Neutrino Oscillations in Matter

PhD Candidacy Exam

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April 19, 2016

OUTLINE

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 Neutrino Oscillations
 Why Do Neutrinos Oscillate
- Matter Effect
 Matter Interaction
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 Stimulated Neutrino Oscillations
- Understanding Stimulated Oscillations Hamiltonian, and Basis Single Frequency Matter Profile Two-frequency Matter Profile
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OVERVIEW

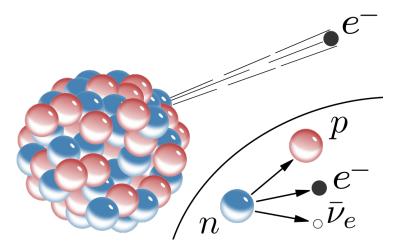
Introduction
What are Neutrinos
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Matter Effect

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WHAT ARE NEUTRINOS?



Beta decay and antineutrino production. Source: Beta_Decay@Wikipedia

WHAT ARE NEUTRINOS?

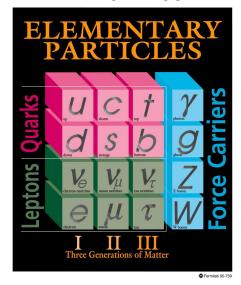
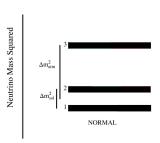


Table of elementary particles. Source: Fermilab

Neutrinos are

- ► fermions.
- electrically neutral,
- ► light.



Adapted from Olga Mena & Stephen Parke (2004)

WHAT ARE NEUTRINOS?

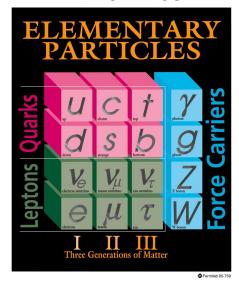
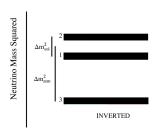


Table of elementary particles. Source: Fermilab

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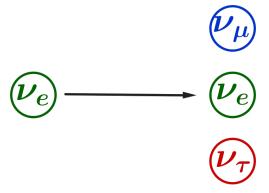
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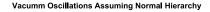
WHAT IS NEUTRINO OSCILLATION?

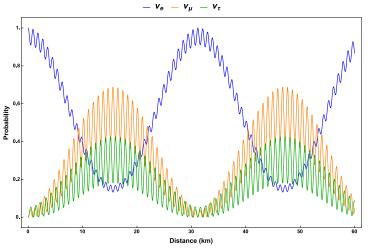




Neutrino Oscillations

WHAT IS NEUTRINO OSCILLATION?





Probabilities of finding neutrinos to be in each flavor.

WHY DO NEUTRINOS OSCILLATE?

Equation of Motion

$$i\partial_x\ket{\Psi}=\hat{\mathbf{H}}\ket{\Psi}$$

► Basis: Hamiltonian diagonalized basis/mass basis/propagation basis, {|\(\nu_1\), |\(\nu_2\)\}.

Þ

$$H=-rac{\omega_{
m v}}{2}\sigma_{3}, \qquad ext{where} \ \omega_{
m v}=rac{\delta m^{2}}{2E}=rac{m_{2}^{2}-m_{1}^{2}}{2E}.$$

► The system can be solved given initial condition of the amplitudes of the two eigenstates $(\langle \nu_1 | \Psi(0) \rangle, \langle \nu_2 | \Psi(0) \rangle)^T$,

$$\begin{pmatrix} \langle \nu_1 | \Psi(x) \rangle \\ \langle \nu_2 | \Psi(x) \rangle \end{pmatrix} = \begin{pmatrix} \langle \nu_1 | \Psi(0) \rangle \exp{(i\omega_v x/2)} \\ \langle \nu_2 | \Psi(0) \rangle \exp{(-i\omega_v x/2)} \end{pmatrix}$$

WHY DO NEUTRINOS OSCILLATE?

