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Bibliography Report

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**Making Flips With Quadrotors In Constrained
Environments**

Jury

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Abstract

Within the rapidly growing aerial robotics market, one of the most substantial challenges in the quadrotor community is performing aggressive maneuvers, especially multi-flip maneuvers. A proper physical definition of the issue is not addressed by the current approaches in the field and several key aspects of this maneuver are still overlooked. It can be shown, in particular, that making a flip with a quadrotor means crossing the parallel singularity of the dynamic model. The aim of the master thesis is to explore the possibility of defining aggressive trajectories for quadrotors on the basis of their dynamic model degeneracy analysis and to adapt various strategies to control the robot in a closed loop. In addition, the possibility to perform the aggressive maneuver in constrained environments will also be investigated. Therefore, the analysis will be extended from the previous studied to create general feasible trajectories that will allow quadrotors to perform aggressive flip maneuvers while passing through a constrained environment and while guaranteeing a satisfactory degree of robustness to the uncertainties of the dynamic model.

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Notations

b	Thrust factor
l	horizontal distance: From the center of the propeller to the CoG
Ω	Spinning speed of a propeller
C_{R_m}	Rolling moment coefficient
C_T	Thrust coefficient
H	Hub force
CT	Continuous time
DT	Discrete time

Abbreviations

IGE In Ground Effect
OGE Out of Ground Effect

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Introduction

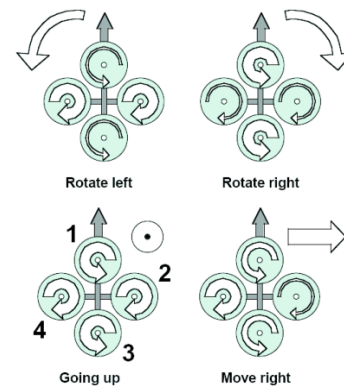
The aim of this section is to provide a general summary of the robotic platform that is used for this master thesis and to illustrate the main objective of the research work. In specific, in the sections below, quadrotors and parallel robots are briefly presented.

The quadrotor platform

A quadrotor is a type of unmanned aerial vehicle (UAV) with four rotors and six degrees of freedom. Typically, drones have a small size and low inertia which allows it to be controlled by simple flight control systems. It is typically designed in a cross-configuration such that the electronics are held in the center of the platform and the rotors are placed at the borders. An example of a real quadrotor, namely the DJI Phantom, is shown in fig. 1a. The quadrotor is typically built in a way such that a pair of opposite rotors rotate clockwise, whereas the other pair of rotors rotates in counter-clockwise. The attitude and the position of the drone are controlled by changing the spinning speed of the rotors. An example is shown in figure 1b.



(a) A DJI Phantom quadcopter (UAV)¹



(b) Typical quadrotor configuration. The width of the arrows is proportional to the angular speed of the propellers.[1]

Figure 1: A commercial quadrotor platform, with a typical quadrotor configuration.

The distinctive mechanical design of the quadrotor permits the actuation system to control all of the six degrees of freedom, even though it is under-actuated. This is due to the fact that the rotational and translational dynamics are tightly coupled. Thus, all the translational and rotational motions can be carried off by properly controlling the magnitude and direction of the spinning speed of the rotors.

¹https://en.wikipedia.org/wiki/Quadcopter#/media/File:Quadcopter_camera_drone_in_flight.jpg, accessed on 01/08/2021.

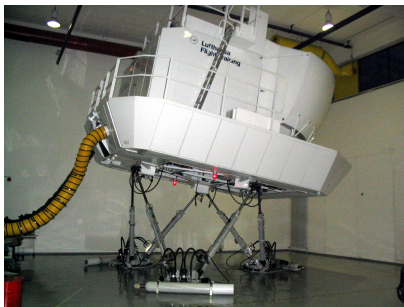
Over the last few years, quadrotors have gained a large popularity in academia and in the industry. This is due to several reasons, such as:

1. Quadrotors are very simple to design and they can be easily assembled using relatively cheap components.
2. As quadrotors became more and more affordable and dependable, the number of quadrotors real-world applications has grown significantly. They are being used for aerial photography, agriculture, surveillance, inspection tasks, in addition to many other uses as well.
3. Quadrotors are quite agile and maneuverable during flight. Especially when compared to other types of unmanned aerial vehicles (UAVs).

However, on the main challenges in the quadrotors community is the capability to design control and planning methods that will allow the quadrotors to carry out aggressive maneuvers. The fast dynamics associated with typically small dimensions of such agile quadrotors, in addition to several aerodynamic effects that will become important during aggressive flight maneuvers, are just a few of the main problems that are faced during the system control design. Moreover, accurate tracking of the provided trajectory is a very big issue in the case of aggressive maneuvers when the rotors are commanded high speeds and accelerations, which will cause rotors to become saturated and may also cause delays.

Parallel manipulators

A parallel manipulator is a mechanical system that consists of two connected platforms, the fixed platform and the moving platform. The latter is linked to the fixed platform thanks to at least two serial chains that are working in parallel. When compared to serial manipulators, parallel manipulators are more accurate and rigid. In addition, the ability to install the motors next to the fixed platform is a very important feature for parallel manipulators. Moreover, parallel manipulators can be used in a wide variety of applications that demand precision and high payload combined with high speed.[3]



(a) Gough-Stewart used for a flight-simulator application.²



(b) The "PAR4" 4 degrees of freedom, high-speed, parallel robot prototype.³

Figure 2: Two examples of parallel robots.

¹https://en.wikipedia.org/wiki/Stewart_platform#/media/File:Simulator-flight-compartment.jpeg, accessed on 01/08/2021.

²https://en.wikipedia.org/wiki/Parallel_manipulator#/media/File:Prototype_robot_parallel%20PAR4.jpg, accessed on 01/08/2021.

