

# **Golden-Phase Prime-Shift Operators**

Gauge-Twisted Arithmetic Dynamics  
and Spectral Probes of Prime Structure

Enter Yourname

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## Abstract

In this work we introduce a class of arithmetic operators obtained by equipping prime-shift dynamics on  $\ell^2(\mathbb{N})$  with nontrivial  $U(1)$  gauge phases. These phases are not added ad hoc, but arise naturally from potential differences assigned to arithmetic states, producing a gauge-covariant prime-shift operator with an almost symmetric forward–backward structure.

A central role is played by *golden-phase gauges*, in which the phase potential is constructed using the golden ratio  $\varphi$  or Fibonacci-based coordinates. The golden ratio is distinguished by its extremal irrationality and optimal equidistribution properties, making it a natural candidate for probing hidden coherence in arithmetic structures. In this sense, the golden phase is interpreted not as an explanatory mechanism for prime distribution, but as a spectral probe: a maximally nonresonant gauge against which arithmetic interference patterns may be tested.

We show that the resulting golden-phase prime-shift operator admits a natural interpretation as a gauge-covariant adjacency operator on an arithmetic graph with edges  $n \leftrightarrow pn$ . On a dense domain of finitely supported states, the operator is formally symmetric, with phase factors arranged in conjugate pairs along inverse transitions. This structure closely parallels discrete gauge theories on lattices, while remaining intrinsically arithmetic.

Beyond its formal definition, the operator provides a framework for numerical and spectral experimentation. By comparing the spectra of truncated operators with and without golden-phase twisting, one may test for phase-sensitive correlations in prime-induced dynamics. Possible outcomes range from complete phase decorrelation to the emergence of stable interference patterns, each carrying distinct implications for operator-based approaches to prime structure.

The construction is compatible with the operator-first philosophy of Quansistor Field Mathematics (QFM) and may be viewed as a gauge refinement of previous prime-shift Hamiltonians. While no claims of explaining prime distribution are made, the framework establishes a falsifiable and extensible setting in which arithmetic dynamics, gauge structure, and spectral analysis interact in a controlled and conceptually transparent manner.

# 1 Motivation and Conceptual Context

The distribution of prime numbers occupies a singular position in mathematics. On the one hand, primes are elementary objects: the indivisible atoms of arithmetic. On the other hand, their global behavior resists direct description, oscillating between apparent randomness and deep structural regularity. This tension has motivated a long history of attempts to reinterpret prime numbers through analytic, geometric, and spectral frameworks.

A recurring theme in these efforts is the idea that prime structure may be more naturally accessed through operators rather than explicit formulas. From the Hilbert–Pólya conjecture to modern operator-theoretic approaches, the underlying intuition is that primes are not merely static data points, but manifestations of a deeper arithmetic dynamics whose invariants are spectral in nature.

In this work we pursue this operator-first perspective, but with a specific focus: the role of *phase* as a probe of arithmetic coherence.

## 1.1 Phase as a Diagnostic Rather Than an Explanation

The introduction of phases in arithmetic contexts is often viewed with suspicion, as phases may suggest hidden periodicity or numerological constructions. Here we adopt a different stance. The phase introduced in this paper is not intended to explain the distribution of primes. Instead, it functions as a diagnostic tool: a controlled perturbation designed to test whether prime-induced dynamics exhibits coherent interference under maximally nonresonant rotations.

In physical systems, phase coherence reveals itself spectrally. If a system is truly random with respect to a given gauge, phase twisting produces no persistent spectral signal. If, however, hidden structure is present, interference patterns may survive truncation, averaging, and regularization. Our approach applies this logic to arithmetic dynamics.

## 1.2 Why the Golden Ratio

The golden ratio

$$\varphi = \frac{1 + \sqrt{5}}{2}$$

is distinguished among irrational numbers by its extremal Diophantine properties. It is the most poorly approximable irrational by rationals, and sequences generated by rotations modulo  $\varphi$  exhibit near-optimal equidistribution.

This property underlies the appearance of the golden angle in phyllotaxis, quasi-crystals, and low-discrepancy sampling schemes. In all these contexts,  $\varphi$  does not impose structure; rather, it minimizes spurious alignment.

For this reason, the golden ratio is an ideal candidate for an arithmetic gauge. If prime-induced dynamics displays phase coherence even under golden-phase twisting, such coherence cannot be attributed to accidental resonance. Conversely, the absence of any

persistent spectral effect provides evidence that the dynamics is phase-blind with respect to optimal irrational rotations.

### 1.3 Arithmetic Dynamics as a Graph

The natural habitat for the operators studied here is not the real line but the discrete set  $\mathbb{N}$ , equipped with a multiplicative adjacency structure. Each prime  $p$  generates transitions

$$n \longleftrightarrow pn,$$

defining an infinite, directed arithmetic graph. Functions in  $\ell^2(\mathbb{N})$  may be viewed as quantum states supported on this graph, while prime-shift operators act as weighted adjacency operators.

This perspective emphasizes that prime multiplication is not merely an arithmetic operation, but a dynamical move within a structured state space. The addition of phase to these transitions corresponds to assigning a local geometric structure to the arithmetic graph, analogous to a gauge field on a lattice.

### 1.4 Gauge Structure and Symmetry

The prime-shift operators considered in this paper are equipped with  $U(1)$  gauge phases derived from potential differences between arithmetic states. This ensures that forward and backward transitions carry conjugate phases, producing an operator that is formally symmetric on natural dense domains.

Importantly, these phases are not arbitrary decorations. They arise from a single potential function on  $\mathbb{N}$ , guaranteeing consistency across composite paths in the arithmetic graph. As in lattice gauge theory, such a construction preserves local symmetry while allowing global interference effects.

### 1.5 Relation to Operator-Based Arithmetic Programs

The present work aligns with a broader operator-theoretic program in arithmetic, including the Hilbert–Pólya conjecture and more recent prime-shift Hamiltonians. Within the framework of Quansistor Field Mathematics (QFM), arithmetic operators are interpreted as fundamental dynamical entities, with spectra encoding invariant information about number-theoretic structure.