Flavor basis

Neutrino wave function in flavor basis $\{|\nu_{\rm e}\rangle\,, |\nu_{\mu}\rangle\}$ is related to state in energy basis $\{|\nu_{\rm 1}\rangle\,, |\nu_{\rm 2}\rangle\}$ through

$$\begin{pmatrix} |\nu_{\rm e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} = \begin{pmatrix} \cos\theta_{\rm v} & \sin\theta_{\rm v} \\ -\sin\theta_{\rm v} & \cos\theta_{\rm v} \end{pmatrix} \begin{pmatrix} |\nu_{1}\rangle \\ |\nu_{2}\rangle \end{pmatrix}$$

 $\theta_v\text{:}$ vacuum mixing angle

WHY DO NEUTRINOS OSCILLATE?

Flavor basis

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 $\theta_{\rm v}$: vacuum mixing angle

Hamiltonian H

Mass basis

$$\begin{split} \frac{\omega_{v}}{2} \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix} & \frac{\omega_{v}}{2} \begin{pmatrix} -\cos 2\theta_{v} & \sin 2\theta_{v} \\ \sin 2\theta_{v} & \cos 2\theta_{v} \end{pmatrix} \\ = -\frac{\omega_{v}}{2} \boldsymbol{\sigma}_{3} & = \frac{\omega_{v}}{2} \left(-\cos 2\theta_{v} \boldsymbol{\sigma}_{3} + \sin 2\theta_{v} \boldsymbol{\sigma}_{1} \right) \end{split}$$

NATURE OF NEUTRINO OSCILLATION

Transition Probability

$$P(|\nu_{\rm e}\rangle \rightarrow |\nu_{\mu}\rangle) = \sin^2(2\theta_{\rm v})\sin^2(\omega_{\rm v}x/2)$$

- $\omega_{\rm v} = (m_2^2 m_1^2)/2E$ determines oscillation wavelength.
- ▶ Mixing angle θ_v determines flavor oscillation amplitude.

OVERVIEW

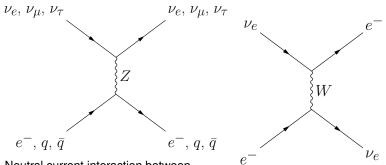
Introduction

Matter Effect
Matter Interaction
MSW Effect
Solar Neutrino Problem
Stimulated Neutrino Oscillations

Understanding Stimulated Oscillations

Summary & Future Work

MATTER INTERACTION



Neutral current interaction between $\nu_{\rm e},\, \nu_{\mu},\, \nu_{\tau},\,$ and e^- , quarks and antiquarks.

Charged current interaction between $\nu_{\rm e}$ and e^-

MATTER INTERACTION

Hamiltonian with matter interaction in flavor basis ($\omega_{\rm v} = \delta m^2/2E$):

$$\mathbf{H} = \begin{array}{cc} \frac{\omega_{\mathrm{v}}}{2} \begin{pmatrix} -\cos 2\theta_{\mathrm{v}} & \sin 2\theta_{\mathrm{v}} \\ \sin 2\theta_{\mathrm{v}} & \cos 2\theta_{\mathrm{v}} \end{pmatrix} & \pm \sqrt{2}G_{\mathrm{F}}n_{\mathrm{e}}(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

- ► Vacuum Hamiltonian
- ▶ Matter interaction

MATTER INTERACTION

Hamiltonian with matter interaction in flavor basis ($\omega_{\rm v} = \delta m^2/2E$):

$$\mathbf{H} = \frac{\omega_{\mathbf{v}}}{2} \left(-\cos 2\theta_{\mathbf{v}} \boldsymbol{\sigma}_{3} + \sin 2\theta_{\mathbf{v}} \boldsymbol{\sigma}_{1} \right) + \frac{\lambda(x)}{2} \boldsymbol{\sigma}_{3}$$

- ► Vacuum Hamiltonian
- ► Matter interaction
- $\lambda(x) = \sqrt{2}G_{\rm F}n_{\rm e}(x)$

