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# Quasi-structure in the dual Hessian for distributed MPC with non-delayed couplings

Emil Klintberg  $\cdot$  Sebastien Gros

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Abstract Bla bla bla

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## 1 Introduction

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1.1 Motivation: Distributed MPC with non-delayed couplings

We consider M discrete time linear systems given by:

$$x_{1,i+1} = A_{1,i}x_{1,i} + B_{1,i}u_{1,i}$$
(1a)

$$x_{M,i+1} = A_{M,i} x_{M,i} + B_{M,i} u_{M,i}$$
 (1c)

where  $x_{k,i} \in \mathbb{R}^{p_k}$  and  $u_{k,i} \in \mathbb{R}^{m_k}$  represents the state and input of system k. We also assume state and input constraints:

$$x_{k,i} \in X_{k,i} \subseteq \mathbb{R}^{p_k}$$
 (2a)

$$u_{k,i} \in U_{k,i} \subseteq \mathbb{R}^{m_k}$$
 (2b)

Emil Klintberg first address

Tel.: +123-45-678910 Fax: +123-45-678910

E-mail: fauthor@example.com

Sebastien Gros second address

Moreover, we assume non-delayed couplings between the systems:

$$h_i(x_{1,i}, u_{1,i}, \dots, x_{M,i}, u_{M,i}) = 0, \quad i = 0, \dots, N-1$$
 (3)

This means that we can state the MPC problem over the horizon N as:

$$\min_{x_{k,i}, u_{k,i}} \sum_{k=1}^{M} \left( \sum_{i=0}^{N-1} l(x_{k,i}, u_{k,i}) + l_f(x_{k,N}) \right)$$
 (4a)

s.t. 
$$x_{k+1,i} = A_{k,i} x_{k,i} + B_{k,i} u_{k,i}$$
 (4b)

$$h_k(x_{k,1}, u_{k,1}, \dots, x_{k,M}, u_{k,M}) = 0$$
 (4c)

$$x_{k,i} \in X_{k,i}, \quad u_{k,i} \in U_{k,i} \tag{4d}$$

Furthermore, let us assume that the non-delayed couplings (4c) are linear, constraints (4d) are polyhedral and that the objective function (4a) is quadratic, and introduce the following notation:  $z_{k,i} = [x_{k,i}^T \ u_{k,i}^T]^T \in \mathbb{R}^{n_k}$ . The MPC problem can then be stated as:

$$\min_{z} \sum_{k=1}^{M} \sum_{i=0}^{N} \frac{1}{2} z_{k,i}^{T} H_{k,i} z_{k,i} + c_{k,i}^{T} z_{k,i}$$
 (5a)

s.t. 
$$\sum_{k=1}^{M} F_{k,i} z_{k,i} = e_i$$
 (5b)

$$C_{k,i}z_{k,i} + D_{k,i+1}z_{k,i+1} = d_{k,i} (5c)$$

$$G_{k,i}z_{k,i} \le f_{k,i} \tag{5d}$$

where  $H_{k,i} \in \mathbb{S}_{++}^{n_k \times n_k}$ ,  $C_{k,i} \in \mathbb{R}^{l_k \times n_k}$ ,  $D_{k,i+1} \in \mathbb{R}^{l_k \times n_k}$  and  $d_{k,i} \in \mathbb{R}^{l_k}$  form the dynamics,  $F_{k,i} \in \mathbb{R}^{r_i \times n_k}$  and  $e_i \in \mathbb{R}^{r_i}$  yield the coupling constraints,  $G_{k,i} \in \mathbb{R}^{t_k \times n_k}$  and  $f_{k,i} \in \mathbb{R}^{t_k}$  form the local constraints.

