

Stimulated Neutrino Flavor Conversions and Rabi Oscillations

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OUTLINE

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 - What are Neutrinos
 - Neutrino Oscillations
 - Why Do Neutrinos Oscillate
2. Matter Effect
 - Interactions with Matter
 - MSW Effect
 - Stimulated Neutrino Oscillations
3. Stimulated Neutrino Flavor Transitions
 - Rabi Oscillations
 - Single Frequency Matter Profile and Rabi Oscillations
4. Single Frequency Matter Potential Revisited
 - Basis and Formalism
 - Multifrequency Matter Potential
5. Summary

OVERVIEW

Introduction

What are Neutrinos

Neutrino Oscillations

Why Do Neutrinos Oscillate

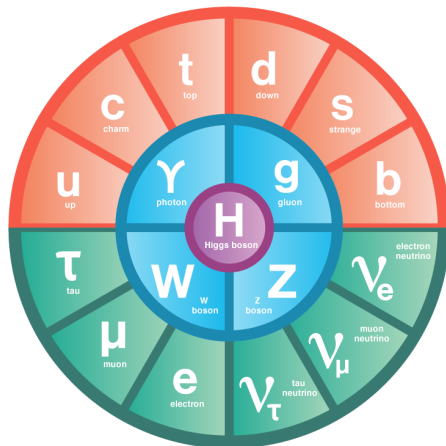
Matter Effect

Stimulated Neutrino Flavor Transitions

Single Frequency Matter Potential Revisited

Summary

WHAT ARE NEUTRINOS?

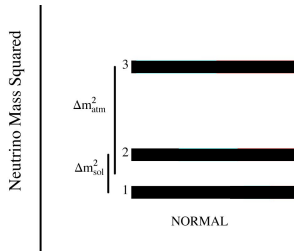


Elementary particles.

Source: symmetrymagazine.org

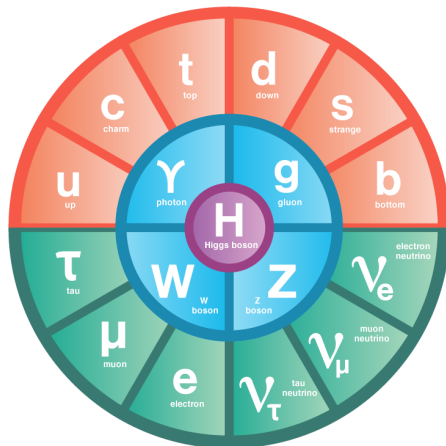
Neutrinos are

- fermions,
- electrically neutral,
- light.



Adapted from Olga Mena & Stephen Parke (2004)

WHAT ARE NEUTRINOS?

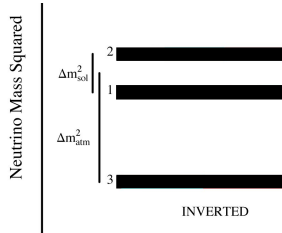


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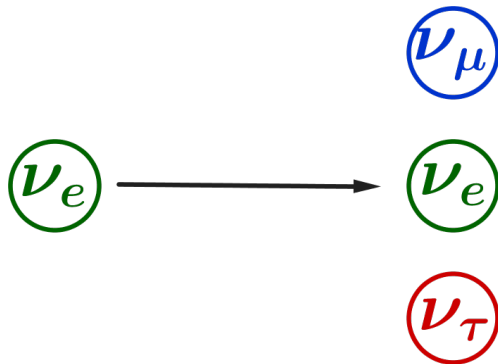
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WHAT ARE NEUTRINO OSCILLATIONS?

Neutrino Oscillations

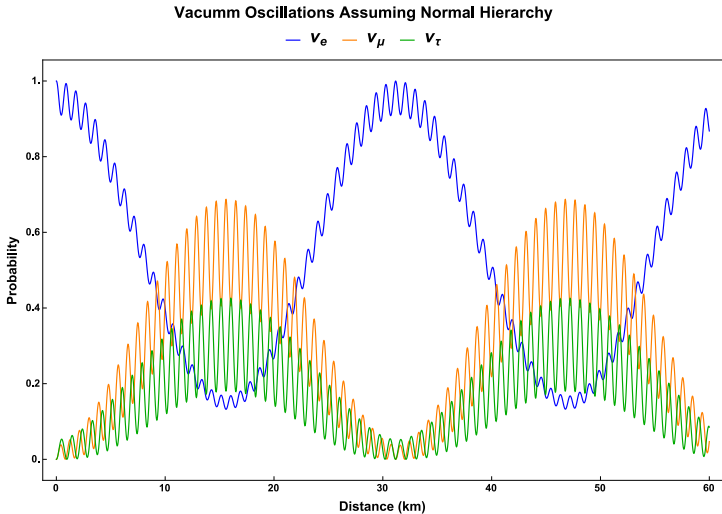
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Neutrino Flavor Conversions



Neutrino Oscillations

WHAT ARE NEUTRINO OSCILLATIONS?



Probabilities of finding neutrinos to be in each flavor.

WHY DO NEUTRINOS OSCILLATE?

$$\begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_\nu & \sin \theta_\nu \\ -\sin \theta_\nu & \cos \theta_\nu \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

Equation of Motion

$$i\partial_x \begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \mathbf{H} \begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix}$$

WHY DO NEUTRINOS OSCILLATE?

$$\begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

Equation of Motion

$$i\partial_x \begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \mathbf{H} \begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix}$$

$$\mathbf{H} = \frac{\omega_v}{2} (-\cos 2\theta_v \sigma_3 + \sin 2\theta_v \sigma_1)$$

- Oscillation frequency:

$$\omega_v = \frac{\delta m^2}{2E} = \frac{m_2^2 - m_1^2}{2E}$$

- Mixing angle θ_v

FLAVOR ISOSPIN

Hamiltonian: $\mathbf{H} = -\frac{\vec{\sigma}}{2} \cdot \vec{H}$

Flavor isospin: $\vec{s} = \Psi^\dagger \frac{\vec{\sigma}}{2} \Psi$

Equation of motion

$$\dot{\vec{s}} = \vec{s} \times \vec{H}$$

electron flavor



muon flavor



FLAVOR ISOSPIN

Hamiltonian: $\mathbf{H} = -\frac{\vec{\sigma}}{2} \cdot \vec{H}$

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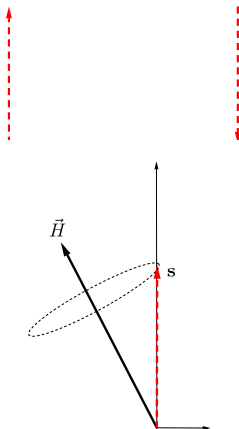
Vacuum oscillation Hamiltonian

$$\frac{\omega_v}{2} (-\cos 2\theta_v \sigma_3 + \sin 2\theta_v \sigma_1)$$

$$\rightarrow \cos 2\theta_v \begin{pmatrix} 0 \\ 0 \\ \omega_v \end{pmatrix} - \sin 2\theta_v \begin{pmatrix} \omega_v \\ 0 \\ 0 \end{pmatrix}$$

electron flavor

muon flavor



OVERVIEW

Introduction

Matter Effect

- Interactions with Matter

- MSW Effect

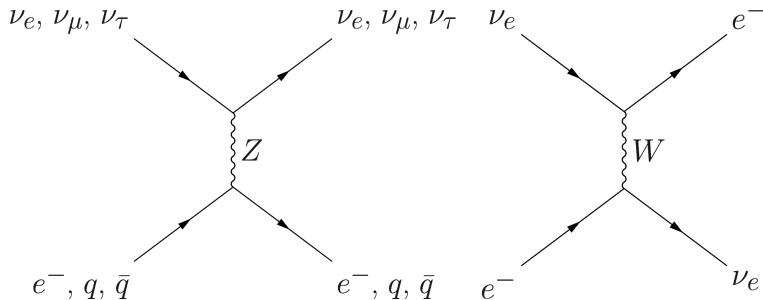
- Stimulated Neutrino Oscillations

Stimulated Neutrino Flavor Transitions

Single Frequency Matter Potential Revisited

Summary

INTERACTIONS WITH MATTER



Neutral current interaction between ν_e, ν_μ, ν_τ , and e^- , quarks etc.

