

Asymptotics Problem 8.6: Complete Pedagogical Solution

Boundary Layers When $p(x)$ Vanishes at Both Boundaries

Problem 1. For the o.d.e. $\varepsilon y'' + \sin(x)y' + \sin(2x)y = 0$, $0 \leq x \leq \pi$, with $y(0) = \pi$, $y(\pi) = 0$, consider asymptotic expansions for $\varepsilon \rightarrow 0$ with a boundary layer at either end of the interval, and find one or more asymptotic expansions for the solution $y(x)$ to leading order.

Solution: Step-by-Step Atomic Breakdown

Step 1: Identify the Problem Structure and Classification

Strategy: We have a singularly perturbed second-order linear ODE of the general form (Lecture Notes §6.2, Eq. (340)):

$$\varepsilon y'' + p(x)y' + q(x)y = 0, \quad y(0) = \alpha, \quad y(1) = \beta,$$

where in our case:

$$\begin{aligned} p(x) &= \sin(x) \\ q(x) &= \sin(2x) = 2 \sin(x) \cos(x) \\ \text{Domain: } &[0, \pi] \\ \alpha &= \pi, \quad \beta = 0 \end{aligned}$$

The key question is: **Where are the boundary layers located?**

Key Concept: From Lecture Notes §6.2.1, the standard theory tells us:

- If $p(x) > 0$ throughout $[a, b]$, the boundary layer is at the **left** endpoint ($x = a$)
- If $p(x) < 0$ throughout $[a, b]$, the boundary layer is at the **right** endpoint ($x = b$)
- If $p(x_0) = 0$ for some x_0 , special analysis is required

Step 2: Analyze Where $p(x) = \sin(x)$ Vanishes

What we observe:

$$p(x) = \sin(x) = 0 \quad \text{at} \quad x = 0 \quad \text{and} \quad x = \pi$$

Warning: The coefficient $p(x) = \sin(x)$ vanishes at **both** boundaries of the domain $[0, \pi]$! This means the standard theory (which assumes $p(x) \neq 0$ on $[0, 1]$, see Lecture Notes Eq. (341)) does not directly apply. We cannot immediately determine the boundary layer location from the sign of $p(x)$ at the boundaries.

Justification: Why is this problematic? Recall from Lecture Notes §6.2.1 that boundary layer location depends on the sign of $p(x)$:

- When $p(x_0) > 0$: The inner solution has form $Y_0(X) = A + Be^{-p_0 X}$, which decays as $X \rightarrow +\infty$ only if $p_0 > 0$
- When $p(x_0) < 0$: The exponential $e^{-p_0 X}$ grows as $X \rightarrow +\infty$, preventing matching

When $p(x_0) = 0$, the exponential character changes completely, and we need a different dominant balance analysis.

Strategy: Since we cannot determine the boundary layer location *a priori*, we must:

1. Compute the outer solution (valid away from any boundary layers)
2. Try an inner solution at $x = 0$ and check if matching is possible
3. Try an inner solution at $x = \pi$ and check if matching is possible
4. Based on which matchings succeed, determine the actual structure

This is the systematic workflow from Lecture Notes §6.2.3.

Step 3: Compute the Outer Solution

What we do: Set $\varepsilon = 0$ in the ODE to find the leading-order outer solution.

Technique: The outer expansion is $y = y_0 + \varepsilon y_1 + \dots$. At leading order (ε^0):

$$\sin(x)y'_0 + \sin(2x)y_0 = 0$$

This is a first-order linear ODE (we've lost the highest derivative!).

Step 3a: Solving the Reduced Equation

Simplify using $\sin(2x) = 2\sin(x)\cos(x)$:

$$\sin(x)y'_0 + 2\sin(x)\cos(x)y_0 = 0$$

For $x \neq 0, \pi$ (where $\sin(x) \neq 0$), divide by $\sin(x)$:

$$y'_0 + 2\cos(x)y_0 = 0$$

Technique: This is separable:

$$\frac{dy_0}{y_0} = -2\cos(x)dx$$

Integrate both sides:

$$\ln|y_0| = -2\sin(x) + C$$

Therefore:

$$y_0(x) = a e^{-2\sin(x)}$$

where a is an arbitrary constant.

