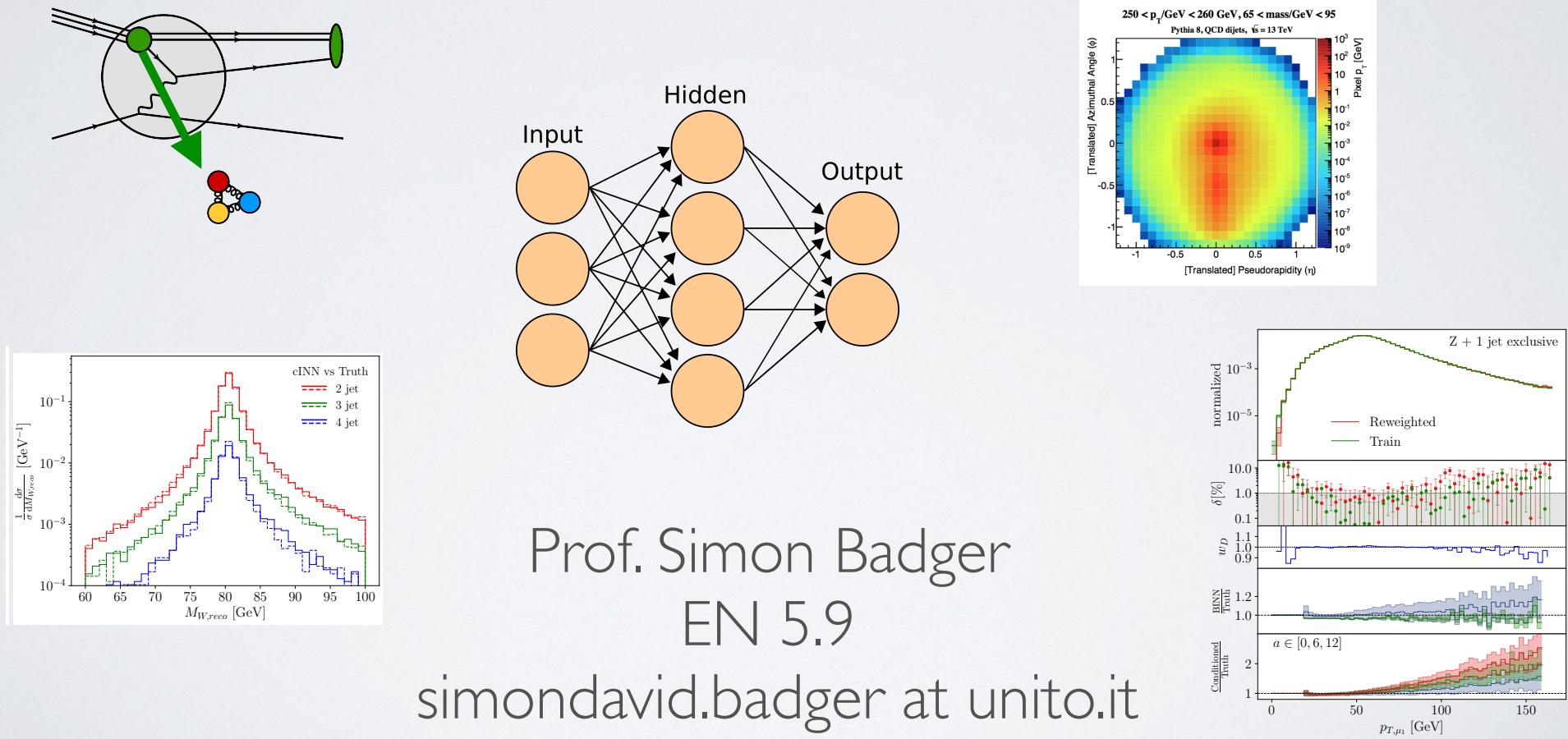


Machine Learning for applied and high energy physics



Course overview

some motivation

High energy collider phenomenology

From field theory to experiment

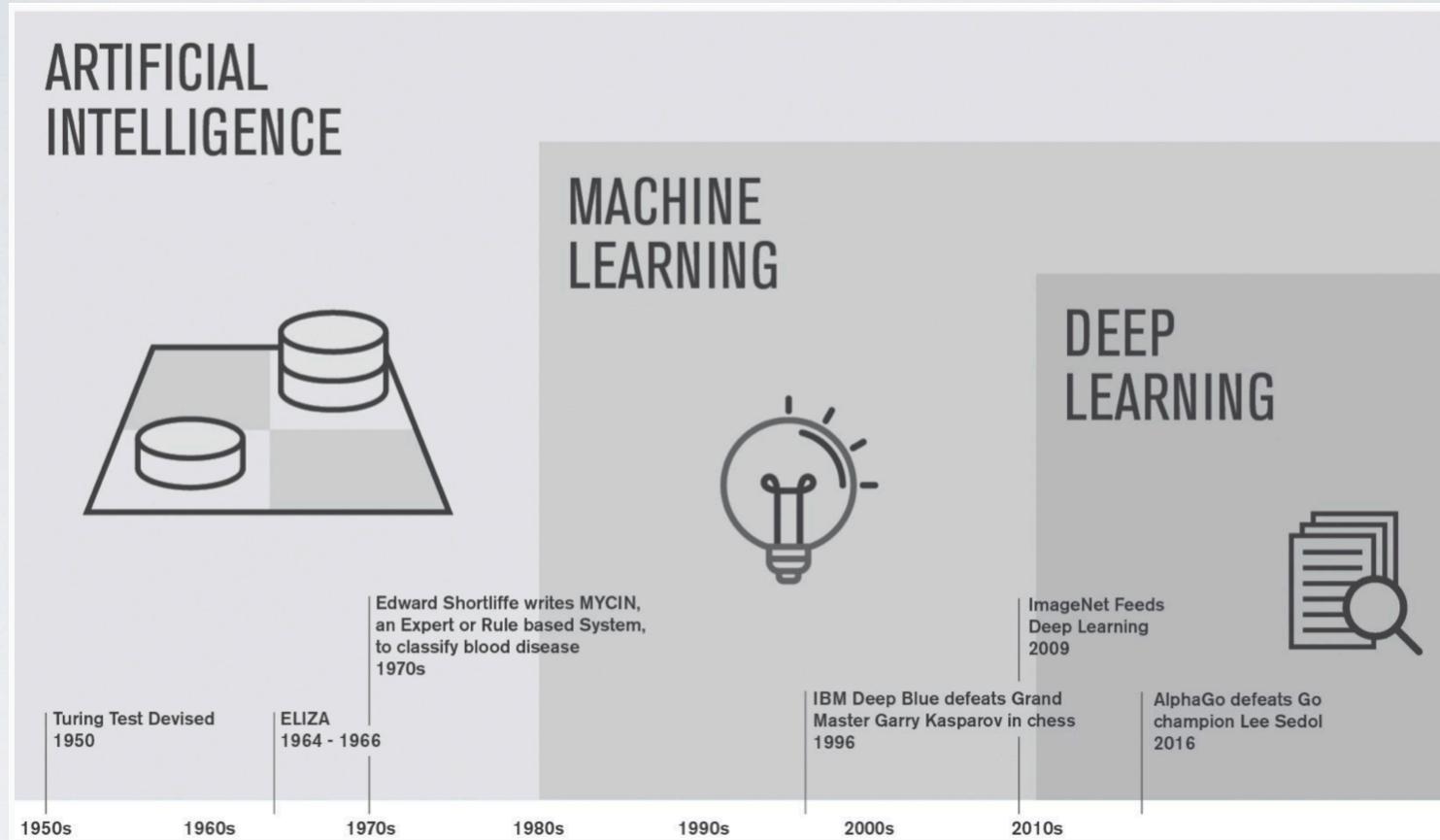
Simulations

The high cost of numerical integration

Machine Learning for fundamental physics

How can ML help us understand the laws of Nature

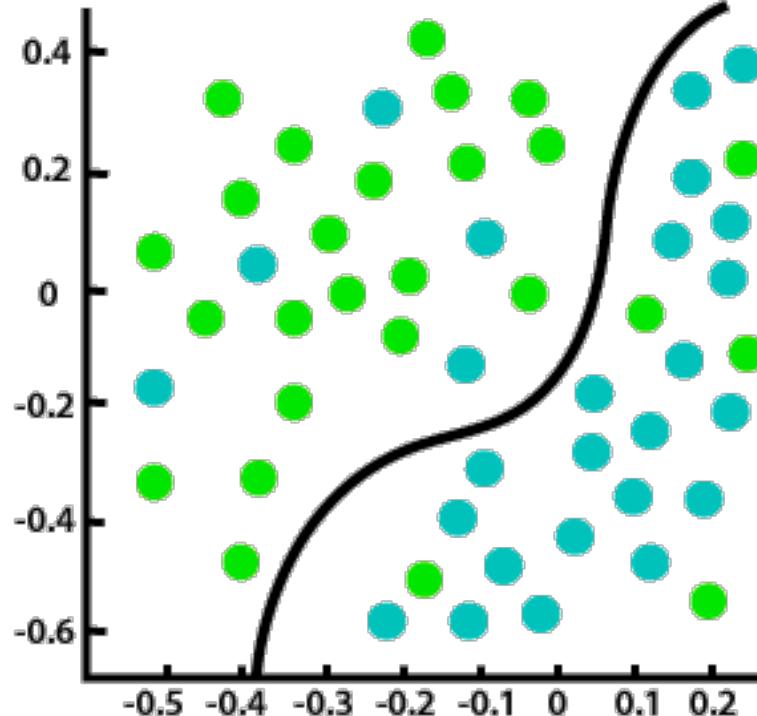
ML algorithms have been around for a while