Charged current interaction between ν_e and e^-

MATTER INTERACTION

Hamiltonian with matter interaction in flavor basis ($\omega_v = \delta m^2/2E$):

$$\mathbf{H} = \frac{\omega_v}{2} (-\cos 2\theta_v \sigma_3 + \sin 2\theta_v \sigma_1) + \frac{\lambda(x)}{2} \sigma_3$$

► Vacuum Hamiltonian

► Matter interaction

► $\lambda(x) = \sqrt{2}G_F n_e(x)$

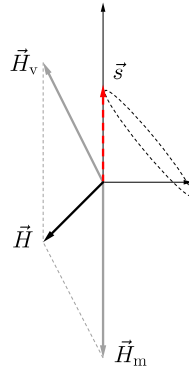
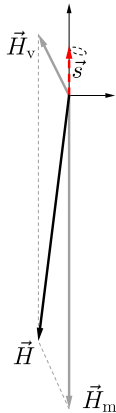
MSW EFFECT

$$\mathbf{H} = \frac{\omega_v}{2} (-\cos 2\theta_v \boldsymbol{\sigma}_3 + \sin 2\theta_v \boldsymbol{\sigma}_1) + \frac{\lambda(x)}{2} \boldsymbol{\sigma}_3$$

$$\rightarrow \omega_v \begin{pmatrix} -\sin 2\theta_v & \\ & 0 \\ & & \cos 2\theta_v \end{pmatrix} + \begin{pmatrix} 0 & \\ & 0 \\ & & -\lambda(x) \end{pmatrix}$$

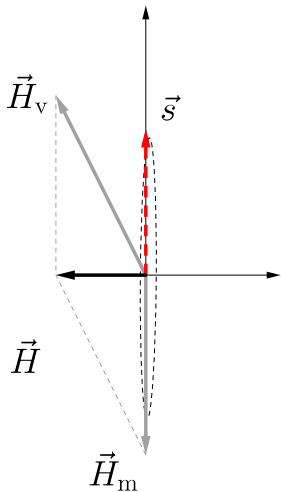
$$\rightarrow \vec{H}_v + \vec{H}_m(x)$$

MSW EFFECT

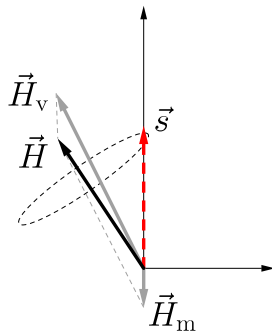


Large density

MSW EFFECT



MSW Resonance

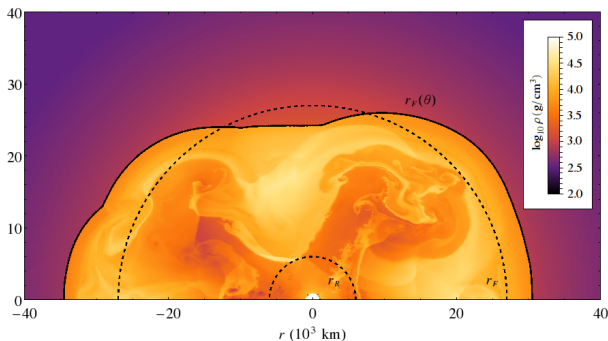


Low density

MORE COMPLICATED MATTER EFFECT

Why Do We Care

Astrophysical environments: supernovae etc



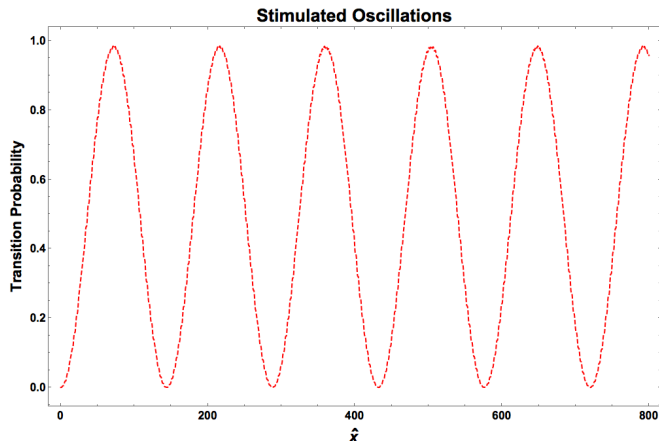
Supernova shock and turbulence. E. Borriello, et al (2014)

Matter Profile

$$\begin{aligned}n_e(x) &= n_0 + \delta n \cos(kx) \\ \Rightarrow \lambda(x) &= \lambda_0 + \delta \lambda \cos(kx)\end{aligned}$$

STIMULATED NEUTRINO OSCILLATIONS

P. Krastev and A. Smirnov (1989); J. Kneller et al (2013);
K. Patton et al (2014);



Stimulated oscillations. $\lambda(x) = \lambda_0 + A \cos(kx)$ with $\hat{x} = \omega_m x$, $A = 0.1\omega_m$,
 $k = 0.995\omega_m$, $\theta_m = \pi/6$

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RABI OSCILLATIONS

Rabi Oscillation

Hamiltonian

$$-\frac{\omega_m}{2}\sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$

$$E_2 = \frac{\omega_m}{2}$$

$$E_1 = -\frac{\omega_m}{2}$$

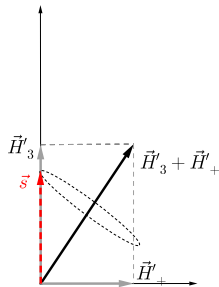
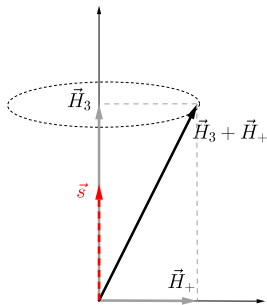
Incoming light



Frequency : k

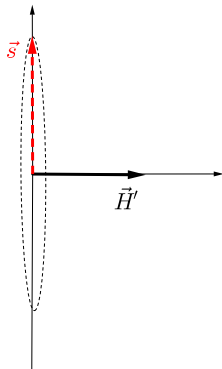
RABI OSCILLATIONS

$$\vec{H}_3 = \omega_m \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \vec{H}_+ = \alpha \begin{pmatrix} \cos(kt) \\ -\sin(kt) \\ 0 \end{pmatrix} \quad \vec{H}'_3 = (\omega_m - k) \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}, \vec{H}'_+ = \alpha \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$



RABI OSCILLATIONS

$$\vec{H}'_3 = (\omega_m - k) \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = 0$$



RABI OSCILLATIONS

Rabi Oscillation

Hamiltonian

$$-\frac{\omega_m}{2}\sigma_3 - \frac{\alpha}{2} \begin{pmatrix} 0 & e^{ikt} \\ e^{-ikt} & 0 \end{pmatrix}$$

Rabi formula

$$P_{1 \rightarrow 2} = \frac{1}{1 + D^2} \sin^2 \left(\frac{\Omega_R}{2} t \right).$$