The golden-phase prime-shift operator introduced here should be viewed as a gauge refinement of these ideas. It neither replaces nor contradicts existing constructions, but extends them by introducing a tunable, conceptually motivated phase degree of freedom.

### 1.6 Philosophical Perspective

At a philosophical level, this work adopts the view that arithmetic structure may not be fully visible in static representations. Just as physical systems reveal their laws through

dynamics and response to perturbation, arithmetic systems may reveal hidden regularities only when probed through carefully chosen operators.

The golden phase plays the role of such a probe. It does not assert order where none exists, nor does it impose symmetry. It merely asks a precise question: whether prime-induced dynamics remains spectrally inert under a maximally nonresonant gauge, or whether subtle coherence persists.

The remainder of this paper develops the formal construction of golden-phase prime-shift operators, analyzes their symmetry and spectral properties, and proposes numerical protocols through which these questions may be tested in a falsifiable manner.

## 2 Prime-Shift Operators on $\ell^2(\mathbb{N})$

In order to introduce gauge phases in a controlled manner, we first establish the untwisted arithmetic operator that forms the backbone of the construction. This operator encodes prime multiplication and division as transitions on  $\mathbb{N}$  and serves as the reference dynamics against which phase effects will later be measured.

### 2.1 The Arithmetic State Space

Let  $\ell^2(\mathbb{N})$  denote the Hilbert space of square-summable complex-valued functions on the natural numbers, equipped with the standard inner product

$$\langle \psi, \phi \rangle = \sum_{n=1}^{\infty} \overline{\psi(n)} \phi(n).$$

Elements of  $\ell^2(\mathbb{N})$  will be interpreted as arithmetic states, with  $\psi(n)$  representing the amplitude assigned to the integer  $n$ .

Throughout this work, we regard  $\mathbb{N}$  not merely as a set, but as a graph whose structure is induced by prime multiplication. Each prime  $p$  defines a directed edge

$$n \longrightarrow pn,$$

with a corresponding inverse edge  $pn \longrightarrow n$  whenever  $p \mid pn$ . This defines an infinite, locally finite arithmetic graph encoding the multiplicative structure of the integers.

### 2.2 Prime-Shift Transitions

The basic dynamical moves on this graph are prime shifts. For each prime  $p$ , we consider the forward shift

$$(T_p^+ \psi)(n) = \psi(pn),$$

and the backward shift

$$(T_p^- \psi)(n) = \mathbf{1}_{p|n} \psi(n/p),$$

where  $\mathbf{1}_{p|n}$  denotes the indicator function of divisibility by  $p$ .

These operators represent the elementary transitions generated by prime multiplication and division. Their action is purely combinatorial and reflects the local structure of the arithmetic graph.

### 2.3 The Untwisted Prime-Shift Operator

We define the untwisted prime-shift operator  $H_0$  by

$$(H_0 \psi)(n) = \sum_{p \in \mathbb{P}} w(p) (\psi(pn) + \mathbf{1}_{p|n} \psi(n/p)),$$

where  $\mathbb{P}$  denotes the set of prime numbers and  $w(p)$  is a real-valued weight function. In many applications one takes

$$w(p) = \frac{1}{p},$$

which reflects the natural logarithmic scaling of prime contributions, though other choices are possible.

The operator  $H_0$  acts as a weighted adjacency operator on the arithmetic graph. Each term in the sum corresponds to a pair of directed edges linking  $n$  and  $pn$ , with symmetric weights assigned to forward and backward transitions.

## 2.4 Formal Symmetry

On the dense subspace  $c_{00}(\mathbb{N})$  of finitely supported functions, the operator  $H_0$  is formally symmetric. Indeed, for  $\psi, \phi \in c_{00}(\mathbb{N})$  one finds

$$\langle \psi, H_0\phi \rangle = \langle H_0\psi, \phi \rangle,$$

as each forward transition  $n \rightarrow pn$  is paired with the corresponding backward transition  $pn \rightarrow n$ .

This symmetry is purely algebraic and does not rely on convergence of the full prime sum. Questions of essential self-adjointness and spectral completeness are subtle and depend on the choice of weights and domains; these issues will be addressed only insofar as they are relevant for numerical and formal analysis.

## 2.5 Regularization and Truncation

Since the sum over primes diverges for typical choices of  $w(p)$ , the operator  $H_0$  is best understood as a formal object or as a family of regularized operators. Common regularization strategies include:

- imposing a prime cutoff  $p \leq P$ ,
- damping weights  $w(p)$  by a smooth factor,
- restricting the state space to  $n \leq N$ .

Such regularizations are not artifacts but essential components of any computational or spectral investigation. Importantly, all constructions introduced in later sections are compatible with these regularized settings.

## 2.6 Role as a Reference Dynamics

The operator  $H_0$  provides a natural reference dynamics for arithmetic systems. It captures the essential multiplicative structure of  $\mathbb{N}$  while remaining free of any phase information. In subsequent sections, this operator will be extended by the introduction of gauge phases on arithmetic transitions.

By comparing the spectral behavior of  $H_0$  with that of its gauge-twisted counterparts, we obtain a controlled framework in which phase sensitivity and interference effects can be meaningfully identified and analyzed.

### 3 U(1) Gauge Phases on Arithmetic Graphs

Having established the untwisted prime-shift operator as a reference dynamics, we now introduce gauge phases on arithmetic transitions. The purpose of this construction is not to modify the underlying combinatorial structure of prime multiplication, but to endow it with a local geometric degree of freedom.

The resulting operators may be viewed as gauge-covariant adjacency operators on the arithmetic graph introduced in the previous section.

#### 3.1 Arithmetic Potentials and Phase Differences

Let

$$\eta : \mathbb{N} \longrightarrow \mathbb{R}$$

be a real-valued function on the arithmetic state space. We interpret  $\eta(n)$  as a potential assigned to the integer  $n$ . From this potential we define phase factors on directed prime transitions by

$$U(n, p) = e^{2\pi i(\eta(pn) - \eta(n))}.$$

This construction assigns a  $U(1)$  phase to each directed edge

$$n \longrightarrow pn,$$

while ensuring that the phase associated with the inverse transition  $pn \longrightarrow n$  is given by the complex conjugate

$$\overline{U(n, p)} = e^{-2\pi i(\eta(pn) - \eta(n))}.$$

In this way, phase information is encoded exclusively through potential differences, guaranteeing consistency across composite paths in the arithmetic graph.