Additionally, to avoid an unnecessary heavy notation at places where we are only dealing with decomposition in space, we introduce the following augmented notations:  $z_k = [z_{k,0}^T \dots, z_{k,N-1}^T]^T \in \mathbb{R}^{n_k}$ . The MPC problem (5) can then be expressed as:

$$\min_{z} \sum_{k=1}^{M} \frac{1}{2} z_k^T H_k z_k + c_k^T z_k$$
 (6a)

s.t. 
$$\sum_{k=1}^{M} F_k z_k = e \tag{6b}$$

$$C_k z_k = d_k \tag{6c}$$

$$G_k z_k \le f_k \tag{6d}$$

where  $H_k \in \mathbb{S}_{++}^{(N+1)n_k \times (N+1)n_k}$ ,  $C_k \in \mathbb{R}^{Nl_k \times (N+1)n_k}$  and  $d_k \in \mathbb{R}^{Nl_k}$ ,  $F_k \in \mathbb{R}^{(N+1)r \times (N+1)n_k}$  and  $e \in \mathbb{R}^{(N+1)r}$ ,  $G_k \in \mathbb{R}^{(N+1)t_k \times (N+1)n_k}$  and  $f_k \in \mathbb{R}^{(N+1)t_k}$ 

and the matrices are accordingly given by:

$$H_k = egin{bmatrix} H_{k,0} & & & & \\ & \ddots & & \\ & H_{k,N} \end{bmatrix},$$
 
$$C_k = egin{bmatrix} C_{k,0} & D_{k,1} & & & \\ & C_{k,1} & D_{k,2} & & & \\ & & \ddots & & & \\ & & C_{k,N-1} & D_{k,N} \end{bmatrix},$$
 
$$A_k = egin{bmatrix} A_{k,0} & & & \\ & \ddots & & \\ & & A_{k,N} \end{bmatrix}.$$

#### 2 Prelimenaries

#### 2.1 Dual decomposition with second-order information

We introduce the dual variables  $\lambda \in \mathbb{R}^{(N+1)r}$  corresponding to the coupling constraints (6b) and define the Lagrange function as

$$\mathcal{L}(z,\lambda) = \sum_{k=1}^{M} (\frac{1}{2} z_k^T H_k z_k + c_k^T z_k) + \lambda^T (\sum_{k=1}^{M} F_k z_k - e)$$
 (8)

Note that  $\mathcal{L}(z,\lambda)$  is separable in z, i.e.

$$\mathcal{L}(z,\lambda) = \sum_{k=1}^{M} \mathcal{L}_k(z_k,\lambda) \tag{9}$$

with

$$\mathcal{L}_{k}(z_{k},\lambda) = \frac{1}{2} z_{k}^{T} H_{k} z_{k} + c_{k}^{T} z_{k} + \lambda^{T} (C_{k} z_{k} - \frac{1}{M} d)$$
 (10)

The dual function  $d(\lambda) = -\min_{z \in \mathcal{Z}} \mathcal{L}(z, \lambda)$  can thus be evaluated in parallel as:

$$d(\lambda) = -\sum_{k=1}^{M} \min_{z_k \in \mathcal{Z}_k} \mathcal{L}_k(z_k, \lambda)$$
(11)

Since (6) is strictly convex, (11) is convex and continuously differentiable, but not twice differentiable. However, the Hessian of  $d(\lambda)$  is a piecewise constant matrix and change with the active-set [5].

The non-smoothness implies that  $d(\lambda)$  is not self-concordant and the solution to the dual problem is hence not easily tracked with Newton's method.

However, if we relax the inequality constraints (6d) with a self-concordant log-barrier, according to:

$$\min_{z} \sum_{k=1}^{M} \frac{1}{2} z_{k}^{T} H_{k} z_{k} + c_{k}^{T} z_{k} - \tau \sum_{i=1}^{(N+1)t_{k}} \log([s_{k}]_{i})$$
 (12a)

s.t. 
$$\sum_{k=1}^{M} F_k z_k = e$$
 (12b)

$$C_k z_k = d_k \tag{12c}$$

$$G_k z_k + s_k = f_k \tag{12d}$$

where  $\tau > 0$  will be referred to as the *barrier parameter*, the resulting *relaxed dual function*  $d(\lambda, \tau)$  is self-concordant [6]. This opens for the possibility of solving a sequence of dual problems  $\{\min_{\lambda} d(\lambda, \tau)\}_{\tau \to 0}$  where each problem is self-concordant and therefore easily solved with Newton's method.