Relative detuning

$$D = \left| \frac{\omega_m - k}{\alpha} \right|.$$

Rabi frequency

$$\Omega_R = \sqrt{\alpha^2 + (\omega_m - k)^2}$$

$$E_2 = \frac{\omega_m}{2}$$

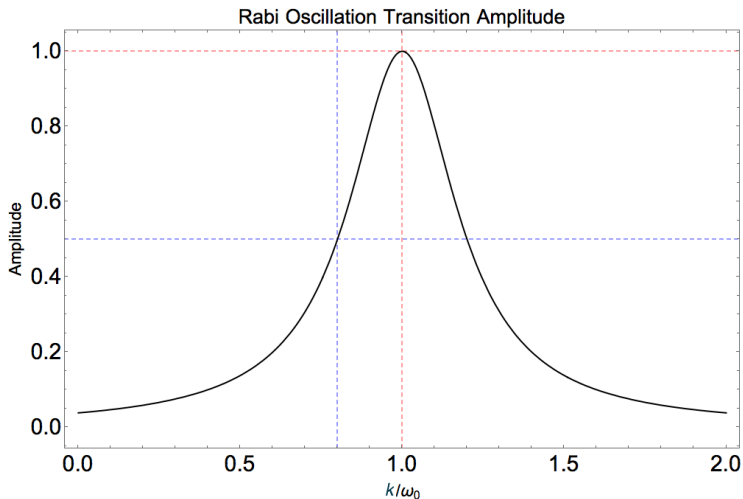
$$E_1 = -\frac{\omega_m}{2}$$

Incoming light



Frequency : k

RABI OSCILLATIONS



Amplitude of Rabi oscillations for different driving field frequency k

HAMILTONIAN IN MATTER BASIS $\begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} \psi_L \\ \psi_H \end{pmatrix}$

Matter Potential

$$\lambda(x) = \lambda_0$$

Basis

Background matter basis (eigen energy basis):

$$\mathbf{H} = \frac{1}{2} (-\omega_m) \sigma_3$$

HAMILTONIAN IN MATTER BASIS $\begin{pmatrix} \psi_e \\ \psi_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} \psi_L \\ \psi_H \end{pmatrix}$

Matter Potential

$$\lambda(x) = \lambda_0 + A \cos(kx)$$

Basis

Background matter basis (eigen energy basis):

$$\mathbf{H} = \frac{1}{2} (-\omega_m + A \cos(kx) \cos 2\theta_m) \sigma_3 - \frac{A \cos(kx)}{2} \sin 2\theta_m \sigma_1$$

HAMILTONIAN IN MATTER BASIS

$$\alpha = \frac{\sin 2\theta_m}{2}$$

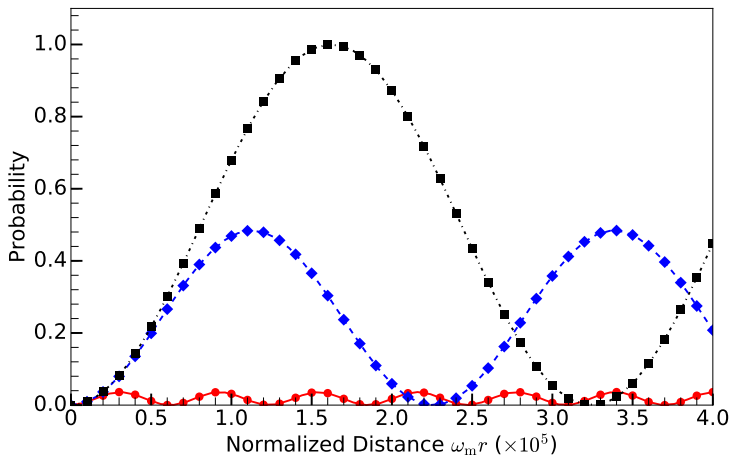
Matter potential frequency

$$k \sim \omega_m$$

$$\begin{aligned} \mathbf{H} &= \frac{1}{2} \left(-\omega_m + \cos 2\theta_m A \cos(kx) \right) \sigma_3 - \frac{\sin 2\theta_m}{2} A \cos(kx) \sigma_1 \\ &\rightarrow \omega_m \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + \alpha \begin{pmatrix} \cos(kx) \\ -\sin(kx) \\ 0 \end{pmatrix} + \alpha \begin{pmatrix} \cos(-kx) \\ -\sin(-kx) \\ 0 \end{pmatrix} \end{aligned}$$

RABI FORMULA WORKS

$$\vec{H} \sim \omega_m \begin{pmatrix} 0 & \\ 0 & 1 \end{pmatrix} + \alpha \begin{pmatrix} \cos(kx) & \\ -\sin(kx) & 0 \end{pmatrix}$$



Lines: Rabi formula

Dots, diamonds, triangles, and squares are for $k = \omega_m$,
 $k = (1 - 2 \times 10^{-5})\omega_m$, and $k = (1 - 10^{-4})\omega_m$ respectively.

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Single Frequency Matter Potential Revisited
Basis and Formalism
Multifrequency Matter Potential

Summary

SINGLE FREQUENCY MATTER POTENTIAL REVISITED

We have been making approximations.

$$\begin{aligned}\mathbf{H} &= \frac{1}{2} \left(-\omega_m + \cos 2\theta_m A \cos(kx) \right) \sigma_3 - \frac{\sin 2\theta_m}{2} A \cos(kx) \sigma_1 \\ &\rightarrow -\frac{\omega_m}{2} \sigma_3 - \frac{A \sin 2\theta_m}{2} \cos(kx) \sigma_1\end{aligned}$$

RABI BASIS

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \boldsymbol{\sigma}_3 - \frac{\delta\lambda(x)}{2} \sin \theta_m \boldsymbol{\sigma}_1.$$

A Better Basis

Define Rabi basis $\{|\tilde{\nu}_L\rangle, |\tilde{\nu}_H\rangle\}$ in which the wave function is related to wave function in background matter basis $\{|\nu_L\rangle, |\nu_H\rangle\}$ through

$$\begin{pmatrix} \psi_L \\ \psi_H \end{pmatrix} = \begin{pmatrix} e^{-i\eta(x)} & 0 \\ 0 & e^{i\eta(x)} \end{pmatrix} \begin{pmatrix} \tilde{\psi}_L \\ \tilde{\psi}_H \end{pmatrix},$$

where

$$\eta(x) - \eta(0) = \frac{\cos 2\theta_m}{2} \int_0^x \delta\lambda(\tau) d\tau.$$

SINGLE FREQUENCY MATTER POTENTIAL

Matter profile

$$\lambda(x) = \lambda_0 + A \sin(kx),$$

Hamiltonian in new basis

$$\tilde{\mathbf{H}} = -\frac{\omega_m}{2}\sigma_3 - \frac{\delta\lambda(x)}{2}\sin 2\theta_m \begin{pmatrix} 0 & e^{2i\eta(x)} \\ e^{-2i\eta(x)} & 0 \end{pmatrix} = -\frac{\omega_m}{2}\sigma_3 + \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

Hamiltonian in New Basis

$$\begin{aligned} h &\equiv -\frac{\delta\lambda(x)}{2}e^{2i\eta(x)} \\ &= \frac{i}{4} \left[\exp \left(ikx + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right. \\ &\quad \left. - \exp \left(-ikx + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right] \end{aligned}$$

SINGLE FREQUENCY MATTER POTENTIAL

Off-diagonal Term in Our System

$$\tilde{\mathbf{H}} = -\frac{\omega_m}{2}\sigma_3 + \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

$$h \propto \left[\exp \left(ikx + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) - \exp \left(-ikx + i \cos 2\theta_m \frac{A}{k} \cos(kx) \right) \right]$$

Jacobi-Anger expansion

$$e^{i\beta \cos(kx)} = \sum_{n=-\infty}^{\infty} i^n J_n(\beta) e^{inkx},$$

where $J_n(\beta)$ are Bessel's functions of the first kind.

