

On Frequency Tuning of Inverted F Antenna

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ABSTRACT

An antenna is a crucial element in the successful design of any radio system. Consequently, the operational characteristic of the system are designed around the directional properties of the antennas. By the wide spread of mobile phones during the last decade, new antenna designs with wider frequency bands have been introduced; single-band devices have developed into multiband and multimode terminals. Although the average terminal size has decreased drastically, the internal antennas have been designed into standard solutions, to meet the new design requirements, taking into consideration the strict limitations set on the energy absorbed by the users of mobile devices. In this paper, a model for frequency tunable ferrite-based Inverted F Antenna is proposed by altering the permeability of the ferrite material.

KEYWORDS

Antennas, Wireless, Communications, frequency tuning.

I. INTRODUCTION

Internal antennas are widely used in mobile wireless communication devices like smart phones, which have been a major incentive of internal antenna research and design during the last decade. As a result, different techniques that can be possibly used for changing the operating frequency of the antenna are introduced. They are mainly categorized as: mechanical techniques, using tunable electronic circuits, and alteration of the properties of the antenna material. In this paper, we will focus on the alteration of the electrical properties (permittivity, permeability, and conductivity) of the antenna material.

The mechanical techniques, which include changing the shape or dimensions of the

antenna, have a disadvantage of lack of reliability to time delay, and short life time due to mechanical movements. On the other hand the alteration of the antenna material does not require a mechanical movement of the antenna as suitable material is incorporated into the antenna such that changing its electrical properties lead to change in the antenna operating frequency [1].

II. ALTERATION OF FERRITES PERMEABILITY

Ferrites are magnetic materials made from a mixture of metal oxides. Due to their electrical characteristics; they can be used in many applications like antennas, electronic and communication devices, inductors, transformers and in many other applications. Ferrites have high permeability at low frequencies, and a dielectric constant of 10 or more, besides its high resistivity. The behavior of ferrites at high frequencies is well studied in many research publications e.g. [2] & [3]. In this section a brief overview of ferrites characteristics is presented.

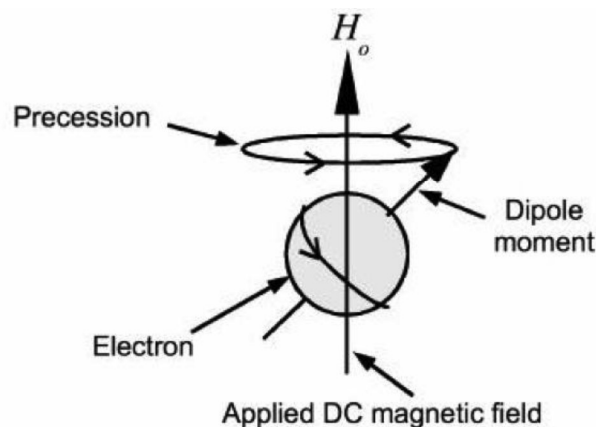


Fig. 1. Spin of electron about the constant field H_0 .

The magnetic dipole moment is primarily responsible for the magnetic properties of the ferrites, where it is produced by the spin of electrons in the material. Fig. 1 depicts this process where a constant magnetic field H_0 , is applied to the ferrite, and the electrons start to spin about the axis of the magnetic field at a frequency of $\omega_0 = 2\pi\gamma H_0$ (called Larmor frequency), where $\gamma = 2.8\text{MHz/Oersted}$. When the magnetic field is intense, the electron spins causing the dipole moments are aligned together and form one large dipole moment. This case is called magnetization saturation M_s , which can be used as an indicator of the density of magnetic flux that is needed to align all the electron spins of the ferrite.