The relaxed dual function is separable and can be evaluated in parallel as

$$d(\lambda, \tau) = -\sum_{k=1}^{N} \min_{z_k \in \mathcal{Z}_k} \left( \mathcal{L}_k(z_k, \lambda) - \tau \sum_{i=1}^{m_k} \log([s_k]_i) \right)$$
(13)

Hence, evaluating (13) involves solving local subproblems of the form

$$\min_{z_k} \quad \frac{1}{2} z_k^T H_k z_k + c_k^T z_k + \lambda^T F_k z_k - \tau \sum_{i=1}^{m_k} \log([s_k]_i)$$
s.t. 
$$C_k z_k = d_k$$

$$G_k z_k + s_k = f_k$$

$$s_k \ge 0$$
(14)

The relaxed dual problem then reads

$$\min_{\lambda} d(\lambda, \tau) \tag{15}$$

from which solution, the solution to (12) can be recovered according to strong duality [2].

Strict convexity also implies that the gradient of  $d(\lambda, \tau)$  is given by the residual of the coupling constraints [1], i.e.

$$\nabla d(\lambda, \tau) = -\sum_{k=1}^{N} F_k z_k^*(\lambda, \tau) + e$$
 (16)

where  $z_k^*(\lambda, \tau) = \arg\min_{z_k \in \mathcal{Z}_k} \mathcal{L}_k(z_k, s_k, \lambda, \tau)$ . The dual Hessian is then given by

$$\nabla^2 d(\lambda, \tau) = -\sum_{k=1}^N F_k \frac{\partial z_k^*(\lambda, \tau)}{\partial \lambda}$$
 (17)

A Newton direction  $\varDelta\lambda$  in the dual space can then be obtained as a solution to the Newton system

$$\nabla^2 d(\lambda, \tau) \Delta \lambda + \nabla d(\lambda, \tau) = 0 \tag{18}$$

#### 3 Structure in the dual Hessian

#### 3.1 Dual Hessian

By introducing  $y_k = \tau/s_k \in \mathbb{R}^{m_k}$ , the KKT conditions to (14) are given by

$$r_{k}(w_{k}^{*}, \lambda, \tau) = \begin{bmatrix} r_{Dk}(w_{k}^{*}, \lambda) \\ r_{Ek}(w_{k}^{*}) \\ r_{Ik}(w_{k}^{*}) \\ r_{Ck}(w_{k}^{*}, \tau) \end{bmatrix} = 0$$
 (19a)

$$s_k^* > 0, \quad y_k^* > 0 \tag{19b}$$

where we use the notation  $w_k = [z_k^T, \mu_k^T, y_k^T, s_k^T]^T$  for the local primal-dual variables and  $r_k(w_k, \lambda, \tau)$  is given by:

$$r_{D_k}(w_k, \lambda) = H_k z_k + c_k + F_k^T \lambda + C_k^T \mu_k + G_k^T y_k$$
 (20a)

$$r_{Ek}(w_k) = C_k z_k - d_k \tag{20b}$$

$$r_{Ik}(w_k) = G_k z_k + s_k - f_k \tag{20c}$$

$$r_{Ck}(w_k, \tau) = Y_k s_k - \tau \mathbf{1} \tag{20d}$$

As we have seen in previous section, in order to form the dual Hessian we need to compute  $\frac{\partial z_k^*}{\partial \lambda}$ , where  $z_k^* = z_k^*(\lambda)$  is the optimal primal solution and is hence fulfilling (19). By differentiating (19), the following linear system is obtained:

$$\begin{bmatrix} H_k & C_k^T & G_k^T & 0 \\ C_k & 0 & 0 & 0 \\ G_k & 0 & 0 & I \\ 0 & 0 & S_k & Y_k \end{bmatrix} \begin{bmatrix} \frac{\partial z_k^*}{\partial \lambda} \\ \frac{\partial \mu_k^*}{\partial \lambda} \\ \frac{\partial y_k}{\partial \lambda} \\ \frac{\partial z_k^*}{\partial \lambda} \end{bmatrix} = - \begin{bmatrix} F_k^T \\ 0 \\ 0 \\ 0 \end{bmatrix}$$
 (21)