SINGLE FREQUENCY MATTER POTENTIAL

Scaled Quantities

Characteristic scale: ω_m

- ▶ $\hat{A} = A/\omega_m$
- ▶ $\hat{k} = k/\omega_m$
- ▶ $\hat{x} = \omega_m x$
- ▶ $\hat{h} = h/\omega_m$

SINGLE FREQUENCY MATTER POTENTIAL

Rotation Wave Approximation

The off-diagonal element of Hamiltonian

$$\tilde{\mathbf{H}} = -\frac{\omega_m}{2}\sigma_3 + \sum_{n=-\infty}^{\infty} \begin{pmatrix} 0 & \frac{1}{2}\hat{B}_n e^{i(n\hat{k})\hat{x}} \\ \frac{1}{2}\hat{B}_n^* e^{-i(n\hat{k})\hat{x}} & 0 \end{pmatrix}$$

where $\hat{B}_n = -(-i)^n n \hat{k} \tan 2\theta_m J_n(\hat{A} \cos 2\theta_m / \hat{k})$.

RABI OSCILLATIONS WITH MULTIFREQUENCIES

$$\vec{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 x) \\ -\sin(k_1 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} \cos(k_2 x) \\ -\sin(k_2 x) \\ 0 \end{pmatrix}$$

Corotating frame of the second frequency,

$$\vec{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m - k_2 \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 - k_2 x) \\ -\sin(k_1 - k_2 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

Energy gap in this frame becomes the length of the vector

$$\begin{pmatrix} 0 \\ 0 \\ \omega_m - k_2 \end{pmatrix} + \alpha_2 \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

RABI OSCILLATIONS WITH MULTIFREQUENCIES

Energy gap in this frame becomes the length of the vector

$$\sqrt{(\omega_m - k_2)^2 + \alpha_2^2} \rightarrow \omega_m - k_2 + \frac{1}{2} \frac{\alpha_2^2}{\omega_m - k_2}$$

Relative detuning

$$D' = \left| \frac{\omega_m - k_1}{\alpha_1} + \frac{\alpha_2^2}{2\alpha_1(\omega_m - k_2)} \right|$$

RABI OSCILLATIONS WITH MULTIFREQUENCIES

Consider $k_1 = \omega_m$

$$D' = \left| \frac{\alpha_2^2}{2\alpha_1(\omega_m - k_2)} \right|$$

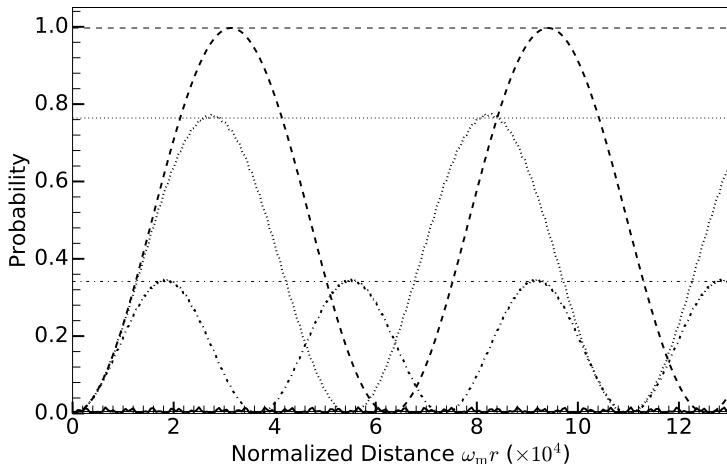
Amplitude reduces from 1 to 1/2 if

$$D' = 1.$$

Two driving frequencies k_1 , and k_2 , with amplitude α_1 , and α_2
Destruction effect: $k_1 = \omega_m$,

$$|\alpha_2| \gg \sqrt{2\omega_m |\alpha_1(k_2 - \omega_m)|} \equiv \alpha_{2,c}$$

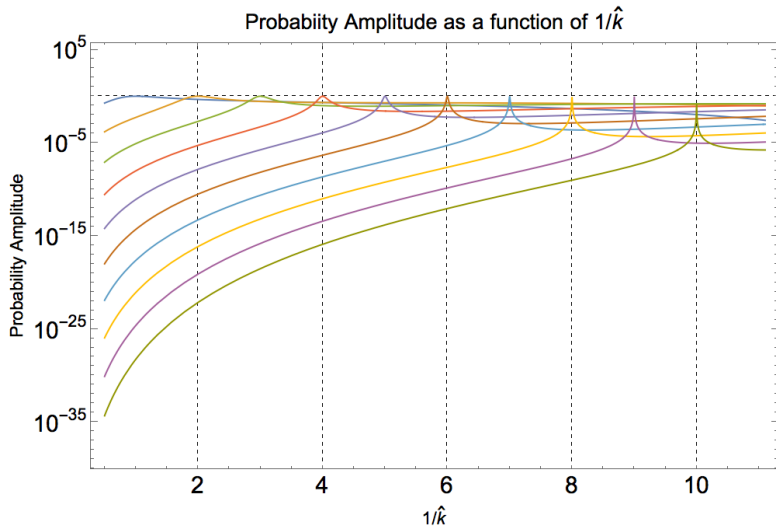
RABI OSCILLATIONS WITH MULTIFREQUENCIES



Grid lines: amplitude predicted using $1/(1 + D'^2)$

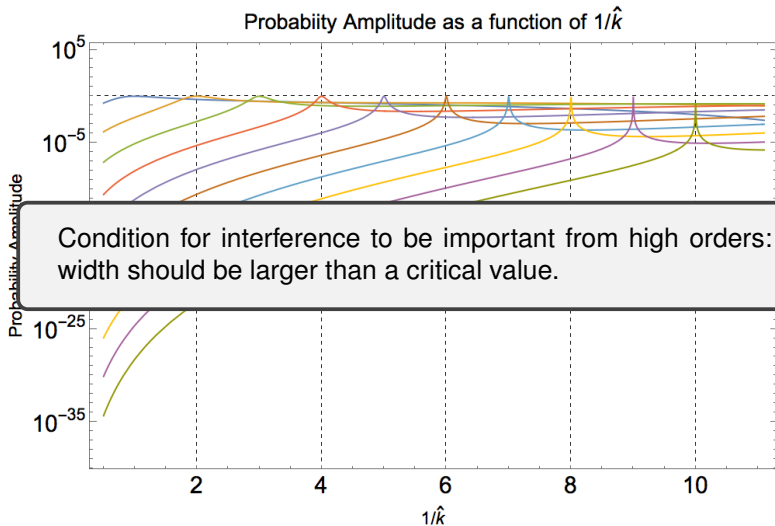
α_2, k_1 values			
Dashed	dotted	dash-dotted	solid
$10^{-2} \omega_m, 10 \omega_m$	$10^{-2} \omega_m, 10^{-1} \omega_m$	$5.0 \times 10^{-2} \omega_m, 10 \omega_m$	$5 \times 10^{-2} \omega_m, 10^{-1} \omega_m$

SINGLE FREQUENCY MATTER POTENTIAL



Resonances of different $n = 1/\hat{k}$. Width becomes extremely narrow for high orders.

SINGLE FREQUENCY MATTER POTENTIAL



Resonances of different $n = 1/\hat{k}$. Width becomes extremely narrow for high orders.

