### Signals and Systems

Lecturer: Dr. Lin Xu, Dr. Xiran Cai

Email: xulin1@shanghaitech.edu.cn

caixr@shanghaitech.edu.cn

Office: 3-428(Xu), 3-438(Cai) SIST

Tel: 20684449(Xu), 20684431(Cai)

**ShanghaiTech University** 



# Chapter 3: Fourier Series Representation of Periodic signals

- ☐ The response of LTI systems to complex exponentials
- **□** Fourier series representation of continuous periodic signals
- **□** Convergence of the Fourier series
- **☐** Properties of continuous-time Fourier series
- **□** Fourier series representation of discrete –time periodic signals
- Properties of discrete FS
- **□** Fourier series and LTI systems



#### Recall Chapter 2

☐ Objective: characterization of a LTI system

$$x(t) \longrightarrow \boxed{\qquad \qquad } y(t)$$

 $\square x(t)$  is considered as linear combinations of a basis signal  $\delta(t)$ 

$$x(t) = \int_{-\infty}^{\infty} x(\tau)\delta(t-\tau) d\tau \quad \to \quad y(t) = \int_{-\infty}^{\infty} x(\tau)h(t-\tau) d\tau$$

- $\square$   $\delta(t)$  is not the only one. In general, a basic signal should satisfy
  - It can be used to construct a broad and useful class of signals
  - The response of an LTI system to the basic signal is simple



#### Continuous-time

$$e^{st}$$
  $\longrightarrow$   $UTI$   $\longrightarrow$   $U(t) = ?$ 

$$y(t) = \int_{-\infty}^{\infty} h(\tau)e^{s(t-\tau)}d\tau = e^{st} \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau$$

Let 
$$\int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau = H(s) \rightarrow y(t) = H(s)e^{st}$$

- $e^{st}$  is an eigenfunction of the system
- For a specific value s, H(s) is the corresponding eigenvalue



#### Discrete-time

$$z^n \longrightarrow \boxed{\text{LTI}} \longrightarrow y[n] = ?$$

$$y[n] = \sum_{k=-\infty}^{\infty} h[k]z^{n-k} = z^n \sum_{k=-\infty}^{\infty} h[k]z^{-k}$$

Let 
$$H[z] = \sum_{k=-\infty}^{\infty} h[k]z^{-k}$$
  $\rightarrow y[n] = H[z]z^n$ 

- $z^n$  is an eigenfunction of the system
- For a specific value z, H[z] is the corresponding eigenvalue



#### Continuous-time

$$e^{st} \longrightarrow \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau \, e^{st} = H(s)e^{st}$$

If 
$$x(t) = a_1 e^{s_1 t} + a_2 e^{s_2 t} + a_3 e^{s_3 t}$$
  $y(t) = ?$ 

$$y(t) = a_1 H(s_1) e^{s_1 t} + a_2 H(s_2) e^{s_2 t} + a_1 H(s_3) e^{s_3 t}$$

Generally, if 
$$x(t) = \sum_{k} a_k e^{s_k t}$$

$$y(t) = \sum_{k} a_k H(s_k) e^{s_k t}$$



#### Discrete-time

$$z^n \longrightarrow \underbrace{\text{LTI}}_{k=-\infty} h[k] z^{-k} z^n = H[z] z^n$$

If 
$$x[n] = \sum_{k} a_k Z_k^n$$

$$y[n] = \sum_{k} a_k H(z_k) Z_k^n$$



#### **Examples**

For a LTI system y(t) = x(t - 3), determine H(s)

#### Solution 1:

$$let x(t) = e^{st}, y(t) = e^{s(t-3)} = e^{-3s}e^{st}$$
$$\therefore H(s) = e^{-3s}$$

#### Solution 2:

$$H(s) = \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau = \int_{-\infty}^{\infty} \delta(\tau - 3)e^{-s\tau}d\tau = e^{-3s}$$



