



Neutrino Flavor Conversions in Dense Medium: Matter Stimulation, Dispersion Relation, and Neutrino Halo

PhD Defense

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Lei Ma

Outline for Section 1

1. Neutrino Oscillations

1.1 What are Neutrinos

1.2 Neutrino Oscillations

1.3 Why Do Neutrinos Oscillate

1.4 Matter Effect

1.4.1 Matter Interaction

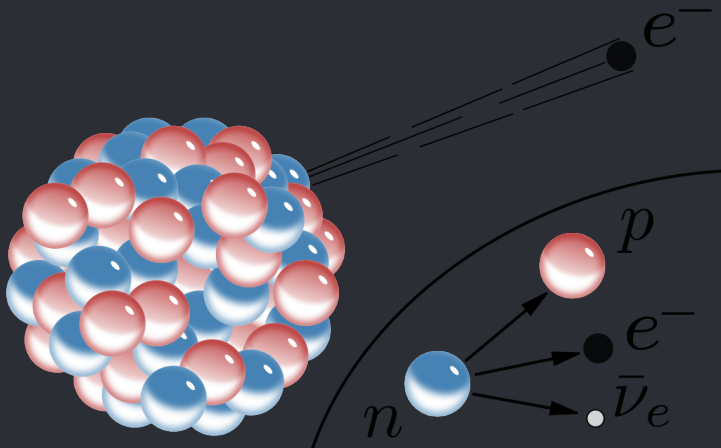
1.4.2 MSW Effect

1.5 Stimulated Neutrino Oscillations

1.6 Two-frequency Matter Profile

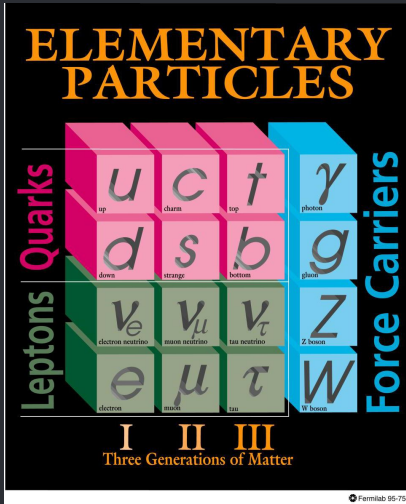
2. Summary

What are Neutrinos?



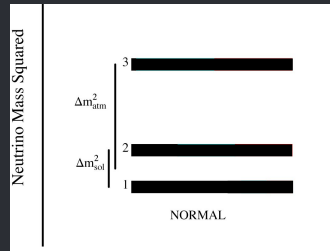
Beta decay and antineutrino production. Source: Beta_Decay@Wikipedia

What are Neutrinos?



Neutrinos are

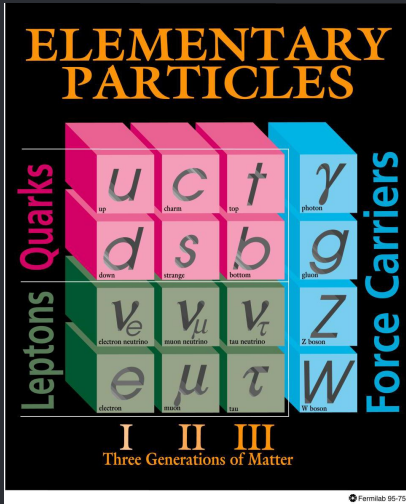
- fermions,
- electrically neutral,
- light.



Adapted from Olga Mena & Stephen Parke (2004)

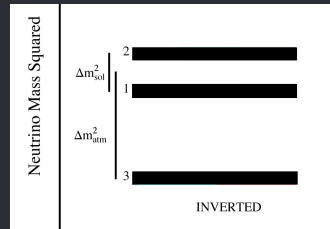
Table of elementary particles. Source: Fermilab

What are Neutrinos?



Neutrinos are

- fermions,
- electrically neutral,
- light.