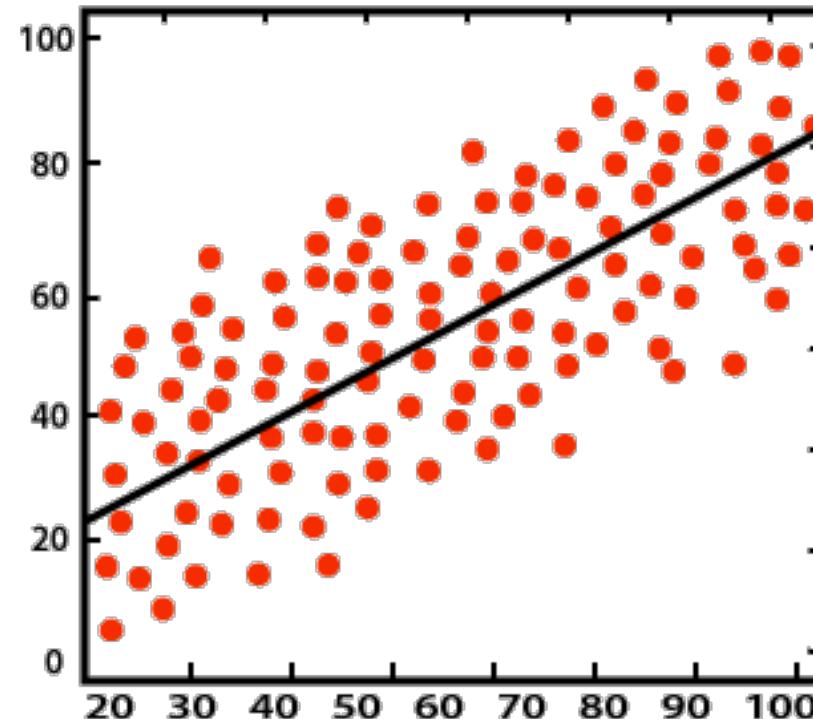
also in HEP

e.g. Finding Gluon Jets With a Neural Trigger,
Lonnblad, Petersen, Rognvaldsson (1990)

Supervised Learning



Classification



Regression

“Learn from examples”

the world of ML techniques can be a bit intimidating



Course overview

structure

- Introduction to collider simulations [1]
What are the aims? What are the challenges?
- Overview of Machine Learning in HEP [1]
What techniques are being used? Look at the [Living Review](#). (Pandas dataframes)
- Non-linear regression: Neural Networks [2]
Perceptron, analytic structure, training and backpropagation
- (Deep) Neural Networks [1]
Optimised training methods, TensorFlow examples

Course overview

structure

- Logistic Regression for Classification [2]
Cross-entropy, Soft-Max, ROC curves
- Application: Supersymmetry searches [1]
signal/background discrimination
- Application: Optimising collider simulations [2]
Phase space integration, Amplitude Neural Networks, Neural network errors
- Convolutional Neural Networks [1]
2d image processing
- Application: Jet substructure and Jet image classification [1]
Jet algorithms, CNNs for boosted Higgs searches

Course overview

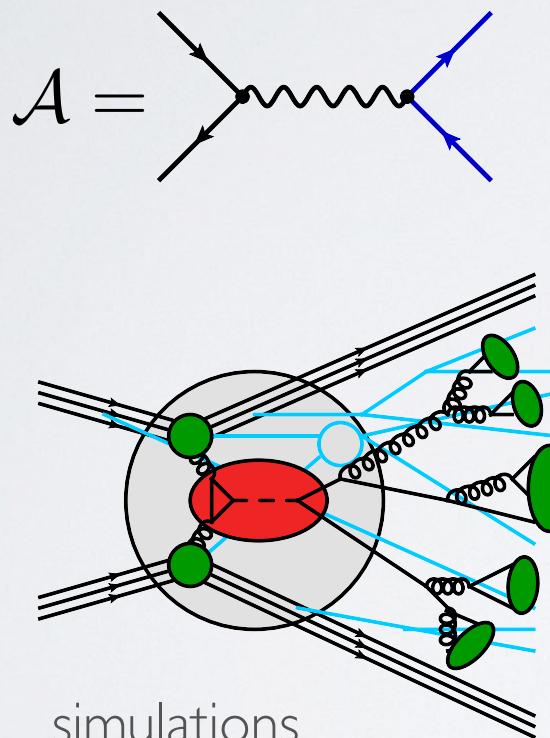
Exam

- Project based assessment
- Demonstrate understanding of course elements
 - choose the right tool for the problem at hand
 - technical skills: demonstrate ability to extend code from a given examples
- 15 minute presentation
 - physics background
 - methods
 - results

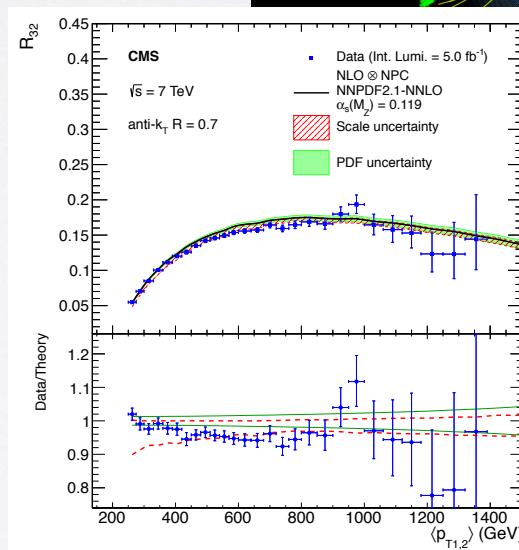
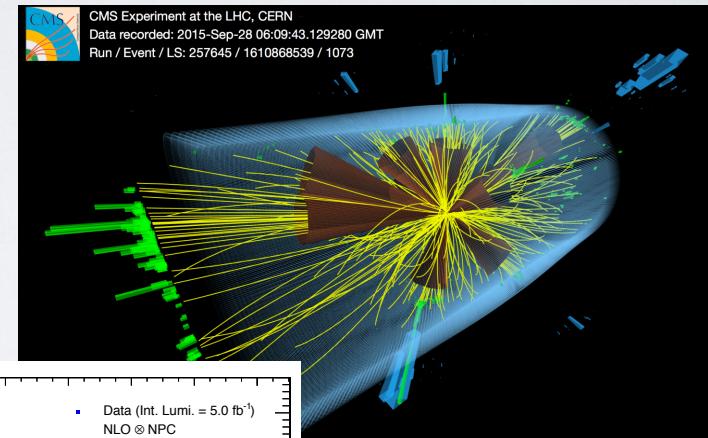
NB: no requirement to be original research - many examples in the 'Living Review' will be too much work