#### **Examples**

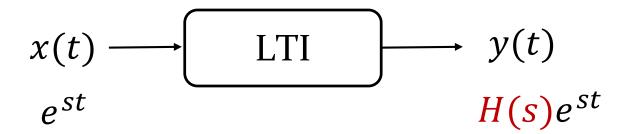
For a LTI system 
$$y(t) = x(t-3)$$
  
If  $x(t) = \cos(4t) + \cos(7t)$ ,  $y(t) = ?$   
Solution 1:  $y(t) = \cos(4(t-3)) + \cos(7(t-3))$   
Solution 2:  $x(t) = \frac{1}{2}e^{j4t} + \frac{1}{2}e^{-j4t} + \frac{1}{2}e^{j7t} + \frac{1}{2}e^{-j7t}$   
 $y(t) = \frac{1}{2}H(j4)e^{j4t} + \frac{1}{2}H(-j4)e^{-j4t} + \frac{1}{2}H(j7)e^{j7t} + \frac{1}{2}H(-j7)e^{-j7t}$   
 $H(s) = e^{-3s} = \frac{1}{2}e^{-j12}e^{j4t} + \frac{1}{2}e^{j12}e^{-j4t} + \frac{1}{2}e^{-j21}e^{j7t} + \frac{1}{2}e^{j21}e^{-j7t}$   
 $= \frac{1}{2}e^{j4(t-3)} + \frac{1}{2}e^{-j4(t-3)} + \frac{1}{2}e^{j7(t-3)} + \frac{1}{2}e^{-j7(t-3)}$ 

# Chapter 3: Fourier Series Representation of Periodic signals

- ☐ The response of LTI systems to complex exponentials
- **□** Fourier series representation of continuous periodic signals
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- Properties of discrete FS
- **□** Fourier series and LTI systems



#### <u>Recall</u>



- $\square$  Decompose x(t) into linear combinations of basis signals, which should satisfy
  - It can be used to construct a broad and useful class of signals
  - The response of an LTI system to the basic signal is simple
- ☐ Complex exponentials are eigenfunctions of a LTI system
- $\square$  Can we represent x(t) as linear combinations of complex exponentials?



#### Linear combination of harmonically related complex exponentials

 $\Box$  Harmonically related complex exponentials (consider  $e^{st}$  with s purely imaginary)

$$\emptyset_k(t) = e^{jk\omega_0 t} = e^{jk(2\pi/T_0)t}, k = 0, \pm 1, \pm 2, ...$$

For any  $k \neq 0$ , fundamental frequency  $|k|\omega_0$ ; fundamental period  $\frac{2\pi}{|k|\omega_0} = \frac{T_0}{|k|}$ 

 $\square$  Linear combination of  $\emptyset_k(t)$  is also periodic

$$\sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{jk(2\pi/T_0)t} = x(t)$$

$$x(t) \text{ is periodic}$$



#### Linear combination of harmonically related complex exponentials

 $\square$  Can any x(t) (periodic) be decomposed as Linear combination of  $\emptyset_k(t)$ 

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$
? Yes for most periodic signals

Because  $e^{jk\omega_0t}$  are orthonormal:  $\langle e^{jk_1\omega_0t}, e^{jk_2\omega_0t} \rangle = 0$ 

$$\langle e^{jk_1\omega_0 t}, e^{jk_2\omega_0 t} \rangle = \frac{1}{T} \int_0^T e^{jk_1\omega_0 t} e^{-jk_2\omega_0 t} dt = \begin{cases} 1, k_1 = k_2 \\ 0, k_1 \neq k_2 \end{cases}$$



#### Linear combination of harmonically related complex exponentials

☐ Fourier Series representation

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$

- $\square$   $\omega_0$  is the fundamental frequency
- $\Box$  For  $a_k e^{jk\omega_0 t}$ 
  - > k = 0: DC component
  - $> k = \pm 1$ : fundamental (first harmonic) components
  - $> k = \pm N$ : Nth harmonic components



#### Linear combination of harmonically related complex exponentials

☐ An example

If 
$$x(t) = \sum_{k=-3}^{3} a_k e^{jk2\pi t}$$

And 
$$a_0 = 1$$
,  $a_1 = a_{-1} = 1/4$ ,  $a_2 = a_{-2} = 1/2$ ,  $a_3 = a_{-3} = 1/3$ 

$$x(t) = 1 + \frac{1}{4} \left( e^{j2\pi t} + e^{-j2\pi t} \right) + \frac{1}{2} \left( e^{j4\pi t} + e^{-j4\pi t} \right) + \frac{1}{3} \left( e^{j6\pi t} + e^{-j6\pi t} \right)$$