MSW EFFECT

Hamiltonian in Vacuum

$$\mathbf{H}_{ ext{vacuum}} = rac{\omega_{ ext{v}}\cos 2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_3 + rac{\omega_{ ext{v}}\sin 2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_1$$

$$egin{aligned} \mathbf{H} &= rac{\lambda(x) - \omega_{ ext{v}}\cos2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_{3} + rac{\omega_{ ext{v}}\sin2 heta_{ ext{v}}}{2}oldsymbol{\sigma}_{1} \ &= rac{\omega_{ ext{m}}(x)\cos2 heta_{ ext{m}}(x)}{2}oldsymbol{\sigma}_{3} + rac{\omega_{ ext{m}}(x)\sin2 heta_{ ext{m}}(x)}{2}oldsymbol{\sigma}_{1}, \end{aligned}$$

where

$$\begin{split} \omega_{\rm m}(x) &= \sqrt{\left(\lambda(x) - \omega_{\rm v}\cos 2\theta_{\rm v}\right)^2 + \omega_{\rm v}^2\sin^2 2\theta_{\rm v}},\\ \tan 2\theta_{\rm m}(x) &= \frac{\omega_{\rm v}\sin 2\theta_{\rm v}}{\omega_{\rm v}\cos 2\theta_{\rm v} - \lambda(x)}. \end{split}$$

MSW EFFECT

Constant matter profile λ_0 as an example,

Significance of $\theta_{\rm m}$

Define matter basis $\{|\nu_L\rangle, |\nu_H\rangle\}$

$$\begin{pmatrix} |\nu_{\rm e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} = \begin{pmatrix} \cos\theta_{\rm m} & \sin\theta_{\rm m} \\ -\sin\theta_{\rm m} & \cos\theta_{\rm m} \end{pmatrix} \begin{pmatrix} |\nu_{\rm L}\rangle \\ |\nu_{\rm H}\rangle \end{pmatrix}$$

In matter basis

$$\mathbf{H}_{\mathsf{matter-basis}} = -rac{\omega_{\mathsf{m}}}{2}oldsymbol{\sigma_{\mathsf{3}}}$$

MSW RESONANCE

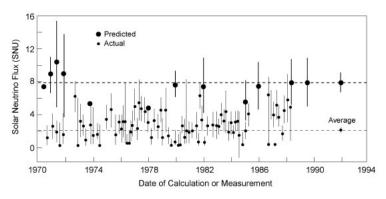
Hamiltonian with Matter Potential

$$\begin{split} \mathbf{H} &= \frac{\lambda(x) - \omega_{v} \cos 2\theta_{v}}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{v} \sin 2\theta_{v}}{2} \boldsymbol{\sigma}_{1} \\ &= \frac{\omega_{m}(x) \cos 2\theta_{m}(x)}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{m}(x) \sin 2\theta_{m}(x)}{2} \boldsymbol{\sigma}_{1} \\ &\tan 2\theta_{m}(x) = \frac{\omega_{v} \sin 2\theta_{v}}{\omega_{v} \cos 2\theta_{v} - \lambda(x)}. \\ \begin{pmatrix} |\nu_{e}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} &= \begin{pmatrix} \cos \theta_{m} & \sin \theta_{m} \\ -\sin \theta_{m} & \cos \theta_{m} \end{pmatrix} \begin{pmatrix} |\nu_{L}\rangle \\ |\nu_{H}\rangle \end{pmatrix} \end{split}$$

Transition Probability

$$P(|\nu_{\rm e}\rangle \to |\nu_{\mu}\rangle) = \sin^2(2\theta_{\rm m})\sin^2(\omega_{\rm m}x)$$

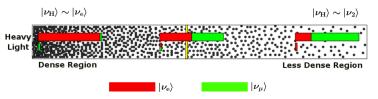
SOLAR NEUTRINO PROBLEM



Chlorine detector (Homestake experiment) results and theory predictions. SNU: 1 event for 10^{36} target atoms per second. Kenneth R. Lang (2010)

