

Quasi-structure in the dual Hessian for distributed MPC with non-delayed couplings

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Received: date / Accepted: date

Abstract Bla bla bla

Keywords First keyword · Second keyword · More

1 Introduction

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1.1 Motivation: Distributed MPC with non-delayed couplings

We consider M discrete time linear systems given by:

$$x_{1,i+1} = A_{1,i}x_{1,i} + B_{1,i}u_{1,i} \quad (1a)$$

$$\vdots \quad (1b)$$

$$x_{M,i+1} = A_{M,i}x_{M,i} + B_{M,i}u_{M,i} \quad (1c)$$

where $x_{k,i} \in \mathbb{R}^{p_k}$ and $u_{k,i} \in \mathbb{R}^{m_k}$ represents the state and input of system k . We also assume state and input constraints:

$$x_{k,i} \in X_{k,i} \subseteq \mathbb{R}^{p_k} \quad (2a)$$

$$u_{k,i} \in U_{k,i} \subseteq \mathbb{R}^{m_k} \quad (2b)$$

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Moreover, we assume non-delayed couplings between the systems:

$$h_i(x_{1,i}, u_{1,i}, \dots, x_{M,i}, u_{M,i}) = 0, \quad i = 0, \dots, N-1 \quad (3)$$

This means that we can state the MPC problem over the horizon N as:

$$\min_{x_{k,i}, u_{k,i}} \sum_{k=1}^M \left(\sum_{i=0}^{N-1} l(x_{k,i}, u_{k,i}) + l_f(x_{k,N}) \right) \quad (4a)$$

$$\text{s.t.} \quad x_{k+1,i} = A_{k,i}x_{k,i} + B_{k,i}u_{k,i} \quad (4b)$$

$$h_k(x_{k,1}, u_{k,1}, \dots, x_{k,M}, u_{k,M}) = 0 \quad (4c)$$

$$x_{k,i} \in X_{k,i}, \quad u_{k,i} \in U_{k,i} \quad (4d)$$

Furthermore, let us assume that the non-delayed couplings (4c) are linear, constraints (4d) are polyhedral and that the objective function (4a) is quadratic, and introduce the following notation: $z_{k,i} = [x_{k,i}^T \quad u_{k,i}^T]^T \in \mathbb{R}^{n_k}$. The MPC problem can then be stated as:

$$\min_z \sum_{k=1}^M \sum_{i=0}^N \frac{1}{2} z_{k,i}^T H_{k,i} z_{k,i} + c_{k,i}^T z_{k,i} \quad (5a)$$

$$\text{s.t.} \quad \sum_{k=1}^M F_{k,i} z_{k,i} = e_i \quad (5b)$$

$$C_{k,i} z_{k,i} + D_{k,i+1} z_{k,i+1} = d_{k,i} \quad (5c)$$

$$G_{k,i} z_{k,i} \leq f_{k,i} \quad (5d)$$

where $H_{k,i} \in \mathbb{S}_{++}^{n_k \times n_k}$, $C_{k,i} \in \mathbb{R}^{l_k \times n_k}$, $D_{k,i+1} \in \mathbb{R}^{l_k \times n_k}$ and $d_{k,i} \in \mathbb{R}^{l_k}$ form the dynamics, $F_{k,i} \in \mathbb{R}^{r_i \times n_k}$ and $e_i \in \mathbb{R}^{r_i}$ yield the coupling constraints, $G_{k,i} \in \mathbb{R}^{t_k \times n_k}$ and $f_{k,i} \in \mathbb{R}^{t_k}$ form the local constraints.