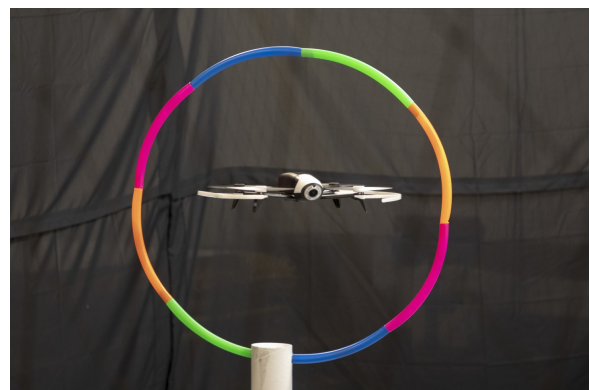
However, parallel manipulators are subject to singularities, which can lead to big problems in the robot workspace in case they were not handled correctly. Thus, the study of the singular configurations of parallel manipulators is very important. Because, even just before reaching a singularity, the performance of the parallel manipulator will decrease dramatically. Moreover, the robot may lose the ability of moving in a certain direction, gain uncontrollable motions and the mechanism could even break. The main difference between serial and parallel manipulators is that singularity configurations may also appear inside the robot workspace (depending on the dimensions of the robot) and not just at the boundaries of the robot workspace, which can significantly decrease the area of the robot workspace. As a result, many works have been developed by robotics researchers in order to allow parallel manipulator manipulators to safely cross these singularities by using trajectory planning and specific control methods.

The goal of this thesis

This master thesis lies at the intersection of parallel robotics and aerial robotics. The two fields may seem very different from each other. However, quadrotors can be seen as a particular case of a parallel manipulator. In fact, a parallel manipulator is made up of a wrench system, applied by the robot limbs on the moving platform. And, this wrench system will define the motion of the moving platform. In the same manner, each propeller in a quadrotor can be considered as limb of a parallel robot and the moving platform to be controlled can be considered as the body of the drone. Specifically, the goal of this master thesis is to study a distinct class of aggressive maneuvers for quadrotors, namely multi-flip maneuvers. By doing multi-flip maneuvers, full rotations around one or more axes of the body of the quadrotor can be done. In addition, the quadrotor must also do the flips in a constrained environment.



(a) Quadrotor performing a triple flip.[4]



(b) Quadrotor going through a loop.¹

Figure 3: Representation of the issues to be tackled in this master thesis.

¹<https://newatlas.com/drones/muscle-signals-drone-control/#gallery:2>, accessed on 01/08/2021.

Outline of the work

The rest of the bibliography is structured as follows:

- Chapter 1** is devoted to introduce the system modeling of quadrotors. Specifically, a simplified dynamic model of the quadrotor will be presented by using Euler-Lagrange formalism. Then, moving on from the simple dynamic model, a more detailed dynamic model will be presented by using the Newton-Euler formalism. Finally, the state-space model of the quadrotor will also be derived.
- Chapter 2** provides an overview of state of the art in quadrotor control in addition to introducing the different potential control methods that can be used during the master thesis in order to properly control the quadrotor.
- Chapter 3** provides detailed explanations of how multi-flip maneuvers can be handled. Then, the link between a quadrotor performing a flip and a parallel robot crossing a singularity will be explained. In the end, a literature review is provided in order to show how the problem is tackled by different researches.
- Chapter 4** is devoted to trajectory optimization. By using trajectory optimization, it will be possible to create feasible trajectories for quadrotors to perform the aggressive maneuvers in constrained environments.

System Modeling

The goal of this chapter is to present the dynamic model of the quadrotor. The mathematical notation and the physics of the quadrotor platform are expressed using the Newton-Euler formalism. Then, the state-space model that will be coded on the controller of the quadrotor will be derived.

1.1 Concepts and Generalities

The dynamic model of the quadrotor will be derived based on the following assumptions:

- The quadrotor has a rigid structure.
- The quadrotor has a symmetrical structure.
- The center of gravity (CoG) and the fixed frame at the center of the body are assumed to be coincident.
- The propellers of the quadrotor are assumed to be rigid.
- The thrust and drag forces are assumed to be proportional to the square of the spinning speed of each propeller.

The helicopter is a complex mechanical system, it gathers many physical effects from the domain of mechanics and aerodynamics [5]. Thus, all the significant effects including the gyroscopic effects must be considered in the modeling of the quadrotor. A small list of the most important effects that a helicopter is subject to [6] are briefly described in table 1.1:

Table 1.1: The main physical effects that the helicopter is subject to.

Effect	Source	formulation
Aerodynamic effects	Rotation of propeller	$C\Omega^2$
	Flapping of blades	
	Change in propeller spinning speed	
Inertial counter torques		$J\dot{\Omega}$
Gravitational effect	Position of the center of mass	
Gyroscopic effects	Orientation change of the rigid body	$I\dot{\theta}\psi$
	Orientation change of the propeller plane	$J\Omega_r\dot{\theta}, \dot{\phi}$
Friction	All helicopter motions	$C\dot{\phi}, \dot{\theta}, \dot{\psi}$

1.2 Modelling with Euler-Lagrange Formalism

The dynamics of the rotation of a simple quadrotor are modeled using the Euler-Lagrange Formalism in this section. A fixed frame E for the world frame and body fixed frame B for the quadrotor are considered as represented in figure 1.1. The orientation of the quadrotor frame in space is provided by a rotation R from B to E , where $R \in SO3$ is a 3×3 rotation matrix.

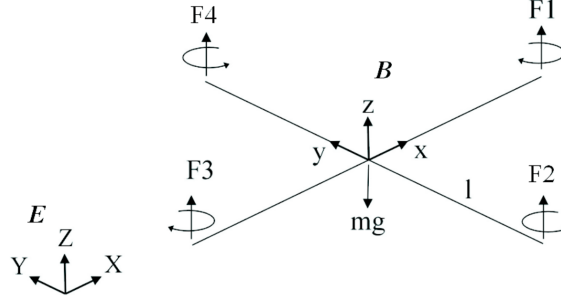


Figure 1.1: Coordinate system of a simple quadrotor. [1]

1.2.1 Kinematics

For any point of the body frame of the quadrotor expressed in the fixed world frame, the following can be written (c : \cos , s : \sin):

$$\begin{cases} r_X = (c\psi c\theta)x + (c\psi s\theta s\phi - s\psi c\phi)y + (c\psi s\theta c\phi + s\psi s\phi)z \\ r_Y = (s\psi c\theta)x + (s\psi s\theta s\phi + c\psi c\phi)y + (s\psi s\theta c\phi - c\psi s\phi)z \\ r_Z = (-s\theta)x + (c\theta s\phi)y + (c\theta c\phi)z \end{cases} \quad (1.1)$$