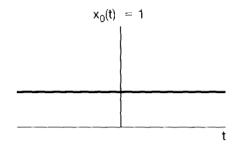
$$= 1 + \frac{1}{2}\cos 2\pi t + \cos 4\pi t + \frac{2}{3}\cos 6\pi t$$

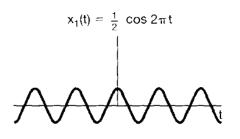


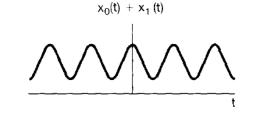
#### Linear combination of harmonically related complex exponentials

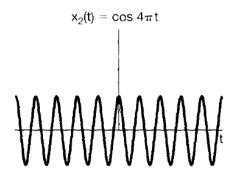
☐ An example

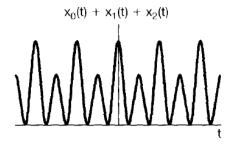
$$1 + \frac{1}{2}\cos 2\pi t + \cos 4\pi t + \frac{2}{3}\cos 6\pi t$$

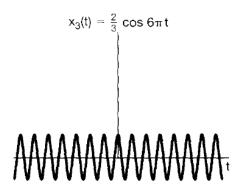


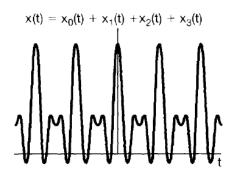
















https://www.youtube.com/watch?v=cUD1gMAl6W4





#### Linear combination of harmonically related complex exponentials

☐ Real signal

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \qquad x^*(t) = \sum_{k=-\infty}^{\infty} a_k^* e^{-jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_{-k}^* e^{jk\omega_0 t}$$

Real 
$$\Rightarrow x(t) = x^*(t) \Rightarrow a_k = a_{-k}^*$$
, or  $a_k^* = a_{-k}$  (Conjugate symmetry)

☐ Alternative form of Fourier Series for real signal

$$x(t) = a_0 + \sum_{k=1}^{\infty} \left[ a_k e^{jk\omega_0 t} + a_{-k} e^{-jk\omega_0 t} \right]$$

$$= a_0 + \sum_{k=1}^{\infty} 2\mathcal{R}e \left[ a_k e^{jk\omega_0 t} \right] = a_0 + 2\sum_{k=1}^{\infty} A_k \cos(k\omega_0 t + \theta_k)$$

$$a_k = A_k e^{j\theta_k}$$

#### Determine the Fourier Series Representation

$$\int_{0}^{T} x(t)e^{-jn\omega_{0}t}dt = \int_{0}^{T} \sum_{k=-\infty}^{\infty} a_{k}e^{jk\omega_{0}t} e^{-jn\omega_{0}t}dt$$

$$= \sum_{k=-\infty}^{\infty} a_{k} \left[ \int_{0}^{T} e^{j(k-n)\omega_{0}t}dt \right] = Ta_{n}$$

$$\therefore a_n = \frac{1}{T} \int_0^T x(t) e^{-jn\omega_0 t} dt$$

$$a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$$



#### Fourier Series pair

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t}$$
 Synthesis equation

$$a_k = \frac{1}{T} \int_T x(t)e^{-jk\omega_0 t} dt$$
 Analysis equation

 $\square$   $a_k$ : Fourier Series coefficients or spectral coefficients of x(t)

$$a_0 = \frac{1}{T} \int_T x(t) dt$$



#### Determine the Fourier Series Representation

$$x(t) = \sin \omega_0 t$$

$$\sin \omega_0 t = \frac{1}{2j} e^{j\omega_0 t} - \frac{1}{2j} e^{-j\omega_0 t}$$

$$\therefore a_1 = \frac{1}{2j} \qquad a_{-1} = -\frac{1}{2j} \qquad a_k = 0, \text{ for } k \neq \pm 1$$



#### Determine the Fourier Series Representation

$$x(t) = 1 + \sin \omega_0 t + 2 \cos \omega_0 t + \cos \left( 2\omega_0 t + \frac{\pi}{4} \right)$$

$$x(t) = 1 + \frac{1}{2j} \left[ e^{j\omega_0 t} - e^{-j\omega_0 t} \right] + \left[ e^{j\omega_0 t} + e^{-j\omega_0 t} \right]$$