In conjunction with the constant magnetic field H_0 , a tiny magnitude alternating magnetic field at a radio frequency ω , can be used to produce a forced spin of magnetic moment. In this case, the Polder permeability tensor [4] is produced such as

$$[\mu] = \begin{bmatrix} \mu & j\delta & 0 \\ -j\delta & \mu & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (1)$$

where μ and δ are given as:-

$$\mu = \mu' - j\mu'' = 1 + \frac{\omega_m(\omega_0 + j\omega_L)}{(\omega_0 + j\omega_L)^2 - \omega^2} \quad (2)$$

$$\delta = \frac{-\omega_m\omega_0}{(\omega_0 + j\omega_L)^2 - \omega^2} \quad (3)$$

where $\omega_m = 8\pi^2\gamma M_s$, and $\omega_L = \pi\gamma\Delta H$. Ferrites are commercially obtainable including nickel, lithium, and magnesium ferrites, with different shapes including bars and sheets. The magnetization saturations of these ferrites have ranges from $300 \text{ Gauss} \leq 4\pi M_s \leq 5000 \text{ Gauss}$, and a resonance width from $10 \text{ Oersted} \leq \Delta H \leq 900 \text{ Oersted}$. As we are dealing with radio frequencies, the frequencies f_0, f_m , and f_L are measured in MHz, where $f_0 = \gamma H_0$, $f_m = 8\gamma\pi M_s$, and $f_L = \pi\gamma\Delta H/2$.

If the polarization of the applied magnetic field B and the density of the magnetic flux H is circular, where (+) denotes clockwise and (-) denotes anticlockwise polarizations, then B can be expressed in terms of H as:-

$$\begin{bmatrix} B_+ \\ B_- \\ B_z \end{bmatrix} = \mu_0 \begin{bmatrix} \mu_+ & 0 & 0 \\ 0 & \mu_- & 0 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} H_+ \\ H_- \\ H_z \end{bmatrix} \quad (4)$$

where $\mu_+ = (\mu + \delta)$, and $\mu_- = (\mu - \delta)$. Fig. 2 depicts a characteristic performance of μ_+ and μ_- versus the frequency of the RF magnetic field. It can be seen that the clockwise waves resonates at f_0 where μ'' exhibits a significant peak which can be observed at that frequency. This indicates a significant reduction in the magnitude of the clockwise waves, in contradiction with case of anticlockwise waves, where there is no attenuation. Moreover, the waves do not propagate if the frequency of the RF magnetic field lies in the range $f_0 < f < f_0 + f_m$ as the real part (μ') of the ferrites permeability is negative in this case.

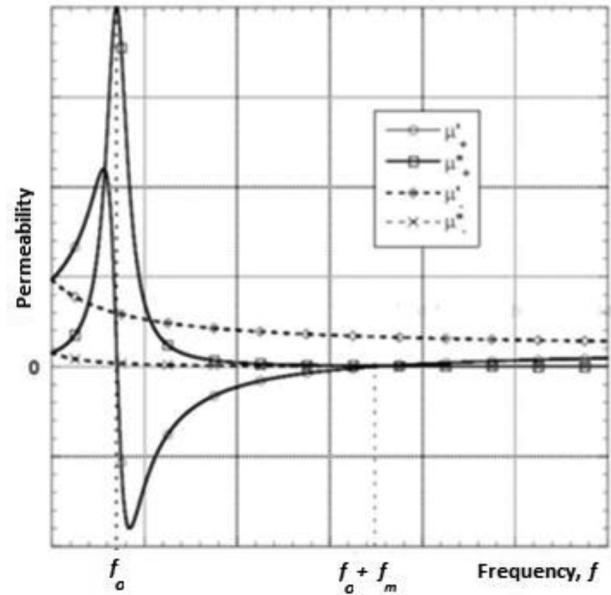


Fig. 2 Ferrites permeability as a function of the frequency of RF magnetic field.

This change in the magnetic field intensity when a radio frequency magnetic field propagates through the ferrites can be employed in designing frequency tunable antennas. In order to investigate this case, the effective permeability μ_{eff} can be used to estimate the range of frequency tuning of the ferrite material. The magnitude of μ_{eff} depends on the direction of the static field H_0 and on the radio frequency magnetic field H . Here, we have three scenarios: parallel bias with $\mu_{eff}=1$,

transverse bias with $\mu_{\text{eff}} = (\mu^2 - \delta^2)/\mu$, and longitudinal bias with $\mu_{\text{eff}} = \mu \pm \delta$. In the transverse bias case μ_{eff} will be negative, and the waves will not propagate when the frequency of the RF magnetic field lies in the range $\sqrt{f_0(f_0 + f_m)} < f < (f_0 + f_m)$.