Additionally, to avoid an unnecessary heavy notation at places where we are only dealing with decomposition in space, we introduce the following augmented notations: $z_k = [z_{k,0}^T \dots z_{k,N-1}^T]^T \in \mathbb{R}^{n_k}$. The MPC problem (5) can then be expressed as:

$$\min_z \sum_{k=1}^M \frac{1}{2} z_k^T H_k z_k + c_k^T z_k \quad (6a)$$

$$\text{s.t.} \quad \sum_{k=1}^M F_k z_k = e \quad (6b)$$

$$C_k z_k = d_k \quad (6c)$$

$$G_k z_k \leq f_k \quad (6d)$$

where $H_k \in \mathbb{S}_{++}^{(N+1)n_k \times (N+1)n_k}$, $C_k \in \mathbb{R}^{Nl_k \times (N+1)n_k}$ and $d_k \in \mathbb{R}^{Nl_k}$, $F_k \in \mathbb{R}^{(N+1)r \times (N+1)n_k}$ and $e \in \mathbb{R}^{(N+1)r}$, $G_k \in \mathbb{R}^{(N+1)t_k \times (N+1)n_k}$ and $f_k \in \mathbb{R}^{(N+1)t_k}$

and the matrices are accordingly given by:

$$\begin{aligned}
 H_k &= \begin{bmatrix} H_{k,0} & & \\ & \ddots & \\ & & H_{k,N} \end{bmatrix}, \\
 C_k &= \begin{bmatrix} C_{k,0} & D_{k,1} & & & \\ & C_{k,1} & D_{k,2} & & \\ & & \ddots & \ddots & \\ & & & C_{k,N-1} & D_{k,N} \end{bmatrix}, \\
 A_k &= \begin{bmatrix} A_{k,0} & & \\ & \ddots & \\ & & A_{k,N} \end{bmatrix}.
 \end{aligned}$$

2 Preliminaries

2.1 Dual decomposition with second-order information

We introduce the dual variables $\lambda \in \mathbb{R}^{(N+1)r}$ corresponding to the coupling constraints (6b) and define the Lagrange function as

$$\mathcal{L}(z, \lambda) = \sum_{k=1}^M \left(\frac{1}{2} z_k^T H_k z_k + c_k^T z_k \right) + \lambda^T \left(\sum_{k=1}^M F_k z_k - e \right) \quad (8)$$

Note that $\mathcal{L}(z, \lambda)$ is separable in z , i.e.

$$\mathcal{L}(z, \lambda) = \sum_{k=1}^M \mathcal{L}_k(z_k, \lambda) \quad (9)$$

with

$$\mathcal{L}_k(z_k, \lambda) = \frac{1}{2} z_k^T H_k z_k + c_k^T z_k + \lambda^T \left(C_k z_k - \frac{1}{M} d \right) \quad (10)$$

The dual function $d(\lambda) = -\min_{z \in \mathcal{Z}} \mathcal{L}(z, \lambda)$ can thus be evaluated in parallel as:

$$d(\lambda) = - \sum_{k=1}^M \min_{z_k \in \mathcal{Z}_k} \mathcal{L}_k(z_k, \lambda) \quad (11)$$

Since (6) is strictly convex, (11) is convex and continuously differentiable, but not twice differentiable. However, the Hessian of $d(\lambda)$ is a piecewise constant matrix and change with the active-set [5].

The non-smoothness implies that $d(\lambda)$ is not self-concordant and the solution to the dual problem is hence not easily tracked with Newton's method.

However, if we relax the inequality constraints (6d) with a self-concordant log-barrier, according to:

$$\min_z \sum_{k=1}^M \frac{1}{2} z_k^T H_k z_k + c_k^T z_k - \tau \sum_{i=1}^{(N+1)t_k} \log([s_k]_i) \quad (12a)$$

$$\text{s.t.} \quad \sum_{k=1}^M F_k z_k = e \quad (12b)$$

$$C_k z_k = d_k \quad (12c)$$

$$G_k z_k + s_k = f_k \quad (12d)$$

where $\tau > 0$ will be referred to as the *barrier parameter*, the resulting *relaxed dual function* $d(\lambda, \tau)$ is self-concordant [6]. This opens for the possibility of solving a sequence of dual problems $\{\min_{\lambda} d(\lambda, \tau)\}_{\tau \rightarrow 0}$ where each problem is self-concordant and therefore easily solved with Newton's method.