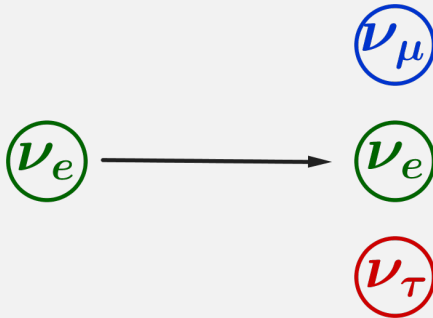


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Table of elementary particles. Source: Fermilab

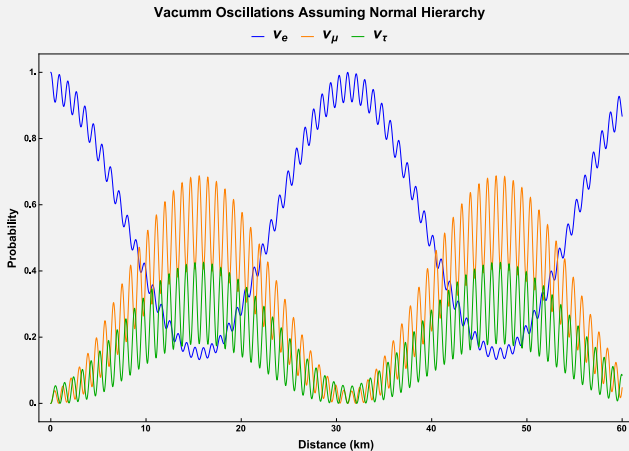
What is Neutrino Oscillation?

Neutrino Oscillation
||
Neutrino Flavor Conversion



Neutrino Oscillations

What is Neutrino Oscillation?



Probabilities of finding neutrinos to be in each flavor.

Why Do Neutrinos Oscillate?

Equation of Motion

$$i\partial_x |\Psi\rangle = \hat{H} |\Psi\rangle$$

- Basis: Hamiltonian diagonalized basis/mass basis/propagation basis, $\{|\nu_1\rangle, |\nu_2\rangle\}$.

-

$$\mathbf{H} = -\frac{\omega_\nu}{2}\sigma_3, \quad \text{where } \omega_\nu = \frac{\delta m^2}{2E} = \frac{m_2^2 - m_1^2}{2E}.$$

- The system can be solved given initial condition of the amplitudes of the two eigenstates $(\langle\nu_1|\Psi(0)\rangle, \langle\nu_2|\Psi(0)\rangle)^T$,

$$\begin{pmatrix} \langle\nu_1|\Psi(x)\rangle \\ \langle\nu_2|\Psi(x)\rangle \end{pmatrix} = \begin{pmatrix} \langle\nu_1|\Psi(0)\rangle \exp(i\omega_\nu x/2) \\ \langle\nu_2|\Psi(0)\rangle \exp(-i\omega_\nu x/2) \end{pmatrix}$$

Why Do Neutrinos Oscillate?

Flavor basis

Neutrino wave function in flavor basis $\{|\nu_e\rangle, |\nu_\mu\rangle\}$ is related to state in energy basis $\{|\nu_1\rangle, |\nu_2\rangle\}$ through

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$

θ_v : vacuum mixing angle

Why Do Neutrinos Oscillate?

Flavor basis

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θ_v : vacuum mixing angle

Hamiltonian **H**

Mass basis

$$\frac{\omega_v}{2} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \\ = -\frac{\omega_v}{2} \sigma_3$$

Flavor basis

$$\frac{\omega_v}{2} \begin{pmatrix} -\cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & \cos 2\theta_v \end{pmatrix} \\ = \frac{\omega_v}{2} (-\cos 2\theta_v \sigma_3 + \sin 2\theta_v \sigma_1)$$

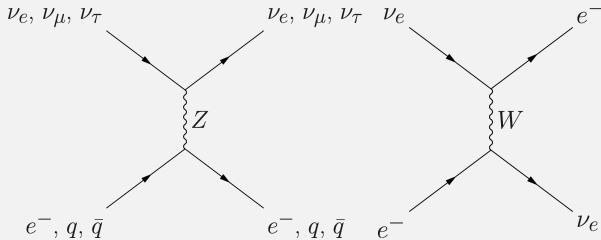
Nature of Neutrino Oscillation

Transition Probability

$$P(|\nu_e\rangle \rightarrow |\nu_\mu\rangle) = \sin^2(2\theta_\nu) \sin^2(\omega_\nu x/2)$$

- $\omega_\nu = (m_2^2 - m_1^2)/2E$ determines oscillation wavelength.
- Mixing angle θ_ν determines flavor oscillation amplitude.

Matter Interaction



Neutral current interaction

between ν_e, ν_μ, ν_τ ,
and e^-, q, \bar{q} .