I. Introduction to collider physics

amplitudes



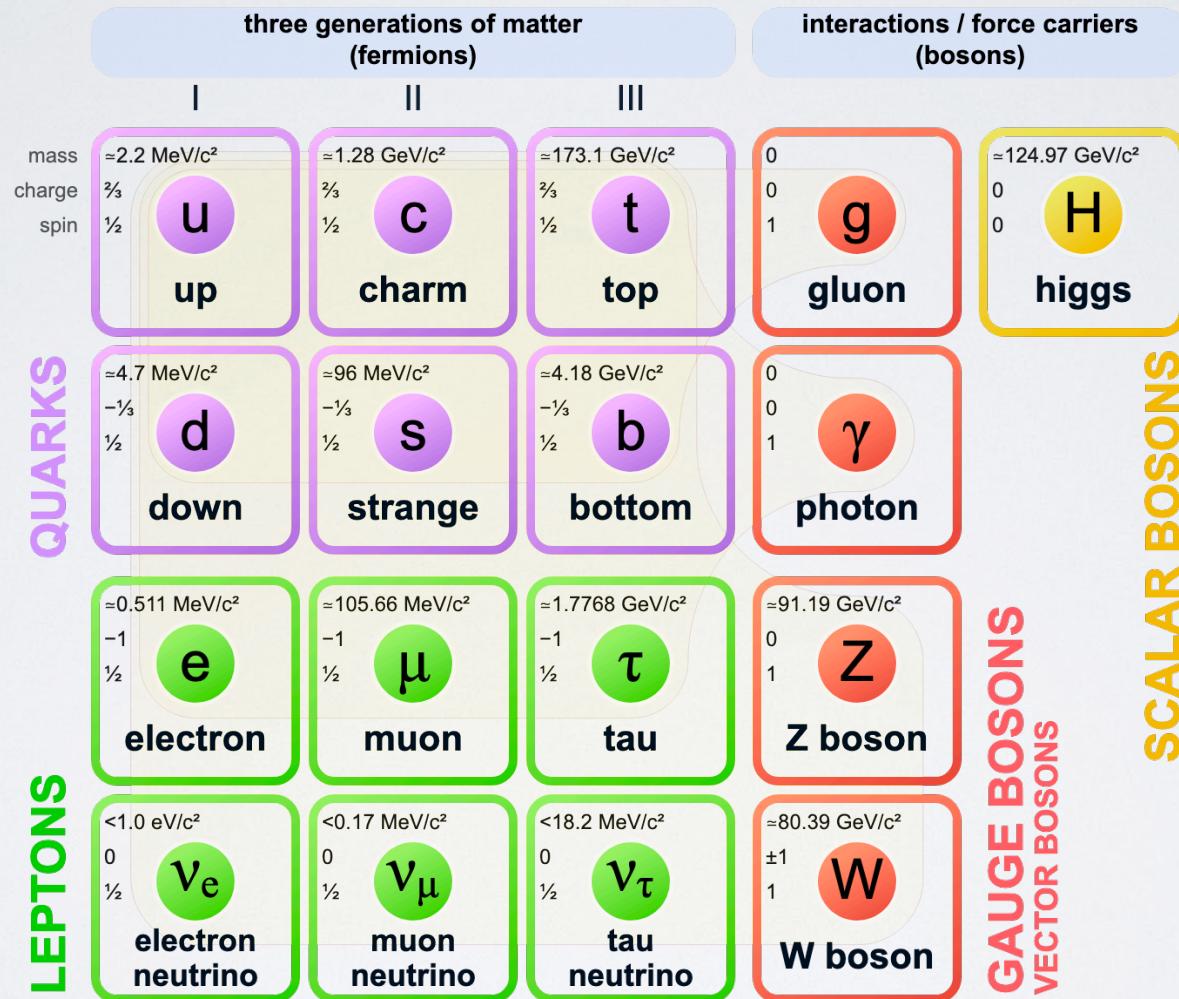
detectors



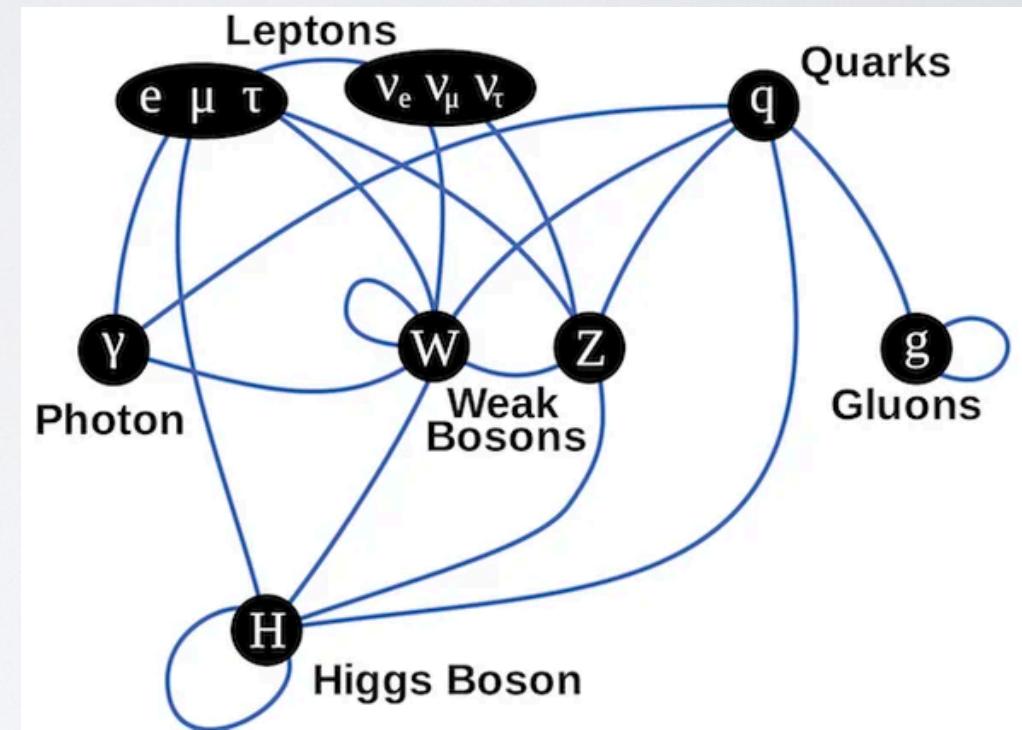
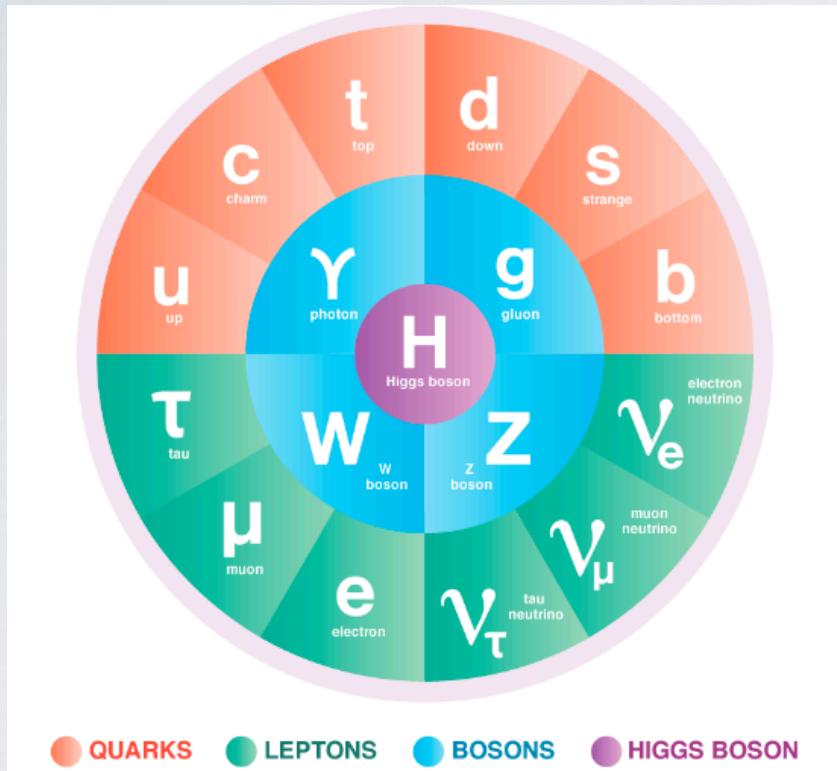
phenomenology

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

Standard Model of Elementary Particles

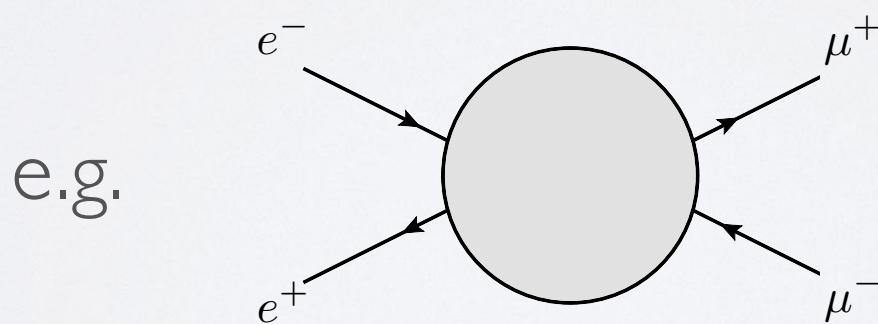


$$SU(3)_c \times SU(2)_L \times U(1)_Y$$



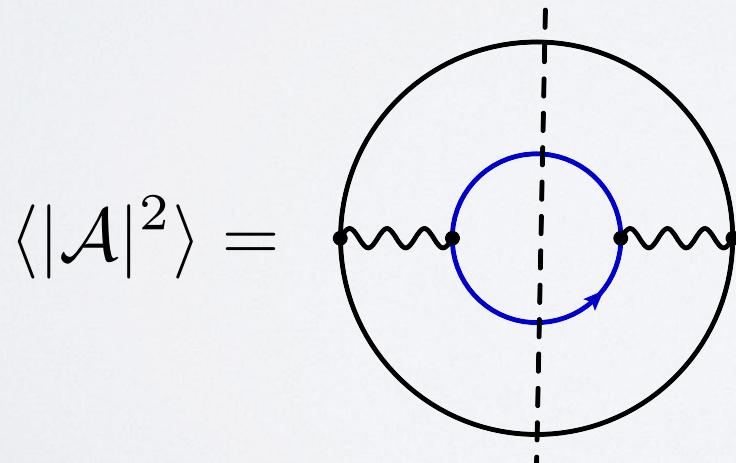
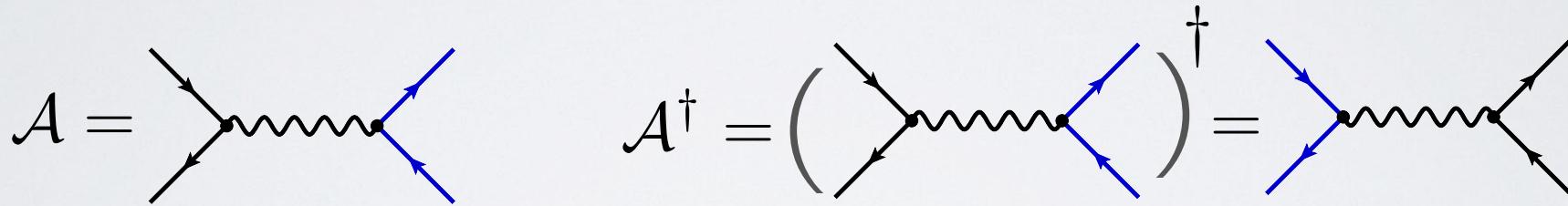
What do theorists see?
scattering cross sections

$$\sigma = \frac{1}{\text{flux}} \int \delta^{(4)}(Q - \sum_i p_i) \prod_i d^4 p \delta^{(+)}(p_i^2 - m_i^2) \langle |\mathcal{A}|^2 \rangle$$