#### 3.2 Gauge Transformations

Given a potential function  $\eta$ , we define a unitary gauge transformation  $G_\eta$  on  $\ell^2(\mathbb{N})$  by

$$(G_\eta \psi)(n) = e^{2\pi i\eta(n)} \psi(n).$$

Two potentials  $\eta$  and  $\eta'$  related by

$$\eta'(n) = \eta(n) + c$$

for a constant  $c \in \mathbb{R}$  generate identical phase factors and therefore define the same gauge. More generally, the construction depends only on the equivalence class of  $\eta$  modulo additive constants.

Under a gauge transformation, the untwisted prime-shift operator  $H_0$  transforms

formally as

$$G_\eta^{-1} H_0 G_\eta,$$

producing an operator whose action includes the phase factors  $U(n, p)$ . This observation provides a direct link between gauge phases and unitary conjugation at the formal level.

### 3.3 Gauge-Covariant Prime-Shift Operators

Using the phase factors defined above, we introduce the gauge-covariant prime-shift operator

$$(H_\eta \psi)(n) = \sum_{p \in \mathbb{P}} w(p) (U(n, p) \psi(pn) + \mathbf{1}_{p|n} \overline{U(n/p, p)} \psi(n/p)).$$

Explicitly,

$$(H_\eta \psi)(n) = \sum_{p \in \mathbb{P}} w(p) \left( e^{2\pi i(\eta(pn) - \eta(n))} \psi(pn) + \mathbf{1}_{p|n} e^{2\pi i(\eta(n/p) - \eta(n))} \psi(n/p) \right).$$

This operator differs from  $H_0$  only by the presence of conjugate phase factors on forward and backward transitions. The underlying adjacency structure of the arithmetic graph remains unchanged.

### 3.4 Formal Symmetry

On the dense domain  $c_0(\mathbb{N})$  of finitely supported functions, the operator  $H_\eta$  is formally symmetric. Indeed, for  $\psi, \phi \in c_0(\mathbb{N})$ , each term corresponding to a transition  $n \leftrightarrow pn$  contributes complex conjugate phases when evaluated in the inner product, yielding

$$\langle \psi, H_\eta \phi \rangle = \langle H_\eta \psi, \phi \rangle.$$

This symmetry is a direct consequence of constructing phases from potential differences. No additional constraints on  $\eta$  are required beyond real-valuedness.

### 3.5 Interpretation as a Discrete Gauge Theory

The structure introduced above closely parallels lattice gauge theory in mathematical physics. The arithmetic graph plays the role of a lattice, prime multiplication defines adjacency, and the phase factors  $U(n, p)$  act as  $U(1)$  link variables.

However, unlike conventional lattices, the arithmetic graph is intrinsically multiplicative and nonuniform. Gauge phases therefore probe arithmetic structure rather than spatial geometry. In this sense, the construction defines a genuinely arithmetical gauge theory, in which phase coherence and interference reflect properties of prime-induced dynamics.

### 3.6 Preparation for Golden-Phase Gauges

The framework developed in this section is fully general and places no restrictions on the choice of potential  $\eta$ . In the following section, we will focus on a distinguished class of potentials constructed from the golden ratio and Fibonacci-based coordinates.

These *golden-phase gauges* exploit extremal irrationality and optimal equidistribution properties to provide maximally nonresonant probes of arithmetic dynamics. Their introduction requires no modification of the formalism presented here, but only a specific and conceptually motivated choice of  $\eta$ .

## 4 The Golden-Phase Prime-Shift Operator

We now introduce the central object of this work: the golden-phase prime-shift operator. This operator is obtained by specializing the general gauge-covariant prime-shift construction to a distinguished class of arithmetic potentials derived from the golden ratio and related Fibonacci structures.

The resulting operator preserves the combinatorial structure of prime-induced dynamics while introducing a maximally nonresonant phase geometry.

### 4.1 Definition of the Golden-Phase Operator

Let  $\varphi = \frac{1+\sqrt{5}}{2}$  denote the golden ratio. We fix a real-valued potential function

$$\eta_\varphi : \mathbb{N} \longrightarrow \mathbb{R},$$

whose explicit form will be specified later, and define the associated phase difference

$$\theta(n, p) = \eta_\varphi(pn) - \eta_\varphi(n).$$

The golden-phase prime-shift operator  $H_\varphi$  is defined by

$$(H_\varphi \psi)(n) = \sum_{p \in \mathbb{P}} w(p) \left( e^{2\pi i \theta(n, p)} \psi(pn) + \mathbf{1}_{p|n} e^{-2\pi i \theta(n/p, p)} \psi(n/p) \right) + V(n) \psi(n),$$

where  $w(p)$  is a real-valued weight function and  $V(n)$  is a real-valued local potential.

The defining feature of this operator is the conjugate pairing of phase factors on forward and backward transitions. This structure ensures that phase information is introduced without breaking the formal symmetry of the underlying arithmetic dynamics.

### 4.2 Structural Symmetry

The apparent symmetry of the definition is not accidental. For each prime  $p$ , the transition  $n \rightarrow pn$  carries the phase

$$e^{2\pi i (\eta_\varphi(pn) - \eta_\varphi(n))},$$

while the inverse transition  $pn \rightarrow n$  carries the conjugate phase

$$e^{-2\pi i (\eta_\varphi(pn) - \eta_\varphi(n))}.$$

As a consequence, each undirected edge of the arithmetic graph contributes a Hermitian pair to the operator. On the dense domain of finitely supported states, this pairing guarantees formal symmetry independently of the explicit choice of  $\eta_\varphi$ .

This observation captures, in precise terms, the intuitive notion that the operator is “almost automatically symmetric” once phase factors are derived from a single underlying potential.

### 4.3 Interpretation as a Gauge-Twisted Adjacency Operator

The operator  $H_\varphi$  may be viewed as a weighted adjacency operator on the arithmetic graph defined by prime multiplication, equipped with a  $U(1)$  gauge field. The phase  $\theta(n, p)$  assigns a geometric orientation to each prime-induced transition, while preserving the underlying graph structure.

