DETECTION OF CNB II

Stodolsky, L. (1975). Speculations on Detection of the "Neutrino Sea." Physical Review Letters, 34(2), 110–112.

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OUTLINE

- · Neutrino Electron Interaction
- · Possible Detection Methods
- · Summary

NEUTRINO ELECTRON INTERACTION ——

WEAK INTERACTION

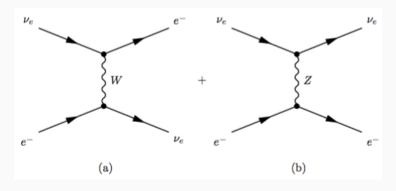


Figure: Neutrino-electron interaction

Effective Lagrangian

$$\begin{split} \mathscr{L}_{\text{eff}} &= -\frac{\mathsf{G}_{\text{F}}}{\sqrt{2}} \{ [\bar{\nu}_{\text{e}} \gamma^{\rho} (1 - \gamma^5) \mathbf{e}] [\bar{\mathbf{e}} \gamma_{\rho} (1 - \gamma^5) \nu_{\text{e}}] \\ &+ [\bar{\nu}_{\text{e}} \gamma^{\rho} (1 - \gamma^5) \nu_{\text{e}}] [\bar{\mathbf{e}} \gamma_{\rho} (\mathbf{g}_{\text{V}}^{\text{l}} - \mathbf{g}_{\text{A}}^{\text{l}} \gamma^5) \mathbf{e}] \} \end{split}$$

- · Red: charged current which exchanges the charge;
- · Blue: neutral current

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FIERZ TRANSFORMATION

In the context of weak interaction, for a Lagrangian,

$$\mathscr{L}(\Psi_{1},\Psi_{2},\Psi_{3},\Psi_{4}) = [\bar{\Psi}_{1}\gamma^{\mu}(1-\gamma^{5})\Psi_{2}][\bar{\Psi}_{3}\gamma_{\mu}(1-\gamma^{5})\Psi_{4}],$$

exchange the field Ψ_2 and Ψ_4 doesn't change the result. The Lagrangian is called V-A theory because people define

$$\begin{split} \mathscr{L}^{V}(\Psi_1,\Psi_2,\Psi_3,\Psi_4) &= [\bar{\Psi}_1\gamma^{\mu}\Psi_2][\bar{\Psi}_3\gamma_{\mu}\Psi_4],\\ \mathscr{L}^{A}(\Psi_1,\Psi_2,\Psi_3,\Psi_4) &= [\bar{\Psi}_1\gamma^{\mu}\gamma^5\Psi_2][\bar{\Psi}_3\gamma_{\mu}\gamma^5\Psi_4]. \end{split}$$

Fierz Transformed Lagrangian

$$\begin{split} \mathscr{L}_{\text{eff}} &= -\frac{\mathsf{G}_{\text{F}}}{\sqrt{2}} \{ [\bar{\nu}_{\text{e}} \gamma^{\rho} (1 - \gamma^5) \mathbf{e}] [\bar{\mathbf{e}} \gamma_{\rho} (1 - \gamma^5) \nu_{\text{e}}] \\ &+ [\bar{\nu}_{\text{e}} \gamma^{\rho} (1 - \gamma^5) \nu_{\text{e}}] [\bar{\mathbf{e}} \gamma_{\rho} (\mathbf{g}_{\text{V}}^{\text{l}} - \mathbf{g}_{\text{A}}^{\text{l}} \gamma^5) \mathbf{e}] \} \end{split}$$

is transformed to

$$\mathscr{L}_{\text{eff}} = -\frac{\mathsf{G}_{\text{F}}}{\sqrt{2}} [\bar{\nu}_{\text{e}} \gamma^{\rho} (1 - \gamma^5) \nu_{\text{e}}] [\bar{\mathsf{e}} \gamma_{\rho} ((1 + \mathsf{g}_{\text{V}}^{\text{l}}) - (1 + \mathsf{g}_{\text{A}}^{\text{l}}) \gamma^5) \mathsf{e}]$$

What's Good about This New Lagrangian

$$\mathscr{L}_{\text{eff}} = -\frac{\mathsf{G}_{\mathsf{F}}}{\sqrt{2}} [\bar{\nu}_{\mathsf{e}} \gamma^{\rho} (1 - \gamma^5) \nu_{\mathsf{e}}] [\bar{\mathsf{e}} \gamma_{\rho} ((1 + \mathsf{g}_{\mathsf{V}}^{\mathsf{l}}) - (1 + \mathsf{g}_{\mathsf{A}}^{\mathsf{l}}) \gamma^5) \mathsf{e}]$$

