Neutrino Oscillations in Matter

PhD Candidacy Exam

Lei Ma **Supervisor**: Huaiyu Duan

> Department of Physics UNM

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OUTLINE

- Introduction
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 Neutrino Oscillations
 Why Oscillations
- Matter Effect
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- 3. Understanding Stimulated Oscillations Hamiltonian, and Basis Single Frequency Matter Profile
- 4. Summary & Future Work

OVERVIEW

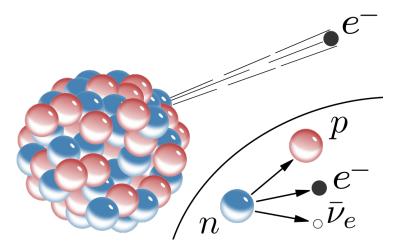
Introduction
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Understanding Stimulated Oscillations

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WHAT ARE NEUTRINOS?



Beta decay and antineutrino production. Source: Beta_Decay@Wikipedia

WHAT ARE NEUTRINOS?

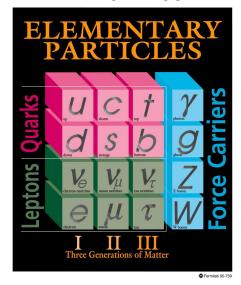
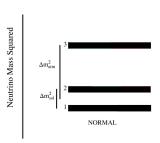


Table of elementary particles. Source: Fermilab

Neutrinos are

- ► fermions.
- electrically neutral,
- ► light.



Adapted from Olga Mena & Stephen Parke (2004)

WHAT ARE NEUTRINOS?

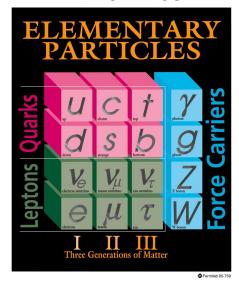
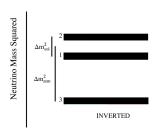


Table of elementary particles. Source: Fermilab

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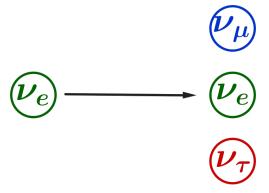
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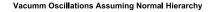
WHAT IS NEUTRINO OSCILLATION?

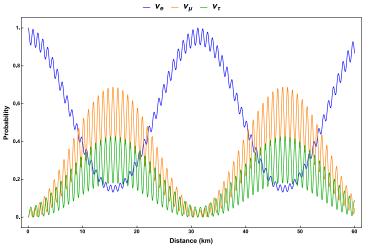




Neutrino Oscillations

WHAT IS NEUTRINO OSCILLATION?





Probabilities of finding neutrinos to be in each flavors.

WHY DO NEUTRINOS OSCILLATE?

Equation of Motion

$$i\partial_x\ket{\Psi}=\hat{\mathbf{H}}\ket{\Psi}$$

► Basis: Hamiltonian diagonalized basis/mass basis/propagation basis, {|\(\nu_1\), |\(\nu_2\)\}.

Þ

$$H=-rac{\omega_{
m v}}{2}\sigma_{3}, \qquad ext{where} \ \omega_{
m v}=rac{\delta m^{2}}{2E}=rac{m_{2}^{2}-m_{1}^{2}}{2E}.$$

► The system can be solved given initial condition of the amplitudes of the two eigenstates $(\langle \nu_1 | \Psi(0) \rangle, \langle \nu_2 | \Psi(0) \rangle)^T$,

$$\begin{pmatrix} \langle \nu_1 | \Psi(x) \rangle \\ \langle \nu_2 | \Psi(x) \rangle \end{pmatrix} = \begin{pmatrix} \langle \nu_1 | \Psi(0) \rangle \exp{(i\omega_v x/2)} \\ \langle \nu_2 | \Psi(0) \rangle \exp{(-i\omega_v x/2)} \end{pmatrix}$$

WHY DO NEUTRINOS OSCILLATE?

Flavor basis

Neutrino wave function in flavor basis $\{|\nu_{\rm e}\rangle\,, |\nu_{\mu}\rangle\}$ is related to state in energy basis $\{|\nu_{\rm 1}\rangle\,, |\nu_{\rm 2}\rangle\}$ through

$$\begin{pmatrix} |\nu_{\rm e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} = \begin{pmatrix} \cos\theta_{\rm v} & \sin\theta_{\rm v} \\ -\sin\theta_{\rm v} & \cos\theta_{\rm v} \end{pmatrix} \begin{pmatrix} |\nu_{1}\rangle \\ |\nu_{2}\rangle \end{pmatrix}$$

 $\theta_v\text{:}$ vacuum mixing angle

WHY DO NEUTRINOS OSCILLATE?