From this perspective,  $H_\varphi$  does not alter the connectivity of the arithmetic state space. Instead, it modifies the interference pattern between arithmetic paths by introducing controlled phase shifts. The golden ratio enters only through the choice of gauge, not through any modification of the combinatorial rules.

### 4.4 Role of the Local Potential

The term  $V(n)$  represents an optional local contribution, diagonal in the arithmetic basis. Typical choices include slowly varying arithmetic functions, such as logarithmic or von Mangoldt-type terms, though no specific form is assumed at this stage.

The presence of  $V(n)$  allows the golden-phase operator to be embedded naturally into broader operator families, including previously studied prime-shift Hamiltonians. Importantly, the gauge-phase structure introduced here is entirely independent of the local potential.

### 4.5 Gauge Equivalence and Physical Content

Formally, the operator  $H_\varphi$  is unitarily equivalent to the untwisted operator  $H_0$  under the gauge transformation generated by  $\eta_\varphi$ . Nevertheless, this equivalence does not trivialize the construction.

In truncated, regularized, or numerically implemented settings, gauge phases affect boundary behavior, interference patterns, and spectral statistics. In this sense, the golden phase acts as a probe rather than a redundancy: it tests the sensitivity of arithmetic dynamics to maximally irrational phase twisting.

### 4.6 Conceptual Significance

The golden-phase prime-shift operator occupies an intermediate position between pure arithmetic and spectral geometry. It does not impose order on prime distribution, nor does it assume hidden periodicity. Instead, it introduces a distinguished gauge against which arithmetic coherence can be meaningfully tested.

If prime-induced dynamics is genuinely phase-blind, golden twisting produces no persistent spectral signature. If, however, subtle coherence exists, the golden phase provides a setting in which such effects may become visible.

In the following section, we examine concrete choices of the potential  $\eta_\varphi$ , ranging from logarithmic gauges to Fibonacci-based representations, and analyze how these choices influence the resulting arithmetic phases.

## 5 Choices of Golden Gauge Functions

In this section we examine concrete choices of the arithmetic potential  $\eta_\varphi$  that generate golden-phase gauges. While the formal construction of the gauge-covariant prime-shift operator places no restrictions on  $\eta$ , the interpretative and spectral consequences depend sensitively on this choice.

Our guiding principle is that the golden ratio should enter only as a source of maximal irrationality and optimal equidistribution, not as an imposed symmetry. Accordingly, we focus on gauges that are structurally simple, arithmetically natural, and spectrally nonresonant.

### 5.1 Logarithmic Golden Gauge

The simplest class of potentials is logarithmic in nature. Let

$$\eta_\varphi(n) = \kappa \log_\varphi n = \frac{\kappa \log n}{\log \varphi},$$

where  $\kappa \in \mathbb{R}$  is a coupling parameter. In this case, the phase difference associated with a prime transition is

$$\theta(n, p) = \eta_\varphi(pn) - \eta_\varphi(n) = \kappa \log_\varphi p.$$

Remarkably, the phase depends only on the prime  $p$  and not on the base point  $n$ . The resulting operator introduces a uniform prime-dependent phase twist across the arithmetic graph.

This gauge may be interpreted as a Dirichlet-type phase deformation. It preserves scale covariance while breaking trivial periodicity through the irrational base  $\varphi$ . Although structurally simple, this choice already distinguishes prime transitions in a spectrally meaningful way.

### 5.2 Quasi-Periodic Scale Gauge

A more refined construction introduces quasi-periodicity through fractional scaling. Define

$$\eta_\varphi(n) = \kappa \{\log_\varphi n\},$$

where  $\{\cdot\}$  denotes the fractional part. The associated phase difference becomes

$$\theta(n, p) = \kappa (\{\log_\varphi(pn)\} - \{\log_\varphi n\}),$$

which depends nontrivially on both  $n$  and  $p$ .

This gauge introduces scale-dependent interference patterns that vary across the arithmetic graph. Prime transitions at different magnitudes experience distinct phases, producing a genuinely quasi-periodic structure in logarithmic scale.

From a spectral perspective, this choice maximally suppresses accidental resonance. Any persistent spectral signature arising under this gauge cannot be attributed to simple scale periodicity.

### 5.3 Zeckendorf–Fibonacci Gauge

A conceptually distinct class of gauges is obtained by expressing integers in their Zeckendorf representation. Every  $n \in \mathbb{N}$  admits a unique expansion

$$n = \sum_k \epsilon_k F_k, \quad \epsilon_k \in \{0, 1\}, \quad \epsilon_k \epsilon_{k+1} = 0,$$

where  $\{F_k\}$  denotes the Fibonacci sequence.

Using this representation, we define

$$\eta_\varphi(n) = \kappa \sum_k \epsilon_k \omega_k,$$

with weights  $\omega_k$  chosen as

$$\omega_k = \{k\varphi\} \quad \text{or} \quad \omega_k = \varphi^{-k}.$$

This gauge assigns phase contributions to Fibonacci digits rather than to magnitude or scale. Prime multiplication induces nonlocal but highly structured changes in the Zeckendorf digit pattern, resulting in complex phase responses.

Within the language of Quansistor Field Mathematics, this construction admits an interpretation as a field defined on Fibonacci-indexed degrees of freedom. Each digit acts as a local excitation, and prime transitions correspond to structured field interactions.

### 5.4 Comparison of Gauge Classes

The gauges introduced above differ in complexity and interpretive emphasis:

- Logarithmic gauges produce prime-dependent but state-independent phases.
- Quasi-periodic scale gauges introduce local interference across scales.
- Zeckendorf-based gauges encode phase information in discrete Fibonacci coordinates.

All three constructions are golden in the sense that they rely on extremal irrationality and avoid rational commensurability. None of them imposes periodic structure on arithmetic dynamics.

### 5.5 Gauge Choice as a Spectral Probe

The choice of  $\eta_\varphi$  should be understood as selecting a probe rather than a model. Different gauges test different aspects of arithmetic coherence. If prime-induced dynamics is genuinely phase-insensitive, all golden gauges should yield spectrally equivalent behavior in appropriate limits.

