1 Rayleigh-Benard Convection

Rayleigh-Benard convection is described by the non-dimensionalized equations

$$\nabla \cdot \boldsymbol{u} = 0,$$

$$\frac{1}{\Pr} (\partial_t \boldsymbol{u} + \boldsymbol{u} \cdot \nabla \boldsymbol{u}) = -\nabla p + \operatorname{Ra}\theta \boldsymbol{e}_z + \nabla^2 \boldsymbol{u},$$

$$\partial_t \theta + \boldsymbol{u} \cdot \nabla \theta - u_z = \nabla^2 \theta,$$
(1)

with the boundary conditions

$$u_z = \partial_z u_x = \partial_z u_y = \theta = 0$$
 at $z = 0, 1$. (2)

We can rewrite this as a system of first order (in z) PDEs:

$$\frac{1}{\Pr} (\partial_t u + u \partial_x u + v \partial_y u + w du) = -\partial_x p + \partial_x^2 u + \partial_y^2 u + \partial_z du,
\frac{1}{\Pr} (\partial_t v + u \partial_x v + v \partial_y v + w dv) = -\partial_y p + \partial_x^2 v + \partial_y^2 v + \partial_z dv,
\frac{1}{\Pr} (\partial_t w + u \partial_x w + v \partial_y w + w \partial_z w) = -\partial_z p + \operatorname{Ra}\theta +
+ \partial_x^2 w + \partial_y^2 w - (\partial_x du + \partial_y dv),
\partial_t \theta + u \partial_x \theta + v \partial_y \theta + w d\theta - w = \partial_x^2 \theta + \partial_y^2 \theta + \partial_z d\theta,
\partial_z u = du,
\partial_z v = dv,
\partial_z \theta = d\theta.$$
(3)

Separating out the linear and nonlinear terms gives

$$\frac{1}{\Pr}\partial_{t}u + \partial_{y}v + \partial_{z}w = 0,$$

$$\frac{1}{\Pr}\partial_{t}u + \partial_{x}p - \partial_{x}^{2}u - \partial_{y}^{2}u - \partial_{z}du = -\frac{1}{\Pr}\left(u\partial_{x}u + v\partial_{y}u + wdu\right),$$

$$\frac{1}{\Pr}\partial_{t}v + \partial_{y}p - \partial_{x}^{2}v - \partial_{y}^{2}v - \partial_{z}dv = -\frac{1}{\Pr}\left(u\partial_{x}v + v\partial_{y}v + wdv\right),$$

$$\frac{1}{\Pr}\partial_{t}w + \partial_{z}p - \operatorname{Ra}\theta - \partial_{x}^{2}w - \partial_{y}^{2}w + (\partial_{x}du + \partial_{y}dv)$$

$$= -\frac{1}{\Pr}\left(u\partial_{x}w + v\partial_{y}w + w\partial_{z}w\right),$$

$$\partial_{t}\theta - w - \partial_{x}^{2}\theta - \partial_{y}^{2}\theta - \partial_{z}d\theta = -u\partial_{x}\theta - v\partial_{y}\theta - wd\theta,$$

$$\partial_{z}u - du = 0,$$

$$\partial_{z}v - dv = 0,$$

$$\partial_{z}\theta - d\theta = 0.$$
(4)

The matrix version is:

$$L_{0} = \begin{pmatrix} 0 & \partial_{x} & \partial_{y} & 0 & 0 & 0 & 0 & 0 \\ \partial_{x} & -\partial_{x}^{2} - \partial_{y}^{2} & 0 & 0 & 0 & 0 & 0 & 0 \\ \partial_{y} & 0 & -\partial_{x}^{2} - \partial_{y}^{2} & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -\partial_{x}^{2} - \partial_{y}^{2} & -\operatorname{Ra} & \partial_{x} & \partial_{y} & 0 \\ 0 & 0 & 0 & -1 & -\partial_{x}^{2} - \partial_{y}^{2} & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 \end{pmatrix}. (6)$$