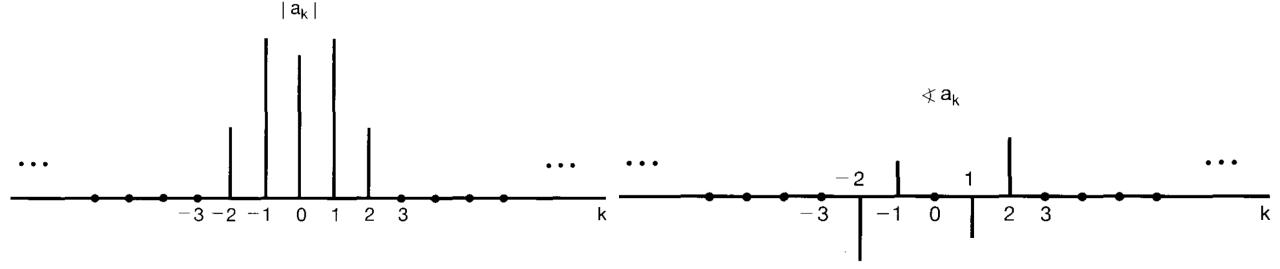
$$+ \frac{1}{2} \left( e^{j(2\omega_0 t + \pi/4)} + e^{-j(2\omega_0 t + \pi/4)} \right)$$



#### Determine the Fourier Series Representation

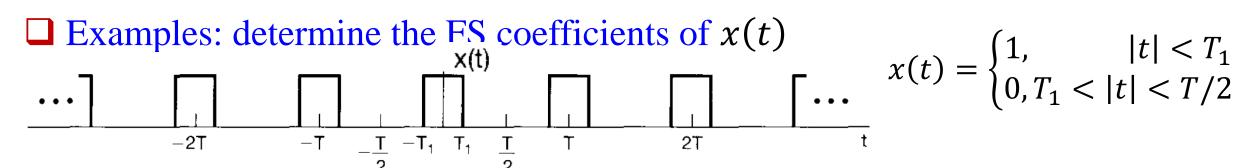
$$x(t) = 1 + \left(1 + \frac{1}{2j}\right)e^{j\omega_0 t} + \left(1 - \frac{1}{2j}\right)e^{-j\omega_0 t} + \frac{1}{2}e^{j\pi/4}e^{j2\omega_0 t} + \frac{1}{2}e^{-j\pi/4}e^{-j2\omega_0 t}$$

$$a_0 \quad a_1 \quad a_{-1} \quad a_2 \quad a_{-2}$$





#### Determine the Fourier Series Representation



$$a_0 = \frac{1}{T} \int_{-T/2}^{T/2} x(t) dt = \frac{1}{T} \int_{-T_1}^{T_1} 1 dt = \frac{T_1}{T}$$

$$a_{k} = \frac{1}{T} \int_{-T_{1}}^{T_{1}} e^{-jk\omega_{0}t} dt = -\frac{1}{jk\omega_{0}T} e^{-jk\omega_{0}t} \Big|_{-T_{1}}^{T_{1}} = \frac{2}{k\omega_{0}T} \left[ \frac{e^{jk\omega_{0}T_{1}} - e^{-jk\omega_{0}T_{1}}}{2j} \right]$$

$$= \frac{2\sin(k\omega_0 T_1)}{k\omega_0 T} = \frac{\sin(k\omega_0 T_1)}{k\pi} = \frac{2T_1}{T} \frac{\sin(k\omega_0 T_1)}{k\omega_0 T_1} = \frac{2\pi}{k\omega_0 T_1}$$

$$= \frac{2\sin(k\omega_0 T_1)}{k\omega_0 T_1} = \frac{\sin(k\omega_0 T_1)}{k\pi} = \frac{2T_1}{T} \frac{\sin(k\omega_0 T_1)}{k\omega_0 T_1}$$

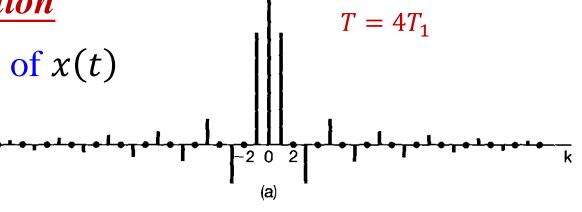
$$= \frac{2\sin(k\omega_0 T_1)}{k\omega_0 T_1} = \frac{\sin(k\omega_0 T_1)}{k\pi} = \frac{2\pi}{T} \frac{\sin(k\omega_0 T_1)}{k\omega_0 T_1}$$