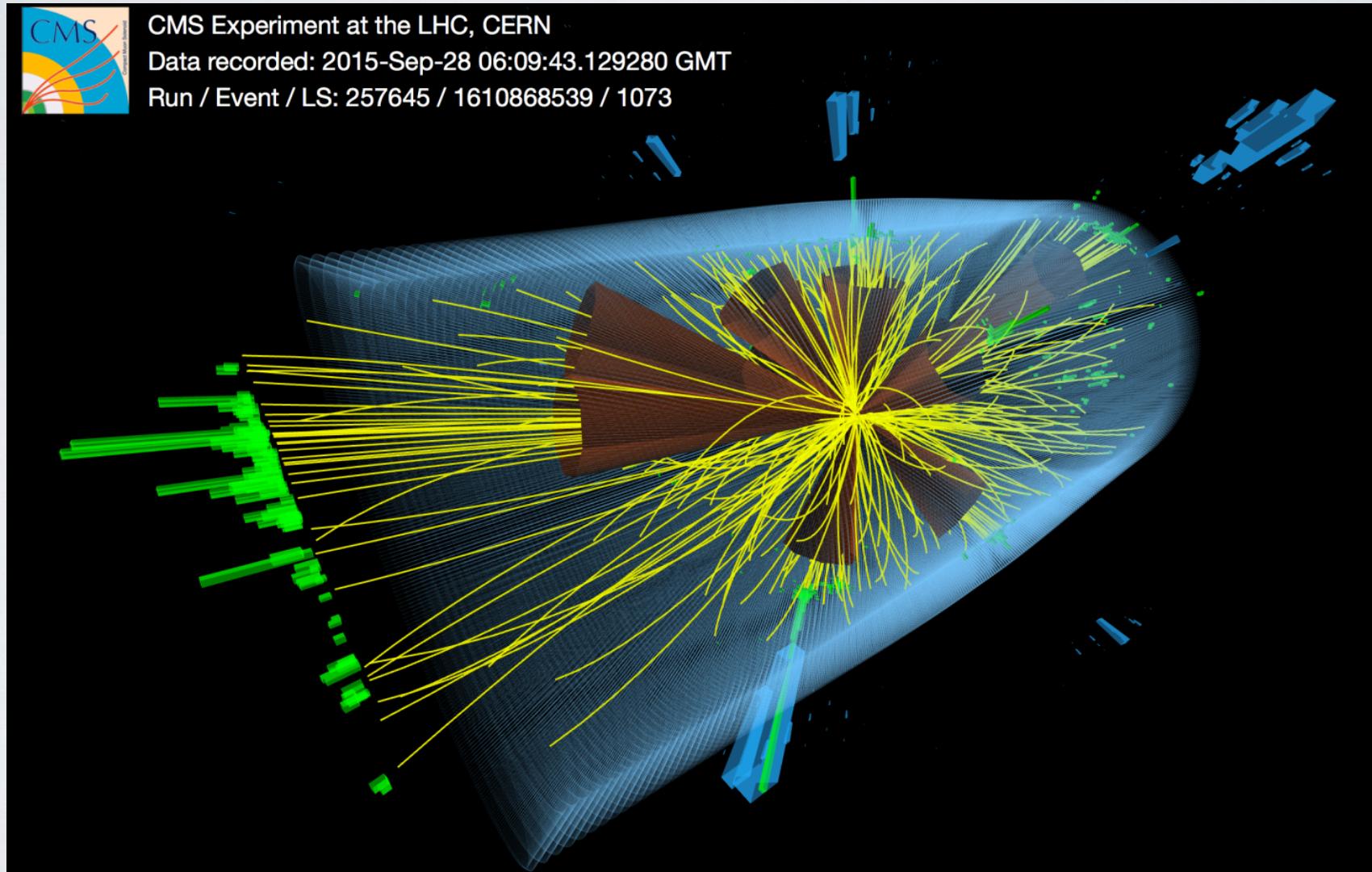
Scattering amplitudes: Probabilities from the Lagrangian

$$\langle |\mathcal{A}|^2 \rangle = \sum_{\text{spins}} \mathcal{A}^\dagger \mathcal{A}$$



diagrammatic
representation of the
squared amplitude

What do we see in the experiments?



What do we see in the experiments?

tracking detectors

calorimeters

exclusive particle identification

(see π , p , J/ψ NOT u,d,c,g)

unstable particle decay

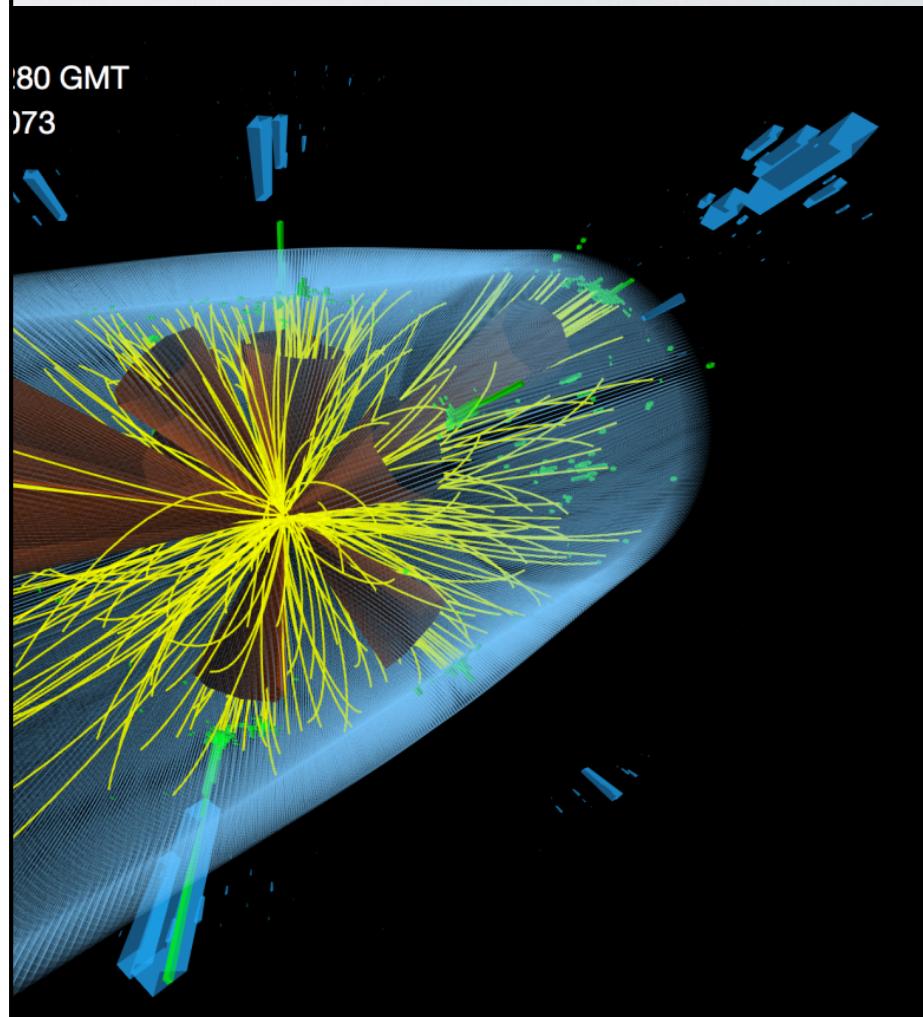
(see e,μ,b NOT W,Z,H,t)

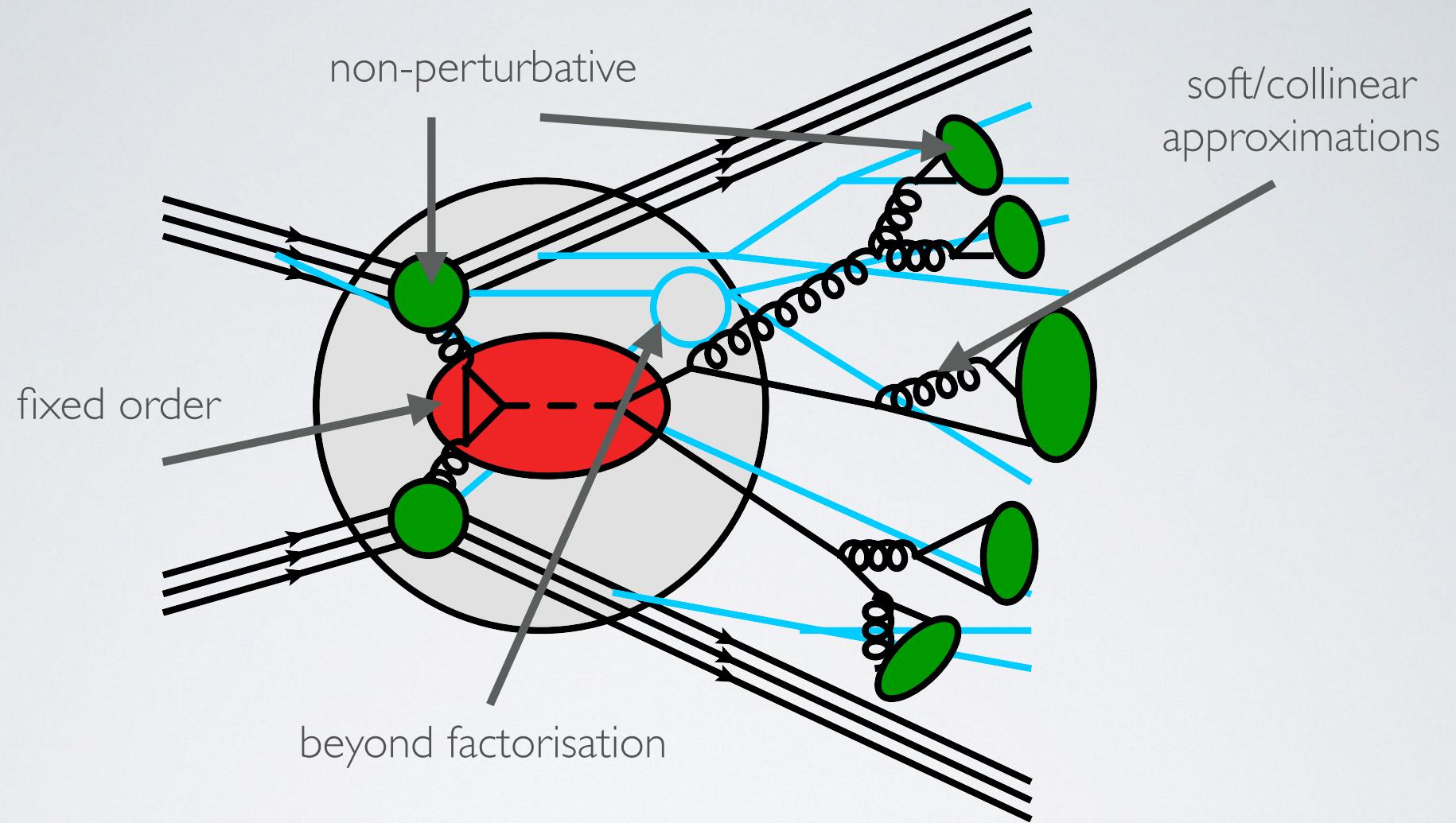
there will be missing energy

(can't see ν)

precise physical dimensions!