MSW EFFECT AND SOLAR NEUTRINOS

$$\begin{split} \mathbf{H} &= \frac{\lambda(\mathbf{x}) - \omega_{\mathrm{v}} \cos 2\theta_{\mathrm{v}}}{2} \boldsymbol{\sigma}_{3} + \frac{\omega_{\mathrm{v}} \sin 2\theta_{\mathrm{v}}}{2} \boldsymbol{\sigma}_{1} \\ \begin{pmatrix} |\nu_{\mathrm{L}}\rangle \\ |\nu_{\mathrm{H}}\rangle \end{pmatrix} &= \begin{pmatrix} \cos \theta_{\mathrm{m}} & -\sin \theta_{\mathrm{m}} \\ \sin \theta_{\mathrm{m}} & \cos \theta_{\mathrm{m}} \end{pmatrix} \begin{pmatrix} |\nu_{\mathrm{e}}\rangle \\ |\nu_{\mu}\rangle \end{pmatrix} \\ \mathbf{H}_{\mathrm{matter-basis}} &= -\frac{\omega_{\mathrm{m}}}{2} \boldsymbol{\sigma}_{3} \end{split}$$



Yellow bar is the resonance point. Red: $|\nu_e\rangle$. Green: $|\nu_{\mu}\rangle$. Adapted from Smirnov, 2003.

MSW EFFECT

Suppose
$$\omega_{\mathrm{v}} = (m_2^2 - m_1^2)/2E < 0,$$

$$\mathbf{H} = \begin{bmatrix} -\frac{\omega_{\mathrm{v}}}{2} \begin{pmatrix} -\cos 2\theta_{\mathrm{v}} & \sin 2\theta_{\mathrm{v}} \\ \sin 2\theta_{\mathrm{v}} & \cos 2\theta_{\mathrm{v}} \end{pmatrix} & +\sqrt{2}G_{\mathrm{F}}n_{\mathrm{e}}(x) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \end{bmatrix}$$

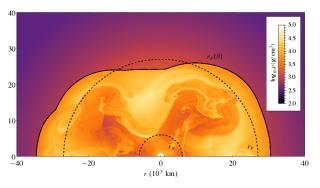
$$\downarrow$$

$$\mathbf{H} = \begin{pmatrix} -\frac{\omega_{\mathrm{v}}}{2} \cos 2\theta_{\mathrm{v}} + \frac{\lambda(x)}{2} \end{pmatrix} \boldsymbol{\sigma}_3 - \frac{\omega_{\mathrm{v}}}{2} \sin 2\theta_{\mathrm{v}} \boldsymbol{\sigma}_1$$

SUPERNOVA MATTER DENSITY PROFILE

Why Do We Care

Astrophysical environments: supernovae, accretion disks etc



Supernova shock and turbulence. E. Borriello, et al (2014)

$$\Delta n_e(r) = \sum_n c_n \sin(k_n r + \phi_n)$$

STIMULATED NEUTRINO OSCILLATIONS

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x)$$

Basis

Background matter basis: Hamiltonian is diagonalized with only background matter profile λ_0 ,

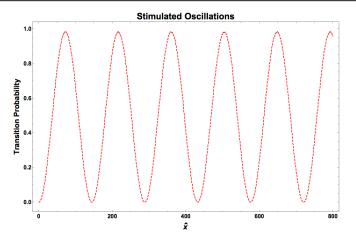
$$H_{background} = -\frac{\omega_m}{2} \sigma_3.$$

Hamiltonian

$$H = \frac{1}{2} \left(-\omega_m + \frac{\delta \lambda(x)}{2} \cos 2\theta_m \right) \sigma_3 - \frac{\delta \lambda(x)}{2} \sin \theta_m \sigma_1.$$

STIMULATED NEUTRINO OSCILLATIONS

P. Krastev and A. Smirnov (1989); J. Kneller et al (2013); K. Patton et al (2014);



Stimulated oscillations. $\lambda(x) = \lambda_0 + A\sin(kx)$ with $\hat{x} = \omega_m x$, $A = 0.1\omega_m$, $k = 0.995\omega_m$, $\theta_m = \pi/6$