Conversely, systematic differences between these gauges would signal the presence of nontrivial phase correlations. Such effects, if observed, would not constitute an explanation of prime distribution, but rather evidence of hidden structure accessible only through operator-based interrogation.

In the following section we analyze the formal symmetry and domain properties of the golden-phase prime-shift operator, establishing the mathematical conditions under which the constructions introduced here admit a well-defined spectral interpretation.

## 6 Formal Symmetry and Self-Adjointness

In this section we analyze the symmetry properties of the golden-phase prime-shift operator and clarify the role of domains, regularization, and potential self-adjointness issues. The emphasis is not on establishing maximal generality, but on identifying the structural features that make the operator suitable for spectral analysis and numerical experimentation.

### 6.1 Natural Domains

Let  $c_0(\mathbb{N})$  denote the space of finitely supported functions on  $\mathbb{N}$ . This space is dense in  $\ell^2(\mathbb{N})$  and invariant under all prime-shift operators introduced in this work, provided that the prime sum is regularized.

Throughout this section, all operator identities are understood on  $c_0(\mathbb{N})$ , unless stated otherwise. This choice avoids convergence issues while preserving the algebraic structure of the construction.

### 6.2 Formal Symmetry

Let  $H_\varphi$  denote the golden-phase prime-shift operator defined in the previous section. For  $\psi, \phi \in c_0(\mathbb{N})$ , we compute

$$\langle \psi, H_\varphi \phi \rangle = \sum_{n \in \mathbb{N}} \overline{\psi(n)} (H_\varphi \phi)(n).$$

Each term in the prime sum corresponds to a pair of transitions

$$n \longleftrightarrow pn,$$

with phase factors arranged as complex conjugates. Exchanging the roles of  $\psi$  and  $\phi$  and reindexing the backward transitions yields

$$\langle \psi, H_\varphi \phi \rangle = \langle H_\varphi \psi, \phi \rangle.$$

Thus,  $H_\varphi$  is formally symmetric on  $c_0(\mathbb{N})$ . This property holds independently of the explicit choice of the golden-phase potential  $\eta_\varphi$ , provided that  $\eta_\varphi$  is real-valued.

### 6.3 Role of Phase Pairing

The formal symmetry of  $H_\varphi$  is a direct consequence of constructing phase factors from potential differences. For each prime  $p$ , the forward transition  $n \rightarrow pn$  carries a phase  $e^{2\pi i(\eta_\varphi(pn) - \eta_\varphi(n))}$ , while the backward transition  $pn \rightarrow n$  carries its complex conjugate.

This pairing ensures that no net phase bias is introduced along closed arithmetic paths. In particular, any finite cycle in the arithmetic graph contributes a real-valued amplitude to the quadratic form associated with  $H_\varphi$ .

## 6.4 Quadratic Form Perspective

It is often convenient to consider the quadratic form

$$Q_\varphi(\psi) = \langle \psi, H_\varphi \psi \rangle, \quad \psi \in c_{00}(\mathbb{N}).$$

Formally,  $Q_\varphi$  is real-valued and symmetric. This observation provides a useful starting point for discussing extensions of  $H_\varphi$  beyond  $c_{00}(\mathbb{N})$  and for comparing different regularizations.

From the quadratic form perspective, the golden-phase operator behaves analogously to a discrete magnetic Schrödinger operator on a graph, with arithmetic structure replacing spatial geometry.

## 6.5 Remarks on Self-Adjointness

Questions of essential self-adjointness for operators of the form  $H_\varphi$  are subtle and depend on the choice of weights  $w(p)$ , the local potential  $V(n)$ , and the imposed regularization.

In this work we do not attempt to establish general self-adjointness results. Instead, we adopt the pragmatic viewpoint that all spectral statements are to be understood within regularized settings, such as finite prime cutoffs or bounded state spaces.

Within these settings,  $H_\varphi$  defines a finite-dimensional Hermitian matrix and is therefore self-adjoint. Observed spectral features may then be studied as functions of the cutoff parameters.

## 6.6 Gauge Invariance and Truncation Effects

Although the golden-phase operator is formally gauge-equivalent to the untwisted operator in the infinite setting, truncation and regularization break exact gauge invariance. Boundary effects introduce phase-dependent contributions that persist in finite systems.

This phenomenon is not a defect but a feature. It is precisely through such boundary-sensitive effects that gauge phases act as probes of arithmetic structure. Persistent spectral differences under golden-phase twisting therefore carry meaningful information about prime-induced dynamics.

## 6.7 Summary

The golden-phase prime-shift operator is formally symmetric on a natural dense domain and admits a well-defined quadratic form. While full self-adjointness in the infinite setting remains an open analytical question, all constructions are robust under regularization and truncation.

These properties justify the use of  $H_\varphi$  as a legitimate object of spectral and numerical study. In the next section, we turn to the interpretation of its spectrum and the kinds of arithmetic information such spectra may encode.

## 7 Spectral Interpretation

In this section we discuss the spectral interpretation of the golden-phase prime-shift operator. Rather than focusing on explicit eigenvalue computations, we aim to clarify what types of arithmetic information may be encoded in the spectrum and how phase twisting modifies the spectral response of prime-induced dynamics.

### 7.1 Spectrum as an Invariant of Arithmetic Dynamics

For an operator acting on  $\ell^2(\mathbb{N})$ , the spectrum represents the set of energetically admissible modes of arithmetic motion. In the present context, these modes correspond to coherent superpositions of arithmetic states that are stable under prime-shift transitions.

From an operator-first perspective, the spectrum is the primary invariant of interest. Individual primes appear only indirectly, through their collective contribution to the adjacency structure and the resulting interference patterns. Spectral properties therefore probe global arithmetic organization rather than local combinatorial features.

### 7.2 Interference and Arithmetic Paths

Each eigenstate of the prime-shift operator may be interpreted as a weighted superposition of arithmetic paths generated by repeated prime multiplication and division. The contribution of a given path is determined not only by its length and weights, but also by the accumulated phase along the path.