(need fully differential theory predictions)





numerical **event simulations** using Monte Carlo integration

'Monte Carlo Event Generators'

More on scattering at high energies

Factorisation at hadron colliders (e.g. LHC)

$$\hat{\sigma} = \frac{1}{\text{flux}} \int \delta^{(4)}(Q - \sum_i p_i) \prod_i d^4 p \delta^{(+)}(p_i^2 - m_i^2) \langle |\mathcal{A}|^2 \rangle$$

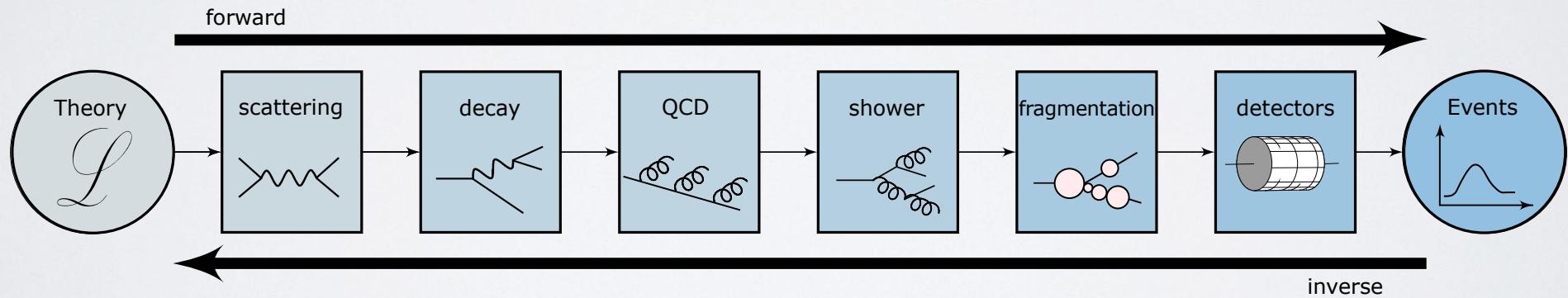
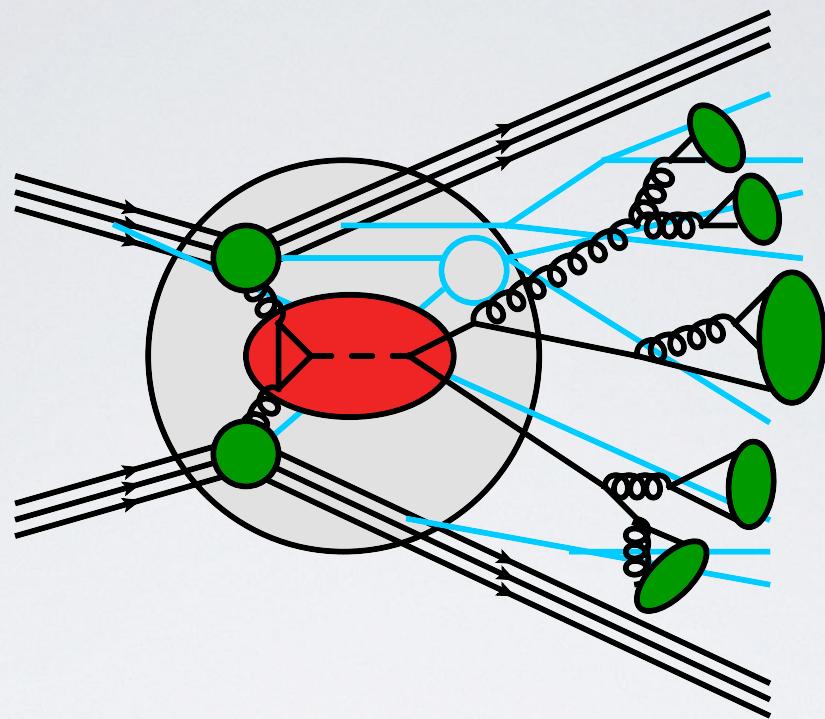
$$\sigma \sim \int dx_1 dx_2 f_i(x_1, Q^2) f_j(x_2, Q^2) \hat{\sigma}(ij \rightarrow \{p\}) \Theta + \mathcal{O}(Q/\Lambda)$$

PDF
parton distribution function

observable
(very broad definition)

beyond leading
order factorisation

The diagram illustrates the decomposition of the cross-section σ into three components. Three arrows point from the terms in the equation to their respective labels: one arrow points from $f_i(x_1, Q^2) f_j(x_2, Q^2)$ to 'PDF parton distribution function'; another arrow points from $\hat{\sigma}(ij \rightarrow \{p\})$ to 'observable (very broad definition)'; and a third arrow points from $\Theta + \mathcal{O}(Q/\Lambda)$ to 'beyond leading order factorisation'.



what precision do we need?

determination of
SM parameters

a lot!

general searches for
BSM resonances

less...

e.g. Drell-Yan with ATLAS (2019)

Many experimental measurements approach 1% precision

Data	
$\sigma(W^+ \rightarrow \mu^+\nu)$ [pb]	3110 ± 0.5 (stat.) ± 28 (syst.) ± 59 (lumi.)
$\sigma(W^- \rightarrow \mu^-\bar{\nu})$ [pb]	2137 ± 0.4 (stat.) ± 21 (syst.) ± 41 (lumi.)
Sum [pb]	5247 ± 0.6 (stat.) ± 49 (syst.) ± 100 (lumi.)
Ratio	1.4558 ± 0.0004 (stat.) ± 0.0040 (syst.)

QCD is the largest coupling and so dominates the perturbative expansion

$$\alpha_s(M_z^2) = 0.117 \pm 0.0009$$

PDG world average (2021)

$$\text{c.f. } \alpha^{(-1)}(0) = 137.035999150(33)$$

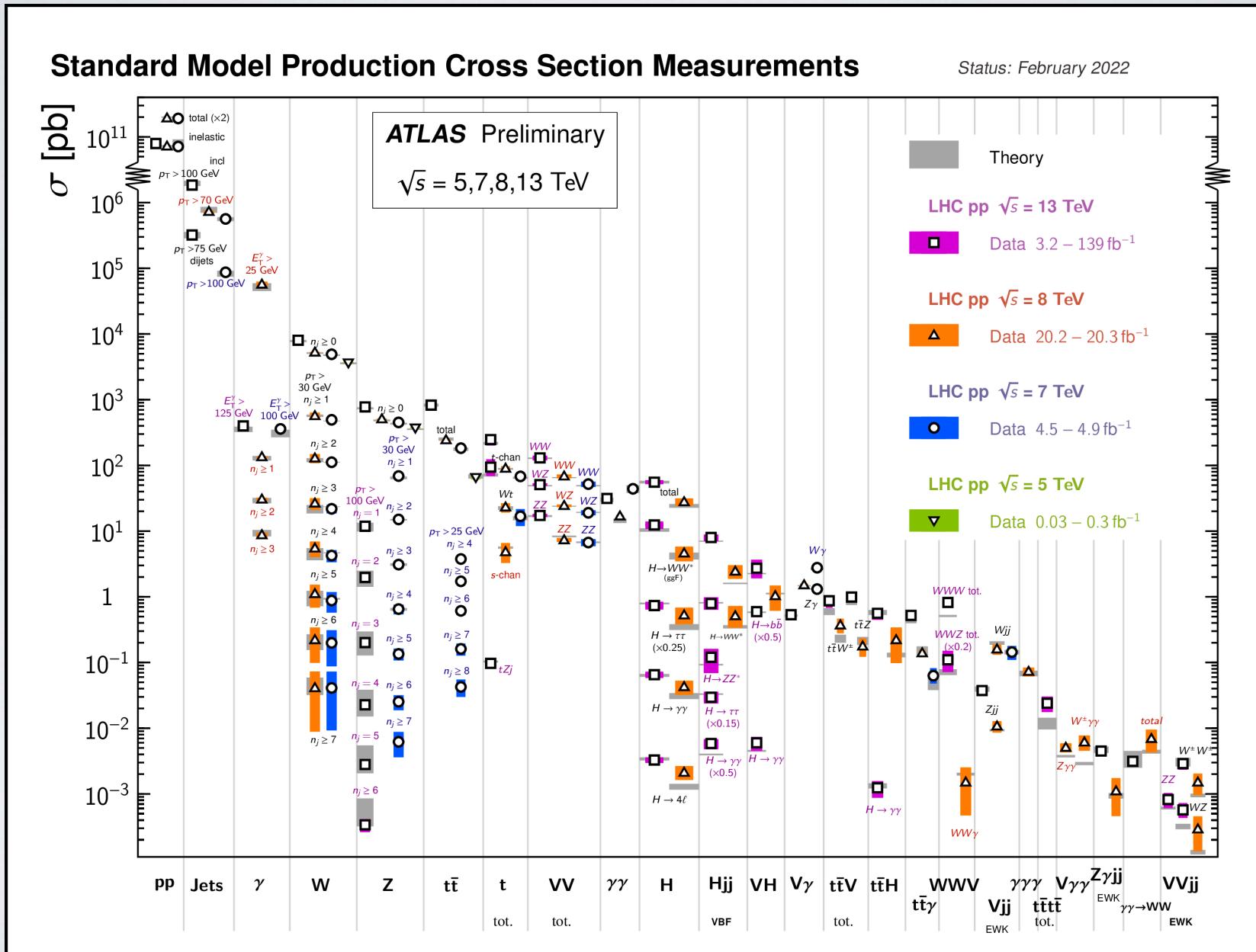
$$d\sigma = d\sigma^{\text{LO}} + \alpha_s d\sigma^{\text{NLO}} + \alpha_s^2 d\sigma^{\text{NNLO}}$$