If block elimination of (21) is used, the normal equations can be formed [7],

$$\Lambda_k \frac{\partial \mu_k^*}{\partial \lambda} = -C_k \Phi_k^{-1} F_k^T \tag{22a}$$

$$\Phi_k \frac{\partial z_k^*}{\partial \lambda} = -F_k^T - C_k^T \frac{\partial \mu_k^*}{\partial \lambda}$$
 (22b)

$$\frac{\partial s_k^*}{\partial \lambda} = -G_k \frac{\partial z_k^*}{\partial \lambda} \tag{22c}$$

$$\frac{\partial s_{k}^{*}}{\partial \lambda} = -G_{k} \frac{\partial z_{k}^{*}}{\partial \lambda}$$

$$\frac{\partial y_{k}^{*}}{\partial \lambda} = -S^{-1} Y \frac{\partial s_{k}^{*}}{\partial \lambda}$$
(22c)

where  $\Phi_k = H_k + G_k^T S_k^{-1} Y_k G_k \in \mathbb{S}_{++}^{(N+1)n_k \times (N+1)n_k}$  and  $\Lambda_k = C_k \Phi_k^{-1} C_k^T \in \mathbb{S}_{++}^{Nl_k \times Nl_k}$ . By using (22a) and (22b), it can be obtained that

$$F_k \frac{\partial z_k}{\partial \lambda} = -F_k (\Phi_k^{-1} - \Phi_k^{-1} C_k^T \Lambda_k^{-1} C_k \Phi_k^{-1}) F_k^T$$
 (23)

This means that the dual Hessian (17) can be written as

$$\nabla^2 d(\lambda, \tau) = \sum_{k=1}^M F_k (\Phi_k^{-1} - \Phi_k^{-1} C_k^T \Lambda_k^{-1} C_k \Phi_k^{-1}) F_k^T$$
 (24)

Let us now turn our attention to the structure coming from the time domain. We have already seen that  $C_k$  is block bidiagonal, and it can trivially be realized that  $\Phi_k$  is block diagonal. This means that  $\Lambda_k$  has a block tridiagonal structure given by:

$$\Lambda_{k} = \begin{bmatrix}
\Lambda_{11} & \Lambda_{12} & & & \\
\Lambda_{12}^{T} & \Lambda_{22} & \ddots & & \\
& \ddots & \ddots & \Lambda_{N-1,N} \\
& & \Lambda_{N-1,N}^{T} & \Lambda_{N,N}
\end{bmatrix} \in \mathbb{S}_{++}^{Nl_{k} \times Nl_{k}}$$
(25)

where

$$\Lambda_{i,i} = C_{k,i-1} \Phi_{k,i-1}^{-1} C_{k,i-1}^T + D_{k,i} \Phi_{k,i}^{-1} D_{k,i}^T \in \mathbb{R}^{l_k, \times l_k}$$
(26a)

$$\Lambda_{i,i+1} = D_{k,i} \Phi_{k,i}^{-1} C_{k,i}^T \in \mathbb{R}^{l_k \times l_k}$$

$$\tag{26b}$$

Accordingly, all matrices in (24) are banded except  $\Lambda_k^{-1}$  which in general is dense.