Thus, the velocities can be derived by differentiation 1.1, and the squared magnitude of the squared velocity can be expressed as follows for any point:

$$v^2 = v_X^2 + v_Y^2 + v_Z^2 \quad (1.2)$$

1.2.2 Energy

Assuming that the matrix of inertia is diagonal, then from equation 1.2, the expression of the kinetics energy can be calculated:

$$T = \frac{1}{2}I_{xx}(\dot{\phi} - \dot{\psi}s\theta)^2 + \frac{1}{2}I_{yy}(\dot{\theta}c\phi + \dot{\psi}s\phi c\theta)^2 + \frac{1}{2}I_{zz}(\dot{\theta}s\phi - \dot{\psi}c\phi)^2 \quad (1.3)$$

Using the formula of the potential energy, equation 1.3 can be expressed in the fixed world frame as:

$$V = \int x dm(x)(-gs\theta) + \int y dm(y)(gs\phi c\theta) + \int z dm(z)(gc\phi c\theta) \quad (1.4)$$

1.2.3 Equation of Motion

By using the Euler-Lagrange formalism:

$$L = T - V \quad , \quad \Gamma_i = \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} \quad (1.5)$$

Where L , Γ_i and \dot{q}_i are the Lagrangian, the generalized forces and the generalized coordinates respectively. Thus, the equations of motion can be expressed as follows:

$$\begin{cases} I_{xx}\ddot{\phi} = \dot{\theta}\dot{\psi}(I_{yy} - I_{zz}) \\ I_{yy}\ddot{\theta} = \dot{\phi}\dot{\psi}(I_{zz} - I_{xx}) \\ I_{zz}\ddot{\psi} = \dot{\phi}\dot{\theta}(I_{xx} - I_{yy}) \end{cases} \quad (1.6)$$

Moreover, the torques that are nonconservative and acting on the quadrotor, are due to two different causes. First, it is due to thrust of each rotor pairs in figure 1.1:

$$\begin{cases} \tau_x = bl(\Omega_4^2 - \Omega_2^2) \\ \tau_y = bl(\Omega_3^2 - \Omega_1^2) \\ \tau_z = bl(\Omega_1^2 - \Omega_2^2 + \Omega_3^2 - \Omega_4^2) \end{cases} \quad (1.7)$$

Second, it is also due to the gyroscopic effect which is the result of the rotation of the propellers:

$$\begin{cases} \tau'_x = J_r\omega_y(\Omega_1 + \Omega_3 - \Omega_2 - \Omega_4) \\ \tau'_y = J_r\omega_x(\Omega_2 + \Omega_4 - \Omega_1 - \Omega_3) \end{cases} \quad (1.8)$$

1.2.4 The Derived Dynamic Model

The dynamic model of the quadrotor which describes the rotations of roll, pitch and yaw consists of three terms:

1. The actuator torques.
2. The gyroscopic effects that are due to the rotation of the rigid body.
3. The gyroscopic effects that are due to rotation of the propeller that is coupled with the rotation of the body.

Thus, the dynamic model of the quadrotor is:

$$\begin{cases} I_{xx}\ddot{\phi} = \dot{\theta}\dot{\psi}(I_{yy} - I_{zz}) - J\dot{\theta}\Omega_r + \tau_x \\ I_{yy}\ddot{\theta} = \dot{\phi}\dot{\psi}(I_{zz} - I_{xx}) + J\dot{\phi}\Omega_r + \tau_y \\ I_{zz}\ddot{\psi} = \dot{\phi}\dot{\theta}(I_{xx} - I_{yy}) + \tau_z \end{cases} \quad (1.9)$$

1.2.5 Rotor Dynamics

DC motors are used to drive the rotors of a quadrotor. So, the established equations of a DC motor are the following:

$$\begin{cases} L \frac{di}{dt} = u - R_{mot}i - k_e \omega_m \\ J_m \frac{d\omega_m}{dt} = \tau_m - \tau_d \end{cases} \quad (1.10)$$

Since small motors are used in which they also have little inductance, then the second order equation of the DC motor dynamics is given by:

$$J_m \frac{d\omega_m}{dt} = -\frac{k_m^2}{R_{mot}} \omega_m - \tau_d + \frac{k_m}{R_{mot}} u \quad (1.11)$$

When the gearbox and the propeller models are introduced, then equation 1.12 becomes:

$$\begin{cases} \dot{\omega}_m = -\frac{1}{\tau} \omega_m - \frac{d}{\eta r^3 J_t} \omega_m^2 + \frac{1}{k_m \tau} u \\ \frac{1}{\tau} = \frac{k_m^2}{R J_t} \end{cases} \quad (1.12)$$

Moreover, linearization of equation 1.12 can be done around an operation point $\dot{\omega}_0$ to the form $\dot{\omega}_m = -A\omega_m + Bu + C$ with:

$$A = \left(\frac{1}{\tau} + \frac{2d\omega_0}{\eta r^3 J_t} \right), \quad B = \left(\frac{1}{k_m \tau} \right), \quad C = \left(\frac{d\omega_0^2}{\eta r^3 J_t} \right) \quad (1.13)$$

1.3 Modeling with Newton-Euler Formalism

The model above was derived in succession as shown in papers [7, 8, 9]. The dynamic equations below includes contain rolling moments R_m , hub forces H and various aerodynamic effects. Thus, this is a more realistic dynamic model, especially when the quadrotor flies in a forward manner. With the previous versions of the dynamic model, it was required to to moderately tune the control parameters in order to have experiments that are successful.

The dynamic model expressed in the Newton-Euler formalism of a rigid body that is subject to external forces acting on the center of mass is expressed as follows [10]:

$$\begin{bmatrix} mI_{3 \times 3} & 0 \\ 0 & I \end{bmatrix} \begin{bmatrix} \dot{V} \\ \dot{\omega} \end{bmatrix} + \begin{bmatrix} \omega \times mV \\ \omega \times I\omega \end{bmatrix} = \begin{bmatrix} F \\ \tau \end{bmatrix} \quad (1.14)$$

Considering a fixed world frame E and a fixed body frame B on the quadrotor as shown in figure 1.2. Then, by the use of the Euler angles, the orientation of the rigid body of the quadrotor in space is expressed by a rotation R from B to E , where $R \in SO3$ is a rotation matrix.