~10-30 % **~1-10 %**

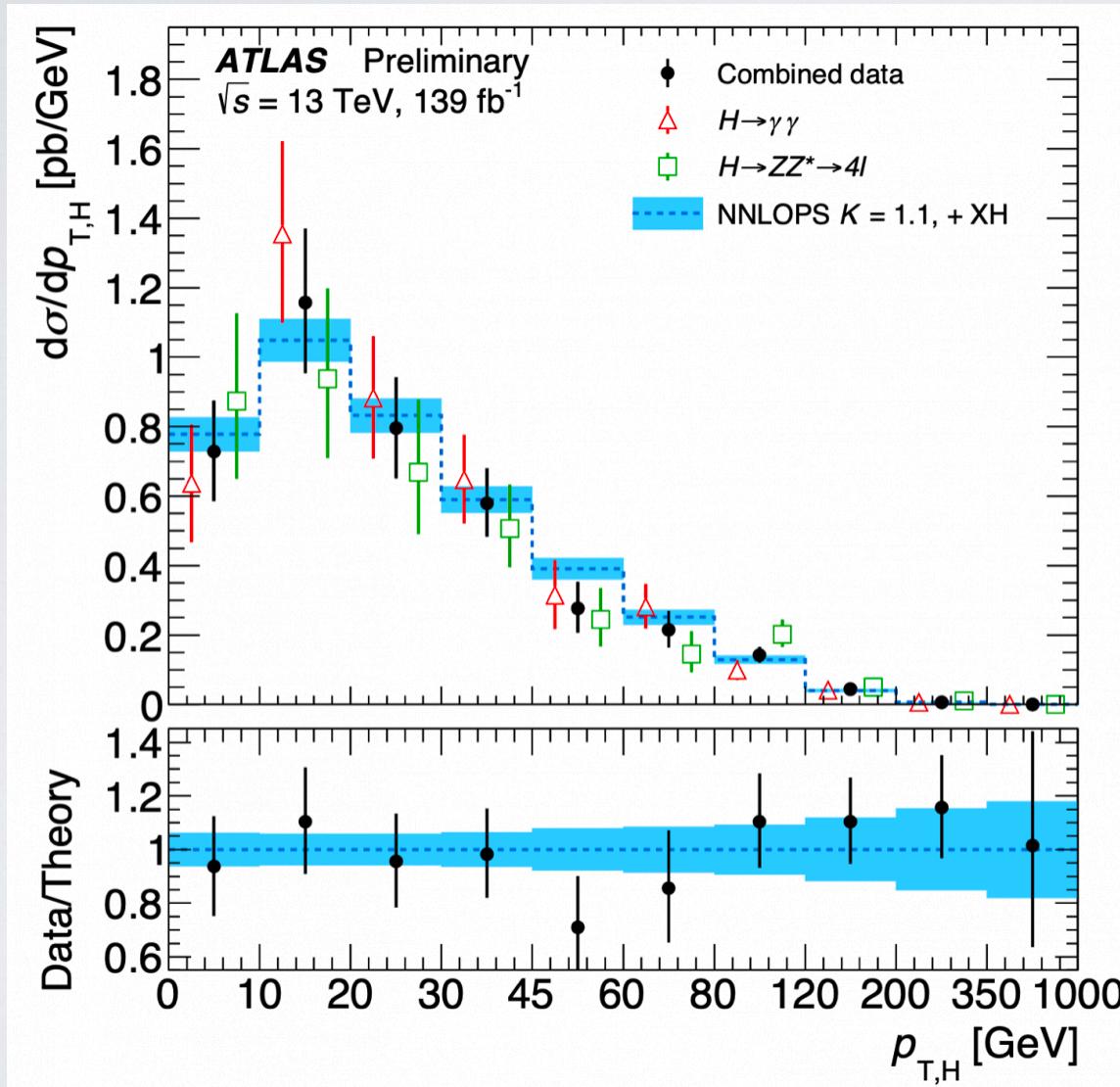
remember:
RG improved PT
is asymptotic

differentially many effects play a role: EW corrections, mass effects etc.

Current Standard Model Tests

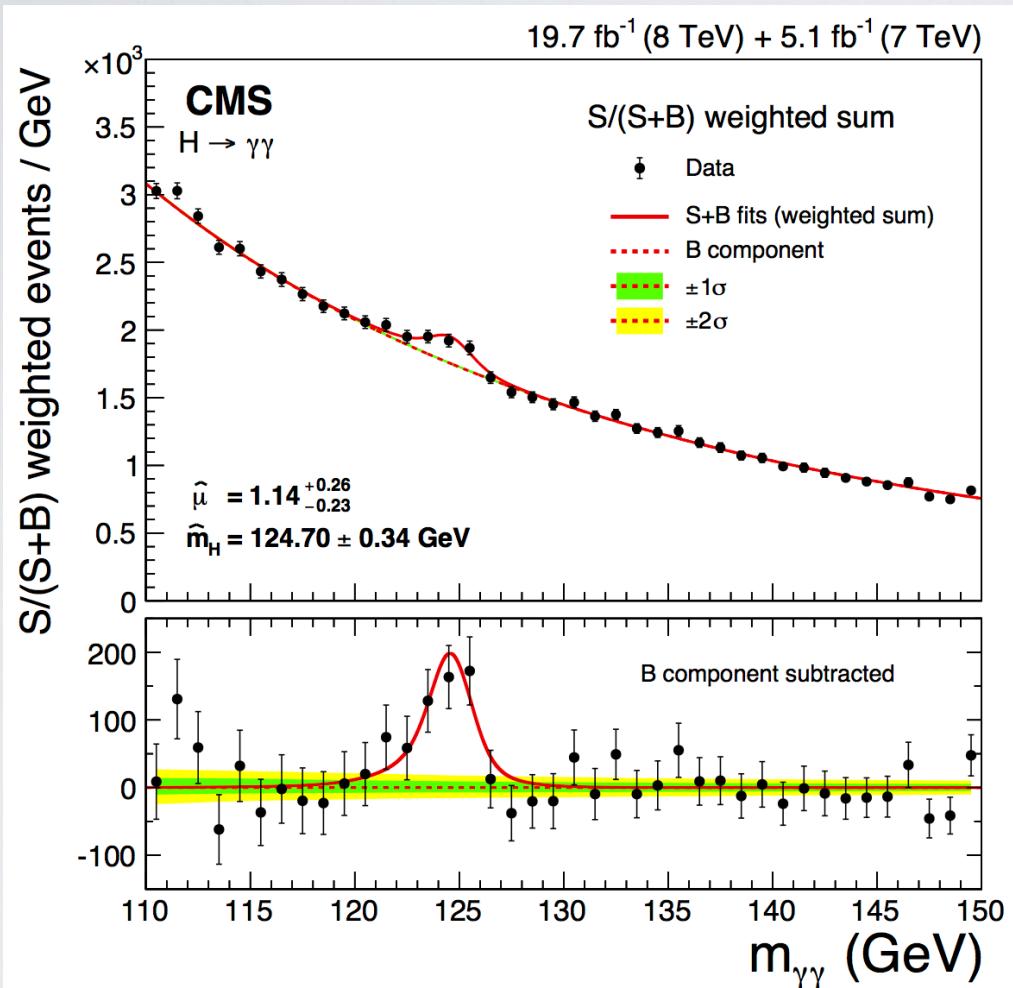


Differential tests offer more information



area under the curve is
the cross section
(Sezione d'urto)