OVERVIEW

Introduction

Matter Effect

Understanding Stimulated Oscillations Hamiltonian, and Basis Single Frequency Matter Profile Two-frequency Matter Profile

Summary & Future Work

Understanding Stimulated Oscillations

Matter profile

$$\lambda(x) = \lambda_0 + A\sin(kx),$$

Hamiltonian in Background Matter Basis

$$\mathbf{H} = \frac{1}{2} \left(-\omega_{m} + \frac{\delta \lambda(\mathbf{x})}{\delta \lambda(\mathbf{x})} \cos 2\theta_{m} \right) \boldsymbol{\sigma}_{3} - \frac{\frac{\delta \lambda(\mathbf{x})}{2}}{2} \sin \theta_{m} \boldsymbol{\sigma}_{1}.$$

A Better Basis

Define new basis $\{|\tilde{\nu}_L\rangle,|\tilde{\nu}_H\rangle\}$ is related to background matter basis $\{|\nu_L\rangle,|\nu_H\rangle\}$ through

$$\begin{pmatrix} |\nu_L\rangle \\ |\nu_H\rangle \end{pmatrix} = \begin{pmatrix} e^{-i\eta(x)} & 0 \\ 0 & e^{i\eta(x)} \end{pmatrix} \begin{pmatrix} |\tilde{\nu}_L\rangle \\ |\tilde{\nu}_H\rangle \end{pmatrix},$$

where

$$\eta(x) - \eta(0) = -\frac{\omega_{\rm m}}{2}x + \frac{\cos 2\theta_{\rm m}}{2} \int_0^x \frac{\delta \lambda(\tau)}{d\tau}.$$

Hamiltonian in new basis

$$\widetilde{\mathbf{H}} = -\frac{\delta \lambda(\mathbf{x})}{2} \sin 2\theta_{\mathrm{m}} \begin{pmatrix} 0 & e^{2i\eta(\mathbf{x})} \\ e^{-2i\eta(\mathbf{x})} & 0 \end{pmatrix} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

Hamiltonian in New Basis

$$\begin{split} h &\equiv -\frac{\delta\lambda(x)}{2}e^{2i\eta(x)} \\ &= \frac{i}{4}\left[\exp\left(i(k+\omega_{\rm m})x + i\cos2\theta_{\rm m}\frac{A}{k}\cos(kx)\right) \\ &-\exp\left(i(-k+\omega_{\rm m})x + i\cos2\theta_{\rm m}\frac{A}{k}\cos(kx)\right)\right] \end{split}$$

RABI OSCILLATION

Rabi Oscillation

Hamiltonian

$$\begin{pmatrix} -\omega_0/2 & \alpha\omega_0e^{i\omega x} \\ \alpha\omega_0e^{-ikx} & \omega_0/2 \end{pmatrix},$$

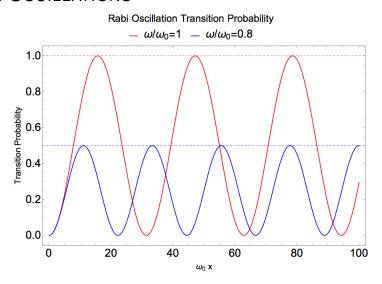
$$E_2 = \frac{\omega_0}{2}$$

Incoming light

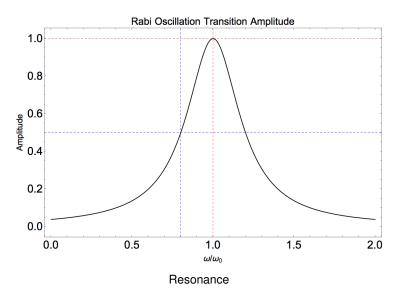
$$E_1 = -\frac{\omega_0}{2}$$

Frequency : ω

RABI OSCILLATIONS



RABI OSCILLATIONS



Off-diagonal Term in Our System

$$\widetilde{\mathbf{H}} = \begin{pmatrix} 0 & h \\ h^* & 0 \end{pmatrix}$$

$$h \propto \left[\exp \left(i(k + \omega_{\rm m})x + i\cos 2\theta_{\rm m} \frac{A}{k}\cos(kx) \right) - \exp \left(i(-k + \omega_{\rm m})x + i\cos 2\theta_{\rm m} \frac{A}{k}\cos(kx) \right) \right]$$