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 $\theta_{\rm v}$: vacuum mixing angle

Hamiltonian H

Mass basis

$$\begin{split} \frac{\omega_{v}}{2} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} & \frac{\omega_{v}}{2} \begin{pmatrix} -\cos 2\theta_{v} & \sin 2\theta_{v} \\ \sin 2\theta_{v} & \cos 2\theta_{v} \end{pmatrix} \\ = -\frac{\omega_{v}}{2} \boldsymbol{\sigma}_{3} & = \frac{\omega_{v}}{2} \left(-\cos 2\theta_{v} \boldsymbol{\sigma}_{3} + \sin 2\theta_{v} \boldsymbol{\sigma}_{1} \right) \end{split}$$

NATURE OF NEUTRINO OSCILLATION

Transition Probability

$$P(|\nu_{\rm e}\rangle \rightarrow |\nu_{\mu}\rangle) = \sin^2(2\theta_{\rm v})\sin^2(\omega_{\rm v}x/2)$$

- $\omega_{\rm v} = (m_2^2 m_1^2)/2E$ determines oscillation wavelength.
- ▶ Mixing angle θ_v determines flavor oscillation amplitude.

OVERVIEW

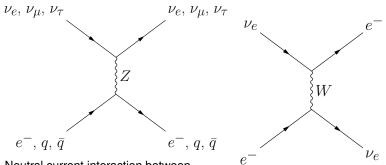
Introduction

Matter Effect
Matter Interaction
MSW Effect
Solar Neutrino Problem
Stimulated Neutrino Oscillations

Understanding Stimulated Oscillations

Summary & Future Work

MATTER INTERACTION



Neutral current interaction between $\nu_{\rm e},\, \nu_{\mu},\, \nu_{\tau},\,$ and e^- , quarks and antiquarks.

Charged current interaction between $\nu_{\rm e}$ and e^-

MATTER INTERACTION

Hamiltonian with matter interaction in flavor basis ($\omega_{\rm v} = \delta m^2/2E$):

$$\mathbf{H} = \begin{array}{cc} \frac{\omega_{\mathrm{v}}}{2} \begin{pmatrix} -\cos 2\theta_{\mathrm{v}} & \sin 2\theta_{\mathrm{v}} \\ \sin 2\theta_{\mathrm{v}} & \cos 2\theta_{\mathrm{v}} \end{pmatrix} & \pm \sqrt{2}G_{\mathrm{F}}n_{\mathrm{e}}(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

- ► Vacuum Hamiltonian
- ▶ Matter interaction

MATTER INTERACTION

Hamiltonian with matter interaction in flavor basis ($\omega_{\rm v} = \delta m^2/2E$):

$$\mathbf{H} = \frac{\omega_{\mathbf{v}}}{2} \left(-\cos 2\theta_{\mathbf{v}} \boldsymbol{\sigma}_{3} + \sin 2\theta_{\mathbf{v}} \boldsymbol{\sigma}_{1} \right) + \frac{\lambda(x)}{2} \boldsymbol{\sigma}_{3}$$

- ► Vacuum Hamiltonian
- ► Matter interaction
- $\lambda(x) = \sqrt{2}G_{\rm F}n_{\rm e}(x)$

MSW EFFECT

Hamiltonian in Vacuum

$$\mathbf{H}_{ ext{vacuum}} = rac{\omega_{ ext{v}}\cos 2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_3 + rac{\omega_{ ext{v}}\sin 2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_1$$

$$egin{aligned} \mathbf{H} &= rac{\lambda(x) - \omega_{ ext{v}}\cos2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_{3} + rac{\omega_{ ext{v}}\sin2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_{1} \ &= rac{\omega_{ ext{m}}(x)\cos2 heta_{ ext{m}}(x)}{2}oldsymbol{\sigma}_{3} + rac{\omega_{ ext{m}}(x)\sin2 heta_{ ext{m}}(x)}{2}oldsymbol{\sigma}_{1}, \end{aligned}$$

where

$$\begin{split} \omega_{\rm m}(x) &= \sqrt{\left(\lambda(x) - \omega_{\rm v}\cos 2\theta_{\rm v}\right)^2 + \omega_{\rm v}^2\sin^2 2\theta_{\rm v}},\\ \tan 2\theta_{\rm m}(x) &= \frac{\omega_{\rm v}\sin 2\theta_{\rm v}}{\omega_{\rm v}\cos 2\theta_{\rm v} - \lambda(x)}. \end{split}$$