The relaxed dual function is separable and can be evaluated in parallel as

$$d(\lambda, \tau) = - \sum_{k=1}^N \min_{z_k \in \mathcal{Z}_k} \left(\mathcal{L}_k(z_k, \lambda) - \tau \sum_{i=1}^{m_k} \log([s_k]_i) \right) \quad (13)$$

Hence, evaluating (13) involves solving local subproblems of the form

$$\begin{aligned} \min_{z_k} \quad & \frac{1}{2} z_k^T H_k z_k + c_k^T z_k + \lambda^T F_k z_k - \tau \sum_{i=1}^{m_k} \log([s_k]_i) \\ \text{s.t.} \quad & C_k z_k = d_k \\ & G_k z_k + s_k = f_k \\ & s_k \geq 0 \end{aligned} \quad (14)$$

The relaxed dual problem then reads

$$\min_{\lambda} d(\lambda, \tau) \quad (15)$$

from which solution, the solution to (12) can be recovered according to strong duality [2].

Strict convexity also implies that the gradient of $d(\lambda, \tau)$ is given by the residual of the coupling constraints [1], i.e.

$$\nabla d(\lambda, \tau) = - \sum_{k=1}^N F_k z_k^*(\lambda, \tau) + e \quad (16)$$

where $z_k^*(\lambda, \tau) = \arg \min_{z_k \in \mathcal{Z}_k} \mathcal{L}_k(z_k, s_k, \lambda, \tau)$. The dual Hessian is then given by

$$\nabla^2 d(\lambda, \tau) = - \sum_{k=1}^N F_k \frac{\partial z_k^*(\lambda, \tau)}{\partial \lambda} \quad (17)$$

A Newton direction $\Delta \lambda$ in the dual space can then be obtained as a solution to the Newton system

$$\nabla^2 d(\lambda, \tau) \Delta \lambda + \nabla d(\lambda, \tau) = 0 \quad (18)$$

3 Structure in the dual Hessian

3.1 Dual Hessian

By introducing $y_k = \tau/s_k \in \mathbb{R}^{m_k}$, the KKT conditions to (14) are given by

$$r_k(w_k^*, \lambda, \tau) = \begin{bmatrix} r_{Dk}(w_k^*, \lambda) \\ r_{Ek}(w_k^*) \\ r_{Ik}(w_k^*) \\ r_{Ck}(w_k^*, \tau) \end{bmatrix} = 0 \quad (19a)$$

$$s_k^* > 0, \quad y_k^* > 0 \quad (19b)$$

where we use the notation $w_k = [z_k^T, \mu_k^T, y_k^T, s_k^T]^T$ for the local primal-dual variables and $r_k(w_k, \lambda, \tau)$ is given by:

$$r_{Dk}(w_k, \lambda) = H_k z_k + c_k + F_k^T \lambda + C_k^T \mu_k + G_k^T y_k \quad (20a)$$

$$r_{Ek}(w_k) = C_k z_k - d_k \quad (20b)$$

$$r_{Ik}(w_k) = G_k z_k + s_k - f_k \quad (20c)$$

$$r_{Ck}(w_k, \tau) = Y_k s_k - \tau \mathbf{1} \quad (20d)$$

As we have seen in previous section, in order to form the dual Hessian we need to compute $\frac{\partial z_k^*}{\partial \lambda}$, where $z_k^* = z_k^*(\lambda)$ is the optimal primal solution and is hence fulfilling (19). By differentiating (19), the following linear system is obtained:

$$\begin{bmatrix} H_k & C_k^T & G_k^T & 0 \\ C_k & 0 & 0 & 0 \\ G_k & 0 & 0 & I \\ 0 & 0 & S_k & Y_k \end{bmatrix} \begin{bmatrix} \frac{\partial z_k^*}{\partial \lambda} \\ \frac{\partial \mu_k^*}{\partial \lambda} \\ \frac{\partial y_k^*}{\partial \lambda} \\ \frac{\partial s_k^*}{\partial \lambda} \end{bmatrix} = - \begin{bmatrix} F_k^T \\ 0 \\ 0 \\ 0 \end{bmatrix} \quad (21)$$