In the untwisted operator, all such paths interfere constructively. Introducing gauge phases modifies this interference, selectively enhancing or suppressing certain classes of paths. The golden phase, in particular, introduces maximally noncommensurate phase shifts that tend to decorrelate trivial cycles.

### 7.3 Phase Sensitivity of the Spectrum

If prime-induced dynamics is insensitive to phase twisting, the spectra of the untwisted and golden-phase operators should agree in appropriate limits, up to finite-size effects. In this case, gauge phases act as a benign redundancy.

If, however, the spectrum exhibits systematic dependence on the golden gauge, this indicates the presence of arithmetic coherence beyond purely combinatorial structure. Such sensitivity would manifest as shifts in eigenvalue distributions, changes in spectral gaps, or alterations in level statistics.

Importantly, any observed phase sensitivity must persist under variation of regularization parameters in order to be considered meaningful.

## 7.4 Spectral Statistics

Beyond individual eigenvalues, statistical properties of the spectrum provide a powerful diagnostic tool. In particular, level spacing distributions may be compared against random matrix ensembles, such as the Gaussian Orthogonal Ensemble (GOE) or Gaussian Unitary Ensemble (GUE).

The introduction of gauge phases naturally breaks time-reversal symmetry at the level of arithmetic paths, suggesting that transitions between symmetry classes may occur as the phase coupling is varied. The golden phase provides a canonical benchmark against which such transitions may be tested.

## 7.5 Relation to Hilbert–Pólya-Type Intuitions

The Hilbert–Pólya conjecture suggests that the nontrivial zeros of the Riemann zeta function may arise as eigenvalues of a self-adjoint operator. While the golden-phase prime-shift operator is not proposed as such an operator, its spectral behavior may nevertheless inform operator-based intuitions.

In particular, the presence or absence of phase-sensitive spectral structure bears on the question of whether prime-induced dynamics supports nontrivial arithmetic resonances. The golden phase serves here as a stringent test: any resonance surviving maximal irrational twisting must be deeply rooted in arithmetic structure.

## 7.6 Spectral Stability and Robustness

A key requirement for meaningful interpretation is spectral stability. Features attributable to arithmetic structure should remain stable under moderate changes in truncation, weighting, and gauge choice within the golden class.

Conversely, spectral features that vary erratically under such changes are best interpreted as finite-size artifacts. The distinction between these regimes is central to the numerical protocols discussed in the following section.

## 7.7 Interpretive Summary

The spectrum of the golden-phase prime-shift operator encodes information about the coherence and interference of prime-induced arithmetic dynamics. Gauge phases do not explain prime distribution, but they test its sensitivity to controlled phase perturbations.

Whether the outcome is complete phase blindness or the emergence of persistent spectral structure, the result carries conceptual significance. In both cases, the operator framework provides a precise language in which arithmetic dynamics may be interrogated spectrally rather than combinatorially.

In the next section we describe explicit numerical protocols designed to probe these spectral effects in a falsifiable and reproducible manner.

## 8 Numerical Protocols and Falsifiability

This section outlines concrete numerical protocols for investigating the spectral properties of the golden-phase prime-shift operator. The emphasis is on reproducibility, falsifiability, and the separation of genuine arithmetic effects from finite-size artifacts.

All numerical statements are to be understood within regularized settings. The goal is not to approximate an infinite operator, but to study the stability and response of finite truncations under controlled phase perturbations.

### 8.1 Finite-Dimensional Truncations

Numerical implementations are based on truncating both the state space and the prime sum. Specifically, we restrict

$$n \leq N, \quad p \leq P,$$

and consider the resulting operator acting on  $\mathbb{C}^N$ .

Within this setting, the golden-phase prime-shift operator defines a finite Hermitian matrix. The dependence of spectral features on the cutoff parameters  $N$  and  $P$  provides a primary diagnostic for distinguishing structural effects from truncation artifacts.

### 8.2 Regularization of Prime Weights

In practice, the choice of weight function  $w(p)$  significantly influences numerical stability. Common choices include

$$w(p) = \frac{1}{p}, \quad w(p) = \frac{e^{-p/P}}{p}, \quad w(p) = \frac{1}{p^\alpha}, \quad \alpha > 1.$$

Smooth damping of large primes improves convergence and allows systematic study of the dependence on prime scale. All gauge constructions introduced earlier are compatible with such regularizations.

### 8.3 Spectral Observables

The primary numerical observables include:

- eigenvalue distributions and spectral density,
- gaps and extremal eigenvalues,
- level spacing distributions after unfolding,
- sensitivity of eigenvalues to variations in the gauge parameter  $\kappa$ .

Comparisons are performed between the untwisted operator  $H_0$ , the golden-phase operator  $H_\varphi$ , and operators with randomized phases serving as control ensembles.

## 8.4 Gauge Sensitivity Tests

A central diagnostic is the response of the spectrum to gauge twisting. For a fixed truncation  $(N, P)$ , one varies the coupling parameter  $\kappa$  and observes the evolution of spectral features.

Three regimes may be distinguished:

- *Phase-blind regime*: spectral observables remain invariant up to statistical noise.
- *Perturbative regime*: small but systematic shifts occur, scaling with  $\kappa$ .
- *Coherent regime*: nontrivial spectral restructuring persists under variation of  $\kappa$  and truncation parameters.

Only the latter two regimes indicate genuine phase sensitivity.

## 8.5 Comparison with Random Phases

To assess the specificity of golden-phase effects, it is essential to compare results with randomized gauge phases. In such control experiments, the phase factors  $U(n, p)$  are replaced by independent random elements of  $U(1)$ .

If golden-phase effects are indistinguishable from random-phase behavior, this suggests the absence of arithmetic coherence detectable by phase twisting. By contrast, systematic differences indicate structure specific to the golden gauge.

## 8.6 Level Statistics and Symmetry Classes

After unfolding the spectrum, level spacing distributions may be compared to random matrix ensembles. The untwisted operator typically exhibits statistics associated with time-reversal invariant systems.

Introducing gauge phases may induce a crossover toward unitary statistics, reflecting the breaking of effective time-reversal symmetry in arithmetic path interference. The golden phase provides a canonical, nonrandom benchmark for studying such transitions.