MSW EFFECT

Constant matter profile λ_0 as an example,

Significance of $\theta_{\rm m}$

Define matter basis $\{|\nu_L\rangle, |\nu_H\rangle\}$

$$\begin{pmatrix} |\nu_{\rm e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} = \begin{pmatrix} \cos\theta_{\rm m} & \sin\theta_{\rm m} \\ -\sin\theta_{\rm m} & \cos\theta_{\rm m} \end{pmatrix} \begin{pmatrix} |\nu_{\rm L}\rangle \\ |\nu_{\rm H}\rangle \end{pmatrix}$$

In matter basis

$$\mathbf{H}_{\mathsf{matter ext{-}basis}} = -rac{\omega_{\mathsf{m}}}{2}oldsymbol{\sigma_3}$$

MSW RESONANCE

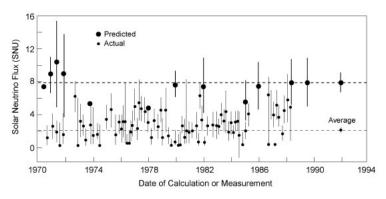
Hamiltonian with Matter Potential

$$\begin{split} \mathbf{H} &= \frac{\lambda(x) - \omega_{v} \cos 2\theta_{v}}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{v} \sin 2\theta_{v}}{2} \boldsymbol{\sigma}_{1} \\ &= \frac{\omega_{m}(x) \cos 2\theta_{m}(x)}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{m}(x) \sin 2\theta_{m}(x)}{2} \boldsymbol{\sigma}_{1} \\ &\tan 2\theta_{m}(x) = \frac{\omega_{v} \sin 2\theta_{v}}{\omega_{v} \cos 2\theta_{v} - \lambda(x)}. \\ \begin{pmatrix} |\nu_{e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} &= \begin{pmatrix} \cos \theta_{m} & \sin \theta_{m} \\ -\sin \theta_{m} & \cos \theta_{m} \end{pmatrix} \begin{pmatrix} |\nu_{L}\rangle \\ |\nu_{H}\rangle \end{pmatrix} \end{split}$$

Transition Probability

$$P(|\nu_{\rm e}\rangle \to |\nu_{\mu}\rangle) = \sin^2(2\theta_{\rm m})\sin^2(\omega_{\rm m}x)$$

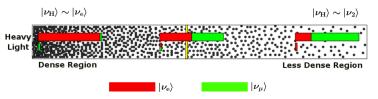
SOLAR NEUTRINO PROBLEM



Chlorine detector (Homestake experiment) results and theory predictions. SNU: 1 event for 10^{36} target atoms per second. Kenneth R. Lang (2010)

MSW EFFECT AND SOLAR NEUTRINOS

$$\begin{split} \mathbf{H} &= \frac{\lambda(\mathbf{x}) - \omega_{\mathrm{v}} \cos 2\theta_{\mathrm{v}}}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{\mathrm{v}} \sin 2\theta_{\mathrm{v}}}{2} \boldsymbol{\sigma}_{1} \\ \begin{pmatrix} |\nu_{\mathrm{L}}\rangle \\ |\nu_{\mathrm{H}}\rangle \end{pmatrix} &= \begin{pmatrix} \cos \theta_{\mathrm{m}} & -\sin \theta_{\mathrm{m}} \\ \sin \theta_{\mathrm{m}} & \cos \theta_{\mathrm{m}} \end{pmatrix} \begin{pmatrix} |\nu_{\mathrm{e}}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} \\ \mathbf{H}_{\mathrm{matter-basis}} &= -\frac{\omega_{\mathrm{m}}}{2} \boldsymbol{\sigma}_{3} \end{split}$$



Yellow bar is the resonance point. Red: $|\nu_e\rangle$. Green: $|\nu_{\mu}\rangle$. Adapted from Smirnov, 2003.

