

# Asymptotics Problem 8.4: Complete Pedagogical Solution

## Boundary Layer with Non-Standard Width

**Problem 1.** Find a first-order uniform expansion as  $\varepsilon \rightarrow 0$  for  $y(x)$  satisfying

$$\varepsilon y'' + x^2 y' - x^3 y = 0, \quad y(0) = \alpha, \quad y(1) = \beta.$$

### Solution: Step-by-Step Atomic Breakdown

#### Step 1: Problem Classification and Initial Analysis

**Strategy:** This is a singularly perturbed second-order linear ODE of the general form

$$\varepsilon y'' + p(x)y' + q(x)y = 0$$

with  $p(x) = x^2$  and  $q(x) = -x^3$ . Our systematic approach follows the workflow from Lecture Notes §6.2.3:

1. Identify candidate locations for boundary layers
2. Determine the boundary layer width  $\delta(\varepsilon)$  via dominant balance
3. Compute outer and inner solutions
4. Match solutions using Prandtl's matching criterion
5. Construct the composite solution

**Justification:** Why is this a boundary layer problem? The coefficient  $\varepsilon$  multiplying the highest derivative  $y''$  is small. As  $\varepsilon \rightarrow 0$ , the order of the ODE effectively reduces from 2 to 1, but we have two boundary conditions. This is the hallmark of a singular perturbation problem where rapid variation (a boundary layer) must occur somewhere to accommodate both boundary conditions.

#### Step 2: Identifying the Boundary Layer Location

**What we examine:** The sign and zeros of the coefficient  $p(x) = x^2$ .

**Key Concept:** From the general theory of boundary layers (Lecture Notes §6.2.1, equations (340)–(354)), for the ODE  $\varepsilon y'' + p(x)y' + q(x)y = 0$ :

- If  $p(x) > 0$  on  $[0, 1]$ : boundary layer at  $x = 0$
- If  $p(x) < 0$  on  $[0, 1]$ : boundary layer at  $x = 1$
- If  $p(x_0) = 0$  for some  $x_0 \in [0, 1]$ : special treatment required

The physical intuition: the sign of  $p(x)$  determines whether information “propagates” from left to right or right to left.

**Analysis of  $p(x) = x^2$ :**

- For  $x > 0$ :  $p(x) = x^2 > 0 \quad \checkmark$

- At  $x = 0$ :  $p(0) = 0$  (vanishes at boundary!)
- At  $x = 1$ :  $p(1) = 1 > 0$  ✓

**Justification:** Since  $p(x) = x^2 > 0$  for all  $x > 0$ , there cannot be a boundary layer in the interior or at  $x = 1$ . The only possible location is  $x = 0$ .

However, this case is special because  $p(0) = 0$ —the coefficient of  $y'$  vanishes exactly at the boundary point where we expect the boundary layer! According to Lecture Notes §6.2.2 (equation (356) and surrounding discussion): “For a boundary point  $x_0$  with  $p(x_0) = 0$ , we can have boundary layers of different width than  $\sim \varepsilon$ .”

This is precisely our situation. We must determine the actual boundary layer width through careful dominant balance analysis.

**Conclusion:** Boundary layer at  $x = 0$  with non-standard width  $\delta \neq O(\varepsilon)$ .

### Step 3: Computing the Outer Solution

**What we do:** In the outer region (away from  $x = 0$ ), the solution varies slowly, so we neglect the  $\varepsilon y''$  term.

**Technique:** Setting  $\varepsilon = 0$  in the original ODE gives the leading-order outer equation:

$$x^2 y'_0 - x^3 y_0 = 0.$$

**Solving the outer equation:**

Divide by  $x^2$  (valid for  $x \neq 0$ , which is the outer region):

$$y'_0 - xy_0 = 0.$$

**Technique:** This is a first-order linear homogeneous ODE. Separate variables:

$$\frac{y'_0}{y_0} = x \implies \frac{dy_0}{y_0} = x dx.$$

Integrate both sides:

$$\ln |y_0| = \frac{x^2}{2} + C'.$$

Exponentiate:

$$y_0(x) = C \exp\left(\frac{x^2}{2}\right),$$

where  $C$  is an arbitrary constant.