The neutral current only processes $\nu_{\mu,\tau} {\rm e}^- \to \nu_{\mu,\tau} {\rm e}^-$ are similar to this

$$\mathscr{L}_{\mathrm{eff},\mu\tau} = -\frac{\mathsf{G}_{\mathsf{F}}}{\sqrt{2}}[\bar{\nu}_{\mathsf{e}}\gamma^{\rho}(1-\gamma^{5})\nu_{\mathsf{e}}][\bar{\mathsf{e}}\gamma_{\rho}(\mathsf{g}_{\mathsf{V}}^{\mathsf{l}}-\mathsf{g}_{\mathsf{A}}^{\mathsf{l}}\gamma^{5})\mathsf{e}]$$

We calculate the $\nu_{\mathrm{e}}\mathrm{e}^{-}
ightarrow \nu_{\mathrm{e}}\mathrm{e}^{-}$ processes only.

SPIN DEPENDENT INTERACTION

Current

Current is

$$\vec{J} = 2\rho \frac{\vec{v}}{\sqrt{1 - v^2}},$$

where v is the velocity of electrons with respect to the CNB.

Interaction Energy

The interaction energies for two different helicity states are,

$$\frac{G_F}{\sqrt{2}} \vec{\sigma} \cdot \vec{J} = \pm \sqrt{2} G_F \rho \frac{v}{\sqrt{1-v^2}}. \label{eq:gf}$$

Number Density

$$ho \propto p_F^3$$
.

Energy Split and Frequency

$$\Delta E = 2\sqrt{2}G_F \rho \frac{\vec{V}}{\sqrt{1 - V^2}} = 0.6 \times 10^{-24} \left(\frac{p_F}{eV}\right)^3 \frac{V}{\sqrt{1 - V^2}} eV$$

What is Fermi Momentum p_F

- \cdot β decay: p_F ≤ 60eV
- · cosmological: $p_F \le 0.75 \times 10^{-2} eV$

ENERGY SPLIT

Energy Split

Energy split is of the order

$$\Delta E \sim 10^{-19} eV$$
 to $10^{-30} eV$

Energy split tells us the frequency.

POSSIBLE DETECTION METHODS

ELECTRONS BEAMS

Electron Beams

- \cdot Electron Beams with equally \pm helicity states.
- · + and states have different energies = two different frequencies.
- · Phase difference between the two states.

How Large is the Phase Difference

$$\begin{split} \phi \sim \Delta E t \sim 2\sqrt{2} G_F \rho z \\ \sim 3 \times 10^{-20} \left(\frac{p_F}{eV}\right)^3 rad/cm \end{split}$$

For one light year, we gain a phase difference of the order $\left(\frac{p_F}{eV}\right)^3$ rad.

FERROMAGNETIC MATERIAL

Ferromagnetic Material

- A big chunk of ferromagnetic material (1 ton) contains 10²⁷ aligned electrons.
- \cdot Torque of order $\left(\frac{p_F}{eV}\right)^3 eV$

Problems

- · a small torque
- · external B field

NEUTRONS

Neutrons

The estimation also works for neutral current only processes. He³ mixed into low temperature He⁴ can retain the spin for a long time. Thus we could use the spin dependent interaction of nuclei to detect CNB.

CONCLUSION

Methods

- · Electrons of equally mixing helicity states : a small phase difference
- · 1 ton of ferromagnetic material : a small torque
- · He³ nuclei : phase difference similar to electrons

REFERENCES

1. Stodolsky, L. (1975). Speculations on Detection of the "Neutrino Sea." Physical Review Letters, 34(2), 110–112.

BACKUPS

Backups

SR Velocity Transformation

$$u' = \frac{u - v}{\sqrt{1 - v^2}}$$

CURRENT

Current

$$\vec{J} = -\frac{i}{2m} \left(\Phi^* \nabla \Phi - \Phi \nabla \Phi^* \right)$$

MAGNETIC MOMENT AND B FIELD

Magnetic Moment and B Field

$$\Delta E = -\vec{\mu} \cdot \vec{B}$$