#### 3.2 Decaying of the dual Hessian

Let us first assume that the problem data is bounded. In other words, let us assume the following:

**Assumption 1** The row and column absolute sums of Jacobians of equality constraints (i.e. (6b) and (6c)) are bounded. Hence:

- 1.  $||C_{k,i}||_{\infty} \leq \gamma$  and  $||C_{k,i}||_{1} \leq \gamma$ 2.  $||D_{k,i}||_{\infty} \leq \gamma$  and  $||D_{k,i}||_{1} \leq \gamma$ 3.  $||F_{k,i}||_{\infty} \leq \gamma$  and  $||F_{k,i}||_{1} \leq \gamma$

It should be observed that Assumption 1 is not by any means restrictive, since any solver would struggle with a problem where it is not fulfilled.

Possible problems with interior-point methods are related to numerical difficulties due to ill-conditioning []. This might occur since elements in  $S_k^{-1}Y_k$ can be close to zero in late iterations when  $\tau$  is small. Our experience is however that this is not a major issue in practice, which is also supported by commercial interior-point implementations []. Moreover, since we assume that (5) is strongly convex, the eigenvalues of  $\Phi_k$  is lower bounded by the smallest eigenvalue of  $H_k$ , even when  $S_k^{-1}Y_k$  is singular. According to this reasoning it is not unrealistic to assume boundedness of  $\Phi_k^{-1}$  which we do next: **Assumption 2** The row and column absolute sums of  $\Phi_k^{-1}$  are bounded. Hence:

$$\|\Phi_{k,i}^{-1}\|_{\infty} = \|\Phi_{k,i}^{-1}\|_{1} \le \gamma_{\Phi_{k}^{-1}}, \quad \forall i$$
 (27)

Inverses of sparse matrices are in general dense, but individual elements are often small in absolute value. Since  $\Lambda_k$  is banded, symmetric and positive definite, we will relay our analysis on the following classical result:

**Lemma 1** If A is Hermitian positive definite and m-banded ( $[A]_{ij} = 0$  if |i-j| > m), the entries of  $A^{-1}$  satisfy the following bound:

$$|[A^{-1}]_{ij}| < K\omega^{|i-j|}, \quad \forall i, j$$

$$(28)$$

where [a,b] is the smallest interval containing the spectrum  $\sigma(A)$  of A,  $K = \max\{a^{-1}, K_0\}$ ,  $K_0 = (1 + \sqrt{\kappa})$ ,  $\omega = q^{1/m}$ ,  $q = \frac{\sqrt{\kappa} - 1}{\sqrt{\kappa} + 1}$ ,  $\kappa = \frac{b}{a}$ .

Proof A proof is given in 
$$[]$$
.

This means that the entries of  $A^{-1}$  are bounded by an exponentially decaying function along each row or column. However, the bound depends on the condition number and the bandwidth of the matrix. Matrices with a high condition number and/or a high bandwidth can accordingly result in a large K and  $\omega \approx 1$ , leading to a slow decay. The opposite, i.e. a low condition number and a small band, would result in a rapid decay.

Observe that due to the block tridiagonal structure of  $\Lambda_k$ , it is  $3l_k$ -banded. Moreover, if we introduce the notation

$$\Lambda_k^{-1} = \begin{bmatrix}
T_{11} & T_{21}^T & \dots & T_{N+1,1}^T \\
T_{21} & T_{22} & \dots & T_{N+1,2}^T \\
\vdots & \vdots & \ddots & \vdots \\
T_{N+1,1} & T_{N+1,2} & \dots & T_{N+1,N+1}
\end{bmatrix}$$
(29)

where  $T_{i,j} \in \mathbb{R}^{l_k \times l_k}$ , we can establish the following proposition:

**Proposition 1** The off-diagonal blocks (i.e.  $T_{i,j}$  where i - j > 0) in  $\Lambda_k^{-1}$  satisfy the following bounds:

$$||T_{i,j}||_{\bullet} \le K_{\Lambda_k} \omega_{\Lambda_k}^{i-j} \tag{30a}$$

(30b)

where  $\bullet$  represents  $\infty$  and 1,  $\sigma_{\min}(\Lambda_k)$  and  $\kappa_{\Lambda_k}$  are the the smallest singular value and the condition number of  $\Lambda_k$  respectively,  $K_{\Lambda_k} = \max\{\sigma_{\min}(\Lambda_k)^{-1}, 1 + \sqrt{\kappa_{\Lambda_k}}\}l_k\omega_{\Lambda_k}^{1/l_k}$  and  $\omega_{\Lambda_k} = \left(\frac{\sqrt{\kappa_{\Lambda_k}}-1}{\sqrt{\kappa_{\Lambda_k}}+1}\right)^{\frac{1}{3}}$ .

*Proof* According to Lemma 1, the element in  $T_{i,j}$  with the largest bound is located in the top-right corner, and is hence element  $[\Lambda_k^{-1}]_{il_k+1,jl_k}$ . By directly applying Lemma 1 it follows that:

$$\max |[T_{ij}]| \le \max \{\sigma_{min}(\Lambda_k)^{-1}, 1 + \sqrt{\kappa_{\Lambda_k}}\} \left(\frac{\sqrt{\kappa_{\Lambda_k}} - 1}{\sqrt{\kappa_{\Lambda_k}} + 1}\right)^{\frac{1}{3l_k}((i-j)l_k + 1)}$$
(31)

where  $\max |[T_{ij}]|$  refers to the maximum absolute value of the components in  $T_{ij}$ . Moreover, since there are  $l_k$  elements in each row or column of a block  $T_{i,j}$ , we obtain the bounds given in (30).

The constants  $K_{\Lambda_k}$  and  $\omega_{\Lambda_k}$  in Proposition 1 depends heavily on the conditioning of  $\Lambda_k$ , which indeed for interior point methods in general can be very high for small values of  $\tau$ . It should however be understood that strong convexity of (5), should improve the worst case conditioning of  $\Phi_k$  and hence also of  $\Lambda_k$ .

To gain some perspective, we have concluded that  $\Lambda_k^{-1}$  is decaying exponentially towards the off-diagonal corners and that all other matrices in (24) are banded. This suggests that also  $\nabla_{\lambda\lambda}^2 d(\lambda,\tau)$  should decay towards the off-diagonal corners. To maintain a simple reasoning, one more stepping stone will be used before arriving at the main results of this section. Therefore, let us introduce the notation:

$$C_k^T \Lambda_k^{-1} C_k = \begin{bmatrix} V_{11} & V_{21}^T & \dots & V_{N+1,1}^T \\ V_{21} & V_{22} & \dots & V_{N+1,2}^T \\ \vdots & \vdots & \ddots & \vdots \\ V_{N+1,1} & V_{N+1,2} & \dots & V_{N+1,N+1} \end{bmatrix}$$
(32)

where  $V_{ij} \in \mathbb{R}^{l_k \times l_k}$ , and look at the decay of  $C_k^T \Lambda_k^{-1} C_k$ .

**Proposition 2** The off diagonal blocks (i.e.  $V_{i,j}$  where i-j>0) of  $C_k^T \Lambda_k^{-1} C_k$  satisfy the following bounds:

$$||V_{i,j}||_{\bullet} \le \gamma^2 \bar{K}_{\Lambda_k} \omega_{\Lambda_k}^{i-j} \tag{33a}$$

(33b)

where  $\bar{K}_{\Lambda_k} = (2 + \omega_{\Lambda_k} + \omega_{\Lambda_k}^{-1})\omega_{\Lambda_k}^{-1}K_{\Lambda_k}$  and  $\bullet$  represents  $\infty$  and 1.