Jacobi-Anger expansion

$$e^{i\beta\cos(kx)} = \sum_{n=-\infty}^{\infty} i^n J_n(\beta) e^{inkx},$$

where $J_n(\beta)$ are Bessel's functions of the first kind.

Scaled Quantities

Characteristic scale: $\omega_{\rm m}$

- $\rightarrow \hat{A} = A/\omega_{\rm m}$
- $\hat{k} = k/\omega_{\rm m}$
- $\hat{\mathbf{x}} = \omega_{\mathrm{m}} \mathbf{x}$
- $\blacktriangleright \hat{h} = h/\omega_{\rm m}$

Rotation Wave Approximation

The off-diagonal element of Hamiltonian

$$\widetilde{\mathbf{H}} = \sum_{n = -\infty}^{\infty} \begin{pmatrix} 0 & \frac{1}{2} \hat{B}_n e^{i(n\hat{k} - 1)\hat{x}} \\ \frac{1}{2} \hat{B}_n^* e^{-i(n\hat{k} - 1)\hat{x}} & 0 \end{pmatrix}$$

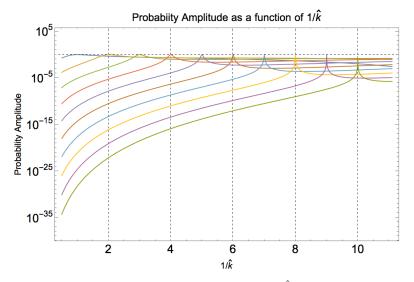
where $\hat{B}_n = -(-i)^n n\hat{k} \tan 2\theta_{\rm m} J_n(\hat{A}\cos 2\theta_{\rm m}/\hat{k})$.

Transition Probability

$$P_{ ext{L}
ightarrow ext{H}}^{(n)} = rac{\left|\hat{B}_{n}
ight|^{2}}{\left|\hat{B}_{n}
ight|^{2} + (n\hat{k} - 1)^{2}} \sin^{2}\left(rac{q^{(n)}}{2}x
ight),$$

where

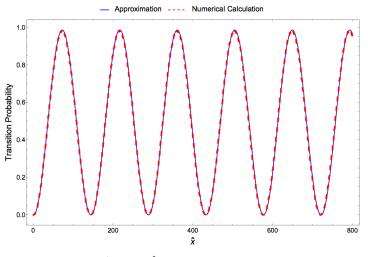
$$q^{(n)}=\sqrt{\left|\Gamma^{(n)}\right|^2+(n\hat{k}-1)^2},\quad ext{frequency of oscillations}$$
 $\Gamma^{(n)}=\left|\hat{B}_n\right|,\quad ext{width of resonance }(n\hat{k} ext{ as parameter})$



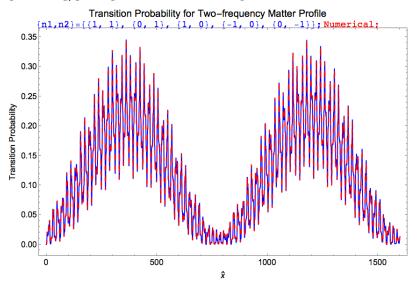
Resonances of different $n = 1/\hat{k}$.

SINGLE FREQUENCY MATTER PROFILE

Stimulated Oscillations



$$\hat{A} = 0.1, \, \hat{k} = 0.995, \, \theta_{\rm m} = \pi/6$$



$$\lambda(x) = \lambda_0 + A_1 \sin(k_1 x) + A_2 \sin(k_2 x)$$
. $\hat{k}_1 = 0.3$, $\hat{k}_2 = 0.7$, $A_1 = A_2 = 0.1$, $\theta_m = \pi/5$.