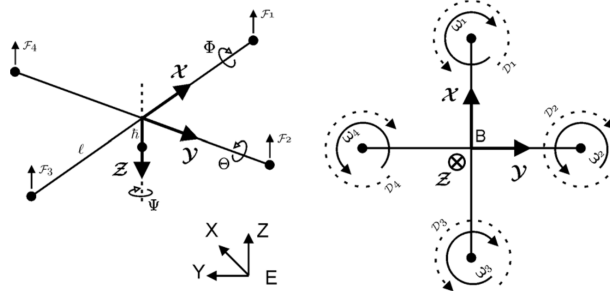


Figure 1.2: Coordinate system of a simple quadrotor. [1]

1.3.1 Aerodynamic Forces and Moments

The aerodynamic forces and moments are computed using a mix of blade element and momentum theory [11]. This is based off of the work of Gary Fay during the project of the Mesicopter [12]. For a simpler readings of the equations provided below, some symbols are recalled:

σ : solidity ration	λ : inflow ratio
a : lift slope	v : induced velocity
μ : rotor advance ration	ρ : air density

Thrust Force

The thrust force is due to all the vertical forces that the blade elements are subject to.

$$\begin{cases} T = C_T \rho A (\Omega R_{rad})^2 \\ \frac{C_T}{\sigma a} = (\frac{1}{6} + \frac{1}{4} \mu^2) \theta_0 - (1 + \mu^2) \frac{\theta_{tw}}{8} - \frac{1}{5} \lambda \end{cases} \quad (1.15)$$

Hub Force

The hub force is due to all the horizontal forces that the blade elements are subject to.

$$\begin{cases} H = C_H \rho A (\Omega R_{rad})^2 \\ \frac{C_H}{\sigma a} = (\frac{1}{4a} \mu \overline{C_d} + \frac{1}{4} \lambda \mu (\theta_0 - \frac{\theta_{tw}}{2})) \end{cases} \quad (1.16)$$

Drag Moment

The drag moment about the rotor shaft is due to the aerodynamic forces that the blade elements are subject to. The horizontal forces that are acting on the rotor are multiplied by the moment arm and integrated over the rotor. The drag moment gives the required power to spin the rotor.

$$\begin{cases} Q = C_Q \rho A (\Omega R_{rad})^2 R_{rad} \\ \frac{C_Q}{\sigma a} = \frac{1}{8a} (1 + \mu^2) \overline{C_d} + \lambda (\frac{1}{6} \theta_0 - \frac{1}{8} \theta_{tw} - \frac{1}{4} \lambda) \end{cases} \quad (1.17)$$

Rolling moment

The rolling moment of a propeller occurs when the blade that is advancing is producing more lift than the blade that is retreating in forwarding flight. It is the integration of the lift of every single section that is acting at a given radius over the entire rotor. The reader should notice that the rolling moment is not the same as the propeller radius, or the overall rolling moment which is due to other effects or the rotation matrix R . So, there should not be any confusion.

$$\begin{cases} R_m = C_{R_m} \rho A (\Omega R_{rad})^2 R_{rad} \\ \frac{C_{R_m}}{\sigma a} = -\mu \left(\frac{1}{6} \theta_0 - \frac{1}{8} \theta_{tw} - \frac{1}{8} \lambda \right) \end{cases} \quad (1.18)$$

Ground Effect

When operating near the ground (at a height equivalent to half the diameter of the rotor), helicopters experience thrust augmentation which is caused by greater efficiency of the rotor. This is linked to a decrease in the velocity of induced airflow. Moreover, this is called Ground Effect. Different approaches to deal with this effect can be found in literature, for example, adaptive techniques can be used [13]. However, the objective is to find a model that is simple and mainly captures the change in the velocity of the induced inflow. Cheeseman [14] states (reached from the images method [15]) that if the power is constant ($T_{OGE} v_{i,OGE} = T_{IGE} v_{i,IGE}$), the generated velocity at the center of the rotor by its image is $\delta v_i = A v_i / 16 \pi z^2$. Cheesman acquired the relation (1.19) by using the assumption that both v_i and δv_i are constant over disk, which results in $v_{i,IGE} = v_i - \delta v_i$.

$$\frac{T_{IGE}}{T_{OGE}} = \frac{1}{1 - \frac{R_{rad}^2}{16z^2}} \quad (1.19)$$

An alternative way to move forward is to consider that the inflow ratio of the inflow ratio is $\lambda_{IGE} = (v_{i,OGE} - \delta v_i - \dot{z}) / \Omega R_{rad}$, where the change of the velocity of the induced inflow is $\delta v_i = v_i / (4z / R_{rad})^2$. Then, the thrust coefficient (1.15) IGE can be rewritten as:

$$\begin{cases} T_{IGE} = C_T^{IGE} \rho A (\Omega R_{rad})^2 \\ \frac{C_T^{IGE}}{\sigma a} = \frac{C_T^{OGE}}{\sigma a} + \frac{\delta v_i}{4 \Omega R_{rad}} \end{cases} \quad (1.20)$$

1.3.2 General Moments and Forces

The motion of the quadrotor is the result of several forces and moments that are originating from different physical effects [1]. In this model, the following effects are considered (with c : \cos , s : \sin).