CASTLE WALL MATTER POTENTIAL

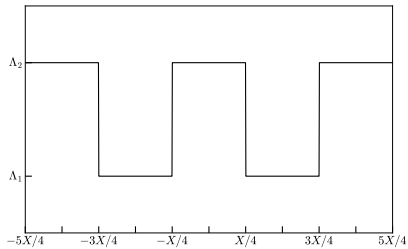
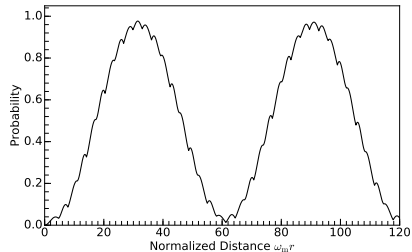


Figure: Castle wall matter profile



OVERVIEW

Introduction

Matter Effect

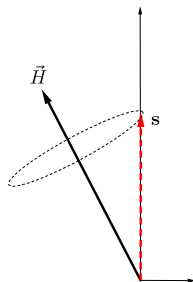
Stimulated Neutrino Flavor Transitions

Single Frequency Matter Potential Revisited

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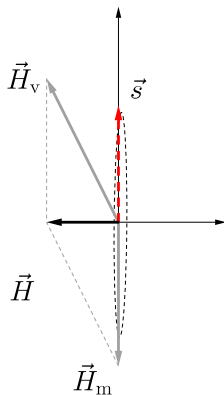
SUMMARY

1. Vacuum oscillations: flavor states are not mass states.
2. MSW resonance: matter potential cancels out the vacuum diagonal elements of the Hamiltonian.
3. Stimulated oscillations: variation in matter profile can cause resonances.
4. Oscillations with two driving fields of different frequencies: large width to destroy the resonance.



SUMMARY

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BACKUP SLIDES

BACKUP SLIDES

WHY DO NEUTRINOS OSCILLATE?

Equation of Motion

$$i\partial_x |\Psi\rangle = \hat{\mathbf{H}} |\Psi\rangle$$

- Basis: Hamiltonian diagonalized basis/mass basis/propagation basis, $\{|\nu_1\rangle, |\nu_2\rangle\}$.

►

$$\mathbf{H} = -\frac{\omega_v}{2}\sigma_3, \quad \text{where } \omega_v = \frac{\delta m^2}{2E} = \frac{m_2^2 - m_1^2}{2E}.$$

- The system can be solved given initial condition of the amplitudes of the two eigenstates $(\langle\nu_1|\Psi(0)\rangle, \langle\nu_2|\Psi(0)\rangle)^T$,

$$\begin{pmatrix} \langle\nu_1|\Psi(x)\rangle \\ \langle\nu_2|\Psi(x)\rangle \end{pmatrix} = \begin{pmatrix} \langle\nu_1|\Psi(0)\rangle \exp(i\omega_v x/2) \\ \langle\nu_2|\Psi(0)\rangle \exp(-i\omega_v x/2) \end{pmatrix}$$

WHY DO NEUTRINOS OSCILLATE?

Flavor basis

Neutrino wave function in flavor basis $\{|\nu_e\rangle, |\nu_\mu\rangle\}$ is related to state in energy basis $\{|\nu_1\rangle, |\nu_2\rangle\}$ through

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$

θ_v : vacuum mixing angle

WHY DO NEUTRINOS OSCILLATE?

Flavor basis

Neutrino wave function in flavor basis $\{|\nu_e\rangle, |\nu_\mu\rangle\}$ is related to state in energy basis $\{|\nu_1\rangle, |\nu_2\rangle\}$ through

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$

θ_v : vacuum mixing angle

Hamiltonian H

Mass basis

$$\begin{aligned} & \frac{\omega_v}{2} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} \\ &= -\frac{\omega_v}{2} \sigma_3 \end{aligned}$$

Flavor basis

$$\begin{aligned} & \frac{\omega_v}{2} \begin{pmatrix} -\cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & \cos 2\theta_v \end{pmatrix} \\ &= \frac{\omega_v}{2} (-\cos 2\theta_v \sigma_3 + \sin 2\theta_v \sigma_1) \end{aligned}$$

NATURE OF NEUTRINO OSCILLATION

Transition Probability

$$P(|\nu_e\rangle \rightarrow |\nu_\mu\rangle) = \sin^2(2\theta_v) \sin^2(\omega_v x/2)$$

- ▶ $\omega_v = (m_2^2 - m_1^2)/2E$ determines oscillation wavelength.
- ▶ Mixing angle θ_v determines flavor oscillation amplitude.

MSW EFFECT

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$

Constant matter profile λ_0 as an example,

Significance of θ_m

Define matter basis (eigenenergy basis) $\{|\nu_L\rangle, |\nu_H\rangle\}$

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix}$$

In matter basis

$$\mathbf{H}_{\text{matter-basis}} = -\frac{\omega_m}{2} \sigma_3$$

MSW EFFECT

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_v & \sin \theta_v \\ -\sin \theta_v & \cos \theta_v \end{pmatrix} \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \end{pmatrix}$$

Constant matter profile λ_0 as an example,

Significance of θ_m

Define matter basis (eigenenergy basis) $\{|\nu_L\rangle, |\nu_H\rangle\}$

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix}$$

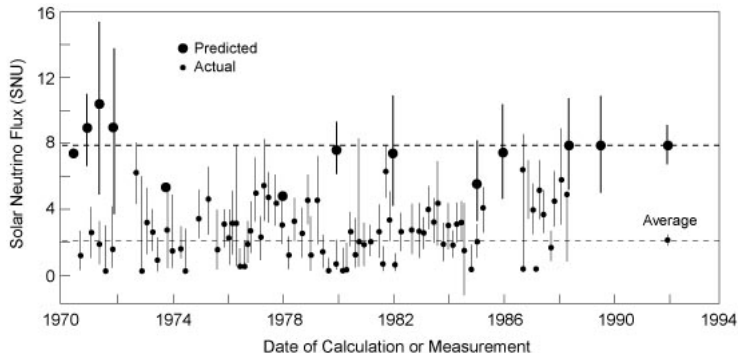
In matter basis

$$\mathbf{H}_{\text{matter-basis}} = -\frac{\omega_m}{2} \sigma_3$$

Transition Probability

$$P(|\nu_e\rangle \rightarrow |\nu_\mu\rangle) = \sin^2(2\theta_m) \sin^2(\omega_m x)$$

SOLAR NEUTRINO PROBLEM



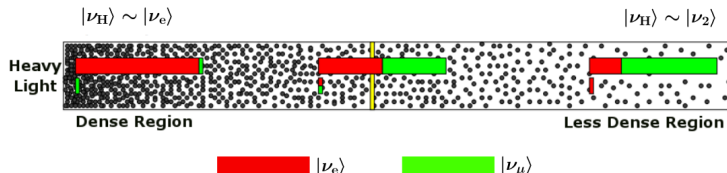
Chlorine detector (Homestake experiment) results and theory predictions.
SNU: 1 event for 10^{36} target atoms per second. Kenneth R. Lang (2010)

MSW EFFECT AND SOLAR NEUTRINOS

$$\mathbf{H} = \frac{\lambda(x) - \omega_\nu \cos 2\theta_\nu}{2} \sigma_3 + \frac{\omega_\nu \sin 2\theta_\nu}{2} \sigma_1$$

$$\begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix} = \begin{pmatrix} \cos \theta_m & -\sin \theta_m \\ \sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \end{pmatrix}$$

$$\mathbf{H}_{\text{matter-basis}} = -\frac{\omega_m}{2} \sigma_3$$



Yellow bar is the resonance point. Red: $|\nu_e\rangle$. Green: $|\nu_\mu\rangle$. Adapted from Smirnov, 2003.