OVERVIEW

Introduction

Matter Effect

Understanding Stimulated Oscillations

Summary & Future Work

SUMMARY & FUTURE WORK

- The fact that neutrino flavor sates are not mass states causes vacuum oscillations.
- ► MSW resonance happens when matter potential cancels out the vacuum diagonal elements of the Hamiltonian.
- ► Even matter profile doesn't match MSW requirement, variation in matter profile can cause resonances.
- Single frequency perturbations in matter profile is a combination of many Rabi oscillations.
- How to understand and calculate systems with multi-frequency matter profile (turbulence).
- ► Combine periodic or even turbulent matter profile with neutrino self-interaction.

ACKNOWLEDGEMENT

I am very thankful to my advisor Professor Huaiyu Duan, and everyone in our group Dr. Sajad Abbar, and Dr. Shashank Shalgar, for all the help in both research and life.

Supported by DOE EPSCoR grant #DE-SC0008142 at UNM.

BACKUP SLIDES

BACKUP SLIDES

PARAMETERS USED FOR VACUUM OSCILLATIONS

$$\begin{array}{l} \theta_{12}=33.36/180\pi;\,\theta_{13}=8.66/180\pi;\,\theta_{23}=40/180*\pi;\,\delta_{cp}=0;\\ m_1^2=0.01;\,m_2^2=m_1^2+0.000079;\,E=1\text{MeV} \end{array}$$

SINGLE FREQUENCY MATTER PROFILE

Why Does It Work?

$$J_n(n \operatorname{sech} \alpha) \sim \frac{e^{-n(\alpha - \tanh \alpha)}}{\sqrt{2\pi n \tanh \alpha}}, \quad ext{for large } n$$

 \Rightarrow

$$\Gamma \propto \hat{B}_n \propto rac{e^{-n(lpha - anh lpha)}}{\sqrt{2\pi n anh lpha}}$$

Small perturbation \Rightarrow Small $\hat{A} \Rightarrow$ Large $\alpha \Rightarrow$ Drops fast at large n.

Matter Profile

$$\lambda(x) = \lambda_0 + \delta\lambda(x), \quad \delta\lambda(x) = A_1 \sin(k_1 x) + A_2 \sin(k_2 x).$$

TWO-FREQUENCY MATTER PROFILE $\hat{h} = \sum_{i=1}^{\infty} \frac{1}{2} \hat{B}_{n} e^{i(n\hat{k}-1)\hat{x}}$,

Hamiltonian Off-diagonal Element

Apply Jacobi-Anger expansion,

$$\hat{h} = \sum_{n_1 = -\infty}^{\infty} \sum_{n_2 = -\infty}^{\infty} \frac{1}{2} \hat{B}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1)\hat{x}},$$

where

$$\hat{B}_{n_1,n_2}(\hat{k}_1,\hat{k}_2)
= -(-i)^{n_1+n_2}(n_1\hat{k}_1 + n_2\hat{k}_2)J_{n_1}\left(\frac{\hat{A}_1\cos 2\theta_{\rm m}}{\hat{k}_1}\right)J_{n_2}\left(\frac{\hat{A}_2\cos 2\theta_{\rm m}}{\hat{k}_2}\right)$$

Which terms are important?

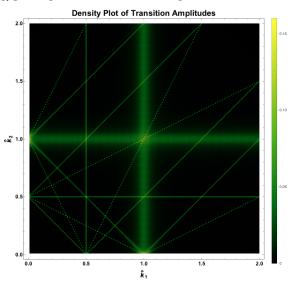
Resonance Lines

There are still resonances, i.e., (almost) zero phases, on lines

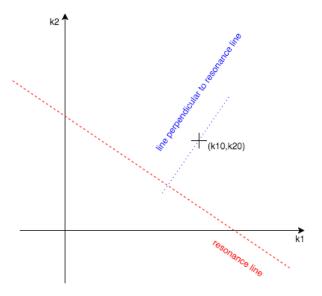
$$n_{1,0}\hat{k}_1 + n_{2,0}\hat{k}_2 - 1 = 0$$

in $\{\hat{k}_1, \hat{k}_2\}$ plane. \Rightarrow Resonance width for each point on resonance lines.