**Applying the boundary condition at  $x = 1$ :**

**Justification:** Since the boundary layer is at  $x = 0$ , the outer solution is valid at  $x = 1$ . We can therefore apply the boundary condition  $y(1) = \beta$  directly to the outer solution:

$$y_0(1) = C \exp\left(\frac{1}{2}\right) = \beta \implies C = \beta e^{-1/2} = \frac{\beta}{\sqrt{e}}.$$

**Final outer solution:**

$$y_0(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right)$$

**Reflection:** Notice that  $y_0(0) = \beta e^{-1/2} = \beta/\sqrt{e} \neq \alpha$  in general. This confirms that the outer solution cannot satisfy both boundary conditions—there must be a boundary layer at  $x = 0$  to transition from  $y(0) = \alpha$  to the outer solution.

## Step 4: Setting Up the Inner Region

**What we do:** Introduce stretched (inner) coordinates to resolve the rapid variation near  $x = 0$ .  
**Technique:** Define the inner variable  $X$  by:

$$x = \delta X, \quad \text{where } \delta = \delta(\varepsilon) \rightarrow 0 \text{ as } \varepsilon \rightarrow 0.$$

Define the inner solution  $Y(X) = y(x) = y(\delta X)$ .

**Justification:** The scaling  $x = \delta X$  “zooms in” on the boundary layer region near  $x = 0$ . When  $x = O(\delta)$ , we have  $X = O(1)$ , so the inner variable  $X$  is order one within the boundary layer. The function  $\delta(\varepsilon)$  represents the boundary layer width—the region where rapid variation occurs.

### Transforming derivatives:

Using the chain rule:

$$\begin{aligned} \frac{dy}{dx} &= \frac{dY}{dX} \cdot \frac{dX}{dx} = \frac{1}{\delta} Y'(X), \\ \frac{d^2y}{dx^2} &= \frac{d}{dx} \left( \frac{1}{\delta} Y' \right) = \frac{1}{\delta} \cdot \frac{dY'}{dX} \cdot \frac{dX}{dx} = \frac{1}{\delta^2} Y''(X). \end{aligned}$$

### Substituting into the original ODE:

The original equation is:

$$\varepsilon y'' + x^2 y' - x^3 y = 0.$$

With  $x = \delta X$  and  $y = Y$ :

$$\begin{aligned} \varepsilon \cdot \frac{1}{\delta^2} Y'' + (\delta X)^2 \cdot \frac{1}{\delta} Y' - (\delta X)^3 Y &= 0 \\ \frac{\varepsilon}{\delta^2} Y'' + \frac{\delta^2 X^2}{\delta} Y' - \delta^3 X^3 Y &= 0 \\ \frac{\varepsilon}{\delta^2} Y'' + \delta X^2 Y' - \delta^3 X^3 Y &= 0. \end{aligned}$$

$$\text{Inner equation: } \frac{\varepsilon}{\delta^2} Y'' + \delta X^2 Y' - \delta^3 X^3 Y = 0$$

## Step 5: Dominant Balance Analysis

**Strategy:** To find the boundary layer width  $\delta$ , we require that the leading terms in the inner equation balance each other as  $\varepsilon \rightarrow 0$ . This is the dominant balance principle (Lecture Notes §2.2.2 and §6.2.2).

### Examining the three terms:

The inner equation has three terms with coefficients:

1.  $Y''$  term: coefficient  $\varepsilon/\delta^2$
2.  $Y'$  term: coefficient  $\delta$
3.  $Y$  term: coefficient  $\delta^3$

**Justification:** We seek  $\delta \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . Compare the magnitudes:

- The  $Y$  term has coefficient  $\delta^3$ .
- The  $Y'$  term has coefficient  $\delta$ .
- Since  $\delta \rightarrow 0$ , we have  $\delta^3 \ll \delta$ , so the  $Y$  term is smaller than the  $Y'$  term.

Therefore, the third term  $(-\delta^3 X^3 Y)$  is subdominant and can be neglected at leading order.