MSW EFFECT

Suppose
$$\omega_{\mathrm{v}} = (m_2^2 - m_1^2)/2E < 0,$$

$$\mathbf{H} = \begin{bmatrix} -\frac{\omega_{\mathrm{v}}}{2} \begin{pmatrix} -\cos 2\theta_{\mathrm{v}} & \sin 2\theta_{\mathrm{v}} \\ \sin 2\theta_{\mathrm{v}} & \cos 2\theta_{\mathrm{v}} \end{pmatrix} & +\sqrt{2}G_{\mathrm{F}}n_{\mathrm{e}}(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \end{bmatrix}$$

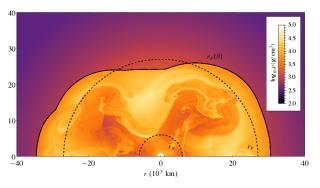
$$\downarrow$$

$$\mathbf{H} = \begin{pmatrix} -\frac{\omega_{\mathrm{v}}}{2} \cos 2\theta_{\mathrm{v}} + \frac{\lambda(x)}{2} \end{pmatrix} \boldsymbol{\sigma}_3 - \frac{\omega_{\mathrm{v}}}{2} \sin 2\theta_{\mathrm{v}} \boldsymbol{\sigma}_1$$

SUPERNOVA MATTER DENSITY PROFILE

Why Do We Care

Astrophysical environments: supernovae, accretion disks etc



Supernova shock and turbulence. E. Borriello, et al (2014)

$$\Delta n_e(r) = \sum_n c_n \sin(k_n r + \phi_n)$$

STIMULATED NEUTRINO OSCILLATIONS

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x)$$

Basis

Background matter basis: Hamiltonian is diagonalized with only background matter profile λ_0 ,

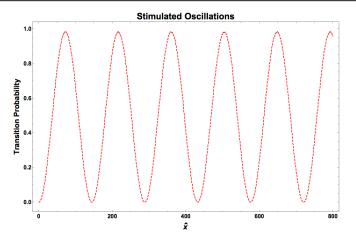
$$H_{background} = -\frac{\omega_m}{2} \sigma_3.$$

Hamiltonian

$$H = \frac{1}{2} \left(-\omega_{m} + \frac{\delta \lambda(x)}{\lambda(x)} \cos 2\theta_{m} \right) \sigma_{3} - \frac{\delta \lambda(x)}{2} \sin \theta_{m} \sigma_{1}.$$

STIMULATED NEUTRINO OSCILLATIONS

P. Krastev and A. Smirnov (1989); J. Kneller et al (2013); K. Patton et al (2014);



Stimulated oscillations. $\lambda(x) = \lambda_0 + A\sin(kx)$ with $\hat{x} = \omega_m x$, $A = 0.1\omega_m$, $k = 0.995\omega_m$, $\theta_m = \pi/6$

OVERVIEW

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Understanding Stimulated Oscillations Hamiltonian, and Basis Single Frequency Matter Profile

Summary & Future Work

Understanding Stimulated Oscillations

Matter profile

$$\lambda(x) = \lambda_0 + A\sin(kx),$$

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} \left(-\omega_{m} + \frac{\delta \lambda(\mathbf{x})}{\delta \lambda(\mathbf{x})} \cos 2\theta_{m} \right) \boldsymbol{\sigma}_{3} - \frac{\frac{\delta \lambda(\mathbf{x})}{2}}{2} \sin \theta_{m} \boldsymbol{\sigma}_{1}.$$

A Better Basis

Define new basis $\{|\tilde{\nu}_L\rangle,|\tilde{\nu}_H\rangle\}$ is related to background matter basis $\{|\nu_L\rangle,|\nu_H\rangle\}$ through

$$\begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix} = \begin{pmatrix} e^{-i\eta(x)} & 0 \\ 0 & e^{i\eta(x)} \end{pmatrix} \begin{pmatrix} |\tilde{\nu}_L\rangle \\ |\tilde{\nu}_H\rangle \end{pmatrix},$$

where

$$\eta(x) - \eta(0) = -\frac{\omega_{\rm m}}{2}x + \frac{\cos 2\theta_{\rm m}}{2} \int_0^x \frac{\delta \lambda(\tau)}{d\tau}.$$

Hamiltonian in new basis

$$\widetilde{\mathbf{H}} = -\frac{\delta \lambda(\mathbf{x})}{2} \sin 2\theta_{\mathrm{m}} \begin{pmatrix} 0 & e^{2i\eta(\mathbf{x})} \\ e^{-2i\eta(\mathbf{x})} & 0 \end{pmatrix} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