$\text{TWO-FREQUENCY MATTER PROFIL}^{\hat{h}} \bar{\bar{E}}^{\sum_{n_1} \sum_{n_2} \frac{1}{2} \hat{h}_{n_1, n_2}(\hat{k}_1, \hat{k}_2) e^{i(n_1 \hat{k}_1 + n_2 \hat{k}_2 - 1) \hat{x}},$



Density plot of transition amplitudes calculated using only one term out of the whole summation in Hamiltonian. $n_1, n_2 \in [-2, 2]$



Resonance line, distance to resonance, and width

Width

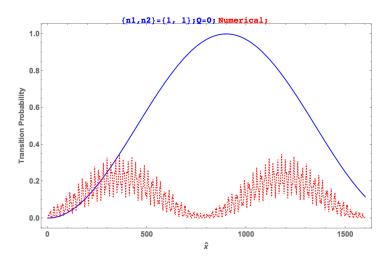
$$\Gamma_2 = rac{\hat{B}_{n_1,n_2}(\hat{k}_{1,\mathrm{intercept}},\hat{k}_{2,\mathrm{intercept}})}{\sqrt{n_1^2 + n_2^2}}.$$

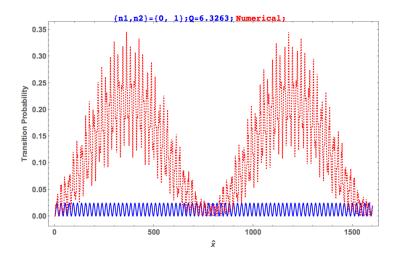
Distance to Resonance Line

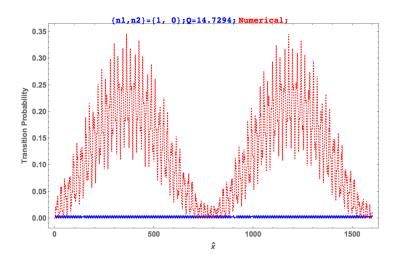
$$d = \frac{|n_1\hat{k}_{10} + n_2\hat{k}_{20} - 1|}{\sqrt{n_1^2 + n_2^2}}.$$

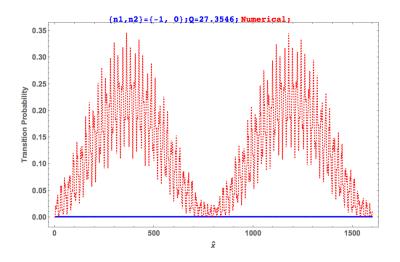
Distance to Resonance Width Ratio

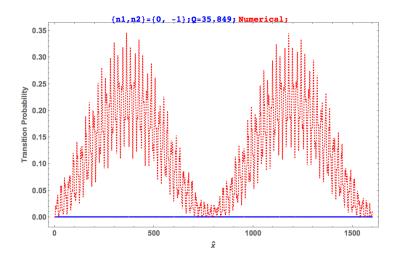
$$Q_2 = \frac{d}{\Gamma_2}.$$

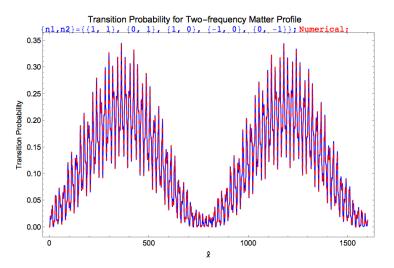












BESSEL'S FUNCTION

$$J_n(\beta) = \sum_{m=0}^{\infty} \frac{(-1)^m}{m! \Gamma(m+n+1)} \left(\frac{\beta}{2}\right)^{2m+n}$$

REFERENCES I