### Balancing the remaining two terms:

**Technique:** For a distinguished limit (where both remaining terms contribute at the same order), we require:

$$\frac{\varepsilon}{\delta^2} \sim \delta.$$

Solving for  $\delta$ :

$$\varepsilon \sim \delta^3 \implies \boxed{\delta = \varepsilon^{1/3}}$$

**Reflection:** This is a **non-standard boundary layer width!** In the typical case where  $p(x_0) \neq 0$  at the boundary, we get  $\delta = \varepsilon$ . Here, because  $p(0) = 0$ , we obtain  $\delta = \varepsilon^{1/3}$ , which is larger than  $\varepsilon$  (since  $\varepsilon^{1/3} \gg \varepsilon$  for small  $\varepsilon$ ).

This confirms the statement from Lecture Notes §6.2.2: when the coefficient of  $y'$  vanishes at the boundary, “we can have boundary layers of different width than  $\sim \varepsilon$ , here  $\sim \sqrt{\varepsilon}$ ” (in some cases) or  $\sim \varepsilon^{1/3}$  (in our case).

**Key Concept:** Verification of dominant balance: With  $\delta = \varepsilon^{1/3}$ :

- $\varepsilon/\delta^2 = \varepsilon/\varepsilon^{2/3} = \varepsilon^{1/3} = \delta \quad \checkmark$  (these balance)
- $\delta^3 = \varepsilon \ll \delta = \varepsilon^{1/3} \quad \checkmark$  (third term is subdominant)

The analysis is self-consistent.

## Step 6: Deriving the Leading-Order Inner Equation

**What we do:** With  $\delta = \varepsilon^{1/3}$ , write the inner equation and identify the leading-order problem. Substituting  $\delta = \varepsilon^{1/3}$ :

$$\varepsilon^{1/3} Y'' + \varepsilon^{1/3} X^2 Y' - \varepsilon X^3 Y = 0.$$

Divide through by  $\varepsilon^{1/3}$ :

$$Y'' + X^2 Y' - \varepsilon^{2/3} X^3 Y = 0.$$

**Technique:** As  $\varepsilon \rightarrow 0$ , the term  $\varepsilon^{2/3} X^3 Y \rightarrow 0$ . At leading order:

$$\boxed{Y_0'' + X^2 Y_0' = 0}$$

This is the leading-order inner equation.

**Justification:** Why can we neglect  $\varepsilon^{2/3} X^3 Y$ ? In the inner region,  $X = O(1)$ , so  $X^3 = O(1)$ . The coefficient  $\varepsilon^{2/3} \rightarrow 0$ , making this term small compared to the  $O(1)$  terms  $Y''$  and  $X^2 Y'$ .

## Step 7: Solving the Inner Equation

**What we solve:**  $Y_0'' + X^2 Y_0' = 0$ .

### Step 7a: First Integration

**Technique:** This is a second-order ODE that can be reduced to first order by the substitution  $P = Y'_0$ :

$$P' + X^2 P = 0.$$

This is a separable first-order linear ODE.

Separate variables:

$$\frac{dP}{P} = -X^2 dX.$$

Integrate:

$$\ln |P| = -\frac{X^3}{3} + C_1.$$

Exponentiate:

$$P = Y'_0 = A \exp\left(-\frac{X^3}{3}\right),$$

where  $A$  is an arbitrary constant.

### Step 7b: Second Integration

**Technique:** Integrate  $Y'_0$  to find  $Y_0$ :

$$Y_0(X) = A \int_0^X \exp\left(-\frac{s^3}{3}\right) ds + B,$$

where  $B$  is another arbitrary constant, and we choose the lower limit of integration as 0 for convenience.

**Justification:** The choice of lower limit 0 is convenient because:

1. At  $X = 0$ , the integral vanishes, giving  $Y_0(0) = B$  directly.
2. This form makes applying the boundary condition at  $x = 0$  (i.e.,  $X = 0$ ) straightforward.