Charged current interaction
between ν_e and e^-

Matter Interaction

Hamiltonian with matter interaction in flavor basis ($\omega_v = \delta m^2/2E$):

$$\mathbf{H} = \frac{\omega_v}{2} \begin{pmatrix} -\cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & \cos 2\theta_v \end{pmatrix} \pm \sqrt{2} G_F n_e(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

- Vacuum Hamiltonian
- Matter interaction

Matter Interaction

Hamiltonian with matter interaction in flavor basis ($\omega_\nu = \delta m^2/2E$):

$$\mathbf{H} = \frac{\omega_\nu}{2} (-\cos 2\theta_\nu \sigma_3 + \sin 2\theta_\nu \sigma_1) + \frac{\lambda(x)}{2} \sigma_3$$

- Vacuum Hamiltonian

- Matter interaction

- $\lambda(x) = \sqrt{2}G_F n_e(x)$

MSW Effect

Hamiltonian in Vacuum

$$\mathbf{H}_{\text{vacuum}} = \frac{\omega_v \cos 2\theta_v}{2} \sigma_3 + \frac{\omega_v \sin 2\theta_v}{2} \sigma_1$$

$$\begin{aligned} \mathbf{H} &= \frac{\lambda(x) - \omega_v \cos 2\theta_v}{2} \sigma_3 + \frac{\omega_v \sin 2\theta_v}{2} \sigma_1 \\ &= \frac{\omega_m(x) \cos 2\theta_m(x)}{2} \sigma_3 + \frac{\omega_m(x) \sin 2\theta_m(x)}{2} \sigma_1, \end{aligned}$$

where

$$\begin{aligned} \omega_m(x) &= \sqrt{(\lambda(x) - \omega_v \cos 2\theta_v)^2 + \omega_v^2 \sin^2 2\theta_v}, \\ \tan 2\theta_m(x) &= \frac{\omega_v \sin 2\theta_v}{\omega_v \cos 2\theta_v - \lambda(x)}. \end{aligned}$$

MSW Effect

Constant matter profile λ_0 as an example,

Significance of θ_m

Define matter basis $\{|\nu_L\rangle, |\nu_H\rangle\}$

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix}$$

In matter basis

$$\mathbf{H}_{\text{matter-basis}} = -\frac{\omega_m}{2} \sigma_3$$

MSW Resonance

$$\begin{aligned}\mathbf{H} &= \frac{\lambda(x) - \omega_v \cos 2\theta_v}{2} \sigma_3 + \frac{\omega_v \sin 2\theta_v}{2} \sigma_1 \\ &= \frac{\omega_m(x) \cos 2\theta_m(x)}{2} \sigma_3 + \frac{\omega_m(x) \sin 2\theta_m(x)}{2} \sigma_1\end{aligned}$$

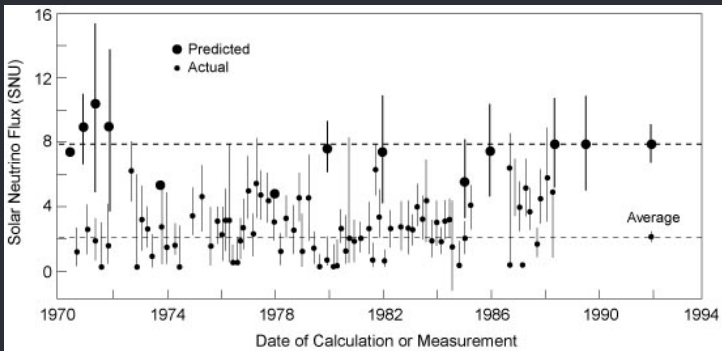
$$\tan 2\theta_m(x) = \frac{\omega_v \sin 2\theta_v}{\omega_v \cos 2\theta_v - \lambda(x)}.$$

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix}$$

Transition Probability

$$P(|\nu_e\rangle \rightarrow |\nu_\mu\rangle) = \sin^2(2\theta_m) \sin^2(\omega_m x)$$

Solar Neutrino Problem



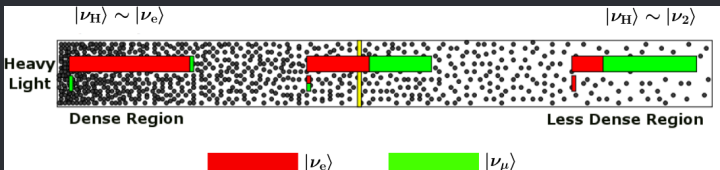
Chlorine detector (Homestake experiment) results and theory predictions. SNU: 1 event for 10^{36} target atoms per second. Kenneth R. Lang (2010)