Justification: Why don't we apply boundary conditions yet? We have only **one** integration constant a , but **two** boundary conditions ($y(0) = \pi$ and $y(\pi) = 0$). The outer solution alone cannot satisfy both boundary conditions — this is the hallmark of a singular perturbation problem. We must determine which boundary condition the outer solution should satisfy based on where the boundary layer is located.

Step 3b: Evaluating the Outer Solution at Boundaries

At $x = 0$:

$$y_0(0) = a e^{-2\sin(0)} = a e^0 = a$$

At $x = \pi$:

$$y_0(\pi) = a e^{-2\sin(\pi)} = a e^0 = a$$

Key Concept: Interestingly, $y_0(0) = y_0(\pi) = a$ because $\sin(0) = \sin(\pi) = 0$. This means the outer solution approaches the **same** constant a at both endpoints. This will be crucial for the matching analysis.

Step 4: Try Inner Solution at $x = 0$

Strategy: Assume there is a boundary layer at $x = 0$. Define the inner variable $X = x/\delta$ where $\delta(\varepsilon) \rightarrow 0$ as $\varepsilon \rightarrow 0$. The boundary layer width δ will be determined by dominant balance.

Step 4a: Transform the ODE

Set $x = \delta X$ and $Y(X) = y(x)$. Then:

$$\begin{aligned}\frac{dy}{dx} &= \frac{1}{\delta} \frac{dY}{dX} = \frac{Y'}{\delta} \\ \frac{d^2y}{dx^2} &= \frac{1}{\delta^2} \frac{d^2Y}{dX^2} = \frac{Y''}{\delta^2}\end{aligned}$$

Substitute into the ODE:

$$\varepsilon \cdot \frac{Y''}{\delta^2} + \sin(\delta X) \cdot \frac{Y'}{\delta} + \sin(2\delta X) \cdot Y = 0$$

Step 4b: Taylor Expand the Coefficients

For small δX :

$$\begin{aligned}\sin(\delta X) &= \delta X - \frac{(\delta X)^3}{6} + \dots \approx \delta X \\ \sin(2\delta X) &= 2\delta X - \frac{(2\delta X)^3}{6} + \dots \approx 2\delta X\end{aligned}$$

The ODE becomes:

$$\begin{aligned}\frac{\varepsilon}{\delta^2} Y'' + \frac{\delta X}{\delta} Y' + 2\delta X \cdot Y &= 0 \\ \frac{\varepsilon}{\delta^2} Y'' + XY' + 2\delta XY &= 0\end{aligned}$$

Justification: Why can we use Taylor expansion? Near $x = 0$, the inner variable $X = O(1)$ corresponds to $x = \delta X = O(\delta)$, which is small. Therefore $\sin(\delta X) \approx \delta X$ is an excellent approximation within the boundary layer.

Step 4c: Dominant Balance Analysis

Technique: We need to determine $\delta(\varepsilon)$ such that the two leading terms balance. Consider the relative sizes:

- First term: $\frac{\varepsilon}{\delta^2} Y''$
- Second term: XY' (coefficient is $O(1)$)
- Third term: $2\delta XY$ (coefficient is $O(\delta)$, subdominant)

For a non-trivial boundary layer, the first two terms must balance:

$$\frac{\varepsilon}{\delta^2} = O(1) \implies \delta^2 = \varepsilon \implies \boxed{\delta = \sqrt{\varepsilon}}$$

Justification: Why this balance? If $\delta \ll \sqrt{\varepsilon}$, then $\varepsilon/\delta^2 \gg 1$ and the Y'' term would dominate alone (no balance). If $\delta \gg \sqrt{\varepsilon}$, then $\varepsilon/\delta^2 \ll 1$ and the Y'' term would be negligible (we'd get the outer equation). Only $\delta = \sqrt{\varepsilon}$ gives a proper balance.

This is consistent with Lecture Notes §6.2.2, Eq. (356): when $p(x_0) = 0$ but $p'(x_0) \neq 0$, the boundary layer width is $\delta = \sqrt{\varepsilon}$, not $\delta = \varepsilon$.