Hamiltonian in New Basis

$$\begin{split} h &\equiv -\frac{\delta\lambda(x)}{2}e^{2i\eta(x)} \\ &= \frac{i}{4}\left[\exp\left(i(k+\omega_{\rm m})x + i\cos2\theta_{\rm m}\frac{A}{k}\cos(kx)\right) \\ &-\exp\left(i(-k+\omega_{\rm m})x + i\cos2\theta_{\rm m}\frac{A}{k}\cos(kx)\right)\right] \end{split}$$

RABI OSCILLATION

Rabi Oscillation

Hamiltonian

$$\begin{pmatrix} -\omega_0/2 & \alpha\omega_0e^{i\omega x} \\ \alpha\omega_0e^{-ikx} & \omega_0/2 \end{pmatrix},$$

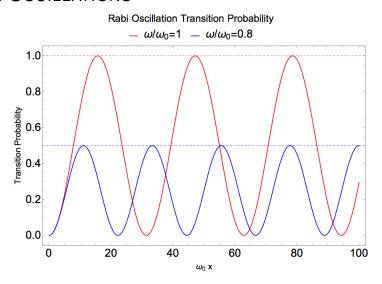
$$E_2 = \frac{\omega_0}{2}$$

Incoming light

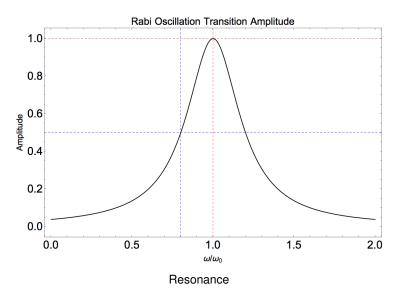
$$E_1 = -\frac{\omega_0}{2}$$

Frequency : ω

RABI OSCILLATIONS



RABI OSCILLATIONS



Off-diagonal Term in Our System

$$\widetilde{\mathbf{H}} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

$$h \propto \left[\exp \left(i(k + \omega_{\rm m})x + i\cos 2\theta_{\rm m} \frac{A}{k}\cos(kx) \right) - \exp \left(i(-k + \omega_{\rm m})x + i\cos 2\theta_{\rm m} \frac{A}{k}\cos(kx) \right) \right]$$

Jacobi-Anger expansion

$$e^{i\beta\cos(kx)} = \sum_{n=-\infty}^{\infty} i^n J_n(\beta) e^{inkx},$$

where $J_n(\beta)$ are Bessel's functions of the first kind.

Scaled Quantities

Characteristic scale: $\omega_{\rm m}$

- $\rightarrow \hat{A} = A/\omega_{\rm m}$
- $\hat{k} = k/\omega_{\rm m}$
- $\hat{\mathbf{x}} = \omega_{\mathrm{m}} \mathbf{x}$
- $\blacktriangleright \hat{h} = h/\omega_{\rm m}$

Rotation Wave Approximation

The off-diagonal element of Hamiltonian

$$\widetilde{\mathbf{H}} = \sum_{n = -\infty}^{\infty} \begin{pmatrix} 0 & \frac{1}{2} \hat{B}_n e^{i(n\hat{k} - 1)\hat{x}} \\ \frac{1}{2} \hat{B}_n^* e^{-i(n\hat{k} - 1)\hat{x}} & 0 \end{pmatrix}$$

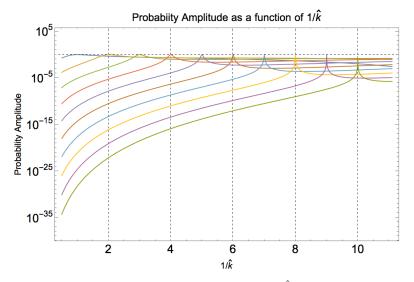
where $\hat{B}_n = -(-i)^n n\hat{k} \tan 2\theta_{\rm m} J_n(\hat{A}\cos 2\theta_{\rm m}/\hat{k})$.

