] + 3\sigma A_{zz}P_{zz}/2}{(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)} \\ &= 1 + \frac{3\sigma A_{zz}P_{zz}/2}{(\sigma_N + 3\sigma)p_f + \sigma_{He}(1 - p_f)} \\ &= 1 + \frac{1}{2}f A_{zz}P_{zz}, \end{aligned} \quad (23)$$

where  $\sigma_1 = \sigma(1 + A_{zz}P_{zz}/2)$  has ben substituted, per Eq. 3, with  $P_B = 0$ . It can be seen that the above result corresponds to Eq. 6.

## 2.5 Overhead

Table 4 summarizes the expected overhead, which sums to 8.6 days. The dominant overhead comes from switching from the polarized to unpolarized state and vice versa, and target anneals. The target will need to be annealed about every other day, and the material replaced once a week. Measurements of the dilution from the unpolarized materials contained in the target, and of the packing fraction due to the granular composition of the target material will be performed with a carbon target.

## 3 Summary

We have investigated the possibility of making high precision measurements of the quasi-elastic tensor asymmetry  $A_{zz}$ . By covering the kinematic range from the QE peak ( $x = 1$ ) up to elastic

Overhead	Number	Time Per (hr)	(hr)
Polarization/depolarization	35	2.0	70.0
Target anneal	13	4.0	52.0
Target T.E. measurement	5	4.0	20.0
Target material change	4	4.0	16.0
Packing Fraction/Dilution runs	6	1.0	6.0
BCM calibration	8	2.0	16.0
Optics	3	4.0	12.0
Linac change	1	8.0	8.0
Momentum/angle change	3	2.0	6.0
			8.6 days

Table 4: Major contributions to the overhead.

scattering ( $x = 2$ ), we expect that this data will provide valuable new insights about the high momentum components of the deuteron wavefunction. We are actively working with several theorists to get state-of-the-art calculations of light cone, virtual nucleon, and six-quark models. Additional calculations are being performed that include final-state interactions, and low  $Q^2$  sensitivity to NN potentials. It is important to note that this is the same kinematic region that has been shown to be correlated with the EMC effect via the  $x > 1$  A/D ( $e, e'$ ) results.

We have found that with 30 days of beam and an additional 8.6 days of overhead,  $A_{zz}$  can be measured with high precision at  $Q^2 = 1.5, 0.7, \text{ and } 0.3 \text{ (GeV}/c)^2$  in Hall C using identical equipment as the upcoming  $b_1$  measurement while being less sensitive to systematic uncertainties. In addition, it will fill a gap in measurements of  $A_{zz}$  between the  $T_{20} \propto A_{zz}$  elastic measurements and the  $b_1 \propto \frac{A_{zz}}{F_1^d}$  deep-inelastic measurements.

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