Step 4d: Leading-Order Inner Equation at $x = 0$

With $\delta = \sqrt{\varepsilon}$, the inner equation becomes:

$$Y'' + XY' + 2\sqrt{\varepsilon}XY = 0$$

At leading order (neglecting the $O(\sqrt{\varepsilon})$ term):

$$Y_0'' + XY_0' = 0$$

This is a second-order ODE with non-constant coefficients.

Step 4e: Solve the Leading-Order Inner Equation

Technique: The equation $Y_0'' + XY_0' = 0$ can be reduced in order. Let $W = Y_0'$:

$$W' + XW = 0 \implies \frac{dW}{W} = -X dX$$

Integrate:

$$\ln|W| = -\frac{X^2}{2} + C_1 \implies W = Y_0' = A \exp\left(-\frac{X^2}{2}\right)$$

Integrate again to find Y_0 :

$$Y_0(X) = A \int_0^X \exp\left(-\frac{s^2}{2}\right) ds + B$$

Justification: Why integrate from 0 to X ? This choice makes the integral vanish at $X = 0$, which simplifies applying the boundary condition. The constant of integration is absorbed into B .

Step 4f: Apply Boundary Condition at $x = 0$

At $x = 0$, we have $X = 0$ and $y(0) = \pi$. Therefore:

$$Y_0(0) = A \int_0^0 \exp\left(-\frac{s^2}{2}\right) ds + B = 0 + B = B$$

The boundary condition $Y_0(0) = \pi$ gives:

$$B = \pi$$

So the inner solution at $x = 0$ is:

$$Y_{0,a}(X) = A \int_0^X \exp\left(-\frac{s^2}{2}\right) ds + \pi$$

Step 4g: Check Matching at $x = 0$

Technique: For matching, we need to check the behavior as $X \rightarrow +\infty$ (moving from the inner region toward the outer region). The Gaussian integral:

$$\int_0^\infty \exp\left(-\frac{s^2}{2}\right) ds = \sqrt{\frac{\pi}{2}}$$

Therefore:

$$\lim_{X \rightarrow +\infty} Y_{0,a}(X) = A \sqrt{\frac{\pi}{2}} + \pi$$

This is a finite constant!

Key Concept: The inner solution approaches a constant as $X \rightarrow +\infty$. This means matching with the outer solution is **possible** in principle. The outer solution also approaches a constant (a) as $x \rightarrow 0$. We will determine both A and a through the matching condition.

Step 5: Try Inner Solution at $x = \pi$

Strategy: Now assume there is a boundary layer at $x = \pi$. Define the inner variable by $x - \pi = \delta X$, or equivalently $x = \pi + \delta X$.

Step 5a: Transform the ODE

Set $x = \pi + \delta X$ and $Y(X) = y(x)$. The derivatives transform the same way:

$$\frac{dy}{dx} = \frac{Y'}{\delta}, \quad \frac{d^2y}{dx^2} = \frac{Y''}{\delta^2}$$

Substitute into the ODE:

$$\frac{\varepsilon}{\delta^2} Y'' + \sin(\pi + \delta X) \cdot \frac{Y'}{\delta} + \sin(2\pi + 2\delta X) \cdot Y = 0$$

Step 5b: Taylor Expand Around $x = \pi$

Use trigonometric identities and Taylor expansion:

$$\begin{aligned}\sin(\pi + \delta X) &= -\sin(\delta X) \approx -\delta X \\ \sin(2\pi + 2\delta X) &= \sin(2\delta X) \approx 2\delta X\end{aligned}$$

The ODE becomes:

$$\begin{aligned}\frac{\varepsilon}{\delta^2} Y'' + \frac{-\delta X}{\delta} Y' + 2\delta X \cdot Y &= 0 \\ \frac{\varepsilon}{\delta^2} Y'' - XY' + 2\delta XY &= 0\end{aligned}$$

Key Concept: Notice the crucial difference: the coefficient of Y' is now $-X$ instead of $+X$! This sign change will have dramatic consequences for whether matching is possible.

Step 5c: Dominant Balance

By the same argument as before, $\delta = \sqrt{\varepsilon}$.