In the saturation phase of the ferrites, the following approximations hold $\mu \cong \mu_0$ and $\cong 4\pi\mu_0\gamma M_r/f$, where M_r indicates the remaining magnetization in the ferrite after the static magnetic field H_0 is turned off [33].

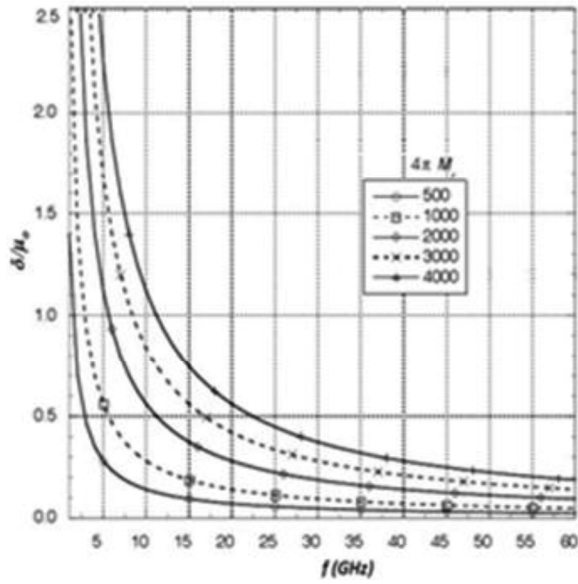


Figure 3. Maximum range of frequency tuning of a ferrite antenna vs. frequency of RF magnetic field [1].

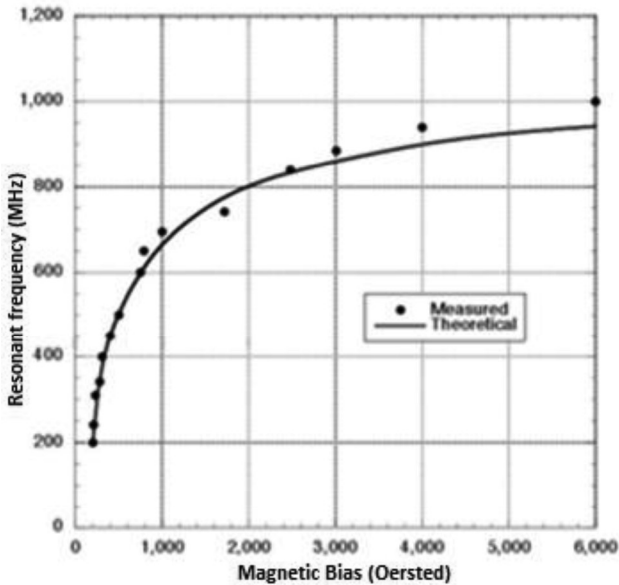


Figure 4. Resonant frequency vs. the static magnetic field H_0 (Magnetic bias), for a rectangular ferrite micro-strip batch antenna [5].

Figure 3 illustrates δ versus the frequency of the RF magnetic field. As the maximum range of frequency tuning of the ferrite antenna (including micro-strip and dielectric resonator antennas) depends on δ , it can be seen from Figure 3 that the maximum frequency tuning range declines by increasing the frequency of the RF magnetic field. The static field H_0 can be produced either by a permanent magnet or by an electromagnet with dc current. Therefore, the resonant frequency of a ferrite micro-strip batch antenna can be controlled by varying the current in the electromagnet that produces the static magnetic field H_0 , as shown in Fig. 4[5].

III. PROPOSED MODEL FOR FERRITE-BASED INVERTED F ANTENNA

In this section a model for a ferrite-based Inverted F antenna is proposed. Figure 5 shows a typical 2.4GHz micro-strip Inverted F antenna, and Fig. 6 shows a typical radiation pattern of the antenna with vertical polarization. The reflection at the feed-point of the antenna in dB is depicted in Fig.7, where it can be seen clearly that the resonant frequency of the antenna occurs at around 2.41GHz.



Fig. 5 Typical 2.4 GHz Inverted F antenna.