**General inner solution:**

$$Y_0(X) = A \int_0^X \exp\left(-\frac{s^3}{3}\right) ds + B$$

### Step 8: Applying the Boundary Condition at $x = 0$

**What we apply:**  $y(0) = \alpha$ , which in inner variables is  $Y_0(0) = \alpha$ .

From the inner solution:

$$Y_0(0) = A \int_0^0 \exp\left(-\frac{s^3}{3}\right) ds + B = 0 + B = B.$$

Therefore:

$$B = \alpha$$

**Inner solution with boundary condition:**

$$Y_0(X) = A \int_0^X \exp\left(-\frac{s^3}{3}\right) ds + \alpha$$

**Reflection:** We have one remaining unknown constant  $A$ , which will be determined by matching the inner and outer solutions in the overlap (intermediate) region.

### Step 9: Asymptotic Matching

**Strategy:** We use **Prandtl's matching rule** (Lecture Notes §6.1.2): the outer limit of the inner solution must equal the inner limit of the outer solution. This is valid here because both limits approach constants, making the matching straightforward.

#### Step 9a: Inner Limit of the Outer Solution

**Technique:** The “inner limit” means taking  $x \rightarrow 0^+$  in the outer solution:

$$\lim_{x \rightarrow 0^+} y_0(x) = \lim_{x \rightarrow 0^+} \beta \exp\left(\frac{x^2 - 1}{2}\right) = \beta \exp\left(\frac{0 - 1}{2}\right) = \beta e^{-1/2} = \frac{\beta}{\sqrt{e}}.$$

### Step 9b: Outer Limit of the Inner Solution

**Technique:** The “outer limit” means taking  $X \rightarrow +\infty$  in the inner solution (since  $X = x/\varepsilon^{1/3} \rightarrow \infty$  as we leave the boundary layer):

$$\lim_{X \rightarrow \infty} Y_0(X) = \lim_{X \rightarrow \infty} \left[ A \int_0^X \exp\left(-\frac{s^3}{3}\right) ds + \alpha \right] = A \cdot I + \alpha,$$

where we define the integral:

$$I = \int_0^\infty \exp\left(-\frac{s^3}{3}\right) ds.$$

### Step 9c: Evaluating the Integral $I$

**Technique:** This is a generalized Gaussian integral. Using the substitution  $u = s^3/3$ , so  $s = (3u)^{1/3}$  and  $ds = (3u)^{-2/3} du$ :

$$I = \int_0^\infty e^{-u} \cdot (3u)^{-2/3} du = 3^{-2/3} \int_0^\infty u^{-2/3} e^{-u} du = 3^{-2/3} \Gamma\left(\frac{1}{3}\right),$$

where  $\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$  is the gamma function (Lecture Notes §2.6.1).

**Justification:** We used the gamma function identity  $\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$  with  $z = 1/3$ :

$$\int_0^\infty u^{1/3-1} e^{-u} du = \int_0^\infty u^{-2/3} e^{-u} du = \Gamma\left(\frac{1}{3}\right).$$

The numerical value is  $\Gamma(1/3) \approx 2.679$ .

**Result:**

$$I = \int_0^\infty \exp\left(-\frac{s^3}{3}\right) ds = 3^{-2/3} \Gamma\left(\frac{1}{3}\right)$$

### Step 9d: Applying Prandtl's Matching Condition

**Technique:** Equate the inner limit of outer and outer limit of inner:

$$\frac{\beta}{\sqrt{e}} = \alpha + A \cdot I.$$

Solve for  $A$ :

$$A = \frac{\beta e^{-1/2} - \alpha}{I} = \frac{\beta/\sqrt{e} - \alpha}{I}.$$

**Matched constant:**

$$A = \boxed{\frac{\beta e^{-1/2} - \alpha}{I} = \frac{\beta e^{-1/2} - \alpha}{3^{-2/3} \Gamma(1/3)}}$$

**Reflection:** The constant  $A$  encodes the “mismatch” between the boundary value  $\alpha$  and the limiting value  $\beta e^{-1/2}$  of the outer solution as  $x \rightarrow 0$ . If  $\alpha = \beta e^{-1/2}$ , then  $A = 0$  and the inner solution reduces to the constant  $\alpha$ —no boundary layer correction is needed!