If block elimination of (21) is used, the *normal equations* can be formed [7],

$$\Lambda_k \frac{\partial \mu_k^*}{\partial \lambda} = -C_k \Phi_k^{-1} F_k^T \quad (22a)$$

$$\Phi_k \frac{\partial z_k^*}{\partial \lambda} = -F_k^T - C_k^T \frac{\partial \mu_k^*}{\partial \lambda} \quad (22b)$$

$$\frac{\partial s_k^*}{\partial \lambda} = -G_k \frac{\partial z_k^*}{\partial \lambda} \quad (22c)$$

$$\frac{\partial y_k^*}{\partial \lambda} = -S^{-1} Y \frac{\partial s_k^*}{\partial \lambda} \quad (22d)$$

where $\Phi_k = H_k + G_k^T S_k^{-1} Y_k G_k \in \mathbb{S}_{++}^{(N+1)n_k \times (N+1)n_k}$ and $\Lambda_k = C_k \Phi_k^{-1} C_k^T \in \mathbb{S}_{++}^{Nl_k \times Nl_k}$. By using (22a) and (22b), it can be obtained that

$$F_k \frac{\partial z_k}{\partial \lambda} = -F_k (\Phi_k^{-1} - \Phi_k^{-1} C_k^T \Lambda_k^{-1} C_k \Phi_k^{-1}) F_k^T \quad (23)$$

This means that the dual Hessian (17) can be written as

$$\nabla^2 d(\lambda, \tau) = \sum_{k=1}^M F_k (\Phi_k^{-1} - \Phi_k^{-1} C_k^T \Lambda_k^{-1} C_k \Phi_k^{-1}) F_k^T \quad (24)$$

Let us now turn our attention to the structure coming from the time domain. We have already seen that C_k is block bidiagonal, and it can trivially be realized that Φ_k is block diagonal. This means that Λ_k has a block tridiagonal structure given by:

$$\Lambda_k = \begin{bmatrix} \Lambda_{11} & \Lambda_{12} & & & \\ \Lambda_{12}^T & \Lambda_{22} & & & \\ & & \ddots & & \\ & & & \ddots & \\ & & & & \Lambda_{N-1,N} \\ & & & \Lambda_{N-1,N}^T & \Lambda_{N,N} \end{bmatrix} \in \mathbb{S}_{++}^{Nl_k \times Nl_k} \quad (25)$$

where

$$\Lambda_{i,i} = C_{k,i-1} \Phi_{k,i-1}^{-1} C_{k,i-1}^T + D_{k,i} \Phi_{k,i}^{-1} D_{k,i}^T \in \mathbb{R}^{l_k \times l_k} \quad (26a)$$

$$\Lambda_{i,i+1} = D_{k,i} \Phi_{k,i}^{-1} C_{k,i}^T \in \mathbb{R}^{l_k \times l_k} \quad (26b)$$

Accordingly, all matrices in (24) are banded except Λ_k^{-1} which in general is dense.

3.2 Decaying of the dual Hessian

Let us first assume that the problem data is bounded. In other words, let us assume the following:

Assumption 1 *The row and column absolute sums of Jacobians of equality constraints (i.e. (6b) and (6c)) are bounded. Hence:*

1. $\|C_{k,i}\|_\infty \leq \gamma$ and $\|C_{k,i}\|_1 \leq \gamma$
2. $\|D_{k,i}\|_\infty \leq \gamma$ and $\|D_{k,i}\|_1 \leq \gamma$
3. $\|F_{k,i}\|_\infty \leq \gamma$ and $\|F_{k,i}\|_1 \leq \gamma$

It should be observed that Assumption 1 is not by any means restrictive, since any solver would struggle with a problem where it is not fulfilled.