MSW EFFECT INVERTED HIERARCHY

Suppose $\omega_v = (m_2^2 - m_1^2)/2E < 0$,

$$\mathbf{H} = -\frac{\omega_v}{2} \begin{pmatrix} -\cos 2\theta_v & \sin 2\theta_v \\ \sin 2\theta_v & \cos 2\theta_v \end{pmatrix} + \sqrt{2}G_F n_e(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$



$$\mathbf{H} = \left(\frac{-\omega_v}{2} \cos 2\theta_v + \frac{\lambda(x)}{2} \right) \sigma_3 - \frac{\omega_v}{2} \sin 2\theta_v \sigma_1$$

HAMILTONIAN

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x)$$

Basis

Background matter basis (eigen energy basis): Hamiltonian is diagonalized with only background matter profile λ_0 ,

$$H_{\text{background}} = -\frac{\omega_m}{2} \sigma_3.$$

Hamiltonian

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin \theta_m \sigma_1.$$

HAMILTONIAN

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} (-\omega_{\text{m}} + \delta\lambda(x) \cos 2\theta_{\text{m}}) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin 2\theta_{\text{m}} \sigma_1.$$

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$$\mathbf{H} = \frac{1}{2} (-\omega_{\text{m}} + \cos 2\theta_{\text{m}} A \cos(kx)) \sigma_3 - \frac{\sin 2\theta_{\text{m}}}{2} A \cos(kx) \sigma_1.$$

HAMILTONIAN

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} (-\omega_m + \delta\lambda(x) \cos 2\theta_m) \sigma_3 - \frac{\delta\lambda(x)}{2} \sin 2\theta_m \sigma_1.$$

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$$\mathbf{H} = \frac{1}{2} \left(-\omega_m + \cos 2\theta_m A \cos(kx) \right) \sigma_3 - \frac{\sin 2\theta_m}{2} A \cos(kx) \sigma_1.$$

RABI OSCILLATIONS

The coupling strength is calculated as

$$\alpha = \langle 1 | \mathbf{d} \cdot \mathbf{E} | 2 \rangle$$

where the electric field is

$$\mathbf{E} = \mathbf{E}_0 \sin(kt).$$

and \mathbf{d} is the dipole moment.

OSCILLATIONS WITH MULTIFREQUENCIES $\alpha = \frac{\sin 2\theta_m}{2} A$

$$\vec{H} = \omega_m \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} + \alpha \begin{pmatrix} \sin(kx) \\ \cos(kx) \\ 0 \end{pmatrix} + \alpha \begin{pmatrix} \sin(kx) \\ -\cos(kx) \\ 0 \end{pmatrix}$$

Dropping off-resonance frequency: Requirement?

OSCILLATIONS WITH MULTIFREQUENCIES

$$\vec{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 x) \\ -\sin(k_1 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} \cos(k_2 x) \\ -\sin(k_2 x) \\ 0 \end{pmatrix}$$

Corotating frame of the second frequency,

$$\vec{H} = \begin{pmatrix} 0 \\ 0 \\ \omega_m - k_2 \end{pmatrix} + \alpha_1 \begin{pmatrix} \cos(k_1 - k_2 x) \\ -\sin(k_1 - k_2 x) \\ 0 \end{pmatrix} + \alpha_2 \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

Energy gap in this frame becomes the length of the vector

$$\begin{pmatrix} 0 \\ 0 \\ \omega_m - k_2 \end{pmatrix} + \alpha_2 \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

OSCILLATIONS WITH MULTIFREQUENCIES

Energy gap in this frame becomes the length of the vector

$$\sqrt{(\omega_m - k_2)^2 + \alpha_2^2} \rightarrow \omega_m - k_2 + \frac{1}{2} \frac{\alpha_2^2}{\omega_m - k_2}$$

Relative detuning

$$D' = \left| \frac{\omega_m - k_1}{\alpha_1} + \frac{\alpha_2^2}{2\alpha_1(\omega_m - k_2)} \right|$$

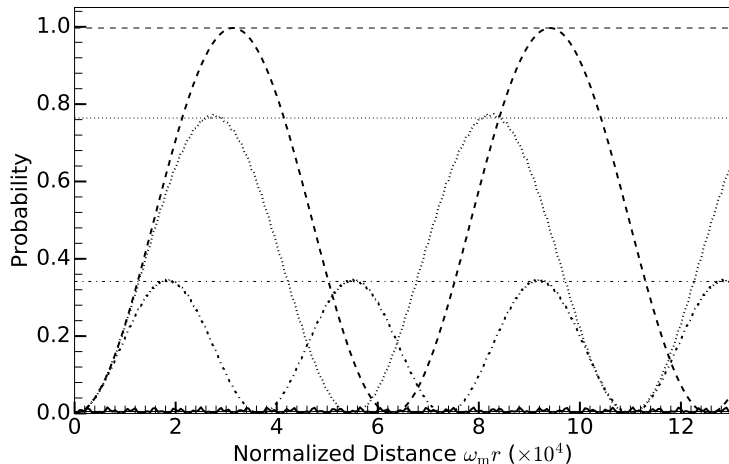
OSCILLATIONS WITH MULTIFREQUENCIES

$$D' = \left| \frac{\omega_m - k_1}{\alpha_1} + \frac{\alpha_2^2}{2\alpha_1(\omega_m - k_2)} \right|$$

Two driving frequencies k_1 , and k_2 , with amplitude α_1 , and α_2

Destruction effect: $k_1 = \omega_m$, $|\alpha_2| \gg \sqrt{2\omega_m|\alpha_1(k_2 - \omega_m)|} \equiv \alpha_{2,c}$

OSCILLATIONS WITH MULTIFREQUENCIES

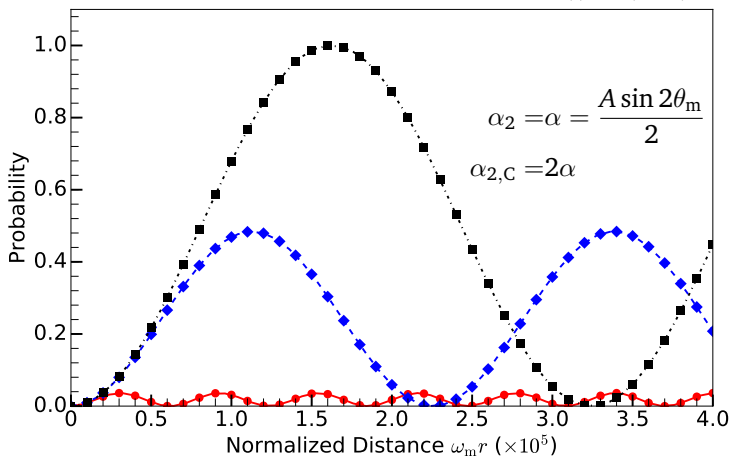


Grid lines: amplitude predicted using $1/(1 + D'^2)$

α_2, k_1 values			
Dashed	dotted	dash-dotted	solid
$10^{-2} \omega_m, 10 \omega_m$	$10^{-2} \omega_m, 10^{-1} \omega_m$	$5.0 \times 10^{-2} \omega_m, 10 \omega_m$	$5 \times 10^{-2} \omega_m, 10^{-1} \omega_m$

RABI FORMULA WORKS

$$\vec{H} \sim \omega_m \begin{pmatrix} 0 & \\ 0 & \\ 1 & \end{pmatrix} + \alpha \begin{pmatrix} \sin(kx) & \\ \cos(kx) & \\ 0 & \end{pmatrix}$$



Lines: Rabi formula

Dots, diamonds, triangles, and squares are for $k = \omega_m$,
 $k = (1 - 2 \times 10^{-5})\omega_m$, and $k = (1 - 10^{-4})\omega_m$ respectively.

PARAMETERS USED FOR VACUUM OSCILLATIONS

$$\theta_{12} = 33.36/180\pi; \theta_{13} = 8.66/180\pi; \theta_{23} = 40/180 * \pi; \delta_{cp} = 0;$$
$$m_1^2 = 0.01; m_2^2 = m_1^2 + 0.000079; E = 1\text{MeV}$$

SINGLE FREQUENCY MATTER PROFILE

Why Does It Work?

$$J_n(n \operatorname{sech} \alpha) \sim \frac{e^{-n(\alpha - \tanh \alpha)}}{\sqrt{2\pi n \tanh \alpha}}, \quad \text{for large } n$$

\Rightarrow

$$\Gamma \propto \hat{B}_n \propto \frac{e^{-n(\alpha - \tanh \alpha)}}{\sqrt{2\pi n \tanh \alpha}}$$

Small perturbation \Rightarrow Small $\hat{A} \Rightarrow$ Large $\alpha \Rightarrow$ Drops fast at large n .