MSW Effect and Solar Neutrinos

$$\mathbf{H} = \frac{\lambda(x) - \omega_\nu \cos 2\theta_\nu}{2} \sigma_3 + \frac{\omega_\nu \sin 2\theta_\nu}{2} \sigma_1$$

$$\begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & -\sin \theta_m \\ \sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix}$$

$$\mathbf{H}_{\text{matter-basis}} = -\frac{\omega_m}{2} \sigma_3$$



Yellow bar is the resonance point. Red: $|\nu_e\rangle$. Green: $|\nu_\mu\rangle$. Adapted from Smirnov, 2003.

MSW Effect

Suppose $\omega_v = (m_2^2 - m_1^2)/2E < 0$,

$$\mathbf{H} = -\frac{\omega_v}{2} \begin{pmatrix} -\cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & \cos 2\theta_v \end{pmatrix} + \sqrt{2}G_F n_e(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

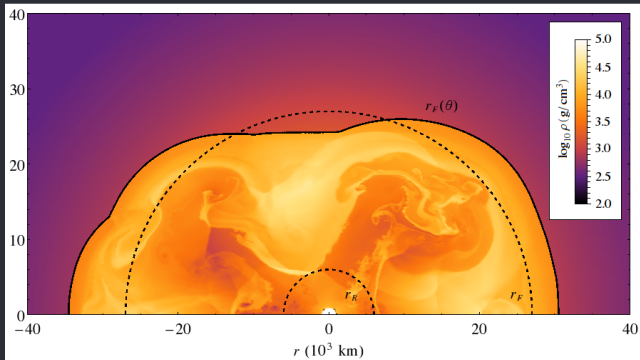


$$\mathbf{H} = \left(\frac{-\omega_v}{2} \cos 2\theta_v + \frac{\lambda(x)}{2} \right) \sigma_3 - \frac{\omega_v}{2} \sin 2\theta_v \sigma_1$$

Supernova Matter Density Profile

Why Do We Care

Astrophysical environments: supernovae, accretion disks etc



Supernova shock and turbulence. E. Borriello, et al (2014)

$$\Delta n(r) = \sum a \sin(kr + t)$$

Stimulated Neutrino Oscillations

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x)$$

Basis

Background matter basis: Hamiltonian is diagonalized with only background matter profile λ_0 ,

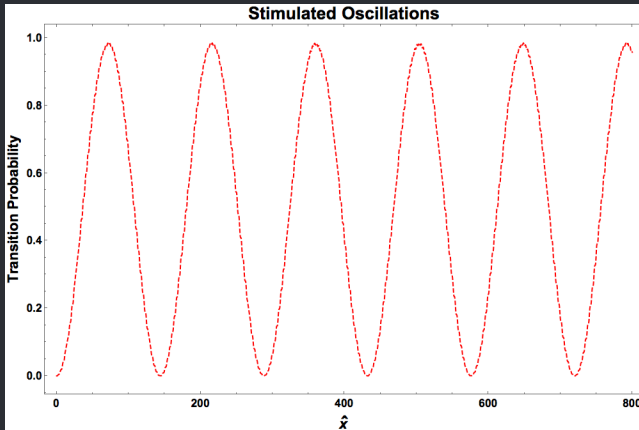
$$H_{\text{background}} = -\frac{\omega_m}{2} \sigma_3.$$

Hamiltonian

$$H = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin \theta_m \sigma_1.$$

Stimulated Neutrino Oscillations

P. Krastev and A. Smirnov (1989); J. Kneller et al (2013);
K. Patton et al (2014);



Stimulated oscillations. $\lambda(x) = \lambda_0 + A \sin(kx)$ with $\hat{x} = \omega_m x$, $A = 0.1 \omega_m$,

Hamiltonian, and Basis, and Rabi Oscillations

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin 2\theta_m \sigma_1.$$

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \cos 2\theta_m A \cos(kx)) \sigma_3 - \frac{\sin 2\theta_m}{2} A \cos(kx) \sigma_1.$$

Hamiltonian, and Basis, and Rabi Oscillations

Rabi Oscillation

Hamiltonian

$$-\frac{\omega_0}{2}\sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$



$$\frac{1}{2} (-\omega_m + \cos 2\theta_m \text{Acos}(kx)) \sigma_3 - \frac{\sin 2\theta_m}{2} \text{Acos}(kx) \sigma_1$$

Hamiltonian, and Basis, and Rabi Oscillations

Rabi Oscillation

Hamiltonian

$$-\frac{\omega_0}{2} \sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$



The transition probability from low energy to high energy is

$$P_{1 \rightarrow 2} = \frac{\alpha^2}{\alpha^2 + (\omega_0 - k)^2} \sin^2 \left(\frac{\Omega_R}{2} t \right),$$

where

$$\Omega_R = \sqrt{\alpha^2 + (\omega_0 - k)^2}$$

is Rabi frequency.

