

Recap

Thrust for cv fixed or moving @ const \underline{u}

$$T = \dot{m} u_e + (p_e - p_a) A_e$$

Equivalent exhaust velocity $u_{eq} \equiv u_e + \frac{p_e - p_a}{\dot{m}} A_e$

such that $T = \dot{m} u_{eq}$

Impulse $I = \int_0^{t_b} \overset{\leftarrow \text{burn time}}{T} dt = u_{eq} \int_0^{t_b} \dot{m} dt = u_{eq} M_p$

Specific impulse $I_{sp} = \frac{I}{M_p g_e} = \frac{u_{eq}}{\underline{g_e}}$ earth gravity @ SL

Typical values for chemical rockets

$200s \leq I_{sp} \leq 500s \rightarrow 2000 \text{ m/s} \rightarrow \text{"cook'n"}$

3. Acceleration

Recall momentum eq. for accelerating cv

$$\frac{d}{dt} \int_{cv} \rho \underline{u}_{xyz} dV + \int_{cs} \rho \underline{u}_{xyz} \underline{u}_{xyz} \cdot d\underline{A} = - \int_{cs} p d\underline{A} + \int_{cv} \rho \underline{g} dV - \underbrace{\int_{cv} \rho \underline{a}_{rel} dV}_{\substack{\text{accel. of cv rel.} \\ \text{to ground XYZ}}} + \underline{F_e}$$

→ Rectilinear trajectory

→ Project in 1 direction, drop xyz subscript

i) $\frac{d}{dt} \int_{cv} \rho \underline{u}_{xyz} dV = 0$ b/c in tanks, $\rho \neq \text{const}$, $\underline{u}_{xyz} \approx 0$
in TC, ρ/\underline{u}_{xyz} vary in space but steady

ii) $\int_{cs} \rho \underline{u}_{xyz} \underline{u}_{xyz} \cdot d\underline{A} = \dot{m} u_e (-\hat{i})$

on exit, $p = p_e + p_a - p_a = p_a + (p_e - p_a)$

iii) $\int_{cs} -p d\underline{A} = -(p_e - p_a) A_e (+\hat{i})$

iv) Project \underline{g} in \hat{i} : $\int_{cv} \rho \underline{g} dV \rightarrow g \cos \theta M(t) (-\hat{i})$

v) $\underline{F_e} = 0 (-\hat{i})$

$$v \dot{v} = - \int_{cv} \rho a_{rel} dV = - \underbrace{a_{rel}}_{\frac{du}{dt}(i)} \underbrace{\int_{cv} \rho dV}_{M(t)} = - \frac{du}{dt}(i) M(t)$$

Reassemble:

$$-u_e \dot{m} = (p_e - p_a) A_e - M(t) g \cos \theta - D - M(t) \frac{du}{dt}$$

$$\underbrace{u_e \dot{m} + (p_e - p_a) A_e}_T = M(t) g \cos \theta + D + M(t) \frac{du}{dt}$$

(Full form)

(A) Simplified version

$$g = D = 0$$

$$\rightarrow \underbrace{u_e \dot{m} + (p_e - p_a) A_e}_{\dot{m} u_{eq}} = M(t) \frac{du}{dt}$$

$$\dot{m} u_{eq} = M(t) \frac{du}{dt}$$

Relate \dot{m} & $M(t)$

$$\underbrace{\dot{m}}_{\substack{< 0 \\ \text{b/c exiting}}} = - \underbrace{\frac{dM(t)}{dt}}_{\substack{< 0 \\ \text{b/c } M \text{ decreases w/time}}} \rightarrow - \frac{dM(t)}{dt} u_{eq} = M(t) \frac{du}{dt}$$

assume constant

$$- \frac{dM(t)}{M(t)} u_{eq} = du$$

$$\rightarrow -u_{eq} \int_0^t \frac{dM(t)}{M(t)} = \int du$$

$$\Delta u = -u_{eq} \ln \frac{M(t)}{M(t=0)} = \boxed{u_{eq} \ln \frac{M_0}{M(t)} = \Delta u} \quad \text{Rocket equation}$$

constant thrust accelerates more as mass ↓

can expand $M(t) = M_0 - \dot{m} t$ assume $\dot{m} = \text{const.}$

$$\Delta u = u_{eq} \ln \frac{M_0}{M_0 - \dot{m} t}$$

If Δu is attained over entire burn, $t \rightarrow t_b$, $M(t_b) = M_s + M_i$

$$\text{Recall } R = \frac{M_0}{M_b} = \frac{M_0}{M_0 - M_p} = \frac{M_0}{M_s + M_i} - M_0 - M_p$$

$$\Delta u(t_b) = u_{eq} \ln R$$

Remarks

- i) It looks like if M_p is large enough, can get $\Delta u(t_b)$ larger than u_{eq} by "any" factor desired (we'll rectify this)
- ii) Since Δu scales w/ u_{eq} , we want u_{eq} as large as possible for fixed \bar{m} , since $\bar{m} = \rho_e u_e t_e$, large u_e requires low ρ_e