Rolling Moments

body gyro effect	$\dot{\theta}\dot{\psi}(I_{yy} - I_{zz})$
propeller gyro effect	$J_r\dot{\theta}\Omega_r$
pitch actuators action	$l(-T_2 + T_4)$
hub moment due to forward flight	$h(\sum_{i=1}^4 H_{yi})$
rolling moment due to forward flight	$(-1)^{i+1} \sum_{i=1}^4 R_{mxi}$

Pitching Moments

body gyro effect	$\dot{\phi}\dot{\psi}(I_{zz} - I_{xx})$
propeller gyro effect	$J_r\dot{\phi}\Omega_r$
pitch actuators action	$l(T_1 - T_3)$
hub moment due to forward flight	$h(\sum_{i=1}^4 H_{xi})$
rolling moment due to side-ward flight	$(-1)^{i+1} \sum_{i=1}^4 R_{myi}$

Yawing Moments

body gyro effect	$\dot{\theta}\dot{\phi}(I_{xx} - I_{yy})$
inertial counter-torque	$J_r\dot{\Omega}_r$
counter-torque unbalance	$(-1)^i \sum_{i=1}^4 Q_i$
hub force unbalance in forward flight	$l(H_{x2} - H_{x4})$
hub force unbalance in sideward flight	$l(-H_{y1} + H_{y3})$

Forces Along z Axis

actuators action	$c\psi c\phi(\sum_{i=1}^4 T_i)$
weight	mg

Forces Along x Axis

actuators action	$(s\psi s\phi + c\psi s\theta c\phi)(\sum_{i=1}^4 T_i)$
hub force in x axis	$-\sum_{i=1}^4 H_{xi}$
friction	$\frac{1}{2}C_x A_c \rho \dot{x} \dot{x} $

Forces Along y Axis

actuators action	$(-c\psi s\phi + s\psi s\theta c\phi)(\sum_{i=1}^4 T_i)$
hub force in y axis	$-\sum_{i=1}^4 H_{yi}$
friction	$\frac{1}{2}C_y A_c \rho \dot{y} \dot{y} $

1.3.3 Equations of Motion

The equations of motion are derived from (1.14) in addition to all the forces and the moments that were listed in subsection 1.3.2.

$$\left\{ \begin{array}{l} I_{xx}\ddot{\phi} = \dot{\theta}\dot{\psi}(I_{yy} - I_{zz}) + J_r\dot{\theta}\Omega_r + l(-T_2 + T_4) - h(\sum_{i=1}^4 H_{yi}) + (-1)^{i+1} \sum_{i=1}^4 R_{mxi} \\ I_{yy}\ddot{\theta} = \dot{\phi}\dot{\psi}(I_{zz} - I_{xx}) - J_r\dot{\phi}\Omega_r + l(T_1 - T_3) + h(\sum_{i=1}^4 H_{xi}) + (-1)^{i+1} \sum_{i=1}^4 R_{mxi} \\ I_{zz}\ddot{\psi} = \dot{\phi}\dot{\psi}(I_{xx} - I_{yy}) + J_r\dot{\Omega}_r + (-1)^i \sum_{i=1}^4 Q_i + l(H_{x2} - H_{x4}) + l(-H_{y1} + H_{y3}) \\ m\ddot{z} = mg - (c\psi c\phi) \sum_{i=1}^4 T_i \\ m\ddot{x} = (s\psi s\phi + c\psi s\theta c\phi) \sum_{i=1}^4 T_i - \sum_{i=1}^4 H_{xi} - \frac{1}{2}C_x A_c \rho \dot{x} |\dot{x}| \\ m\ddot{y} = (-c\psi s\phi + s\psi s\theta c\phi) \sum_{i=1}^4 T_i - \sum_{i=1}^4 H_{yi} - \frac{1}{2}C_y A_c \rho \dot{y} |\dot{y}| \end{array} \right. \quad (1.21)$$

1.4 State-Space Model

The model 1.21 that was developed in subsection 1.3.3 expresses the differential equations of the system. However, for the purpose of control design, it is desirable to reduce the complexity and simplify the model to satisfy the real-time limitations of the embedded control loop. Thus, the thrust and the drag coefficients are assumed to be constant and the hub forces and rolling moments are neglected. As a result, the system can be expressed in state-space form

$\dot{\mathbf{x}} = f(\mathbf{x}, \mathbf{u})$ with \mathbf{x} the state vector and \mathbf{u} the control input vector.

The state vector has the following form:

$$\mathbf{x} = [\phi \ \dot{\phi} \ \theta \ \dot{\theta} \ \psi \ \dot{\psi} \ z \ \dot{z} \ x \ \dot{x} \ y \ \dot{y}]^T \quad (1.22)$$

With,

$$\begin{array}{l|l} x_1 = \phi & x_7 = z \\ x_2 = \dot{x}_1 = \dot{\phi} & x_8 = \dot{x}_7 = \dot{z} \\ x_3 = \theta & x_9 = x \\ x_4 = \dot{x}_3 = \dot{\theta} & x_{10} = \dot{x}_9 = \dot{x} \\ x_5 = \psi & x_{11} = y \\ x_6 = \dot{x}_5 = \dot{\psi} & x_{12} = \dot{x}_{11} = \dot{y} \end{array} \quad (1.23)$$

Moreover, the control input vector has the following form:

$$\mathbf{u} = [u_1 \ u_2 \ u_3 \ u_4]^T \quad (1.24)$$

Where the control inputs are mapped by:

$$\begin{cases} u_1 = b(\Omega_1^2 + \Omega_2^2 + \Omega_3^2 + \Omega_4^2) \\ u_2 = b(-\Omega_2^2 + \Omega_4^2) \\ u_3 = b(\Omega_1^2 - \Omega_3^2) \\ u_4 = d(-\Omega_1^2 + \Omega_2^2 - \Omega_3^2 + \Omega_4^2) \end{cases} \quad (1.25)$$

The transformation matrix between the rate change of the attitude angles $(\dot{\phi}, \dot{\theta}, \dot{\psi})$ and the angular velocities of the body (p, q, r) can be regarded as the identity matrix if the disturbances due to hover flight are small. As a result, the following can be written:

$$(\dot{\phi}, \dot{\theta}, \dot{\psi}) \approx (p, q, r) \quad (1.26)$$

Simulation tests have demonstrated that this assumption is reasonable [1].

