# **Notes**

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March 1, 2025

Here follow the calculation for the fermions 1PI at order  $O(\frac{1}{k})$ 

$$-i\Sigma(p) = -i\int \frac{d^3p}{(2\pi)^3} \gamma^{\mu} \frac{2\pi}{k} \epsilon_{\mu\nu\alpha} \frac{p^{\alpha}}{p^2} \gamma^{\nu} \frac{(\not p + \not q) + m}{(p+q)^2 - m^2} \tag{1}$$

#### **Denominator**

$$\frac{1}{p^2} \frac{1}{(p+q)^2 - m^2} = \int dx \frac{1}{[x((p+q)^2 - m^2) + (1-x)p^2]^2} = \int_0^1 dx \frac{1}{[l^2 - \Delta]^2}$$

$$l = p + xq \; ; \; \Delta = q^2 x(x-1) + xm^2$$

### Numerator

$$-i\frac{2\pi}{k}\gamma^{\mu}\gamma^{\nu}\epsilon_{\mu\nu\alpha}p^{\alpha}(\not\!p+\not\!q+m) = -i\frac{2pi}{k}(g^{\mu\nu}-i\epsilon^{\mu\nu\rho}\gamma_{\rho})\epsilon_{\mu\nu\alpha}p^{\alpha}(\not\!p+\not\!q+m) = \frac{-2\pi}{k}\not\!p(\not\!p+\not\!q+m)$$

### **Together**

$$\frac{-2\pi}{k} \int\limits_{0}^{1} dx \int \frac{d^{3}p}{(2\pi)^{3}} \frac{1}{[l^{2} - \Delta]^{2}} (pp + pq + pm)$$

Changing to l

$$p(p+q+m) = (l-xq)(l-xq+q+m) = l^2 - x lq + lq + lm - xq l - xq q - xq m + x^2 q^2$$

Getting rid of linear contributions in l

$$-\frac{2\pi}{k}\int\limits_{0}^{1}dx\int\frac{d^{3}l}{(2\pi)^{3}}\frac{l^{2}}{[l^{2}-\Delta]^{2}}-\frac{2\pi}{k}\int\limits_{0}^{1}dx[q^{2}x(x-1)-xm\phi]\int\frac{d^{3}l}{(2\pi)^{3}}\frac{1}{[l^{2}-\Delta]^{2}}$$

Wick rotating:

$$-i\frac{2\pi}{k}\int\limits_{0}^{1}dx\int\frac{d^{3}l}{(2\pi)^{3}}\frac{-l^{2}}{[l^{2}+\Delta]^{2}}-i\frac{2\pi}{k}\int\limits_{0}^{1}dx[q^{2}x(x-1)-xm\phi]\int\frac{d^{3}l}{(2\pi)^{3}}\frac{1}{[l^{2}+\Delta]^{2}}$$

NOTE: the first integral clearly shows UV divergences:

$$\int\limits_{a}^{\Lambda}d^{3}l\frac{l^{2}}{[l^{2}+\delta]^{2}}\sim^{a\to\infty}\int\limits_{a}^{\Lambda}dl\frac{l^{4}}{l^{4}}\sim\Lambda$$

still, dimensional regulation is blind to power-law divergences, it only shows logarithmic ones.

Let us us dim-reg.

Knowing:

$$\int \frac{d^d l}{(2\pi)^d} \frac{l^2}{[l^2 + \Delta]^2} = \frac{d}{2} \frac{\Gamma(1 - \frac{d}{2})}{(4\pi)^{d/2}} \frac{1}{\Lambda^{1 - \frac{d}{2}}} \; \; ; \; \; \int \frac{d^d l}{(2\pi)^3} \frac{1}{[l^2 + \Delta]^2} = \frac{\Gamma[2 - \frac{d}{2}]}{(4\pi)^{d/2}} \frac{1}{\Lambda^{2 - \frac{d}{2}}}$$