Transition Probability

$$P_{ ext{L}
ightarrow ext{H}}^{(n)} = rac{\left|\left|\hat{B}_{n}\right|/2
ight|^{2}}{\left|\left|\hat{B}_{n}\right|/2
ight|^{2} + (n\hat{k} - 1)^{2}} \sin^{2}\left(rac{q^{(n)}}{2}x
ight),$$

where

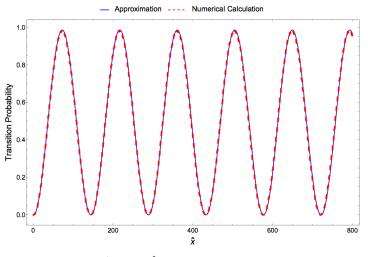
$$q^{(n)}=\sqrt{\left|\Gamma^{(n)}/2\right|^2+(n\hat{k}-1)^2},\quad ext{frequency of oscillations} \ \Gamma^{(n)}=\left|\hat{B}_n\right|,\quad ext{width of resonance }(n\hat{k} ext{ as parameter})$$



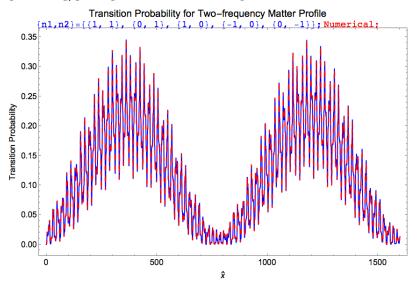
Resonances of different $n = 1/\hat{k}$.

SINGLE FREQUENCY MATTER PROFILE

Stimulated Oscillations



$$\hat{A} = 0.1, \, \hat{k} = 0.995, \, \theta_{\rm m} = \pi/6$$



$$\lambda(x) = \lambda_0 + A_1 \sin(k_1 x) + A_2 \sin(k_2 x)$$
. $\hat{k}_1 = 0.3$, $\hat{k}_2 = 0.7$, $A_1 = A_2 = 0.1$, $\theta_m = \pi/5$.

OVERVIEW

Introduction

Matter Effect

Understanding Stimulated Oscillations

Summary & Future Work

SUMMARY & FUTURE WORK

- The fact that neutrino flavor sates are not mass states causes vacuum oscillations.
- ► MSW resonance happens when matter potential cancels out the vacuum diagonal elements of the Hamiltonian.
- ► Even matter profile doesn't match MSW requirement, variation in matter profile can cause resonances.
- Single frequency perturbations in matter profile is a combination of many Rabi oscillations.
- How to understand and calculate systems with multi-frequency matter profile (turbulence).
- ► Combine periodic or even turbulent matter profile with neutrino self-interaction.

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BACKUP SLIDES

BACKUP SLIDES

PARAMETERS USED FOR VACUUM OSCILLATIONS

$$\begin{array}{l} \theta_{12}=33.36/180\pi;\,\theta_{13}=8.66/180\pi;\,\theta_{23}=40/180*\pi;\,\delta_{cp}=0;\\ m_1^2=0.01;\,m_2^2=m_1^2+0.000079;\,E=1\text{MeV} \end{array}$$

SINGLE FREQUENCY MATTER PROFILE

Why Does It Work?

$$J_n(n \operatorname{sech} \alpha) \sim \frac{e^{-n(\alpha - \tanh \alpha)}}{\sqrt{2\pi n \tanh \alpha}}, \quad ext{for large } n$$

 \Rightarrow

$$\Gamma \propto \hat{B}_n \propto rac{e^{-n(lpha - anh lpha)}}{\sqrt{2\pi n anh lpha}}$$

Small perturbation \Rightarrow Small $\hat{A} \Rightarrow$ Large $\alpha \Rightarrow$ Drops fast at large n.

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x), \quad \delta\lambda(x) = A_1 \sin(k_1 x) + A_2 \sin(k_2 x).$$

TWO-FREQUENCY MATTER PROFILE $\hat{h} = \sum_{i=1}^{\infty} \frac{1}{2} \hat{B}_{n} e^{i(n\hat{k}-1)\hat{x}}$,

Hamiltonian Off-diagonal Element

Apply Jacobi-Anger expansion,

$$\hat{h} = \sum_{n_1 = -\infty}^{\infty} \sum_{n_2 = -\infty}^{\infty} \frac{1}{2} \hat{B}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1)\hat{x}},$$

where

$$\hat{B}_{n_1,n_2}(\hat{k}_1,\hat{k}_2)
= -(-i)^{n_1+n_2}(n_1\hat{k}_1 + n_2\hat{k}_2)J_{n_1}\left(\frac{\hat{A}_1\cos 2\theta_{\rm m}}{\hat{k}_1}\right)J_{n_2}\left(\frac{\hat{A}_2\cos 2\theta_{\rm m}}{\hat{k}_2}\right)$$

Which terms are important?