Let us now introduce the notation:

$$\nabla^2 d(\lambda, \tau) = \begin{bmatrix} W_{11} & W_{21}^T & \dots & W_{N+1,1}^T \\ W_{21} & W_{22} & \dots & W_{N+1,2}^T \\ \vdots & \vdots & \ddots & \vdots \\ W_{N+1,1} & W_{N+1,2} & \dots & W_{N+1,N+1} \end{bmatrix}$$
(34)

where  $W_{i,j} \in \mathbb{R}^{r_i \times r_j}$ . We can now establish a decay towards the off-diagonal corners of  $\nabla^2 d(\lambda, \tau)$ .

**Proposition 3** The off-diagonal blocks (i.e.  $W_{i,j}$  where i-j > 0) of  $\nabla^2 d(\lambda, \tau)$  satisfy the following bounds:

$$||W_{i,j}||_{\bullet} \le \sum_{k=1}^{M} \gamma^4 \gamma_{\Phi_k^{-1}}^2 \bar{K}_{\Lambda_k} \omega_{\Lambda_k}^{i-j}$$
 (35a)

(35b)

where  $\bullet$  represents  $\infty$  and 1.

*Proof* To be done.

Moreover, let us continue by finding a bound on the euclidean distance between  $\nabla^2 d(\lambda, \tau)$  and a band along its diagonal. To do so, we start by recalling Gershgorin's circle theorem:

**Theorem 1** For  $A \in \mathbb{R}^{n \times n}$  with elements  $a_{ij}$ , let  $R_i = \sum_{j \neq i} |a_{ij}|$  be the sum of the absolute values of the non-diagonal entries in the row i. Let  $D(a_{ii}, R_i)$  be the closed disc centered in  $a_{ii}$  with radius  $R_i$ , then every eigenvalue of A lies within at least one of the discs  $D(a_{ii}, R_i)$ .

Proof A proof is given in 
$$[]$$
.

Finally, we can establish our main result:

**Lemma 2** Let  $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}}$  represent the diagonal  $\mathcal{M}$ -block band of  $\nabla^2 d(\lambda, \tau)$ . The following bound holds:

$$\|\lfloor \nabla^2 d(\lambda, \tau) \rceil_{\mathcal{M}} - \nabla^2 d(\lambda, \tau) \|_2 \le (N - \frac{\mathcal{M} - 1}{2}) \sum_{k=1}^{M} \gamma^4 \gamma_{\varPhi_k^{-1}}^2 \tilde{K}_{\Lambda_k} \omega_{\Lambda_k}^{\frac{\mathcal{M} - 1}{2}}$$
(36)

where 
$$\tilde{K}_{\Lambda_k} = \bar{K}_{\Lambda_k} \omega_{\Lambda_k} = (2 + \omega_{\Lambda_k} + \omega_{\Lambda_k}^{-1}) K_{\Lambda_k}$$
.

Proof Since  $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)$  is symmetric, its singular values collapse into the absolute value of its eigenvalues. The problem of finding a bound on the 2-norm is hence reduced to the problem of bounding the magnitude of the largest eigenvalue.

First, observe that all diagonal elements of  $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)$  are zero. This means that every Gershgorin disc will be centered in the origin, and a bound on the largest eigenvalue can be found by finding the largest radius of a Gershgorin disc.

The blocks  $W_{i,j}$  in  $[\nabla^2 d(\lambda,\tau)]_{\mathcal{M}} - \nabla^2 d(\lambda,\tau)$  with the largest bound is located next to the diagonal band of zeros, i.e. where  $i-j=\frac{\mathcal{M}-1}{2}+1$ . Furthermore, there is at least  $N-\frac{\mathcal{M}-1}{2}$  nonzero blocks at each block row. Each block is upper bounded by Proposition X. This means that we can establish a bound on the radius of the Gershgorin discs, and hence the following:

$$\|\lfloor \nabla^2 d(\lambda, \tau) \rceil_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)\|_2 \le (N - \frac{\mathcal{M} - 1}{2}) \sum_{k=1}^{M} \gamma^4 \gamma_{\bar{\Phi}_k^{-1}}^2 \bar{K}_{\Lambda_k} \omega_{\Lambda_k}^{\frac{\mathcal{M} - 1}{2} + 1}$$
(37)

Equation (36) is then directly obtained by introducing  $\tilde{K}_{\Lambda_k} = \bar{K}_{\Lambda_k} \omega_{\Lambda_k}$ .