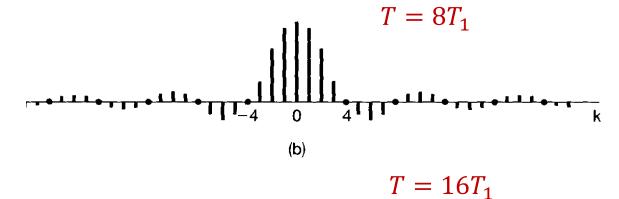


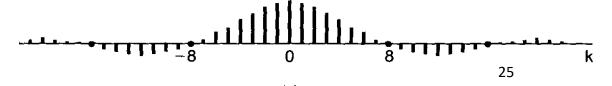
#### Determine the Fourier Series Representation

$$a_k = \frac{2\sin(k\omega_0 T_1)}{k\omega_0 T} = \frac{\sin(k\omega_0 T_1)}{k\pi}$$

$$=\frac{2T_1}{T}\frac{\sin(k\omega_0T_1)}{k\omega_0T_1}, k\neq 0$$







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- Properties of discrete FS
- **□** Fourier series and LTI systems



#### **History**

- Using "trigonometric sum" to describe periodic signal can be tracked back to Babylonians who predicted astronomical events similarly.
- □ L. Euler (in 1748) and Bernoulli (in 1753) used the "normal mode" concept to describe the motion of a vibrating string; though JL Lagrange strongly criticized this concept.
- □ Fourier (in 1807) had found series of harmonically related sinusoids to be useful to describe the temperature distribution through body, and he claimed "any" periodic signal can be represented by such series.
- ☐ Dirichlet (in 1829) provide a precise condition under which a periodic signal can be represented by a Fourier series.



Jean Baptiste Joseph Fourier March 21 1768 - May 16 1830 Born Auxerre, France. Died Paris, France.



#### Convergence problem

- $\square$  Approximate periodic signal x(t) by  $x_N(t) = \sum_{k=-N}^N a_k e^{jk\omega_0 t}$
- ☐ How good the approximation is?

$$e_N(t) = x(t) - x_N(t) = x(t) - \sum_{k=-N}^{N} a_k e^{jk\omega_0 t}$$
  $E_N = \int_T |e_N(t)|^2 dt$ 

- When  $a_k = \frac{1}{T} \int_T x(t) e^{-jk\omega_0 t} dt$ ,  $E_N$  is minimized;
- If  $x_N(t)$  can be expressed as  $\sum_{k=-N}^N a_k e^{jk\omega_0 t}$ ,  $N \to \infty \Rightarrow E_N \to 0$
- ☐ Problem:
  - $a_k$  may be infinite

**Convergence problem!** 

•  $N \to \infty$ ,  $x_N(t)$  may be infinite



#### Two different classes of conditions

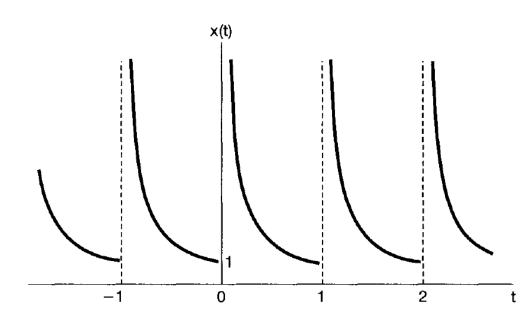
- □ Condition 1: Finite energy condition

  If  $\int_T |x(t)|^2 dt < \infty$ , x(t) can be represented by a FS
  - Guarantees no energy in their difference; FS is not equal to x(t)
- ☐ Condition 2: Dirichlet condition
  - (1) Absolutely integrable  $\int_T |x(t)| dt < \infty$

An example: a periodic signal

$$x(t) = \frac{1}{t}, 0 < t \le 1$$

is not absolutely integrable.