Step 5d: Leading-Order Inner Equation at $x = \pi$

At leading order:

$$Y_0'' - XY_0' = 0$$

Note the **minus sign** compared to the equation at $x = 0$.

Step 5e: Solve the Leading-Order Inner Equation

Technique: Let $W = Y_0'$:

$$W' - XW = 0 \implies \frac{dW}{W} = X dX$$

Integrate:

$$\ln|W| = \frac{X^2}{2} + C_1 \implies W = Y_0' = C \exp\left(\frac{X^2}{2}\right)$$

Integrate again:

$$Y_0(X) = C \int_0^X \exp\left(\frac{s^2}{2}\right) ds + D$$

Step 5f: Apply Boundary Condition at $x = \pi$

At $x = \pi$, we have $X = 0$ and $y(\pi) = 0$. Therefore:

$$Y_0(0) = C \cdot 0 + D = D = 0 \implies \boxed{D = 0}$$

So the inner solution at $x = \pi$ is:

$$Y_{0,b}(X) = C \int_0^X \exp\left(\frac{s^2}{2}\right) ds$$

Step 5g: Check Matching at $x = \pi$

Technique: For matching, we need to check the behavior as $X \rightarrow -\infty$ (moving from the boundary layer at $x = \pi$ toward the interior). Consider:

$$\int_0^X \exp\left(\frac{s^2}{2}\right) ds \quad \text{as } X \rightarrow -\infty$$

Since $\exp(s^2/2) > 0$ for all s , the integral from 0 to $X < 0$ is:

$$\int_0^X = - \int_X^0 \exp\left(\frac{s^2}{2}\right) ds$$

As $X \rightarrow -\infty$, the integrand $\exp(s^2/2) \rightarrow \infty$, so:

$$\left| \int_X^0 \exp\left(\frac{s^2}{2}\right) ds \right| \rightarrow \infty$$

Warning: The inner solution at $x = \pi$ **diverges exponentially** as $X \rightarrow -\infty$:

$$Y_{0,b}(X) \sim C \cdot (\text{exponentially large}) \quad \text{as } X \rightarrow -\infty$$

Unless $C = 0$, this cannot be matched to the outer solution, which is finite.

Justification: Why does this happen? The sign difference in the inner equation changes the Gaussian from $e^{-X^2/2}$ (decaying) to $e^{+X^2/2}$ (growing). The growing Gaussian cannot be integrated to give a finite limit. This is the mathematical manifestation of the physical fact: boundary layers occur where the flow of information (determined by the sign of $p(x)$) is into the boundary, not out of it.

Step 6: Determine the Boundary Layer Structure

Key Concept: Our analysis reveals:

1. At $x = 0$: Inner solution approaches a finite limit as $X \rightarrow +\infty$. **Matching IS possible.**
2. At $x = \pi$: Inner solution diverges as $X \rightarrow -\infty$. **Matching is NOT possible.**

Conclusion: There is a boundary layer **only at** $x = 0$, and **no boundary layer at** $x = \pi$.

Justification: Why no boundary layer at $x = \pi$? Near $x = \pi$, we have $p(x) = \sin(x) \approx -(\pi - x) < 0$ for x slightly less than π . According to the general theory (Lecture Notes §6.2.1), when $p < 0$ near a boundary point, matching fails because the inner solution grows exponentially in the wrong direction. Even though $p(\pi) = 0$ exactly, the **sign of** $p'(\pi) = \cos(\pi) = -1 < 0$ determines that the effective behavior is as if $p < 0$.

Step 7: Apply Boundary Condition to Outer Solution

Since there is no boundary layer at $x = \pi$, the outer solution must satisfy the boundary condition there:

$$y_0(\pi) = 0$$

From Step 3b, $y_0(\pi) = a$. Therefore:

$$a = 0$$

The outer solution is:

$$y_0(x) = 0 \quad \text{for all } x \in (0, \pi]$$

Reflection: This is a remarkable result! The outer solution vanishes identically. The entire non-trivial behavior of the solution is concentrated in the boundary layer at $x = 0$. This happens because the boundary condition at $x = \pi$ forces $a = 0$, and the function $e^{-2\sin(x)}$ (while varying between e^{-2} and 1) is multiplied by zero.