Possible problems with interior-point methods are related to numerical difficulties due to ill-conditioning []. This might occur since elements in $S_k^{-1} Y_k$ can be close to zero in late iterations when τ is small. Our experience is however that this is not a major issue in practice, which is also supported by commercial interior-point implementations []. Moreover, since we assume that (5) is strongly convex, the eigenvalues of Φ_k is lower bounded by the smallest eigenvalue of H_k , even when $S_k^{-1} Y_k$ is singular. According to this reasoning it is not restrictive to assume boundedness of Φ_k^{-1} which we do next:

Assumption 2 *The row and column absolute sums of Φ_k^{-1} are bounded. Hence:*

$$\|\Phi_{k,i}^{-1}\|_\infty = \|\Phi_{k,i}^{-1}\|_1 \leq \gamma_{\Phi^{-1}} \quad (27)$$

Inverses of sparse matrices are in general dense, but individual elements are often small in absolute value. Since Λ_k is banded, symmetric and positive definite, we will relay our analysis on the following classical result:

Lemma 1 *If A is Hermitian positive definite and m -banded ($[A]_{ij} = 0$ if $|i - j| > m$), the entries of A^{-1} satisfy the following bound:*

$$|[A^{-1}]_{ij}| < K\omega^{|i-j|}, \quad \forall i, j \quad (28)$$

where $[a, b]$ is the smallest interval containing the spectrum $\sigma(A)$ of A , $K = \max\{a^{-1}, K_0\}$, $K_0 = (1 + \sqrt{\kappa})$, $\omega = q^{1/m}$, $q = \frac{\sqrt{\kappa}-1}{\sqrt{\kappa}+1}$, $\kappa = \frac{b}{a}$.

Proof A proof is given in []. \square

This means that the entries of A^{-1} are bounded by an exponentially decaying function along each row or column. However, the bound depends on the condition number and the bandwidth of the matrix. Matrices with a high condition number and/or a high bandwidth can accordingly result in a large K and $\omega \approx 1$, leading to a slow decay. The opposite, i.e. a low condition number and a small band, would result in a rapid decay.

Observe that due to the block tridiagonal structure of Λ_k , it is $3l_k$ -banded. Moreover, if we introduce the notation

$$\Lambda_k^{-1} = \begin{bmatrix} T_{11} & T_{21}^T & \cdots & T_{N+1,1}^T \\ T_{21} & T_{22} & \cdots & T_{N+1,2}^T \\ \vdots & \vdots & \ddots & \vdots \\ T_{N+1,1} & T_{N+1,2} & \cdots & T_{N+1,N+1} \end{bmatrix} \quad (29)$$

where $T_{i,j} \in \mathbb{R}^{l_k \times l_k}$, we can establish the following proposition:

Proposition 1 *The off-diagonal blocks (i.e. $T_{i,j}$ where $i - j > 0$) in Λ_k^{-1} satisfy the following bounds:*

$$\|T_{i,j}\|_\infty \leq l_k K_k \omega_k^{(i-j)l_k+1} \quad (30a)$$

$$\|T_{i,j}\|_1 \leq l_k K_k \omega_k^{(i-j)l_k+1} \quad (30b)$$

where K_k and ω_k are defined according to Lemma 1.

Proof According to Lemma 1, the element in $T_{i,j}$ with the largest bound is located in the top-right corner, and is hence element $[\Lambda_k^{-1}]_{il_k+1, jl_k}$. From Lemma 1 it then follows that:

$$\max |[T_{i,j}]| \leq K_k \omega_k^{(i-j)l_k+1} \quad (31)$$

where $\max ||[T_{ij}]||$ refers to the maximum absolute value of the components in T_{ij} . Accordingly, we have that:

$$\|T_{i,j}\|_\infty \leq l_k K_k \omega_k^{(i-j)l_k+1} \quad (32a)$$

$$\|T_{i,j}\|_1 \leq l_k K_k \omega_k^{(i-j)l_k+1} \quad (32b)$$

□

Proposition 1 confirms our intuition that the row and column sum of T_{ij} are decreasing when $|i - j|$ increases. For interior point methods in general, the condition number of A_k can be very high for small values of τ . It should however be understood that strong convexity of (5), should improve the conditioning of Φ_k and hence also of A_k .

