

(Quelques) solutions exactes des équations de Navier-Stokes

Schéma	Hypothèses	Solution	Nom de l'écoulement
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$	$\vec{u} = u(y) \vec{e}_x \quad \text{avec} \quad u(y) = \frac{Uy}{h}$ $p = C^{te}$	Couette plan sans gradient de pression ($\partial p / \partial x = 0$)
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$	$\frac{dp}{dx} = C^{te} \quad \text{et} \quad \vec{u} = u(y) \vec{e}_x$ $\text{avec} \quad u(y) = \frac{1}{2\mu} \frac{dp}{dx} y(h-y)$	Poiseuille plan
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$	$\frac{dp}{dx} = C^{te} \quad \text{et} \quad \vec{u} = u(y) \vec{e}_x$ $\text{avec} \quad u(y) = U \left[\frac{y}{h} - K \frac{y}{h} \left(1 - \frac{y}{h} \right) \right]$ $\text{où} \quad K = \frac{h^2}{2\mu U} \frac{dp}{dx}$	Couette plan avec gradient de pression ($\partial p / \partial x \neq 0$) = "Couette-Poiseuille"
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial \theta} = \vec{0}$	$\frac{dp}{dx} = C^{te} \quad \text{et} \quad \vec{u} = u(r) \vec{e}_x$ $\text{avec} \quad u(r) = -\frac{1}{4\mu} \frac{dp}{dx} (R^2 - r^2)$	Poiseuille cylindrique
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial \theta} = \vec{0}$ $\frac{\partial p}{\partial x} = 0$	$\vec{u} = v(r) \vec{e}_\theta \quad \text{avec} \quad v(r) = Ar + \frac{B}{r}$ $\text{où} \quad A = \frac{\Omega_2 R_2^2 - \Omega_1 R_1^2}{R_2^2 - R_1^2}$ $\text{et} \quad B = \frac{(\Omega_1 - \Omega_2) R_1^2 R_2^2}{R_2^2 - R_1^2}$	Couette cylindrique = "Couette-Taylor"
	$\frac{\partial \vec{u}}{\partial t} = \vec{0}$ $\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$	$\vec{u} = u(y) \vec{e}_x$ $\text{avec} \quad u(y) = \frac{g \sin \alpha}{2\nu} y(2h - y)$ $p(y) = \rho g(h - y) \cos \alpha + P_a$	Film tombant
	$\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$ $\frac{\partial p}{\partial x} = 0$	$\vec{u} = u(y, t) \vec{e}_x$ $\text{avec} \quad u(y, t) = U \left[1 - \operatorname{erf} \left(\frac{y}{2\sqrt{\nu t}} \right) \right]$ $\tau_p = \mu \left. \frac{\partial u}{\partial y} \right _{y=0} = -\rho U \sqrt{\frac{\nu}{\pi t}}$	Premier problème de Stokes
	$\frac{\partial \vec{u}}{\partial x} = \frac{\partial \vec{u}}{\partial z} = \vec{0}$ $\frac{\partial p}{\partial x} = 0$	$\vec{u} = u(y, t) \vec{e}_x$ $u(y, t) = U e^{-y/\delta} \cos(\omega t - \frac{y}{\delta}); \quad \delta = \sqrt{\frac{2\nu}{\omega}}$ $\tau_p = \mu \left. \frac{\partial u}{\partial y} \right _{y=0} = -\frac{\mu U \sqrt{2}}{\delta} \cos(\omega t - \pi/4)$	Second problème de Stokes