SINGLE FREQUENCY MATTER PROFILE

Transition Probability

$$P_{L \rightarrow H}^{(n)} = \frac{|\hat{B}_n|^2}{|\hat{B}_n|^2 + (n\hat{k} - 1)^2} \sin^2 \left(\frac{q^{(n)}}{2} x \right),$$

where

$$q^{(n)} = \sqrt{|\Gamma^{(n)}|^2 + (n\hat{k} - 1)^2}, \quad \text{frequency of oscillations}$$

$$\Gamma^{(n)} = |\hat{B}_n|, \quad \text{width of resonance } (n\hat{k} \text{ as parameter})$$

Resonance conditions

$$\hat{k} \sim \frac{1}{n}$$

TWO-FREQUENCY MATTER PROFILE

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x), \quad \delta\lambda(x) = A_1 \sin(k_1 x) + A_2 \sin(k_2 x).$$

TWO-FREQUENCY MATTER PROFILE $\hat{h} = \sum_{n=-\infty}^{\infty} \frac{1}{2} \hat{B}_n e^{i(n\hat{k}-1)\hat{x}}$,

Hamiltonian Off-diagonal Element

Apply Jacobi-Anger expansion,

$$\hat{h} = \sum_{n_1=-\infty}^{\infty} \sum_{n_2=-\infty}^{\infty} \frac{1}{2} \hat{B}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1)\hat{x}},$$

where

$$\begin{aligned} & \hat{B}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) \\ &= -(-i)^{n_1+n_2} (n_1 \hat{k}_1 + n_2 \hat{k}_2) J_{n_1} \left(\frac{\hat{A}_1 \cos 2\theta_m}{\hat{k}_1} \right) J_{n_2} \left(\frac{\hat{A}_2 \cos 2\theta_m}{\hat{k}_2} \right) \end{aligned}$$

Which terms are important?

SINGLE FREQUENCY MATTER PROFILE REVISITED

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$k_1 = \omega_m$			
n	D	D'_1	$2\pi\omega_m/\Omega_n$
1	0	-	3.2×10^5
-1	10^5	4.8×10^{-6}	3.1
2	1.1×10^9	2.1×10^{-14}	6.3
-2	3.4×10^9	6.9×10^{-15}	2.1

SINGLE FREQUENCY MATTER PROFILE REVISITED

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$k_1 = (1 - 2 \times 10^{-5})\omega_m$			
n	D	D'_1	$2\pi\omega_m/\Omega_n$
1	1	-	2.2×10^5
-1	10^5	1	3.1
2	1.1×10^9	1	6.3
-2	3.4×10^9	1	2.1

SINGLE FREQUENCY MATTER PROFILE REVISITED

Matter profile

$$\lambda(x) = \lambda_0 + A \cos(kx),$$

$$k_1 = (1 - 10^{-4})\omega_m$$

n	D	D'_1	$2\pi\omega_m/\Omega_n$
1	5.2	-	6.2×10^4
-1	10^5	5.2	3.1
2	1.1×10^9	5.2	6.3
-2	3.4×10^9	5.2	2.1

CASTLE WALL MATTER PROFILE

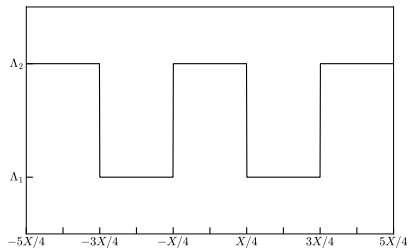
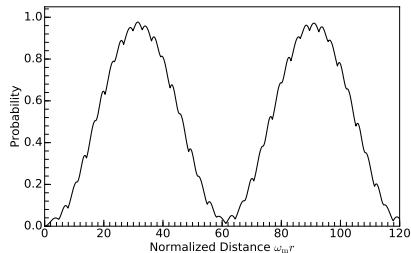


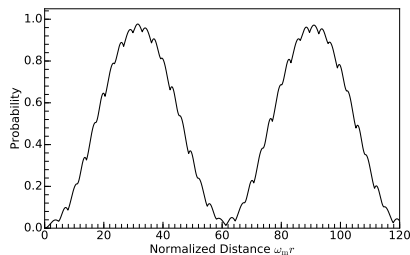
Figure: Castle wall matter profile



CASTLE WALL MATTER PROFILE

Table: Relative detuning of each frequency.

$\{n_1, n_2\}$	D	$D'_{\{1,0\}}$
$\{1, 0\}$	0	-
$\{-1, 0\}$	48	1.0×10^{-2}
$\{0, 1\}$	1.5×10^2	1.1×10^{-3}
$\{2, 0\}$	2.4×10^2	2.0×10^{-4}



TWO-FREQUENCY MATTER PROFILE

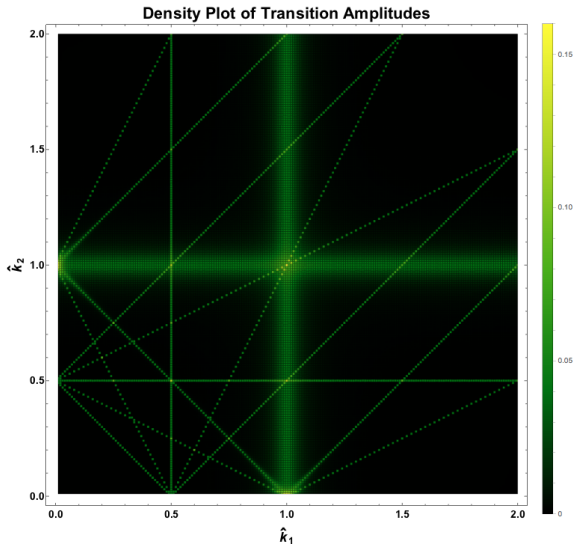
Resonance Lines

There are still resonances, i.e., (almost) zero phases, on lines

$$n_{1,0}\hat{k}_1 + n_{2,0}\hat{k}_2 - 1 = 0$$

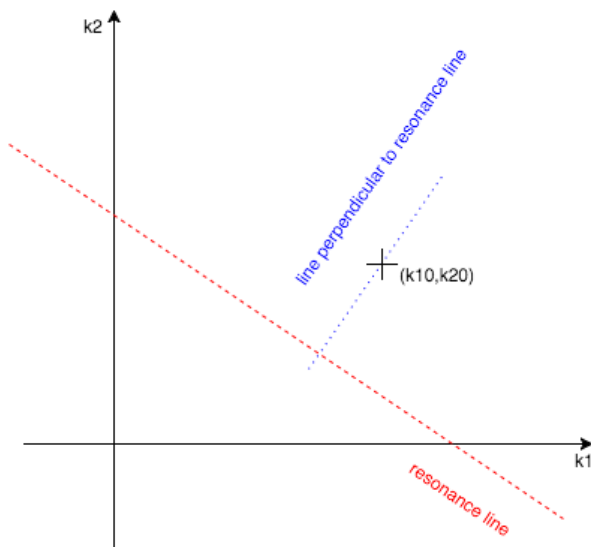
in $\{\hat{k}_1, \hat{k}_2\}$ plane. \Rightarrow Resonance width for each point on resonance lines.