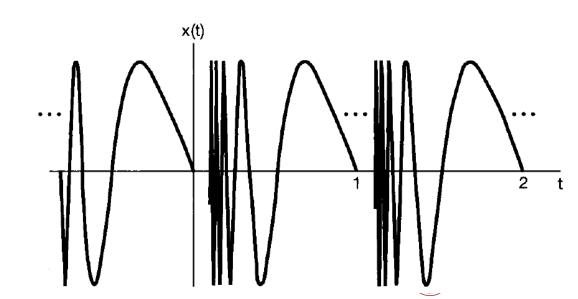
#### Two different classes of conditions

- ☐ Condition 2: Dirichlet condition
  - (2) In any finite interval of time, x(t) is of bounded variation; finite maxima and minima in one period

An example: a periodic signal

$$x(t) = \sin\left(\frac{2\pi}{t}\right), 0 < t \le 1$$

meets (1) but not (2).

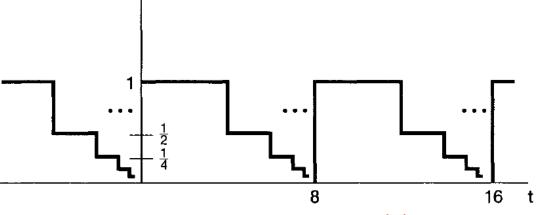


#### Two different classes of conditions

- ☐ Condition 2: Dirichlet condition
  - (3) In any finite interval of time, only a finite number of finite discontinuities

An example: a periodic signal meets (1) and (2) but not (3).

- Dirichlet condition guarantees x(t) equals its Fourier Series representation, except for discontinuous points.
- Three examples are pathological in nature and do not typically arise in practical contexts.



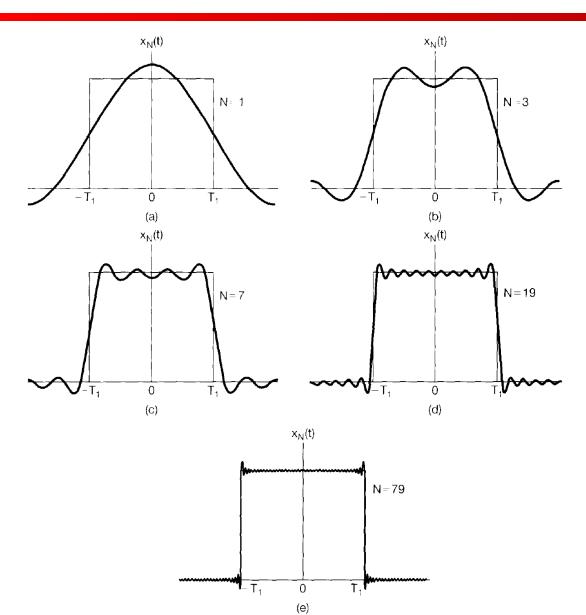
x(t)

#### **Example**

 $\square x(t)$  is a square wave

$$x_N(t) = \sum_{k=-N}^{N} a_k e^{jk\omega_0 t}$$

$$\lim_{N\to\infty} x_N(t_1) = x(t_1)$$



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### Properties of continuous-time FS

☐ Use the notation

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k$$

to signify the paring of a periodic signal with its FS coefficients.

 $\square$  Linearity: if x(t) and y(t) are periodic signals with the same period T

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k$$

$$y(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} b_k \Rightarrow z(t) = Ax(t) + By(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} c_k = Aa_k + Bb_k$$



### Properties of continuous-time FS

☐ Time shifting

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \implies x(t-t_0) \stackrel{\mathcal{FS}}{\longleftrightarrow} e^{-jk\omega_0 t_0} a_k$$

Proof
$$\frac{1}{T} \int_{T} x(t - t_{0}) e^{-jk\omega_{0}t} dt = \frac{1}{T} \int_{T} x(\tau) e^{-jk\omega_{0}(\tau + t_{0})} d\tau$$

$$= e^{-jk\omega_{0}t_{0}} \frac{1}{T} \int_{T} x(\tau) e^{-jk\omega_{0}\tau} d\tau$$

$$= e^{-jk\omega_{0}t_{0}} a_{k}$$



### Properties of continuous-time FS

☐ Time reversal

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \implies y(t) = x(-t) \stackrel{\mathcal{FS}}{\longleftrightarrow} b_k = a_{-k}$$

Proof

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \Longrightarrow x(-t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0(-t)} = \sum_{k=-\infty}^{\infty} a_k e^{j(-k)\omega_0 t}$$
$$= \sum_{m=-\infty}^{\infty} a_{-m} e^{jm\omega_0 t}$$

 $\square$  If x(t) even,  $a_{-k} = a_k$ , if x(t) odd,  $a_{-k} = -a_k$ 



☐ Time scaling

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \implies y(t) = x(\alpha t) \stackrel{\mathcal{FS}}{\longleftrightarrow} b_k = a_k$$

Proof

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \Longrightarrow x(\alpha t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 \alpha t} = \sum_{k=-\infty}^{\infty} a_k e^{jk(\alpha \omega_0)t}$$

FS coefficients the same, but fundamental frequency changed.