Hamiltonian, and Basis, and Rabi Oscillations

Rabi Oscillation

Hamiltonian

$$-\frac{\omega_0}{2}\sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$



The transition probability from low energy to high energy is

$$P_{1 \rightarrow 2} = \frac{1}{1 + D^2} \sin^2 \left(\frac{\Omega_R}{2} t \right),$$

where

$$D = \left| \frac{\omega_0 - k}{\alpha} \right|.$$

Visualizing Rabi Oscillations

$$-\frac{\omega_0}{2}\sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$

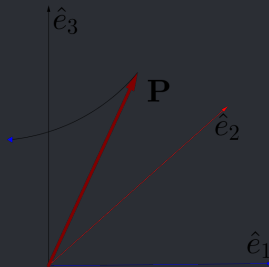
$$-\frac{\omega_0}{2}\sigma_3 - \frac{\alpha}{2} \cos(kt)\sigma_1 + \frac{\alpha}{2} \sin(kt)\sigma_2$$

$$= \begin{pmatrix} \alpha \cos(kt) & -\alpha \sin(kt) & \omega_0 \end{pmatrix} \begin{pmatrix} -\sigma_1/2 \\ -\sigma_2/2 \\ -\sigma_3/2 \end{pmatrix}$$

$$= \vec{H} \cdot (-\vec{\sigma}/2)$$

$$D = \left| \frac{\omega_0 - k}{\alpha} \right|$$

is



Interferences of Rabi Oscillations

$$\begin{pmatrix} 0 \\ 0 \\ \omega_0 \end{pmatrix} + \alpha \begin{pmatrix} \cos(kt) \\ -\sin(kt) \\ 0 \end{pmatrix}$$

$$\begin{aligned} \mathbf{H} &= \frac{1}{2} \left(-\omega_m + \cos 2\theta_m A \cos(kx) \right) \sigma_3 - \frac{\sin 2\theta_m}{2} A \cos(kx) \sigma_1 \\ &\rightarrow -\frac{\omega_m}{2} \sigma_3 - \frac{A \sin 2\theta_m}{2} \cos(kx) \sigma_1 \end{aligned}$$

$$\tilde{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m \end{pmatrix} + \frac{A \sin 2\theta_m}{2} \begin{pmatrix} \cos(kx) \\ -\sin(kx) \\ 0 \end{pmatrix} + \frac{A \sin 2\theta_m}{2} \begin{pmatrix} \cos(-kx) \\ -\sin(-kx) \\ 0 \end{pmatrix}$$

Two frequencies!

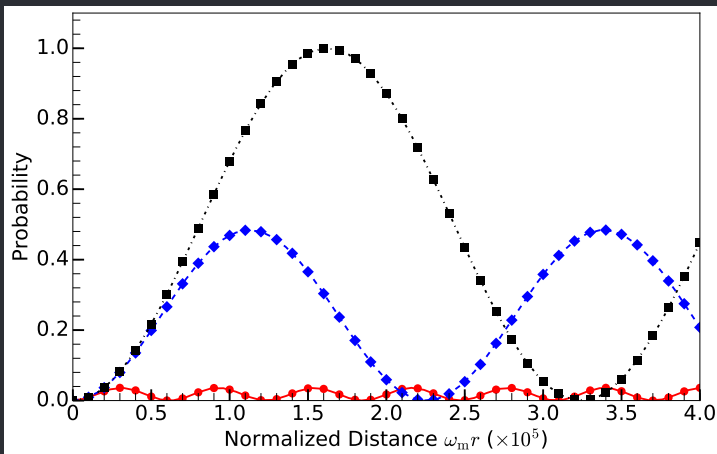
Interferences of Rabi Oscillations

$$\tilde{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 x) \\ -\sin(k_1 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} \cos(k_2 x) \\ -\sin(k_2 x) \\ 0 \end{pmatrix}$$