Finding new particles



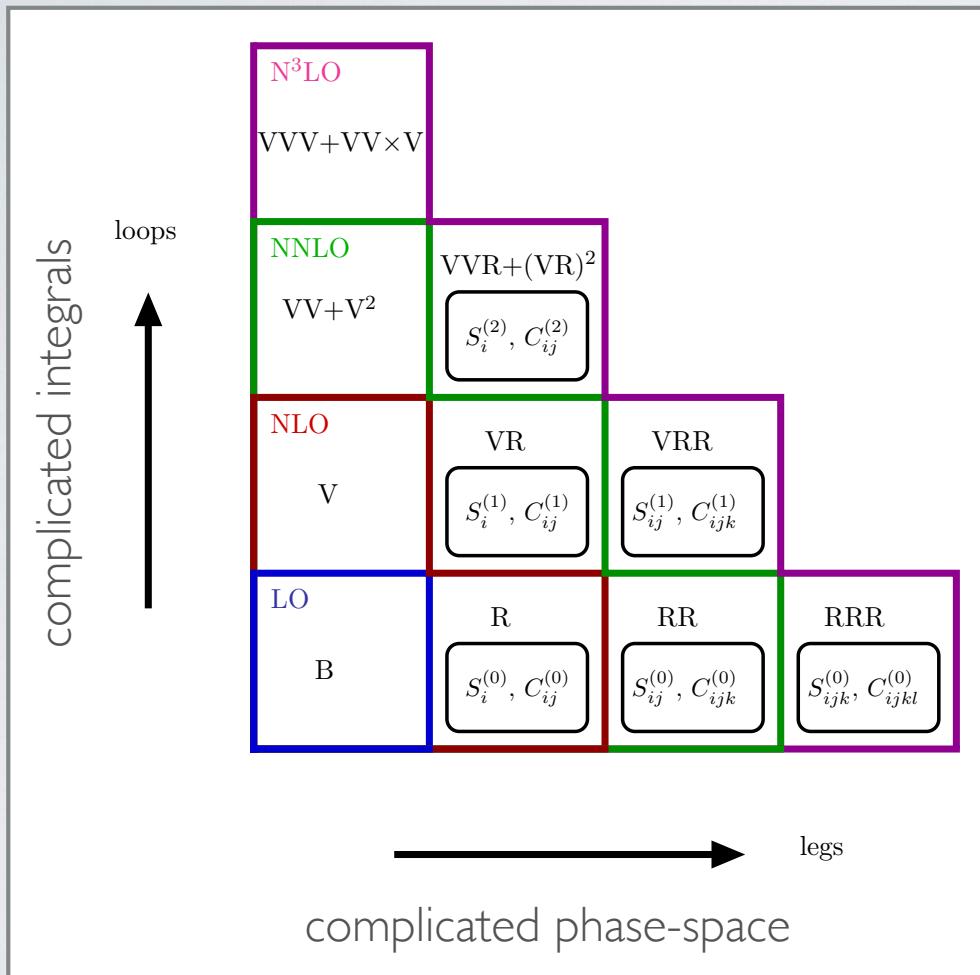
bump hunting is relatively easy

$$\mathcal{A} \sim \frac{1}{p^2 - m_h^2 - im_h\Gamma_h}$$

leads to the
Breit-Wigner distribution

Γ_h = decay width

Precision theory for experiments: quantum corrections from fixed order calculations



$$d\sigma \left(\text{black circle} + X \right) = \text{one loop} + \text{two loops} + \text{three loops} + \dots + \mathcal{O}(\alpha_s^8)$$

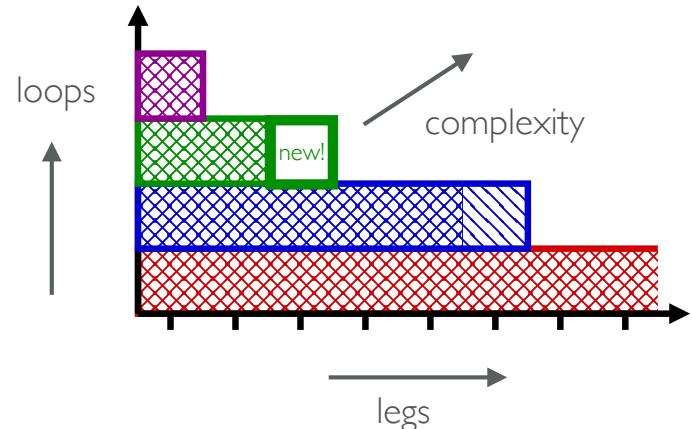
The equation shows the expansion of a differential cross-section $d\sigma$ in terms of coupling constants. It starts with a black circle representing the tree-level contribution, followed by a sum of diagrams involving one loop, two loops, three loops, and higher-order terms, plus a term proportional to $\mathcal{O}(\alpha_s^8)$.

keeping theory in line with experiments takes **years** of dedicated effort

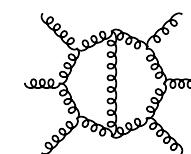
Growing Complexity

loops	1	2	3	4	5
diagrams	5	30	450	50,000	1.5×10^6
year	1973	1974	1980/1993	1997/2005	2016

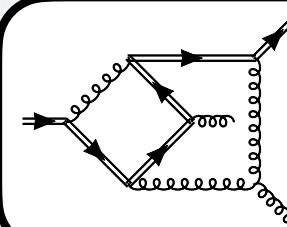
Diagrams contributing to QCD β function up to 5 loops



more scales = more complicated



algebraic complexity
e.g. six-gluon scattering

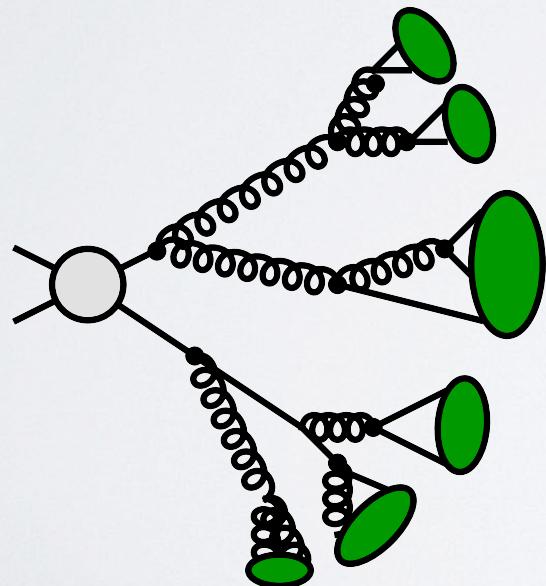


analytic complexity
e.g. $pp \rightarrow tt$

More radiation: e.g. parton showers

The cancellation of IR divergences is connected to the universal singular behaviour of QCD

The probability of a parton emitting a soft or collinear gluon is independent of the hard scattering process



parton showers give a statistical evolution of the hard scattering process until the hadronisation scale is reached

the effect is an explicit re-summation of soft emissions **beyond fixed order** perturbation theory

General purpose event generators (long term projects)

Pythia (1982-)

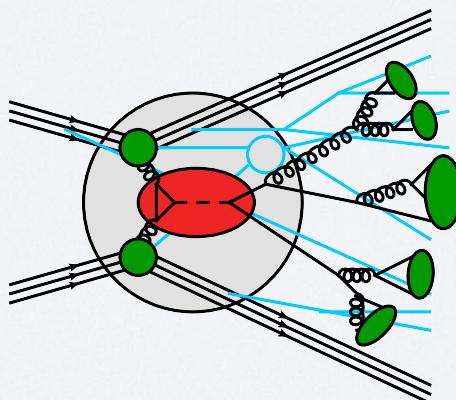
Herwig (1986-)

Sherpa (2002-)

Powheg (2004-)

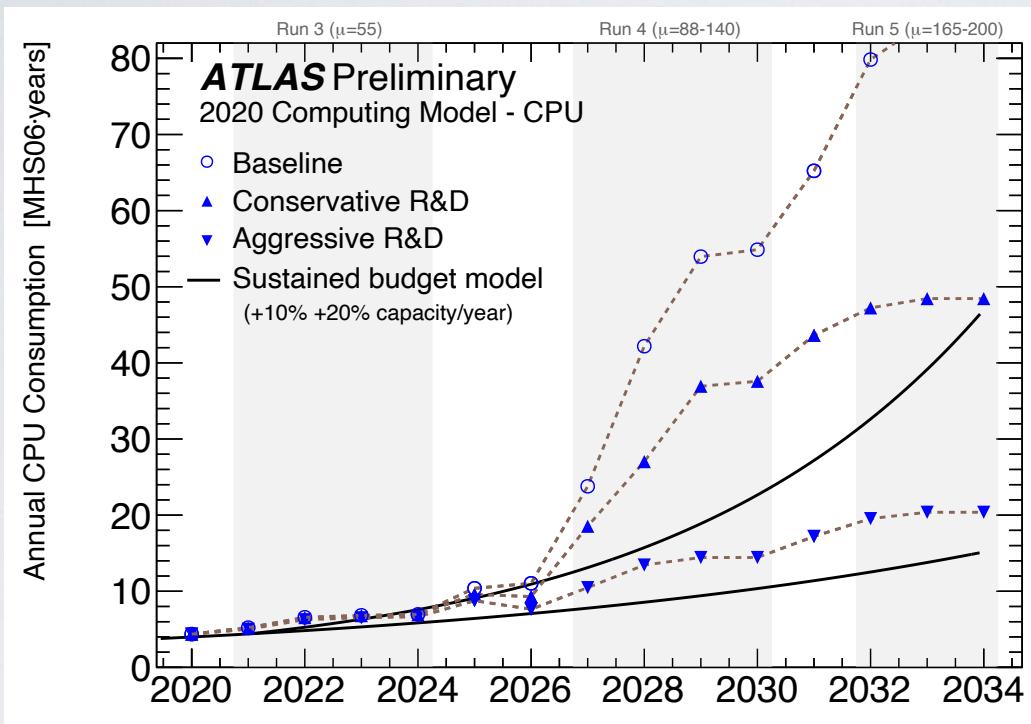
MadGraph (1994-)

collect large numbers
of different models
(hadronisations,
parton showers etc.)
and interface with
specialised codes for
PDFs, amplitudes etc.



what's the computational cost?