### Step 10: Constructing the Composite Solution

**Strategy:** The composite solution is constructed by adding the outer and inner solutions and subtracting their common limit (to avoid double-counting in the overlap region). This follows Lecture Notes §6.1.3 and equation (353):

$$y_c(x) = y_{\text{outer}}(x) + Y_{\text{inner}}\left(\frac{x}{\delta}\right) - (\text{common limit}).$$

**Components:**

$$\text{Outer solution: } y_0(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right)$$

$$\text{Inner solution: } Y_0(X) = \alpha + A \int_0^X \exp\left(-\frac{s^3}{3}\right) ds$$

$$\text{Common limit: } \frac{\beta}{\sqrt{e}} = \beta e^{-1/2}$$

**Technique:** *The common limit is what both solutions approach in the intermediate (matching) region:*

- As  $x \rightarrow 0^+$ :  $y_0(x) \rightarrow \beta e^{-1/2}$
- As  $X \rightarrow \infty$ :  $Y_0(X) \rightarrow \alpha + AI = \beta e^{-1/2}$  (by matching)

**Composite solution formula:**

$$y_c(x) = y_0(x) + Y_0\left(\frac{x}{\varepsilon^{1/3}}\right) - \frac{\beta}{\sqrt{e}}$$

Substituting:

$$y_c(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right) + \alpha + A \int_0^{x/\varepsilon^{1/3}} \exp\left(-\frac{s^3}{3}\right) ds - \frac{\beta}{\sqrt{e}}$$

**Simplifying:**

**Technique:** Using  $A = (\beta e^{-1/2} - \alpha)/I$  and the fact that  $\int_0^\infty e^{-s^3/3} ds = I$ :

$$A \int_0^X e^{-s^3/3} ds = \frac{\beta e^{-1/2} - \alpha}{I} \int_0^X e^{-s^3/3} ds = \left(\frac{\beta}{\sqrt{e}} - \alpha\right) \frac{\int_0^X e^{-s^3/3} ds}{\int_0^\infty e^{-s^3/3} ds}.$$

Therefore, the composite solution can be written as:

$$y_c(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right) + \left(\frac{\beta}{\sqrt{e}} - \alpha\right) \left[ \frac{\int_0^{x\varepsilon^{-1/3}} e^{-s^3/3} ds}{\int_0^\infty e^{-s^3/3} ds} - 1 \right]$$

### Step 11: Final Answer in Standard Form

**First-Order Uniform Expansion:**

$$y_c(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right) + \left(\frac{\beta}{\sqrt{e}} - \alpha\right) \left[ \frac{\int_0^{x\varepsilon^{-1/3}} e^{-s^3/3} ds}{\int_0^\infty e^{-s^3/3} ds} - 1 \right]$$

Equivalently, with explicit integral notation:

$$y_c(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right) + \left(\frac{\beta}{\sqrt{e}} - \alpha\right) \left[ \frac{\int_0^{x/\varepsilon^{1/3}} \exp\left(-\frac{s^3}{3}\right) ds}{\int_0^\infty \exp\left(-\frac{s^3}{3}\right) ds} - 1 \right]$$

$$\text{where } \int_0^\infty \exp\left(-\frac{s^3}{3}\right) ds = 3^{-2/3} \Gamma\left(\frac{1}{3}\right).$$

## Step 12: Verification of the Solution

### Verification 1: Boundary Condition at $x = 0$

**Technique:** As  $x \rightarrow 0^+$ :

- $\beta \exp\left(\frac{x^2-1}{2}\right) \rightarrow \beta e^{-1/2}$
- $x\varepsilon^{-1/3} \rightarrow 0$ , so  $\int_0^{x\varepsilon^{-1/3}} e^{-s^3/3} ds \rightarrow 0$

Therefore:

$$y_c(0) = \frac{\beta}{\sqrt{e}} + \left( \frac{\beta}{\sqrt{e}} - \alpha \right) (0 - 1) = \frac{\beta}{\sqrt{e}} - \frac{\beta}{\sqrt{e}} + \alpha = \alpha \quad \checkmark$$

### Verification 2: Boundary Condition at $x = 1$

**Technique:** At  $x = 1$ , the inner variable  $X = 1/\varepsilon^{1/3} \rightarrow \infty$  as  $\varepsilon \rightarrow 0$ :