No need to set  $d=3-\epsilon$ , since the  $\Gamma$  functions would show no pole for  $\epsilon\to 0$ . It gives finally:

$$\begin{split} -i\Sigma(p) &= -i\frac{3}{4k}\int\limits_{0}^{1}dx\sqrt{q^{2}x(x-1)+xm^{2}} \\ &-i\frac{1}{4k}(\int\limits_{0}^{1}dx\frac{q^{2}x(x-1)}{\sqrt{q^{2}x(x-1)+xm^{2}}} - \int\limits_{0}^{1}dx\frac{xmq}{\sqrt{q^{2}x(x-1)+xm^{2}}}) \end{split}$$

From which:

$$\delta Z_m \cdot m = \Sigma(p = m) = \frac{3}{4k} \int_0^1 dx \ m \cdot x + \frac{1}{4k} \left[ \int_0^1 dx \frac{m^2 x(x - 1)}{mx} - \int_0^1 dx \frac{xm^2}{xm} \right] = 0$$

at order  $O(\frac{1}{k})$ 

# 1 Fermion self energy calculation

Here follow the calculation for the fermions 1PI at order  $O(\frac{1}{k})$ 

$$-i\Sigma(\mathbf{q}) = -i\int \frac{d^3p}{(2\pi)^3} \gamma^{\mu} \frac{2\pi}{k} \epsilon_{\mu\nu\alpha} \frac{p^{\alpha}}{p^2} \gamma^{\nu} \frac{(\mathbf{p} + \mathbf{q}) + m}{(p+q)^2 - m^2}$$
(2)

The **numerator** can be written as:

$$\begin{split} -i\frac{2\pi}{k}\gamma^{\mu}\gamma^{\nu}\epsilon_{\mu\nu\alpha}p^{\alpha}(\not\!p+\not\!q+m) &= -i\frac{2\pi}{k}(g^{\mu\nu}-i\epsilon^{\mu\nu\rho}\gamma_{\rho})\epsilon_{\mu\nu\alpha}p^{\alpha}(\not\!p+\not\!q+m) \\ &= \frac{-2\pi}{k}(2\delta^{\rho}_{\alpha})\gamma_{\rho}p^{\alpha}(\not\!p+\not\!q+m) = \frac{-4\pi}{k}\not\!p(\not\!p+\not\!q+m) \\ &= -\frac{4\pi}{k}(p^2\mathbf{1}+\not\!p(\not\!q+m)) \end{split}$$

Which gives:

$$\begin{split} -i\Sigma(\mathbf{p}) &= \frac{-4\pi}{k} \int \frac{d^3}{(2\pi)^3} \frac{p^2 + \mathbf{p}(\mathbf{p} + m)}{p^2[(p+q)^2 - m^2]} \\ &= -\frac{4\pi}{k} \left[ \int \frac{d^3p}{(2\pi)^3} \frac{\mathbf{1}}{(p+q)^2 - m^2} + \gamma_\mu \left( \int \frac{d^3p}{(2\pi)^3} \frac{p^\mu}{p^2[(p+q)^2 - m^2]} \right) (\mathbf{p} + m) \right] \end{split}$$

Focusing on the ingral between partenthesis: in order the preserve the Lorentz structure, and since the only external momentum is q, it must be:

$$\int \frac{d^3p}{(2\pi)^3} \frac{p^{\mu}}{p^2[(p+q)^2 - m^2]} = A(q^2) \cdot q^{\mu}$$
$$A(q^2) = \frac{1}{q^2} \int \frac{d^3p}{(2\pi)^3} \frac{p \cdot q}{p^2[(p+q)^2 - m^2]}$$

Substituting  $p\cdot q=\frac{1}{2}[(p+q)^2-m^2-p^2-q^2+m^2]$ 

$$A(q^2) = \frac{1}{2q^2} \int \frac{d^3}{(2\pi)^3} \left[ \frac{1}{p^2} - \frac{1}{(p+q)^2 - m^2} + \frac{m^2 - q^2}{p^2[(p+q)^2 - m^2]} \right]$$