To gain some perspective, we have concluded that A_k^{-1} is decaying exponentially towards the off-diagonal corners and that all other matrices in (24) are banded. This suggests that also $\nabla_{\lambda\lambda}^2 d(\lambda, \tau)$ should decay towards the off-diagonal corners. To maintain a simple reasoning, one more stepping stone will be used before arriving at the main results of this section. Therefore, let us introduce the notation:

$$C_k^T A_k^{-1} C_k = \begin{bmatrix} V_{11} & V_{21}^T & \cdots & V_{N+1,1}^T \\ V_{21} & V_{22} & \cdots & V_{N+1,2}^T \\ \vdots & \vdots & \ddots & \vdots \\ V_{N+1,1} & V_{N+1,2} & \cdots & V_{N+1,N+1} \end{bmatrix} \quad (33)$$

where $V_{ij} \in \mathbb{R}^{l_k \times l_k}$, and look at the decay of $C_k^T A_k^{-1} C_k$.

Proposition 2 *The off diagonal blocks (i.e. $V_{i,j}$ where $i - j > 0$) of $C_k^T A_k^{-1} C_k$ satisfy the following bounds:*

$$\|V_{i,j}\|_\infty \leq (2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^2 l_k K_k \omega_k^{(i-j-1)l_k+1} \quad (34a)$$

$$\|V_{i,j}\|_1 \leq (2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^2 l_k K_k \omega_k^{(i-j-1)l_k+1} \quad (34b)$$

Proof Bla bla bla

□

Let us now introduce the notation:

$$\nabla^2 d(\lambda, \tau) = \begin{bmatrix} W_{11} & W_{21}^T & \cdots & W_{N+1,1}^T \\ W_{21} & W_{22} & \cdots & W_{N+1,2}^T \\ \vdots & \vdots & \ddots & \vdots \\ W_{N+1,1} & W_{N+1,2} & \cdots & W_{N+1,N+1} \end{bmatrix} \quad (35)$$

where $W_{i,j} \in \mathbb{R}^{r_i \times r_j}$. We can now establish a decay towards the off-diagonal corners of $\nabla^2 d(\lambda, \tau)$.

Proposition 3 *The off-diagonal blocks (i.e. $W_{i,j}$ where $i-j > 0$) of $\nabla^2 d(\lambda, \tau)$ satisfy the following bounds:*

$$\|W_{i,j}\|_\infty \leq \sum_{k=1}^M (2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^4 \gamma_{\Phi-1}^2 l_k K_k \omega_k^{(i-j-1)l_k+1} \quad (36a)$$

$$\|W_{i,j}\|_1 \leq \sum_{k=1}^M (2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^4 \gamma_{\Phi-1}^2 l_k K_k \omega_k^{(i-j-1)l_k+1} \quad (36b)$$

Proof To be done.

Finally, let us bound the euclidean distance between $\nabla^2 d(\lambda, \tau)$ and a band along its diagonal. To do so, we start by recalling Gershgorin's circle theorem:

Theorem 1 *For $A \in \mathbb{R}^{n \times n}$ with elements a_{ij} , let $R_i = \sum_{j \neq i} |a_{ij}|$ be the sum of the absolute values of the non-diagonal entries in the row i . Let $D(a_{ii}, R_i)$ be the closed disc centered in a_{ii} with radius R_i , then every eigenvalue of A lies within at least one of the discs $D(a_{ii}, R_i)$.*

Proof A proof is given in []. □

Finally, we can establish our main result:

Lemma 2 *Let $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}}$ represent the diagonal \mathcal{M} -block band of $\nabla^2 d(\lambda, \tau)$. The following bound holds:*

$$\|[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)\|_2 \leq \sum_{k=1}^M (N+1-\mathcal{M})(2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^4 \gamma_{\Phi-1}^2 l_k K_k \omega_k^{\frac{\mathcal{M}-1}{2} l_k + 1} \quad (37)$$

where...

Proof Since $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)$ is symmetric, its singular values collapse into the absolute value of its eigenvalues. The problem of finding a bound on the 2-norm is hence reduced to the problem of bounding the magnitude of the largest eigenvalue.

First, observe that all diagonal elements of $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)$ are zero. This means that every Gershgorin disc will be centered in the origin, and a bound on the largest eigenvalue can be found by finding the largest radius of a Gershgorin disc.