Say something about the principal behavior of the bound...

It is well known that the use of an inexact Hessian will degrade the convergence of Newton's method to a linear rate []. However, let us still look at the relative error in the Newton direction that we get from using  $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}}$  instead of  $\nabla^2 d(\lambda, \tau)$ . First, recall the following lemma:

**Lemma 3** If x is a solution to Ax = b and  $\hat{x}$  is a solution to the pertubed system  $(A + F)\hat{x} = b + f$ , then the following bound holds:

$$\frac{\|\hat{x} - x\|_2}{\|x\|_2} \le \kappa(A)(\rho_A - \rho_b) \tag{38}$$

where  $\kappa(A)$  represents the condition number of A,  $\rho_A = \frac{\|F\|_2}{\|A\|_2}$  and  $\rho_b = \frac{\|f\|_2}{\|b\|_2}$ .

*Proof* A proof is given in [].

Accordingly, by combining Lemma 2 and Lemma 3 we can establish the following:

**Lemma 4** If  $\Delta \hat{\lambda}$  is a solution to  $[\nabla^2 d(\lambda, \tau)]_{\mathcal{M}} \Delta \hat{\lambda} = -\nabla d(\lambda, \tau)$ , the following bound on the relative error compared to the true Newton direction  $\Delta \lambda$  holds:

$$\frac{\|\Delta \hat{\lambda} - \Delta \lambda\|_{2}}{\|\Delta \lambda\|_{2}} \le \|\nabla^{2} d(\lambda, \tau)^{-1}\|_{2} (N - \frac{\mathcal{M} - 1}{2}) \sum_{k=1}^{M} \gamma^{4} \gamma_{\Phi_{k}^{-1}}^{2} \tilde{K}_{\Lambda_{k}} \omega_{\Lambda_{k}}^{\frac{\mathcal{M} - 1}{2}}$$
(39)

*Proof* We can view  $\Delta \hat{\lambda}$  as a solution to a perturbed system:

$$\left(\nabla^2 d(\lambda, \tau) + \left(\lfloor \nabla^2 d(\lambda, \tau) \right\rceil_{\mathcal{M}} - \nabla^2 d(\lambda, \tau)\right) \Delta \hat{\lambda} = -\nabla d(\lambda, \tau) \tag{40}$$

Equation (39) follows then directly from Lemma 2 and Lemma 3.  $\Box$ 

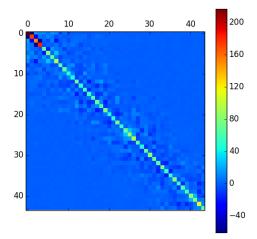
#### 4 Numerical experiments

4.1 Decaying of the dual Hessian

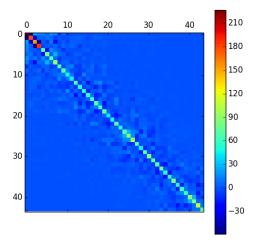
## 4.1.1 A small problem

As an example of a small problem, we use a randomly generated problem with 4 subproblems (i.e. M=4), where each subproblem has N=10, 6 states and 4 controls and hence  $n_k=10$ , and each time instance has 4 coupling constraints.

According to Section 3.2, we expect a slower decay when the conditioning of  $S_k^{-1}Y_k$  gets worse, i.e. when we reduce  $\tau$ . Surprisingly, this is something that does not seem to have a big influence in practice. To illustrate this, Figure 2 shows the dual Hessian for the same problem as in Figure 1, but with  $\tau = 10^{-5}$ .



**Fig. 1**  $\tau = 1$ 



**Fig. 2**  $\tau = 10^{-5}$ 

# 4.1.2 A large problem

# 4.2 Newton steps with inexact Hessian

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