From equations (1.21),(1.22),(1.24), the following expression is obtained after simplification:

$$f(\mathbf{x}, \mathbf{u}) = \begin{pmatrix} \dot{\phi} \\ \dot{\theta}\dot{\psi}a_1 + \dot{\theta}a_2\Omega_r + b_1u_2 \\ \dot{\theta} \\ \dot{\phi}\dot{\psi}a_3 - \dot{\phi}a_4\Omega_r + b_2u_3 \\ \dot{\psi} \\ \dot{\theta}\dot{\psi}a_5 + b_3u_4 \\ \dot{z} \\ g - (\cos \phi \cos \theta) \frac{1}{m}u_1 \\ \dot{x} \\ u_x \frac{1}{m}u_1 \\ \dot{y} \\ u_y \frac{1}{m}u_1 \end{pmatrix} \quad (1.27)$$

With,

$$\begin{array}{l|l} a_1 = (I_{yy} - I_{zz})/I_{xx} & b_1 = l/I_{xx} \\ a_2 = J_r/I_{xx} & b_2 = l/I_{yy} \\ a_3 = (I_{zz} - I_{xx})/I_{yy} & b_3 = l/I_{zz} \\ a_4 = J_r/I_{yy} & \\ a_5 = (I_{xx} - I_{yy})/I_{zz} & \end{array} \quad (1.28)$$

$$\begin{aligned} u_x &= (\cos \phi \sin \theta \cos \psi + \sin \phi \sin \psi) \\ u_y &= (\cos \phi \sin \theta \sin \psi - \sin \phi \cos \psi) \end{aligned} \quad (1.29)$$

It is important to note that that the angles and the derivatives of the angles do not depend on the components of the translation in the system represented by equation (1.27). Contrarily, the translation components depend of the angles. So, the system represented by equation(1.27) can be depicted as two subsystems, the angle subsystem and the translation subsystem as shown in figure

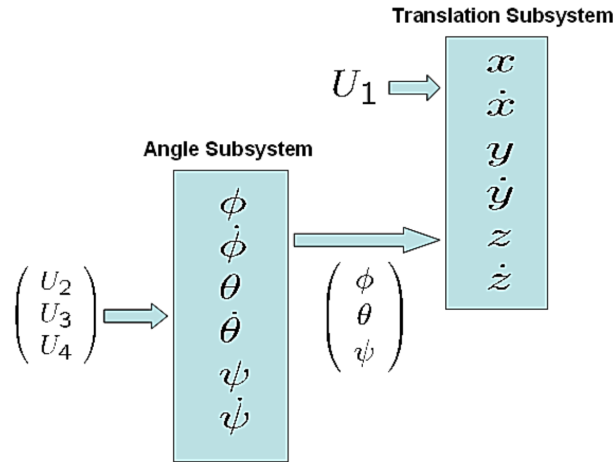


Figure 1.3: Link between the rotation and the translation subsystems.[1]

Control of quadrotors

In the following sections of this chapter, differential flatness, the general control architecture of a quadrotor and different potential control approaches (linear and nonlinear) that can be used to control a quadrotor are explained.

2.1 Differential Flatness

In the quadrotor community, a well-established finding is that the dynamic model of a quadrotor is differentially flat. Moreover, the control design problem in non-linear systems will be considerably simplified. Precisely, a system with state $\mathbf{x} \in \mathbb{R}^n$ and input $\mathbf{u} \in \mathbb{R}^m$ is considered to be *differentially flat* if there exists a set of *flat outputs* $\mathbf{y} \in \mathbb{R}^m$ which have the following form:

$$\mathbf{y} = \mathbf{y}(\mathbf{x}, \mathbf{u}, \dot{\mathbf{u}}, \dots, \mathbf{u}^{(p)}) \quad (2.1)$$

With,

$$\begin{cases} \mathbf{x} = \mathbf{x}(\mathbf{y}, \dot{\mathbf{y}}, \dots, \mathbf{y}^{(q)}) \\ \mathbf{u} = \mathbf{u}(\mathbf{y}, \dot{\mathbf{y}}, \dots, \mathbf{y}^{(r)}) \end{cases} \quad (2.2)$$

Thus, the new set of variables is required to be a function of the state, the input and the derivatives of the input. Moreover, this set should also have the same dimensions as the control input. In this manner, it is possible to rewrite both the state and the input in function of the flat outputs and the derivatives of the flat outputs. This is a very useful property in underactuated systems where $m < n$, such as quadrotors, because, it will allow to generate trajectories in the lower dimensional space m , then this trajectory will be mapped into the full dimensional space n . Another well known example of systems is a car, in which the underactuation is the result of the nonholonomic constraints that are imposed by the wheels. So, for a car, a generated trajectory for (x, y) position of the rear-wheels is enough to specify all the viable trajectories of the system. Formal proofs that the quadrotor system is differentially flat can be found in [16], and [17] for the full model with first-order aerodynamics. The standard choice of flat outputs for the quadrotor is the coordinates of the center of mass and the yaw angle:

$$\mathbf{y} = [x \quad y \quad z \quad \psi]^\top \quad (2.3)$$

Consequently, the problem of generating a feasible trajectory for a quadrotor then trajecing it can be dimensionally decreased from a 6-dimensional space to a 4-dimensional space. By reason of the tight coupling between the rotational and translational dynamics, then defining a trajectory in function of the flat outputs \mathbf{y} is sufficient to properly define the full dynamics \mathbf{x} .

2.2 General Control Architecture

In the last few years, many researches have developed interest in control of quadrotors. As a result, various control approaches have been proposed. The most known control architecture [17] consists of three nested control loops, as shown in figure 2.1, in order to generate the suitable thrust in each actuator to follow the desired signal. This strategy assumes that the attitude dynamics of a quadrotor are much faster than the translational dynamics. Assuming that Euler angles are used to define the attitude and that a navigation module generates the desired trajectory $(\mathbf{r}_d(t), \psi_d(t))$ as shown in section 2.1, then:

Position controller has the objective of driving the errors occurring on the translational dynamics to zero. And, the outputs of this outer loop are the thrust $f = U_1$, which is sent to the motor controller, and the desired attitude $(\theta_d(t), \phi_d(t))$, which corresponds to the reference signal of the attitude controller.

Attitude controller it has the goal of driving the errors occurring on the rotational dynamics to zero. This controller generates the inputs $\boldsymbol{\tau} = [u_2 \quad u_3 \quad u_4]^\top$ that are then sent to the motor controller.

Motor controller This controller receives the control inputs $\boldsymbol{\tau} = [f \quad \boldsymbol{\tau}]^\top$ and maps them into the desired spinning velocities Ω_i for each individual rotor based on equation 1.25. Moreover, low-level control laws are designed and realized in the firmware of the drone to make the convergence from the actual rotations to these desired values.