Step 8: Perform Matching at $x = 0$

Now we match the inner solution at $x = 0$ with the (now known) outer solution $y_0(x) = 0$.

Step 8a: Van Dyke Matching

Technique: Apply Van Dyke's matching rule (Lecture Notes §6.1.3): The outer expansion of the inner solution must equal the inner expansion of the outer solution.

Inner expansion of outer solution:

The outer solution is $y_0(x) = 0$. Express in inner variable $x = \sqrt{\varepsilon}X$:

$$y_0(\sqrt{\varepsilon}X) = 0$$

This is already 0 for all X .

Outer expansion of inner solution:

The inner solution is:

$$Y_{0,a}(X) = A \int_0^X \exp\left(-\frac{s^2}{2}\right) ds + \pi$$

Express in outer variable $X = x/\sqrt{\varepsilon}$ and take $\varepsilon \rightarrow 0$ (which means $X \rightarrow \infty$ for fixed $x > 0$):

$$\lim_{X \rightarrow \infty} Y_{0,a}(X) = A \sqrt{\frac{\pi}{2}} + \pi$$

Step 8b: Matching Condition

For matching:

$$\lim_{X \rightarrow \infty} Y_{0,a}(X) = \lim_{x \rightarrow 0^+} y_0(x)$$

$$A \sqrt{\frac{\pi}{2}} + \pi = 0$$

Solve for A :

$$A = -\frac{\pi}{\sqrt{\pi/2}} = -\pi \sqrt{\frac{2}{\pi}} = -\sqrt{2\pi}$$

$$A = -\sqrt{2\pi}$$

Step 9: Write the Complete Inner Solution

The matched inner solution at $x = 0$ is:

$$Y_0(X) = -\sqrt{2\pi} \int_0^X \exp\left(-\frac{s^2}{2}\right) ds + \pi$$

We can express this using the error function. Recall (Lecture Notes §2.6.2, Eq. (71)):

$$\operatorname{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-t^2} dt$$

With substitution $t = s/\sqrt{2}$:

$$\int_0^X e^{-s^2/2} ds = \sqrt{2} \int_0^{X/\sqrt{2}} e^{-t^2} dt = \sqrt{2} \cdot \frac{\sqrt{\pi}}{2} \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right) = \sqrt{\frac{\pi}{2}} \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right)$$

Therefore:

$$\begin{aligned} Y_0(X) &= -\sqrt{2\pi} \cdot \sqrt{\frac{\pi}{2}} \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right) + \pi \\ &= -\pi \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right) + \pi \\ &= \pi \left[1 - \operatorname{erf}\left(\frac{X}{\sqrt{2}}\right)\right] \end{aligned}$$

Step 10: Write the Composite Solution

Since the outer solution vanishes ($y_0 = 0$), the composite solution equals the inner solution expressed in the original variable x .

Recall $X = x/\sqrt{\varepsilon}$, so:

$$\frac{X}{\sqrt{2}} = \frac{x}{\sqrt{2\varepsilon}}$$

Final Leading-Order Composite Solution:

$$y_c(x) = \pi \left[1 - \operatorname{erf}\left(\frac{x}{\sqrt{2\varepsilon}}\right)\right] = \pi \operatorname{erfc}\left(\frac{x}{\sqrt{2\varepsilon}}\right)$$

where $\operatorname{erfc}(z) = 1 - \operatorname{erf}(z)$ is the complementary error function.

Step 11: Verify the Solution Properties

Technique: Let's verify that our solution satisfies the required properties:

Boundary condition at $x = 0$:

$$y_c(0) = \pi [1 - \operatorname{erf}(0)] = \pi[1 - 0] = \pi \quad \checkmark$$

Boundary condition at $x = \pi$:

As $\varepsilon \rightarrow 0$ with $x = \pi$ fixed:

$$\frac{\pi}{\sqrt{2\varepsilon}} \rightarrow \infty$$

$$\operatorname{erf}\left(\frac{\pi}{\sqrt{2\varepsilon}}\right) \rightarrow 1$$

$$y_c(\pi) \rightarrow \pi[1 - 1] = 0 \quad \checkmark$$

Behavior in the outer region ($x = O(1)$, $x > 0$):