Corotating frame of the second frequency,

$$\tilde{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m - k_2 \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 - k_2 x) \\ -\sin(k_1 - k_2 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

Rabi Formula Works



Dots, diamonds, triangles, and squares are for $k = \omega_m$, $k = (1 - 2 \times 10^{-5})\omega_m$, and $k = (1 - 10^{-4})\omega_m$ respectively.
Lines: Rabi formula

Interferences of Rabi Oscillations

Matter profile with two frequencies,

$$\lambda(x) = \lambda_0 + \frac{\alpha_1}{2} \sin(k_1 x) \sigma_1$$

Understanding Stimulated Oscillations

Matter profile

$$\lambda(x) = \lambda_0 + A \sin(kx),$$

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin \theta_m \sigma_1.$$

A Better Basis

Define new basis $\{|\tilde{\nu}_L\rangle, |\tilde{\nu}_H\rangle\}$ is related to background matter basis $\{|\nu_L\rangle, |\nu_H\rangle\}$ through

$$\begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix} = \begin{pmatrix} e^{-i\eta(x)} & 0 \\ 0 & e^{i\eta(x)} \end{pmatrix} \begin{pmatrix} |\tilde{\nu}_L\rangle \\ |\tilde{\nu}_H\rangle \end{pmatrix},$$

where

Single Frequency Matter Profile

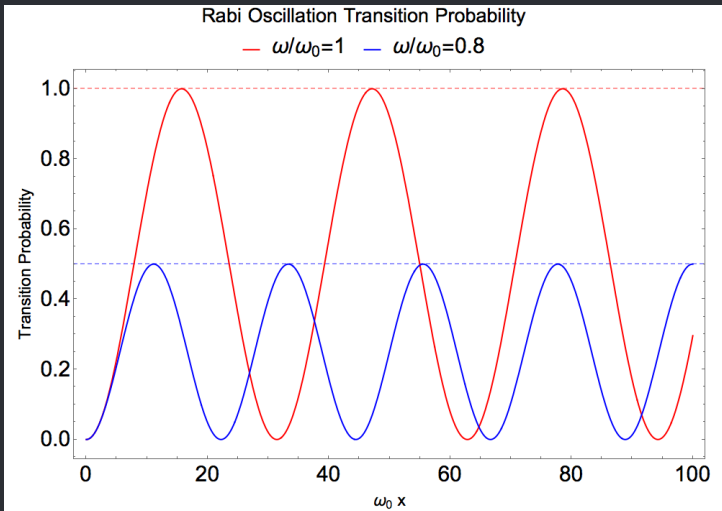
Hamiltonian in new basis

$$\tilde{\mathbf{H}} = -\frac{\delta\lambda(x)}{2} \sin 2\theta_m \begin{pmatrix} 0 & e^{2i\eta(x)} \\ e^{-2i\eta(x)} & 0 \end{pmatrix} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

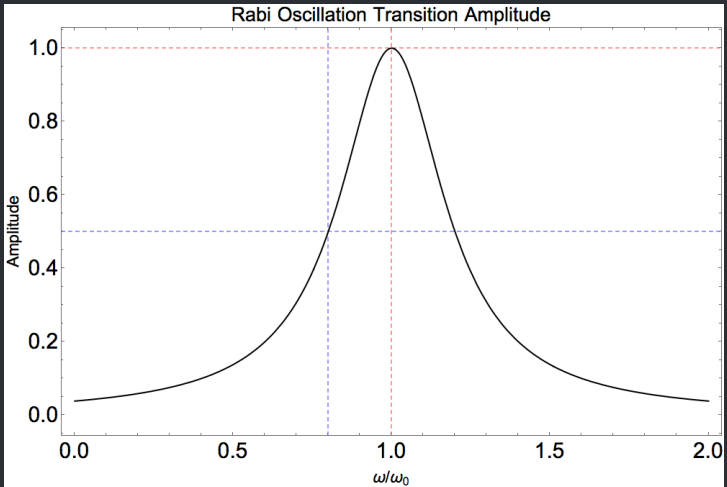
Hamiltonian in New Basis

$$\begin{aligned} h &\equiv -\frac{\delta\lambda(x)}{2} e^{2i\eta(x)} \\ &= \frac{i}{4} \left[\exp \left(i(k + \omega_m)x + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right. \\ &\quad \left. - \exp \left(i(-k + \omega_m)x + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right] \end{aligned}$$

Rabi Oscillations



Rabi Oscillations



Resonance

Single Frequency Matter Profile

Off-diagonal Term in Our System

$$\tilde{\mathbf{H}} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

$$h \propto \left[\exp \left(i(k + \omega_m)x + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) - \exp \left(i(-k + \omega_m)x + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right]$$

Jacobi-Anger expansion

$$e^{i\beta \cos(kx)} = \sum_{n=-\infty}^{\infty} i^n J_n(\beta) e^{inkx},$$

where $J_n(\beta)$ are Bessel's functions of the first kind.