#### ■ Multiplication

$$\begin{array}{ccc}
x(t) & \stackrel{\mathcal{FS}}{\longleftrightarrow} & a_k \\
y(t) & \stackrel{\mathcal{FS}}{\longleftrightarrow} & b_k
\end{array} \implies z(t) = x(t)y(t) & \stackrel{\mathcal{FS}}{\longleftrightarrow} & h_k = \sum_{l=-\infty}^{\infty} a_l b_{k-l}$$

#### ☐ Proof

$$x(t)y(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \sum_{n=-\infty}^{\infty} b_n e^{jn\omega_0 t} = \sum_{k=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} a_k b_n e^{j(k+n)\omega_0 t}$$

$$= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} a_k b_{l-k} e^{jl\omega_0 t} = \sum_{l=-\infty}^{\infty} \sum_{k=-\infty}^{\infty} a_k b_{l-k} e^{jl\omega_0 t}$$

$$= \sum_{k=-\infty}^{\infty} \sum_{l=-\infty}^{\infty} a_k b_{l-k} e^{jl\omega_0 t}$$



□ Conjugation and conjugate symmetry

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \implies z(t) = x^*(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} b_k = a_{-k}^*$$

□ Proof

Proof
$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \quad \therefore x^*(t) = \sum_{k=-\infty}^{\infty} a_k^* e^{-jk\omega_0 t} = \sum_{m=-\infty}^{\infty} a_{-m}^* e^{jm\omega_0 t}$$

- $\square$  If x(t) real,  $a_k^* = a_{-k}$  (conjugate symmetry)  $\Rightarrow |a_k| = |a_{-k}|$ 
  - x(t) real and even  $(a_{-k} = a_k) \Rightarrow a_k = a_k^* \Rightarrow a_k$  real and even
  - x(t) real and odd  $(a_{-k} = -a_k) \Rightarrow a_k = -a_k^* \Rightarrow a_k$  pure imagery and odd
  - x(t) real and odd,  $a_0 = ?$



☐ Differentiation and Integration

$$x(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \Rightarrow \begin{cases} dx(t)/dt \stackrel{\mathcal{FS}}{\longleftrightarrow} jk\omega_0 a_k \\ \int_{-\infty}^t x(\tau)d\tau \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k/(jk\omega_0) \end{cases}$$

Proof

$$\frac{dx(t)}{dt} = \sum_{k=-\infty}^{\infty} a_k \frac{d(e^{jk\omega_0 t})}{dt} = \sum_{k=-\infty}^{\infty} a_k jk\omega_0 e^{jk\omega_0 t}$$

$$\int_{-\infty}^{t} x(\tau)d\tau = \sum_{k=-\infty}^{\infty} a_k \int_{-\infty}^{t} e^{jk\omega_0\tau}d\tau = \sum_{k=-\infty}^{\infty} \frac{a_k}{(jk\omega_0)} e^{jk\omega_0t}$$



☐ Frequency shifting

$$\chi(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_k \Rightarrow e^{jM\omega_0 t} \chi(t) \stackrel{\mathcal{FS}}{\longleftrightarrow} a_{k-M}$$

☐ Proof

$$e^{jM\omega_0 t}x(t) = e^{jM\omega_0 t} \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} = \sum_{k=-\infty}^{\infty} a_k e^{j(k+M)\omega_0 t}$$

$$k + M = l = \sum_{l=-\infty}^{\infty} a_{l-M} e^{jl\omega_0 t}$$



Periodic convolution

$$\begin{array}{ccc}
x(t) & \stackrel{\mathcal{FS}}{\longleftrightarrow} & a_k \\
y(t) & \stackrel{\mathcal{FS}}{\longleftrightarrow} & b_k
\end{array} \implies \int_{T} x(\tau)y(t-\tau)d\tau & \stackrel{\mathcal{FS}}{\longleftrightarrow} & Ta_k b_k$$