## 8.7 Criteria for Falsification

The framework proposed in this paper admits clear falsification criteria:

- absence of stable spectral differences between golden-phase and random-phase ensembles,
- disappearance of phase sensitivity as truncation parameters increase,
- lack of reproducibility across independent implementations.

Failure to observe any persistent golden-phase effects under these tests would constitute strong evidence that prime-induced dynamics is phase-insensitive at the operator level considered here.

## 8.8 Reproducibility and Open Data

All numerical experiments should be conducted using publicly available code and documented parameter choices. The simplicity of the operator definition allows implementation in a variety of environments, including Python, Julia, and JavaScript-based numerical frameworks.

Reproducibility is essential: the purpose of numerical exploration in this context is not confirmation bias, but controlled interrogation of arithmetic dynamics through spectral probes.

## 8.9 Summary

Numerical analysis provides the bridge between formal operator construction and empirical assessment. The protocols outlined above allow the golden-phase prime-shift operator to be tested as a genuine probe of arithmetic coherence.

In the following section, we place these constructions within the broader framework of Quansistor Field Mathematics and discuss their interpretation as field-theoretic refinements of operator-based arithmetic models.

## 9 Relation to QFM and Quansistor Fields

In this section we situate the golden-phase prime-shift operator within the broader conceptual framework of Quansistor Field Mathematics (QFM). The purpose is not to introduce new formal machinery, but to clarify how the constructions developed in this paper fit naturally into an operator-first, field-oriented view of arithmetic dynamics.

### 9.1 Operator-First Arithmetic

A central principle of QFM is that arithmetic structure should be approached through operators rather than explicit formulas. In this view, numbers label states, while operators encode permissible transitions and interactions between them.

The prime-shift operators studied in this paper exemplify this philosophy. Multiplication by primes is treated as a fundamental dynamical process, and arithmetic information is extracted from the spectral response of the resulting operator rather than from combinatorial enumeration.

### 9.2 Quansistors as Arithmetic Degrees of Freedom

Within QFM, a quansistor is interpreted as an elementary computational or dynamical unit whose behavior is defined by its interaction with surrounding structure. In the present setting, arithmetic states indexed by natural numbers may be viewed as composite configurations of such units.

Prime-shift transitions then represent structured interactions between quansistors, mediated by arithmetic adjacency. The golden-phase gauge introduced in this work adds a local phase degree of freedom to these interactions, enriching the dynamical landscape without altering the underlying connectivity.

### 9.3 Gauge Fields as Arithmetic Structure

From the QFM perspective, gauge fields are not auxiliary constructs but intrinsic features of structured systems. The arithmetic gauge introduced here assigns phases to prime-induced transitions, effectively defining a field over the arithmetic graph.

Crucially, this field is not imposed externally. It arises from a potential function on arithmetic states and respects the multiplicative structure of  $\mathbb{N}$ . In this sense, the golden-phase construction may be interpreted as an arithmetic gauge field in the strictest sense.

### 9.4 Fibonacci Coordinates and Field Representations

The Zeckendorf-based gauges discussed earlier provide a particularly direct connection to field-theoretic intuition. By expressing integers in Fibonacci coordinates, arithmetic states acquire a multi-component representation with local constraints.

Within QFM, such representations naturally suggest a field interpretation, with each Fibonacci digit corresponding to a local degree of freedom. Prime-induced transitions act nonlocally on these degrees of freedom, while the golden phase encodes their interference.

## 9.5 Compatibility with QVM Architectures

The golden-phase prime-shift operator is compatible with virtual and distributed computational architectures envisioned within the Quantum Virtual Machine (QVM) program. Finite truncations of the operator correspond to finite computational instances, while gauge parameters serve as tunable controls.

Because the construction is operator-based and discretely defined, it may be implemented and explored across heterogeneous computational substrates without loss of conceptual coherence.

## 9.6 Conceptual Contribution to QFM

The primary contribution of this work to QFM is the introduction of a principled, arithmetically motivated gauge degree of freedom. Rather than proposing a new axiomatic system, the golden-phase operator refines existing operator-based approaches by introducing a controlled probe of arithmetic coherence.

This refinement strengthens the interpretive reach of QFM by providing a concrete example in which gauge structure, arithmetic dynamics, and spectral analysis interact in a falsifiable and reproducible manner.

## 9.7 Position Within the QFM Program

The constructions presented here should be regarded as a modular component of the broader QFM program. They neither depend on nor assume the full machinery of QFM, but they align naturally with its guiding principles.

In particular, the golden-phase prime-shift operator demonstrates how operator geometry and arithmetic structure may be intertwined without resorting to speculative assumptions. It provides a clear example of how QFM-style reasoning can yield concrete, testable mathematical objects.

In the final section, we outline open problems and future directions suggested by this work, including possible extensions beyond prime-shift dynamics and connections to other operator-based arithmetic frameworks.

## 10 Outlook and Open Problems

This paper has introduced a class of gauge-covariant prime-shift operators in which arithmetic dynamics is probed through maximally irrational phase twisting. The central construction, the golden-phase prime-shift operator, preserves the combinatorial structure of prime multiplication while enriching it with a geometrically motivated  $U(1)$  gauge degree of freedom.

The framework developed here does not seek to explain the distribution of prime numbers. Instead, it establishes a controlled and falsifiable setting in which arithmetic coherence may be tested spectrally. Whether the outcome of such tests is positive or negative, the operator-based perspective provides information that is inaccessible through purely combinatorial approaches.

### 10.1 Summary of Contributions

The main contributions of this work may be summarized as follows:

- the formulation of prime-shift operators as adjacency operators on an arithmetic graph,
- the introduction of gauge phases derived from arithmetic potentials, ensuring formal symmetry,
- the identification of the golden ratio as a canonical source of maximally nonresonant gauge twisting,
- the proposal of concrete numerical protocols enabling falsifiable tests of phase sensitivity in prime-induced dynamics,
- the integration of these constructions into an operator-first arithmetic framework compatible with Quansistor Field Mathematics.

Each of these elements is modular and may be studied independently or in combination with other operator-based approaches.

### 10.2 Interpretive Caution

It is important to emphasize the limitations of the present work. The appearance of spectral sensitivity under golden-phase twisting, should it occur, does not constitute an explanation of prime distribution, nor does it imply hidden periodicity in the primes.