Placing it inside the total integral:

$$-i\Sigma(\mathbf{q}) = -\frac{4\pi}{k} \left\{ \int \frac{d^3p}{(2\pi)^3} \frac{\mathbf{1}}{(p+q)^2 - m^2} + \frac{\mathbf{q}(\mathbf{q}+m)}{2q^2} \int \frac{d^3}{(2\pi)^3} \left[ \frac{1}{p^2} - \frac{1}{(p+q)^2 - m^2} + \frac{m^2 - q^2}{p^2[(p+q)^2 - m^2]} \right] \right\}$$

Now, usingin dim-reg and recalling:

$$\int \frac{d^d k}{(2\pi)^d} \frac{1}{(k^2 + \Delta)^n} = \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(n - d/2)}{\Gamma(n)} \left(\frac{1}{\Delta}\right)^{n - d/2}$$

We can evaluate the different integrals inside  $\Sigma(q)$ . (WR: wick rotation, DR: dim-reg)

$$\int \frac{d^3p}{(2\pi)^3} \frac{1}{(p+q)^2 - m^2} = \int \frac{d^3l}{(2\pi)^3} \frac{1}{l^2 - m^2} \stackrel{WR}{=} -i \int \frac{d^3l}{(2\pi)^3} \frac{1}{l^2 + m^2}$$

$$\stackrel{DR}{=} \frac{-i}{8\pi^{3/2}} (-2\sqrt{\pi}) \cdot m = i \frac{m}{4\pi}$$

$$\int \frac{d^3p}{(2\pi)^3} \frac{1}{p^2} \stackrel{DR}{=} 0$$

$$\int \frac{d^3p}{(2\pi)^3} \frac{1}{p^2[(p+q)^2 - m^2]} = \int_0^1 dx \int \frac{d^3k}{(2\pi)^3} \frac{1}{[k^2 - \Delta]^2} \stackrel{WR}{=} i \int_0^1 dx \int \frac{d^3k}{(2\pi)^3} \frac{1}{[k^2 + \Delta]^2}$$

$$\stackrel{DR}{=} i \frac{\sqrt{\pi}}{8\pi\sqrt{\pi}} \int_0^1 dx \frac{1}{(q^2x(x-1) + xm^2)^{\frac{1}{2} + \varepsilon}}$$

Where  $l=p+q; \;\; k=p+xq; \;\; \Delta=q^2x(x-1)+xm^2$ 

**Note:** in the first integral I substituted d=3 directly, whereas in the last one  $d=3-2\varepsilon$ . The reason is to be found in the counter-term calculation at the next page.

This results in:

$$-i\Sigma(\mathbf{q}) = -\frac{4\pi}{k} \left\{ i\frac{m}{4\pi} + \frac{1}{2} \left( \mathbf{1} + \frac{\mathbf{q}m}{q^2} \right) \left[ -i\frac{m}{4\pi} + \frac{i}{8\pi} (m^2 - q^2) \int\limits_0^1 dx \frac{1}{(q^2 x (x-1) + x m^2)^{\frac{1}{2} + \varepsilon}} \right] \right\}$$

## 2 Counter-terms

In the ON-SHELL scheme:

$$\delta Z_m \cdot m = \Sigma(\not p = m) = 0$$

$$\begin{split} \delta Z_{\psi} &= \frac{d}{d \mathbf{q}} \Sigma(\mathbf{q})|_{\mathbf{q}=m} = \ldots = \frac{4\pi}{k} \left\{ \frac{1}{8\pi} \frac{m^2}{q^2} + \frac{d}{d \mathbf{q}} \left[ \frac{1}{8\pi} (m^2 - q^2) \int\limits_0^1 dx \frac{1}{(q^2 x (x-1) + x m^2)^{\frac{1}{2} + \varepsilon}} \right] \right\} \bigg|_{\mathbf{q}=m} \\ &= \frac{1}{2k} - \frac{1}{k} \int\limits_0^1 dx \, \frac{1}{x^{1+2\varepsilon}} = \frac{1}{2k} \left( 1 + \frac{1}{\varepsilon} \right) \end{split}$$