Single Frequency Matter Profile

Scaled Quantities

Characteristic scale: ω_m

- $\hat{A} = A/\omega_m$
- $\hat{k} = k/\omega_m$
- $\hat{x} = \omega_m x$
- $\hat{h} = h/\omega_m$

Single Frequency Matter Profile

Rotation Wave Approximation

The off-diagonal element of Hamiltonian

$$\tilde{\mathbf{H}} = \sum_{n=-\infty}^{\infty} \begin{pmatrix} 0 & \frac{1}{2} \hat{B}_n e^{i(n\hat{k}-1)\hat{x}} \\ \frac{1}{2} \hat{B}_n^* e^{-i(n\hat{k}-1)\hat{x}} & 0 \end{pmatrix}$$

where $\hat{B}_n = -(-i)^n n \hat{k} \tan 2\theta_m J_n(\hat{A} \cos 2\theta_m / \hat{k})$.

Single Frequency Matter Profile

Transition Probability

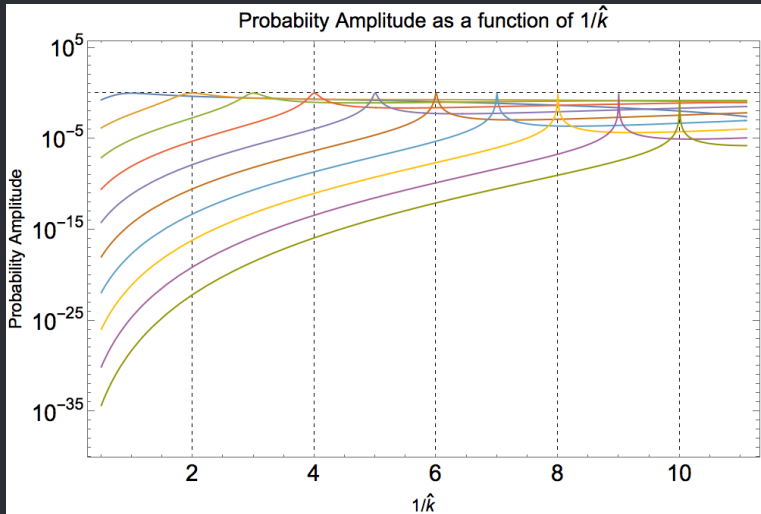
$$P_{L \rightarrow H}^{(n)} = \frac{|\hat{B}_n|^2}{|\hat{B}_n|^2 + (n\hat{k} - 1)^2} \sin^2 \left(\frac{q^{(n)}}{2} x \right),$$

where

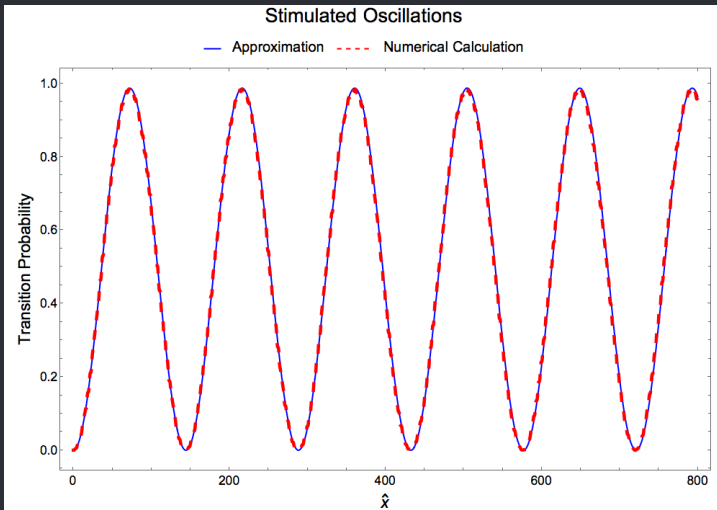
$$q^{(n)} = \sqrt{|\Gamma^{(n)}|^2 + (n\hat{k} - 1)^2}, \quad \text{frequency of oscillations}$$

$$\Gamma^{(n)} = |\hat{B}_n|, \quad \text{width of resonance } (n\hat{k} \text{ as parameter})$$