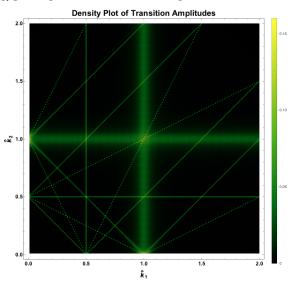
Resonance Lines

There are still resonances, i.e., (almost) zero phases, on lines

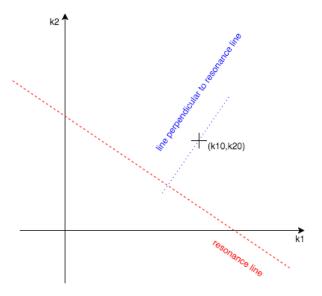
$$n_{1,0}\hat{k}_1 + n_{2,0}\hat{k}_2 - 1 = 0$$

in $\{\hat{k}_1, \hat{k}_2\}$ plane. \Rightarrow Resonance width for each point on resonance lines.

$\text{TWO-FREQUENCY MATTER PROFIL}^{\hat{h}} \bar{\bar{E}}^{\sum_{n_1} \sum_{n_2} \frac{1}{2} \hat{h}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1) \hat{x}},$



Density plot of transition amplitudes calculated using only one term out of the whole summation in Hamiltonian. $n_1, n_2 \in [-2, 2]$



Resonance line, distance to resonance, and width

Width

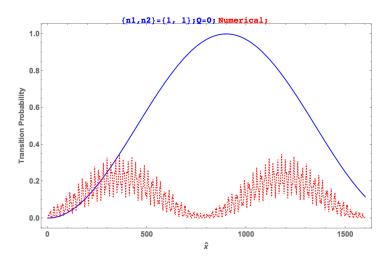
$$\Gamma_2 = rac{\hat{B}_{n_1,n_2}(\hat{k}_{1,\mathrm{intercept}},\hat{k}_{2,\mathrm{intercept}})}{\sqrt{n_1^2 + n_2^2}}.$$

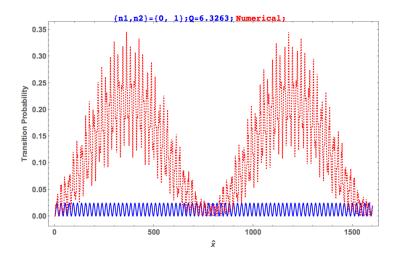
Distance to Resonance Line

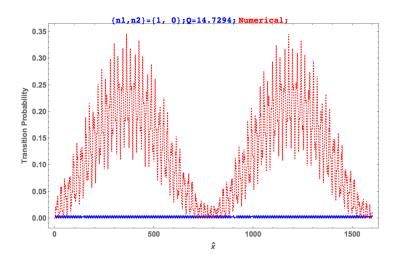
$$d = \frac{|n_1\hat{k}_{10} + n_2\hat{k}_{20} - 1|}{\sqrt{n_1^2 + n_2^2}}.$$

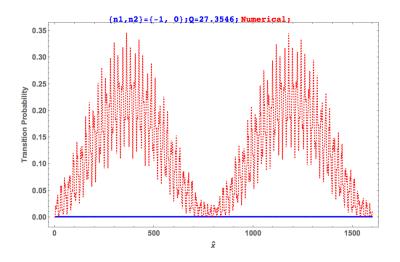
Distance to Resonance Width Ratio

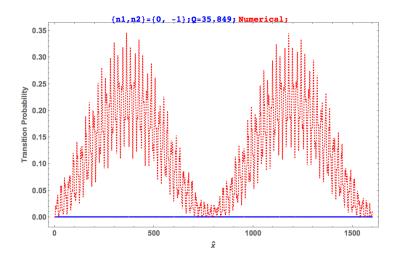
$$Q_2 = \frac{d}{\Gamma_2}.$$

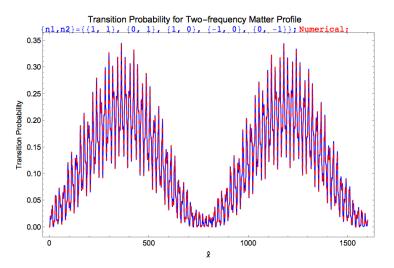












BESSEL'S FUNCTION

$$J_n(\beta) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+n+1)} \left(\frac{\beta}{2}\right)^{2m+n}$$

REFERENCES I