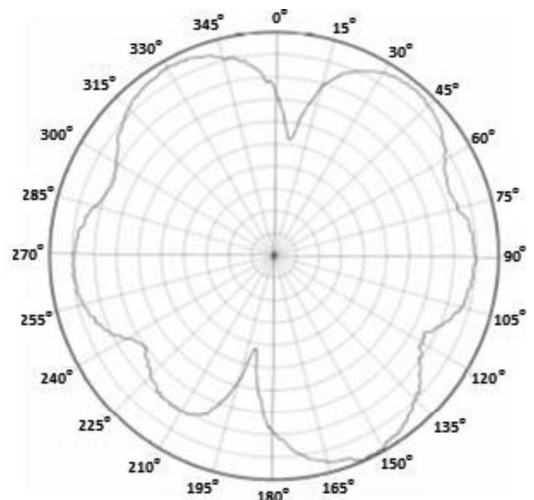


Fig. 6 Typical radiation pattern of 2.4 GHz Inverted F antenna with vertical polarization [6].

Based on the explanations presented in section 1 and 2, the resonant frequency is expected to follow the curve shown in Fig. 8. From this figure it can be seen that the resonant frequency of the inverted F antenna can be significantly shifted from 2.41GHz to 3.15GHz by increasing the bias frequency from 0.2KOersted to 6.0KOersted, with a shift percentage of around 31%. This leads to a significant enhancement in the performance of the antenna and hence the communication systems over wide range of frequencies.

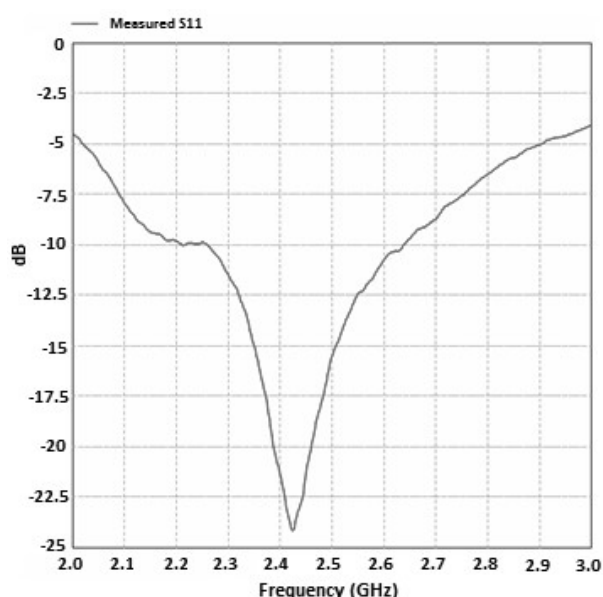


Fig. 7 Reflection at the feed-point of the 2.4GHz Inverted F antenna [6].

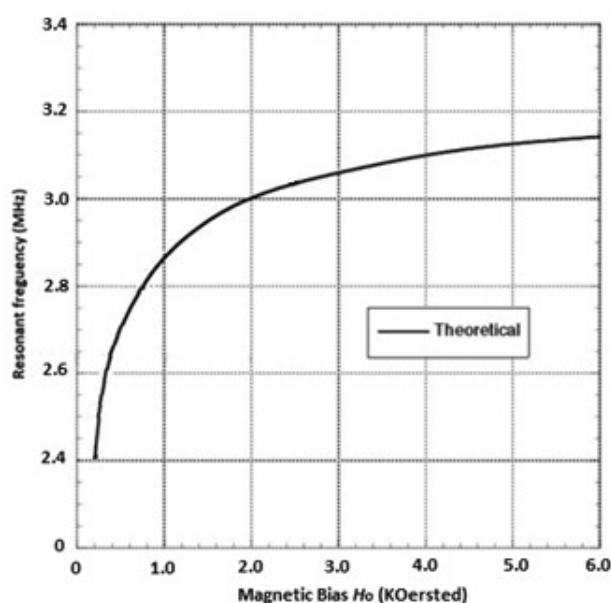


Fig. 8 Resonant frequency vs. the static magnetic field H_0 (Magnetic bias), for 2.4GHz Inverted F antenna.

IV. CONCLUSIONS

In this paper a model for frequency tunable ferrite-based Inverted F antenna is proposed. The resonant frequency of the antenna can be significantly shifted by altering the permeability of the ferrite material in the antenna, and without changing the physical dimensions of the antenna. The resonant frequency shift gives flexibility to operate the antenna over wide range of frequencies which leads to an overall enhancement of the communication systems.

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