TWO-FREQUENCY MATTER PROFILE $\hat{h} \equiv \sum_{n_1} \sum_{n_2} \frac{1}{2} \hat{B}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1)\hat{x}}$



Density plot of transition amplitudes calculated using only one term out of the whole summation in Hamiltonian. $n_1, n_2 \in [-2, 2]$

TWO-FREQUENCY MATTER PROFILE



Resonance line, distance to resonance, and width

TWO-FREQUENCY MATTER PROFILE

Width

$$\Gamma_2 = \frac{\hat{B}_{n_1, n_2}(\hat{k}_{1, \text{intercept}}, \hat{k}_{2, \text{intercept}})}{\sqrt{n_1^2 + n_2^2}}.$$

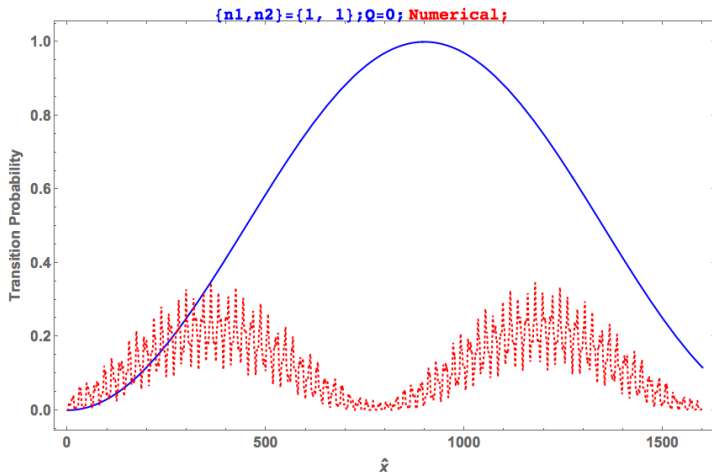
Distance to Resonance Line

$$d = \frac{|n_1 \hat{k}_{10} + n_2 \hat{k}_{20} - 1|}{\sqrt{n_1^2 + n_2^2}}.$$

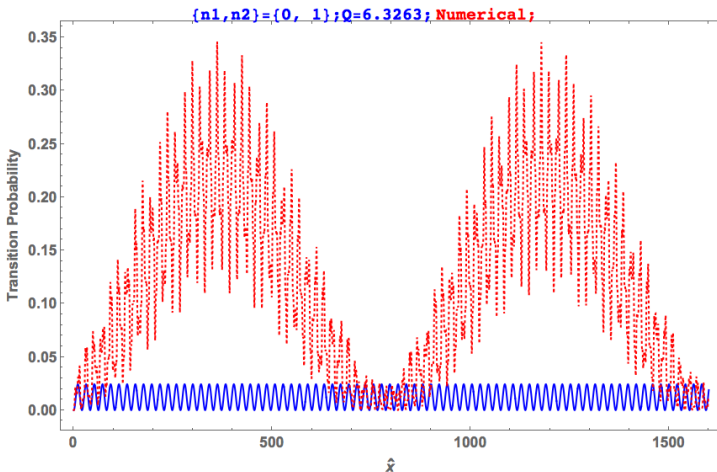
Distance to Resonance Width Ratio

$$Q_2 = \frac{d}{\Gamma_2}.$$

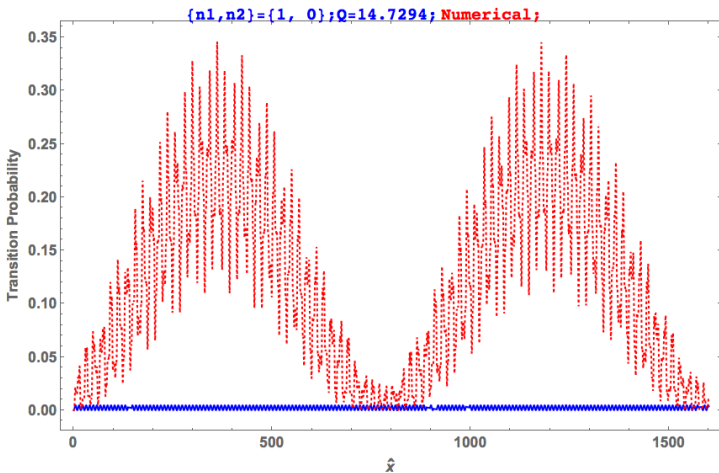
TWO-FREQUENCY MATTER PROFILE



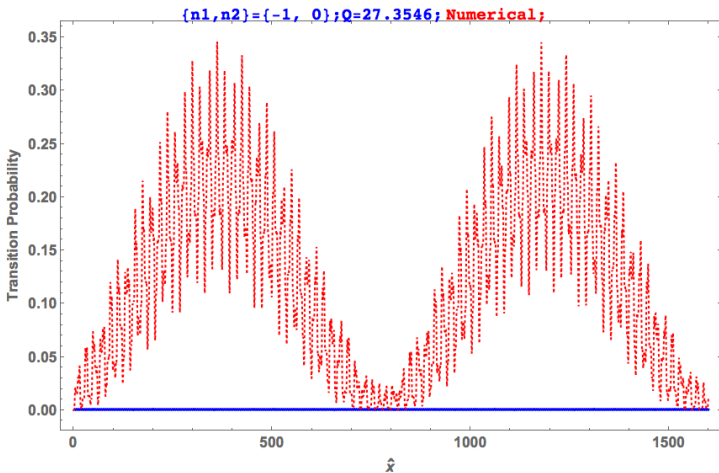
TWO-FREQUENCY MATTER PROFILE



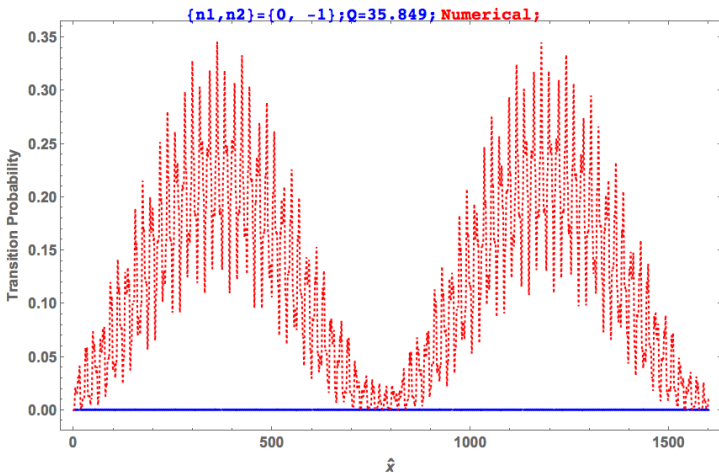
TWO-FREQUENCY MATTER PROFILE



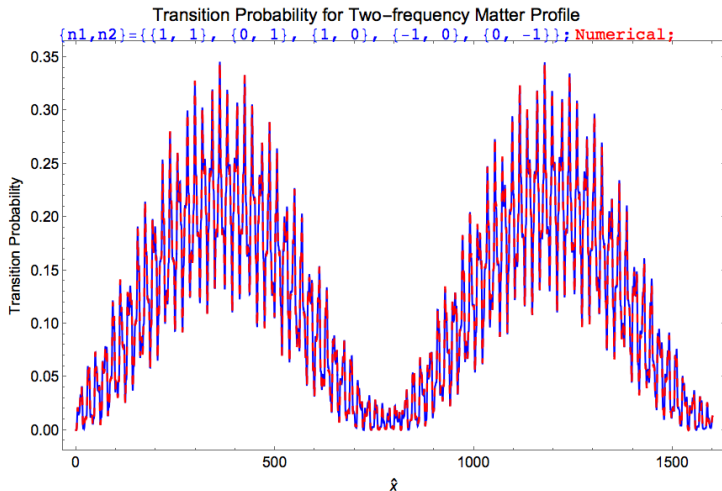
TWO-FREQUENCY MATTER PROFILE



TWO-FREQUENCY MATTER PROFILE



TWO-FREQUENCY MATTER PROFILE



BESSEL'S FUNCTION

$$J_n(\beta) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+n+1)} \left(\frac{\beta}{2}\right)^{2m+n}$$

REFERENCES I