For fixed $x > 0$, as $\varepsilon \rightarrow 0$:

$$\frac{x}{\sqrt{2\varepsilon}} \rightarrow \infty, \quad \text{erf} \rightarrow 1, \quad y_c(x) \rightarrow 0 = y_0(x) \quad \checkmark$$

Boundary layer width:

The transition from $y \approx \pi$ to $y \approx 0$ occurs when $x/\sqrt{2\varepsilon} = O(1)$, i.e., when $x = O(\sqrt{\varepsilon})$. This confirms the boundary layer width is $\delta = \sqrt{\varepsilon}$.

Step 12: Physical and Mathematical Interpretation

Reflection: The solution structure reveals several important features:

1. **Boundary layer only at $x = 0$:** Although $p(x) = \sin(x)$ vanishes at both endpoints, a boundary layer exists only at $x = 0$. This is determined by the **sign of $p'(x)$** :
 - At $x = 0$: $p'(0) = \cos(0) = +1 > 0 \Rightarrow$ boundary layer exists
 - At $x = \pi$: $p'(\pi) = \cos(\pi) = -1 < 0 \Rightarrow$ no boundary layer
2. **Boundary layer width $\sim \sqrt{\varepsilon}$:** When $p(x_0) = 0$ but $p'(x_0) \neq 0$, the boundary layer is **thicker** than the standard $O(\varepsilon)$ width. This is because the coefficient of y' is $O(\delta X) = O(\sqrt{\varepsilon})$ near the boundary, requiring a larger region for the $\varepsilon y''$ term to be important.
3. **Gaussian profile:** The inner solution involves the error function, which arises from integrating a Gaussian. This is characteristic of boundary layers at points where $p(x)$ has a simple zero.
4. **Outer solution vanishes:** The solution is essentially concentrated entirely in the boundary layer. Outside the layer, the solution is asymptotically zero. This happens because the boundary condition $y(\pi) = 0$, combined with the fact that y_0 is constant at both endpoints, forces $y_0 \equiv 0$.
5. **Connection to parabolic cylinder functions:** The inner equation $Y'' + XY' = 0$ is related to parabolic cylinder equations (Lecture Notes §6.2.2). For this specific case, the solution reduces to error functions, but more general cases might require parabolic cylinder functions.

Summary Table

Component	Result
ODE	$\varepsilon y'' + \sin(x)y' + \sin(2x)y = 0$
Domain	$[0, \pi]$
Boundary conditions	$y(0) = \pi, y(\pi) = 0$
Outer solution	$y_0(x) = 0$
Boundary layer location	$x = 0$ only
Boundary layer width	$\delta = \sqrt{\varepsilon}$
Inner variable	$X = x/\sqrt{\varepsilon}$
Inner equation	$Y_0'' + XY_0' = 0$
Inner solution	$Y_0(X) = \pi \left[1 - \operatorname{erf} \left(\frac{X}{\sqrt{2}} \right) \right]$
Composite solution	$y_c(x) = \pi \operatorname{erfc} \left(\frac{x}{\sqrt{2\varepsilon}} \right)$

Connection to Lecture Material

Reflection: This problem illustrates several key concepts from the course:

- **Workflow for boundary layers** (Lecture Notes §6.2.3): We followed the systematic approach of identifying candidate locations, computing inner and outer solutions, and checking matching.
- **Boundary layers of non-standard width** (Lecture Notes §6.2.2): When $p(x_0) = 0$, the dominant balance gives $\delta = \sqrt{\varepsilon}$ rather than $\delta = \varepsilon$.
- **Sign of $p'(x_0)$ determines layer existence:** Even when $p(x_0) = 0$ at both boundaries, only one may have a boundary layer, determined by whether the exponential in the inner solution decays or grows.
- **Error function solutions:** The Gaussian integrals appearing in boundary layer problems (Lecture Notes §2.6.2) lead naturally to error functions in the final solution.
- **Matching determines integration constants:** The constants A in the inner solution and a in the outer solution were both determined by matching conditions and boundary conditions, not independently.