$$F(X) = \begin{pmatrix} 0 \\ -\frac{1}{\Pr} \left(u\partial_x u + v\partial_y u + wdu \right) \\ -\frac{1}{\Pr} \left(u\partial_x v + v\partial_y v + wdv \right) \\ -\frac{1}{\Pr} \left(u\partial_x w + v\partial_y w - w\partial_x u - w\partial_y v \right) \\ -u\partial_x \theta - v\partial_y \theta - wd\theta \\ 0 \\ 0 \\ 0 \end{pmatrix}. \tag{8}$$

The boundary conditions are $w = du = dv = \theta = 0$ at z = 0, 1.

The 2D problem is easier:

$$L_{0} = \begin{pmatrix} 0 & \partial_{x} & 0 & 0 & 0 & 0 \\ \partial_{x} & -\partial_{x}^{2} & 0 & 0 & 0 & 0 \\ 0 & 0 & -\partial_{x}^{2} & -\operatorname{Ra} & \partial_{x} & 0 \\ 0 & 0 & -1 & -\partial_{x}^{2} & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 \end{pmatrix}.$$
 (10)

$$F(X) = \begin{pmatrix} 0 \\ -\frac{1}{\Pr} (u\partial_x u + wdu) \\ -\frac{1}{\Pr} (u\partial_x w - w\partial_x u) \\ -u\partial_x \theta - wd\theta \\ 0 \\ 0 \end{pmatrix}.$$
 (12)

The boundary conditions are $w = du = \theta = 0$ at z = 0, 1.

2 Simplified System

Since we are having problems with the Rayleigh-Benard system, I propose that we consider the following simpler problem that has similar behavior:

$$\partial_t w - \operatorname{Ra}\theta - \nabla^2 w = 0, \tag{13}$$

$$\partial_t \theta - w - \nabla^2 \theta = 0. \tag{14}$$

The boundary conditions are $w = \theta = 0$ at z = 0, 1.

Assuming

$$w, \theta \sim \exp(-i\omega t + imx + inz),$$
 (15)

with $\mathbf{k} = m\mathbf{e}_x + n\mathbf{e}_z$, we have

$$-i\omega w + k^2 w = \operatorname{Ra}\theta,\tag{16}$$

$$-i\omega\theta + k^2\theta = w, (17)$$

so the dispersion relation is

$$(-i\omega + k^2)^2 = \text{Ra},\tag{18}$$

SO

$$\omega = i\sqrt{\text{Ra}} - ik^2. \tag{19}$$

We have instability when the imaginary part of ω is positive, so there is instability only when $k^2 < \sqrt{\text{Ra}}$. In a box of size (1,1), the smallest possible wavenumber is $(0,2\pi)$, so the critical Rayleigh number is $(2\pi)^4 \approx 1559$.

We can rewrite the system in terms of first order (in z) PDEs,

$$\partial_t w - \operatorname{Ra}\theta - \partial_x^2 w - \partial_z dw = 0, \tag{20}$$

$$\partial_t \theta - w - \partial_x^2 \theta - \partial_z d\theta = 0, \tag{21}$$

$$\partial_z w - dw = 0, (22)$$

$$\partial_z \theta - d\theta = 0. (23)$$

Writing this in matrix form, we have

$$L_{0} = \begin{pmatrix} -\partial_{x}^{2} & -\text{Ra} & 0 & 0\\ -1 & -\partial_{x}^{2} & 0 & 0\\ 0 & 0 & -1 & 0\\ 0 & 0 & 0 & -1 \end{pmatrix}, \quad L_{1} = \begin{pmatrix} 0 & 0 & -1 & 0\\ 0 & 0 & 0 & -1\\ 1 & 0 & 0 & 0\\ 0 & 1 & 0 & 0 \end{pmatrix}.$$
 (25)