The blocks $W_{i,j}$ in $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)$ with the largest bound is located next to the diagonal band of zeros, i.e. where $i - j = \frac{\mathcal{M}-1}{2} + 1$. Furthermore, at each block row there are $N + 1 - \mathcal{M}$ nonzero blocks. This means that we can establish a bound on the radius of the Gershgorin discs, and hence the following:

$$\|[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)\|_2 \leq \sum_{k=1}^M (N+1-\mathcal{M})(2 + \omega_k^{l_k} + \omega_k^{-l_k}) \gamma^4 \gamma_{\Phi-1}^2 l_k K_k \omega_k^{\frac{\mathcal{M}-1}{2} l_k + 1} \quad (38)$$

□

Say something about the principal behavior of the bound...

It is well known that the use of an inexact Hessian will degrade the convergence of Newton's method to a linear rate [1]. However, let us still look at the relative error in the Newton direction that we get from using $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}}$ instead of $\nabla^2 d(\lambda, \tau)$. First, recall the following lemma:

Lemma 3 *If x is a solution to $Ax = b$ and \hat{x} is a solution to the perturbed system $(A + F)\hat{x} = b + f$, then the following bound holds:*

$$\frac{\|\hat{x} - x\|_2}{\|x\|_2} \leq \kappa(A)(\rho_A - \rho_b) \quad (39)$$

where $\kappa(A)$ represents the condition number of A , $\rho_A = \frac{\|F\|_2}{\|A\|_2}$ and $\rho_b = \frac{\|f\|_2}{\|b\|_2}$.

Proof A proof is given in [1].

Accordingly, by combining Lemma 2 and Lemma 3 we can establish the following:

Lemma 4

$$\frac{\|\Delta\hat{\lambda} - \Delta\lambda\|_2}{\|\Delta\lambda\|_2} \leq \quad (40)$$

4 Numerical experiments

4.1 Decaying of the dual Hessian

4.1.1 A small problem

As an example of a small problem, we use a randomly generated problem with 4 subproblems (i.e. $M = 4$), where each subproblem has $N = 10$, 6 states and 4 controls and hence $n_k = 10$, and each time instance has 4 coupling constraints.

According to Section 3.2, we expect a slower decay when the conditioning of $S_k^{-1}Y_k$ gets worse, i.e. when we reduce τ . Surprisingly, this is something that does not seem to have a big influence in practice. To illustrate this, Figure 2 shows the dual Hessian for the same problem as in Figure 1, but with $\tau = 10^{-5}$.

4.1.2 A large problem

4.2 Newton steps with inexact Hessian

References

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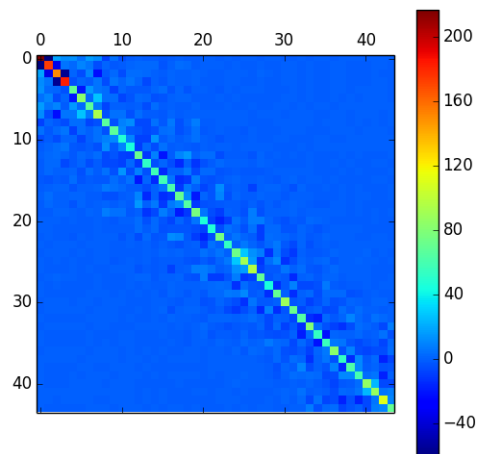


Fig. 1 $\tau = 1$

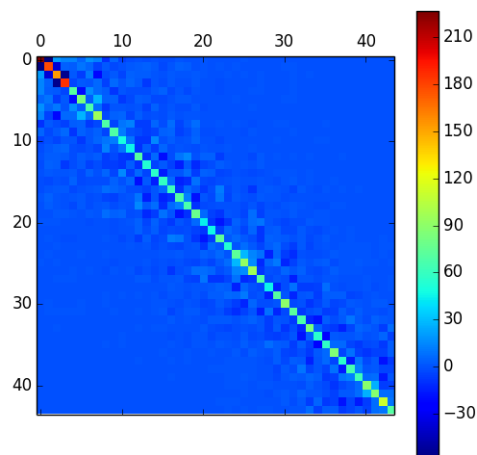


Fig. 2 $\tau = 10^{-5}$

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