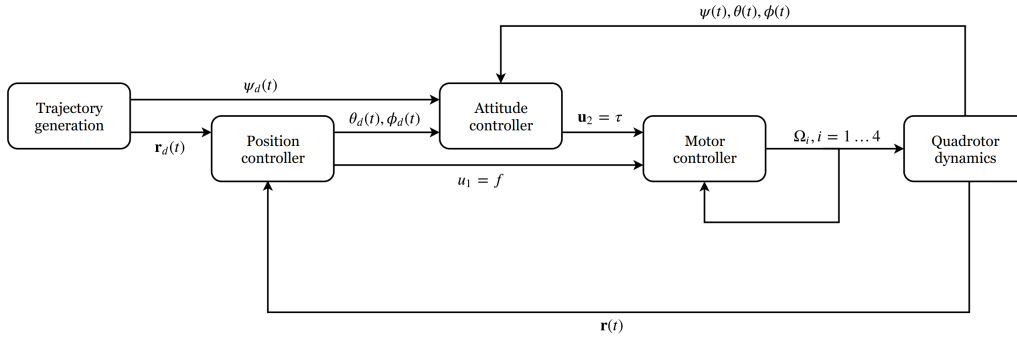


Figure 2.1: General control architecture of a quadrotor.

2.3 General Control Approaches

2.3.1 Method of Linearization

By using extreme assumptions, it is feasible to apply linear control techniques in order to control a quadrotor ([18], [19]). Particularly, this can be made by doing a linearization of the full dynamic model around an equilibrium point $\bar{\mathbf{x}}$ and by using the assumption that the vehicle is only capable of oscillating lightly around the hover point. It is very easy to observe that a feasible equilibrium is provided by a configuration where the center of mass is at a random position $\bar{\mathbf{r}}$ and all the other elements of the state are set to zero. So, the nominal input $\mathbf{U} = \bar{\mathbf{U}}$ to sustain such equilibrium can be assessed as the thrust that is required to compensate the gravity force:

$$\bar{\mathbf{u}} = \begin{bmatrix} f \\ \boldsymbol{\tau} \end{bmatrix} = \begin{bmatrix} mg \\ \mathbf{0}_{3 \times 1} \end{bmatrix} \quad (2.4)$$

At this stage, the complete non-linear dynamics that have the form :

$$\dot{\mathbf{x}} = \bar{\mathbf{f}}(\bar{\mathbf{x}}, \bar{\mathbf{u}}) \quad (2.5)$$

can now be linearized around the hover point $(\bar{\mathbf{x}}, \bar{\mathbf{u}})$ as shown below:

$$\dot{\mathbf{x}} = \left[\frac{\partial \mathbf{f}(\mathbf{x}, \mathbf{u})}{\partial \mathbf{x}} \right]_{(\bar{\mathbf{x}}, \bar{\mathbf{u}})} \mathbf{x} + \left[\frac{\partial \mathbf{f}(\mathbf{x}, \mathbf{u})}{\partial \mathbf{u}} \right]_{(\bar{\mathbf{x}}, \bar{\mathbf{u}})} \mathbf{u} = \mathbf{Ax} + \mathbf{Bu} \quad (2.6)$$

It can be demonstrated that both matrices \mathbf{A} and \mathbf{B} can be used to determine a linear system that is both controllable and observable [18]. Thus, any control technique that is linear can now be used on the quadrotor in order to keep it around a desired equilibrium point, such as optimal LQR/LQG [20, 21] control or simple PD or PID controller [22, 23].

2.3.2 Internal Lyapunov Stability

Before defining the *Lyapunov Direct Method*, the notion of stability will be thoroughly defined first.

Notions of Stability

For a general system without any control input

$$\dot{x}(t) = f(x(t), 0, t) \quad (CT) \quad (2.7)$$

$$\dot{x}(k+1) = f(x(k), 0, k) \quad (DT), \quad (2.8)$$

it is said that a point \bar{x} is called an *equilibrium point* from time t_0 for the continuous system (CT) if $f(\bar{x}, 0, t) = 0, \forall t \geq t_0$. Moreover, in the discrete time (DT) case, the point \bar{x} is an equilibrium point from time k_0 if $f(\bar{x}, 0, k) = 0, \forall k \geq k_0$. If the system begins from state \bar{x} at time t_0 or k_0 , then the system will stay there and will not change with time. It is possible for nonlinear system to have more than one equilibrium point (equilibria). There also exists another class of special solutions in the case of nonlinear systems, these solutions are called *periodic* solution. However, it is outside the scope of this bibliography and interested readers are referred to [24] for more in depth explanation. So, the focus will be on equilibria. It is desired to identify the *stability* of the equilibria in some way. For instance, it is desired to know if, given some small perturbation to the system, the state would either come back to the equilibrium point, remain close to it in some sense, or it diverges.

The most useful notion of stability for an equilibrium point of a nonlinear system is provided by the definition below. Assuming that the equilibrium point is at the origin, because if $\bar{x} \neq 0$, a simple translation can be done to obtain a system that is equivalent with the equilibrium at the origin.

Asymptotic stability A system is said to be *asymptotically stable* about its own equilibrium point at the origin if the following two conditions are satisfied [2]:

1. For any $\epsilon > 0$, $\exists \delta_1 > 0$ such that if $\|x(t_0)\| < \delta_1$, then $\|x(t)\| < \epsilon, \forall t > t_0$.
2. $\exists \delta_2$ such that if $\|x(t_0)\| < \delta_2$, then $x(t) \rightarrow 0$ at $t \rightarrow \infty$.

For the first condition, it is required that the state trajectory should be restricted to a randomly small "ball" that is centered at the equilibrium point and has a radius ϵ , when released from an arbitrary initial condition in a ball that has an adequately small (yet positive) radius δ_1 . This is referred to as *stability in the sense of Lyapunov* (i.s.L.). It is also possible to have stability in the sense of Lyapunov without having asymptotic stability, in that case, it is said that the equilibrium point is *marginally stable*. Moreover, there also exist nonlinear systems that satisfy the second condition without being stable in the sense of Lyapunov. Moreover, an equilibrium point that is *not* stable in the sense of Lyapunov is said to be *unstable*.