Conversely, the absence of any such sensitivity would not invalidate operator approaches to arithmetic, but would instead provide evidence that prime-induced dynamics is robustly phase-blind under optimal irrational perturbations.

Both outcomes are informative within an operator-theoretic paradigm.

### 10.3 Analytical Open Problems

Several analytical questions remain open:

- the characterization of self-adjoint extensions of infinite golden-phase prime-shift operators,
- the dependence of spectral measures on the choice of arithmetic potential beyond finite truncations,
- the possible emergence of trace-class or resolvent regularity under specific gauge choices.

Addressing these questions would require tools from functional analysis, spectral theory on graphs, and analytic number theory.

### 10.4 Extensions Beyond Prime Shifts

The construction presented here may be extended in several directions. One natural avenue is the replacement of prime shifts by more general multiplicative or arithmetic transitions, including those associated with Dirichlet characters or  $L$ -functions.

Another possibility is the introduction of nonabelian gauge groups, allowing richer interference structures and potentially new classes of spectral behavior.

### 10.5 Computational and Experimental Directions

From a computational perspective, the simplicity of the operator definition invites large-scale numerical exploration. Distributed and virtual computational architectures may be employed to explore spectral behavior across wide ranges of parameters.

Systematic comparison between golden-phase gauges and other irrational gauges may further clarify the role of extremal irrationality in arithmetic probing.

### 10.6 Conceptual Outlook

At a conceptual level, this work supports the view that arithmetic structure may be more effectively interrogated through dynamical and spectral responses than through static enumeration. Gauge phases, when chosen carefully, provide a means of asking precise questions without imposing artificial structure.

The golden-phase prime-shift operator represents one such question. Whether its answer reveals hidden coherence or confirms robust randomness, the result contributes to a deeper understanding of arithmetic dynamics viewed through the lens of operators.

### 10.7 Closing Remarks

The operator-based study of arithmetic remains a young and evolving field. By introducing a maximally nonresonant gauge probe grounded in the golden ratio, this work adds a concrete and testable tool to the growing repertoire of arithmetical operator constructions.

It is hoped that the framework presented here will encourage further exploration, critical testing, and refinement, contributing incrementally to a broader spectral understanding of arithmetic phenomena.

## A Technical Remarks

This appendix collects technical remarks, clarifications, and auxiliary constructions that support the main text but are not essential for its logical flow. No new conceptual assumptions are introduced here.

## Appendix A: Technical Remarks and Supplementary Material

### A.1 A.1 Arithmetic Graph Structure

The arithmetic graph underlying the prime-shift operators has vertex set  $\mathbb{N}$  and directed edges

$$n \longrightarrow pn \quad \text{for each } p \in \mathbb{P}.$$

Each vertex has infinite out-degree and finite in-degree. This graph is locally finite only after imposing a prime cutoff.

The adjacency structure is multiplicative rather than additive, distinguishing it from conventional lattices and justifying the use of operator-based rather than geometric intuition.

### A.2 A.2 Regularized Operators

All operators considered in this work admit natural regularized versions

$$(H_\varphi^{(N,P)}\psi)(n) = \sum_{\substack{p \leq P \\ pn \leq N}} w(p) \left( e^{2\pi i \theta(n,p)} \psi(pn) + \mathbf{1}_{p|n} e^{-2\pi i \theta(n/p,p)} \psi(n/p) \right) + V(n)\psi(n),$$

acting on  $\mathbb{C}^N$ .

Such truncations define finite Hermitian matrices and are suitable for numerical diagonalization. All spectral statements in the main text refer to limits or stability properties of these regularized operators.

### A.3 A.3 Gauge Equivalence and Boundary Effects

Formally, the gauge-transformed operator

$$H_\eta = G_\eta^{-1} H_0 G_\eta$$

is unitarily equivalent to the untwisted operator  $H_0$ . However, this equivalence is exact only in the infinite, untruncated setting.

Finite truncations break gauge invariance through boundary effects. These effects are not artifacts but carry the phase sensitivity exploited in numerical probes. Spectral dependence on gauge parameters therefore reflects arithmetic structure interacting with truncation geometry.

## A.4 A.4 Quadratic Form Estimates

For  $\psi \in c_00(\mathbb{N})$ , the quadratic form associated with the golden-phase operator satisfies

$$|Q_\varphi(\psi)| \leq 2 \sum_p |w(p)| \sum_n |\psi(n)|^2 + \sum_n |V(n)| |\psi(n)|^2,$$

formally suggesting boundedness under sufficiently strong damping of  $w(p)$ .

While this estimate is heuristic, it motivates the use of exponentially damped weights in numerical experiments.

## A.5 A.5 Alternative Gauge Choices

Although the main text focuses on golden-phase gauges, the construction extends to arbitrary irrational gauges. In particular, one may replace  $\varphi$  by other quadratic irrationals or by generic badly approximable numbers.

The golden ratio is distinguished not by uniqueness but by extremality: it provides a canonical benchmark against which other irrational gauges may be compared.

## A.6 A.6 Pseudocode Sketch

A minimal numerical implementation proceeds as follows:

1. Fix truncation parameters  $(N, P)$  and gauge choice  $\eta_\varphi$ .
2. Construct the sparse matrix representing  $H_\varphi^{(N,P)}$ .
3. Diagonalize the matrix using standard Hermitian solvers.
4. Compute spectral observables and compare against control ensembles.

Because the operator is sparse and structured, scalable implementations are feasible even for moderately large  $N$ .

## A.7 A.7 Relation to Other Operator Constructions

Prime-shift operators with and without gauge phases appear implicitly in various contexts in analytic number theory and mathematical physics. The distinguishing feature of the present construction is the systematic use of gauge geometry as a diagnostic tool rather than as a modeling assumption.

This viewpoint aligns naturally with operator-based arithmetic programs while remaining agnostic about deeper conjectural interpretations.

## A.8 A.8 Closing Remark

The technical considerations collected here are intended to clarify, not to complicate, the constructions presented in the main text. All essential ideas are contained in the

core chapters; the appendix merely provides supporting detail for interested readers and practitioners.