Single Frequency Matter Profile

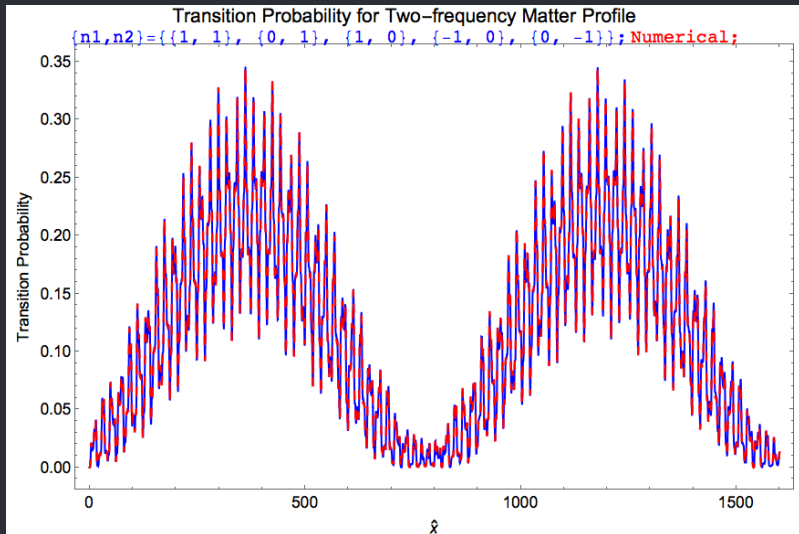


Single Frequency Matter Profile



$$\hat{\lambda} = 0.1, \hat{k} = 0.995, \theta_m = \pi/6$$

Two-frequency Matter Profile



$$\lambda(x) = \lambda_0 + A_1 \sin(k_1 x) + A_2 \sin(k_2 x). \quad \hat{k}_1 = 0.3, \hat{k}_2 = 0.7, A_1 = A_2 = 0.1, \\ \theta_1 = \pi/5$$

Outline for Section 2

1. Neutrino Oscillations

1.1 What are Neutrinos

1.2 Neutrino Oscillations

1.3 Why Do Neutrinos Oscillate

1.4 Matter Effect

1.4.1 Matter Interaction

1.4.2 MSW Effect

1.5 Stimulated Neutrino Oscillations

1.6 Two-frequency Matter Profile

2. Summary

Summary

- The fact that neutrino flavor states are not mass states causes vacuum oscillations.
- MSW resonance happens when matter potential cancels out the vacuum diagonal elements of the Hamiltonian.
- Even matter profile doesn't match MSW requirement, variation in matter profile can cause resonances.
- Single frequency perturbations in matter profile is a combination of many Rabi oscillations.
- How to understand and calculate systems with multi-frequency matter profile (turbulence).
- Combine periodic or even turbulent matter profile with neutrino self-interaction.

Acknowledgement

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Citations

\TeX , \LaTeX , and Beamer

\TeX is a programming language for the typesetting of documents. It was created by Donald Erwin Knuth in the late 1970s and it is documented in *The \TeX book* [1].

In the early 1980s, Leslie Lamport created the initial version of \LaTeX , a high-level language on top of \TeX , which is documented in *\LaTeX : A Document Preparation System* [2]. There exists a healthy ecosystem of packages that extend the base functionality of \LaTeX ; *The \LaTeX Companion* [3] acts as a guide through the ecosystem.

In 2003, Till Tantau created the initial version of Beamer, a \LaTeX package for the creation of presentations. Beamer is documented in the *User's Guide to the Beamer Class* [4].

Bibliography

T_EX, *ΛT_EX*, and Beamer

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- [3] M. Goossens, F. Mittelbach, and A. Samarin. *The ΛT_EX Companion*. Addison-Wesley, 1994.
- [4] Till Tantau. *User's Guide to the Beamer Class Version 3.01*. Available at <http://latex-beamer.sourceforge.net>.
- [5] A. Mertz and W. Slough. Edited by B. Beeton and K. Berry. *Beamer by example* In TUGboat, Vol. 26, No. 1., pp. 68-73.