- $\beta \exp\left(\frac{1-1}{2}\right) = \beta$
- $\int_0^{\varepsilon^{-1/3}} e^{-s^3/3} ds \rightarrow \int_0^{\infty} e^{-s^3/3} ds = I$

Therefore:

$$y_c(1) = \beta + \left( \frac{\beta}{\sqrt{e}} - \alpha \right) \left( \frac{I}{I} - 1 \right) = \beta + \left( \frac{\beta}{\sqrt{e}} - \alpha \right) (0) = \beta \quad \checkmark$$

### Verification 3: Behavior in Outer Region

**Technique:** For fixed  $x > 0$  as  $\varepsilon \rightarrow 0$ :

$$x\varepsilon^{-1/3} \rightarrow \infty, \quad \text{so} \quad \frac{\int_0^{x\varepsilon^{-1/3}} e^{-s^3/3} ds}{I} \rightarrow 1.$$

Therefore:

$$y_c(x) \rightarrow \beta \exp\left(\frac{x^2-1}{2}\right) + \left( \frac{\beta}{\sqrt{e}} - \alpha \right) (1 - 1) = \beta \exp\left(\frac{x^2-1}{2}\right) = y_0(x) \quad \checkmark$$

## Step 13: Physical Interpretation and Key Insights

**Reflection:** This problem illustrates several important asymptotic concepts:

1. **Non-standard boundary layer width:** When the coefficient of  $y'$  vanishes at the boundary ( $p(0) = x^2|_{x=0} = 0$ ), the boundary layer width is determined by dominant balance as  $\delta = \varepsilon^{1/3}$ , not  $\delta = \varepsilon$ . This is because near  $x = 0$ , the  $y'$  term “turns on” gradually (like  $x^2$ ), requiring a wider region for the solution to transition.
2. **The integral  $\int_0^{\infty} e^{-s^3/3} ds$ :** This generalized Gaussian integral appears because the inner equation  $Y'' + X^2 Y' = 0$  has solutions involving  $e^{-X^3/3}$ . The exponent  $X^3/3$  comes from integrating the coefficient  $X^2$ .
3. **Matching via Prandtl's rule:** Both the inner and outer solutions approach constants as they enter the intermediate region, making Prandtl's matching (comparing limits) straightforward. Van Dyke matching would give the same result here.
4. **Composite solution structure:** The additive form  $y_c = y_{\text{outer}} + Y_{\text{inner}} - (\text{common limit})$  ensures that:
  - In the inner region:  $y_{\text{outer}} \approx \text{common limit}$ , so  $y_c \approx Y_{\text{inner}}$
  - In the outer region:  $Y_{\text{inner}} \approx \text{common limit}$ , so  $y_c \approx y_{\text{outer}}$

This is the essence of matched asymptotic expansions from Lecture Notes §6.1.3.

## Summary Table

Quantity	Value
Boundary layer location	$x = 0$
Boundary layer width	$\delta = \varepsilon^{1/3}$
Inner variable	$X = x/\varepsilon^{1/3}$
Outer solution	$y_0(x) = \beta \exp\left(\frac{x^2 - 1}{2}\right)$
Inner equation	$Y_0'' + X^2 Y_0' = 0$
Inner solution	$Y_0(X) = \alpha + A \int_0^X e^{-s^3/3} ds$
Matching constant	$A = \frac{\beta e^{-1/2} - \alpha}{3^{-2/3} \Gamma(1/3)}$
Common limit	$\beta e^{-1/2} = \frac{\beta}{\sqrt{e}}$

## Connection to Lecture Material

**Key Concept:** This problem is a direct application of the boundary layer workflow from Lecture Notes §6.2.3, with the special feature that  $p(x_0) = 0$  at the boundary (case discussed in §6.2.2, equation (356)). The key steps parallel Example 1 on page 61–62 of the lecture notes, which treats  $\varepsilon y'' + x^2 y' - y = 0$ —a closely related problem.

The dominant balance analysis follows §2.2.2, the matching procedure follows §6.1.2 (Prandtl's rule), and the composite solution construction follows §6.1.3 (equation (353)).