LHC MC computing requirement projections

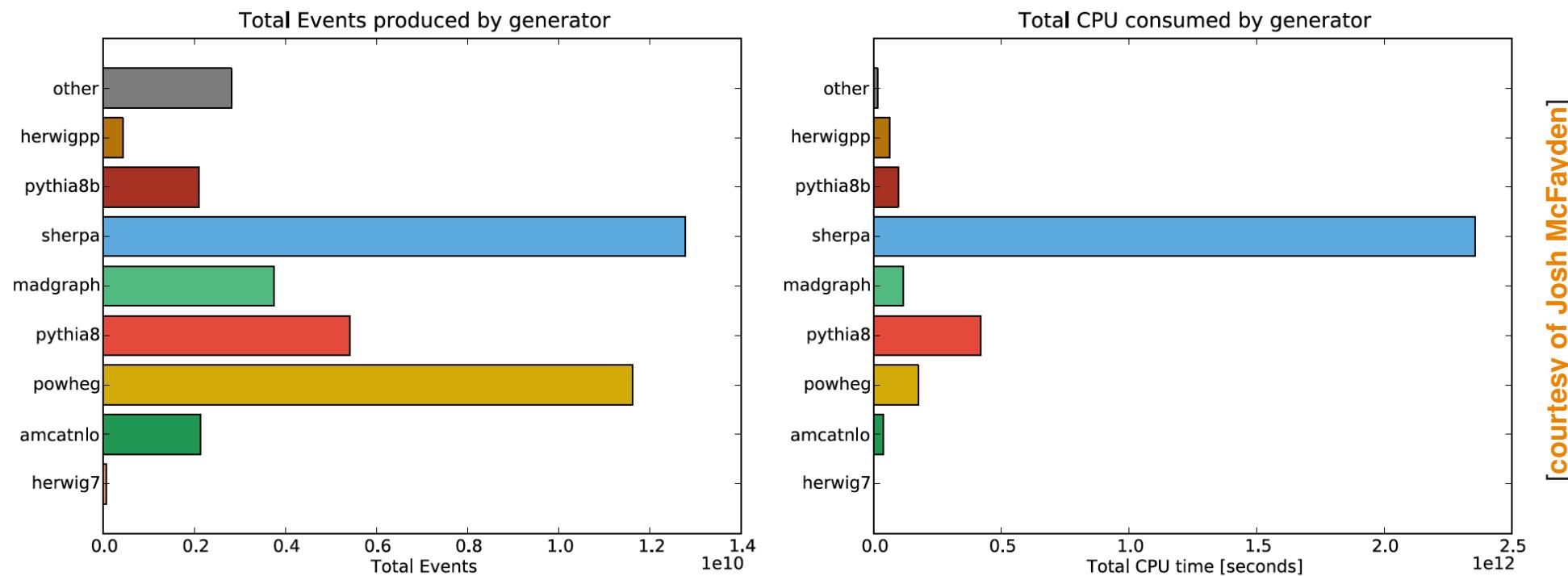


- Moore's Law no longer applies
- HL-LHC could have limitations from MC statistics without new developments

from
Christian Güttschow
at
Taming the accuracy of event generators
workshop CERN, June 2020

Breakdown by generator (bit outdated, but not too bad)

- left plot: does not account for alternative multi-leg setups
- right plot: most CPU spent on high-precision multi-leg calculations
(e.g. for ATLAS: $V + 0, 1, 2j$ @NLO+3, $4j$ @LO and $t\bar{t} + 0, 1j$ @NLO+2, 3, $4j$ @LO)

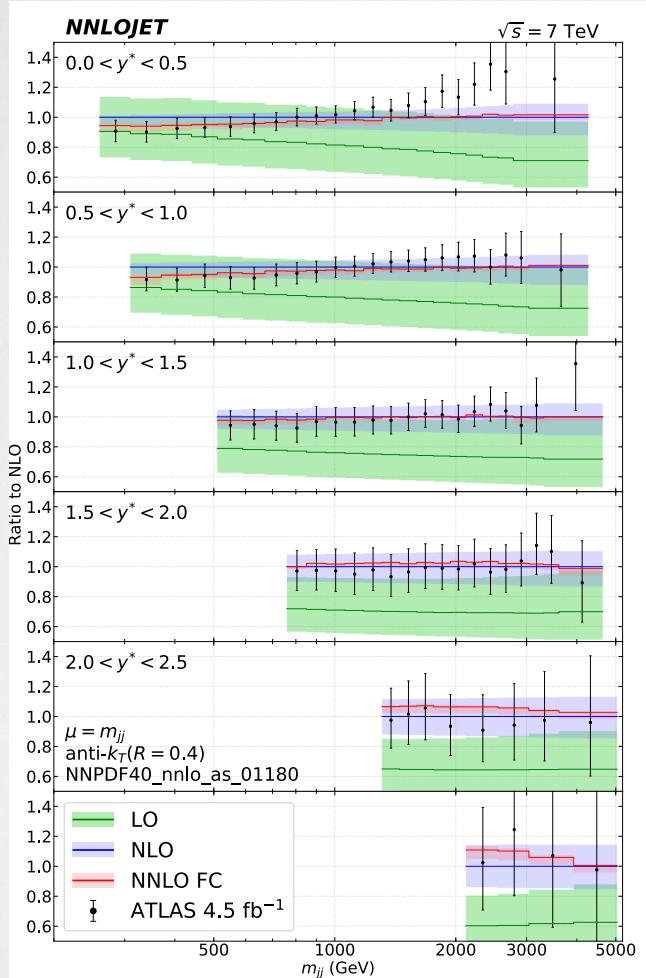


[courtesy of Josh McFayden]

- outlook: CPU spent on expensive setups expected to increase faster than for fast setups

and for NNLO precision?

many are still specialized (private) codes, NNLOjet, MATRIX, etc.



for some of the most expensive
fixed order simulations

~ 200,000 CPU hours

example: full colour di-jets [Chen et al 2204.10173]

Making Precision Predictions

As mentioned before QFT contains divergences in both **UV** and **IR**

observables such as (differential) cross sections must be **inclusive** over the phase-space in order to capture the cancellations from **unresolved** radiation configurations

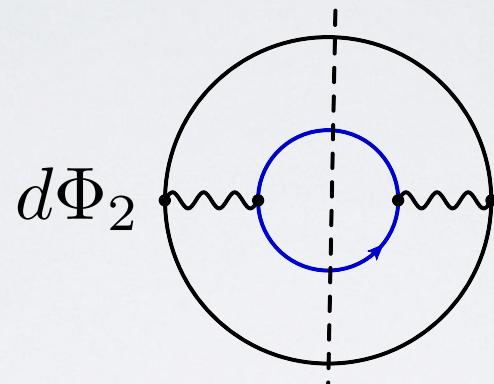
$$\sigma \left(\text{diagram} + X \right) = \int d\Omega_1 \left(\text{diagram} \right)^2 + \int d\Omega_2 \left(\text{diagram} \right)^2 + \int d\Omega_3 \left(\text{diagram} \right)^2 + \dots$$

now consider the perturbative expansion of these amplitudes

Making Precision Predictions

QCD corrections to $e^+e^- \rightarrow q\bar{q}$

LO

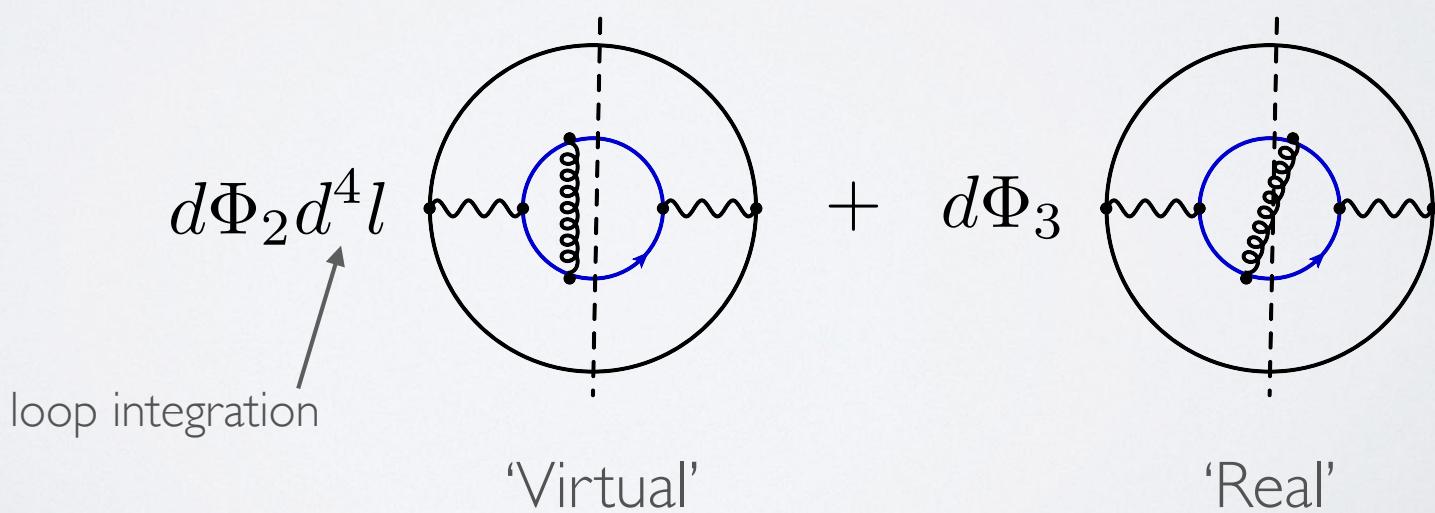


'Born'

IR divergences cancel between real and virtual corrections to physical observables.

Kinoshita-Lee-Nauenberg (KLN) theorem

NLO



Making Precision Predictions

The full calculation involves many steps and techniques...