Lyapunov's Direct Method

General Idea If the following continuous-time system is considered

$$\dot{x}(t) = f(x(t)) \quad (2.9)$$

which has an equilibrium point at the origin ($x=0$). This system is called a time-invariant system because f does not depend explicitly on the time t . In such a system, the stability analysis of the equilibrium point is in general a tedious task. This is because there is no way to write a simple formula which relates the trajectory to the initial state. The main idea of *Lyapunov's direct method* is to establish properties of the equilibrium point (of the nonlinear system in general) by evaluating how a certain carefully chosen scalar function of the changes as the state of the system changes. (The term "*direct*" is used to distinguish this method from the Lyapunov's "*indirect*" method, which tries to establish properties of the equilibrium point by assessing the behavior of the *linearized* system at that point [25].

As an example, a continuous and scalar function $V(x)$ that is equal to 0 at the origin and positive elsewhere is considered in some ball that is enclosing the origin. So, $V(0) = 0$ and $V(x) > 0$ when $x \neq 0$ in this ball. This $V(x)$ be considered as an "*energy*" function. Also, let $\dot{V}(x)$ denote the time derivative of $V(x)$ throughout any trajectory of the system. In other words, $\dot{V}(x)$ varies as $x(t)$ varies proportionally to equation (2.9). If this derivative is negative in all the region (with the exclusion of the origin), then this means that the energy is decreasing with time in a strict manner. Moreover, since the lower bound of the energy is 0, then the energy must go to 0 with time. This means that all trajectories will converge to the origin (zero state). This idea is formalized below.

Lyapunov Functions Assuming V is a continuous map from \mathbb{R}^n to \mathbb{R} . $V(x)$ is called a *locally positive definite* (lpd) function about $x=0$ if

1. $V(0) = 0$.
2. $V(x) > 0$, $0 < \|x\| < r$ for a given r .

Also, the function is called locally *positive semidefinite* (lpsd) if the strict inequality applied on the function in the second condition is changed to $V(x) \geq 0$. The function $V(x)$ is locally *negative definite* (lnd) if $-V(x)$ is lpd. Moreover, $V(x)$ is locally *negative semidefinite* (lnsd) if $-V(x)$ is lpsd. In order to graphically illustrate the locally positive definite function $V(x)$, it is useful to imagine having "contours" of constant V that form (at least in a tiny region about the origin) a set of nested surfaces that are encircling the origin. An example is given in figure 2.2 below:

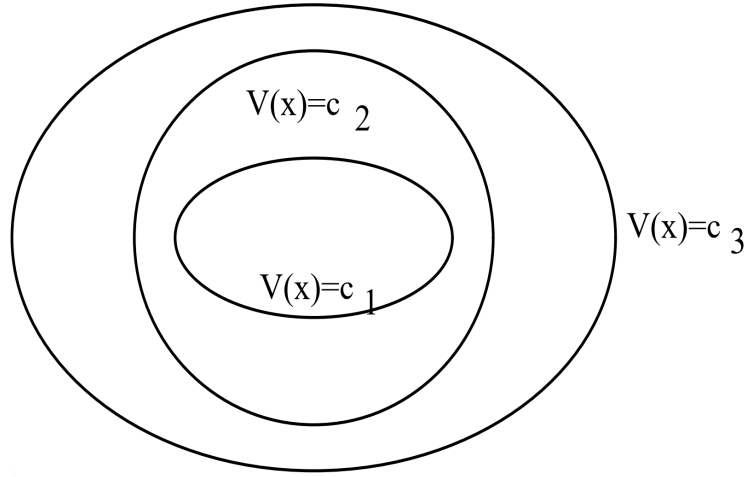


Figure 2.2: Different values for a Lyapunov function, where $c_1 < c_2 < c_3$. [2]

In the continuous time case, the focus in this bibliography will be restricted to $V(x)$ that has first partial derivatives that are continuous. (In the case of discrete time, only continuity will suffice, so differentiability is not required in that case.) The derivative of V with respect to time *along a trajectory of the system* 2.9 is denoted as $\dot{V}(x(t))$. This derivative is expressed as follows:

$$\dot{V}(x(t)) = \frac{dV(x)}{dx} \dot{x} = \frac{dV(x)}{dx} f(x) \quad (2.10)$$

where $\frac{dV(x)}{dx}$ is the *Jacobian* of V with respect to x which contain the partial derivative V with respect to every component of x : $\frac{\partial V}{\partial x_i}$.

Moreover, assuming V is a lpd function (a "candidate Lyapunov function"), and (\dot{V}) is the derivative along trajectories of the system of equation 2.9. Then, V is called a *Lyapunov function* of the system of equation 2.9 if \dot{V} is lnsd [2].

Lyapunov Theorem for Local Stability If a Lyapunov function for the system of equation (2.9) exists, this means that $x = 0$ is a stable equilibrium point in the sense of Lyapunov. Moreover, if $\dot{V} < 0$, $0 < \|x\| < r_1$, for some r_1 , in other words, if \dot{V} is lnd, this means that $x = 0$ is an equilibrium point that is asymptotically stable[2].

Lyapunov Theorem of Global Asymptotic Stability The region in the state space for which the earlier results hold is determined by the region over which $V(x)$ represents a Lyapunov function. It is interesting to determine the "basin of attraction" of an equilibrium point that is asymptotically stable. In other words, the basin of attraction is the set of initial conditions whose following trajectories conclude at the equilibrium point. Thus, an equilibrium point is "*globally asymptotically stable*" (also called asymptotically stable "in the large") if its basin of attraction is the entire state space.

If a function $V(x)$ has the following criteria below,

1. $V(x)$ is positive definite throughout the state-space.
2. $V(x)$ also has the added property that $|V(x)| \rightarrow \infty$ as $\|x\| \rightarrow \infty$.
3. The derivative of V (in other words \dot{V}) is negative definite throughout the state-space.

Then, the equilibrium point at the origin is globally asymptotically stable [2].

Discrete-Time Systems Fundamentally identical results hold for the system

$$x(k+1) = f(x(k)) \tag{2.11}$$

provided that \dot{V} is represented as

$$\dot{V}(x) \triangleq V(f(x)) - V(x),$$

in other words as

$$V(\text{next state}) - V(\text{present state})$$

2.3.3 Model Predictive control

2.3.4 Sliding mode control

2.3.5 Other types of control

Multi-flips maneuver with quadrotors

- 3.1 Quadrotor flip physics
- 3.2 Link to parallel robots
- 3.3 Control approaches for multi-flip maneuvers

Trajectory optimization

Conclusion

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