

# Noether's Theorem

## I. TRANSLATIONAL SYMMETRY IN HAMILTONIAN

Consider some Hamiltonian:

$$H[\phi(\mathbf{r})] = \int \psi(\phi, \partial_\alpha \phi) d^3r, \quad (1)$$

where  $\psi$  is the energy density (including the gradient terms). Let's suppose we now perform a spatial translation  $\delta \mathbf{r}$  on the scalar field  $\phi(\mathbf{r})$ :

$$\phi(\mathbf{r}) \rightarrow \phi(\mathbf{r} - \delta \mathbf{r}) = \phi(\mathbf{r}) - \underbrace{\delta r_\alpha \partial_\alpha \phi}_{\delta \phi}. \quad (2)$$

The energy density must also be translated *via*  $\delta \mathbf{r}$ , *i.e.*

$$\psi(\mathbf{r}) \rightarrow \psi(\mathbf{r} - \delta \mathbf{r}) = \psi(\mathbf{r}) - \underbrace{\delta r_\alpha \partial_\alpha \psi}_{\delta \psi}. \quad (3)$$

Thus the change in the energy density due this spatial translation is:

$$\delta \psi = -\delta r_\alpha \partial_\alpha \psi. \quad (4)$$

However, we also know that:

$$\delta \psi = \left[ \frac{\partial \psi}{\partial \phi} - \partial_\alpha \left( \frac{\partial \psi}{\partial (\partial_\alpha \phi)} \right) \right] \delta \phi + \partial_\alpha \left( \frac{\partial \psi}{\partial (\partial_\alpha \phi)} \delta \phi \right). \quad (5)$$

If  $\phi(\mathbf{r})$  remains in equilibrium before and after the transformation,  $\phi(\mathbf{r})$  must then satisfy Euler-Lagrange equation and the terms inside the square bracket above vanish. Thus we have:

$$\delta \psi = \partial_\alpha \left( \frac{\partial \psi}{\partial (\partial_\alpha \phi)} \delta \phi \right) \quad \text{and} \quad \delta \psi = -\delta r_\alpha \partial_\alpha \psi. \quad (6)$$

In other words,

$$\partial_\alpha \left( \frac{\partial \psi}{\partial (\partial_\alpha \phi)} \delta r_\beta \partial_\beta \phi \right) - \delta r_\beta \partial_\beta \psi = 0 \quad (7)$$

$$\partial_\alpha \underbrace{\left( \frac{\partial \psi}{\partial (\partial_\alpha \phi)} \partial_\beta \phi - \psi \delta_{\alpha\beta} \right)}_{\mathcal{J}_{\alpha\beta}} = 0 \quad (8)$$

$$\mathcal{J}_{\alpha\beta} = \text{constant} \quad (9)$$

We often call  $\mathcal{J}_{\alpha\beta}$  the Noether current.

For a symmetric Landau energy, we have:

$$\psi = \underbrace{\frac{\alpha}{2}\phi^2 + \frac{\beta}{4}\phi^4}_{f(\phi)} + \frac{\kappa}{2}|\nabla\phi|^2, \quad (10)$$

and the Noether current is:

$$\mathcal{J}_{\alpha\beta} = \kappa(\partial_\alpha \phi)(\partial_\beta \phi) - \left[ f(\phi) + \frac{\kappa}{2}|\nabla\phi|^2 \right] \delta_{\alpha\beta} = \text{constant} \quad (11)$$

In particular in one-dimension  $\alpha = \beta = x$ , we get:

$$\mathcal{J}_{xx} = \frac{\kappa}{2} \left( \frac{\partial \phi}{\partial x} \right)^2 - f(\phi) = \text{constant}. \quad (12)$$