☐ Proof

$$\int_{T} x(\tau)y(t-\tau)d\tau = \int_{T} \sum_{k=-\infty}^{\infty} a_{k}e^{jk\omega_{0}\tau} \sum_{n=-\infty}^{\infty} b_{n}e^{jn\omega_{0}(t-\tau)}d\tau$$

$$= \int_{T} \sum_{k=-\infty}^{\infty} \sum_{n=-\infty}^{\infty} a_{k} e^{jk\omega_{0}\tau} b_{n} e^{-jn\omega_{0}\tau} e^{jn\omega_{0}t} d\tau$$

☐ Parseval's relation

$$\frac{1}{T} \int_{T} |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k|^2$$

Proof

$$\frac{1}{T} \int_{T} |x(t)|^{2} dt = \frac{1}{T} \int_{T} x(t) x^{*}(t) dt = \frac{1}{T} \int_{T} x(t) \sum_{k=-\infty}^{\infty} a_{k}^{*} e^{-jk\omega_{0}t} dt$$

$$= \sum_{k=-\infty}^{\infty} a_{k}^{*} \frac{1}{T} \int_{T} x(t) e^{-jk\omega_{0}t} dt$$

$$= \sum_{k=-\infty}^{\infty} a_{k}^{*} a_{k} = \sum_{k=-\infty}^{\infty} |a_{k}|^{2}$$



☐ Parseval's relation

$$\frac{1}{T} \int_{T} |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |a_k|^2$$

$$\frac{1}{T} \int_{T} |a_{k}e^{jk\omega_{0}t}|^{2} dt = \frac{1}{T} \int_{T} |a_{k}|^{2} dt = |a_{k}|^{2}$$

- $\square |a_k|^2$  is the average power in the kth harmonic component of x(t)
- $\Box$  Total average power in x(t) equals the sum of the average powers in all of its harmonic components



## **Properties**

**Property** 



		$x(t)$ Periodic with period T and $y(t)$ fundamental frequency $\omega_0 = 2\pi/T$	$egin{aligned} a_k \ b_k \end{aligned}$
Linearity	3.5.1	Ax(t) + By(t)	$Aa_k + Bb_k$
Time Shifting	3.5.2	$x(t-t_0)$	$a_k e^{-jk\omega_0 t_0} = a_k e^{-jk(2\pi/T)t_0}$
Frequency Shifting		$e^{jM\omega_0t}=e^{jM(2\pi/T)t}x(t)$	$a_{k-M}$
Conjugation	3.5.6	$x^*(t)$	$a_{-k}^*$
Time Reversal	3.5.3	x(-t)	$a_{-k}$
Time Scaling	3.5.4	$x(\alpha t), \alpha > 0$ (periodic with period $T/\alpha$ )	$a_k$
Periodic Convolution		$\int_T x(\tau)y(t-\tau)d\tau$	$Ta_kb_k$
Multiplication	3.5.5	x(t)y(t)	$\sum_{l=-\infty}^{+\infty} a_l b_{k-l}$
Differentiation		$\frac{dx(t)}{dt}$	$jk\omega_0 a_k = jk\frac{2\pi}{T}a_k$
Integration		$\int_{-\infty}^{t} x(t) dt $ (finite valued and periodic only if $a_0 = 0$ )	$\left(\frac{1}{jk\omega_0}\right)a_k = \left(\frac{1}{jk(2\pi/T)}\right)a_k$
			$\begin{cases} a_k = a_{-k}^* \\ \Re e\{a_k\} = \Re e\{a_{-k}\} \end{cases}$
Conjugate Symmetry for Real Signals	3.5.6	x(t) real	$\begin{cases} \mathfrak{G}m\{a_k\} = -\mathfrak{G}m\{a_{-k}\} \\  a_k  =  a_{-k}  \\ \not< a_k = -\not< a_{-k} \end{cases}$
Real and Even Signals	3.5.6	x(t) real and even	$a_k$ real and even
Real and Odd Signals	3.5.6	x(t) real and odd	$a_k$ purely imaginary and odd
Even-Odd Decomposition			$\Re\{a_k\}$
of Real Signals		$\begin{cases} x_e(t) = \mathcal{E}v\{x(t)\} & [x(t) \text{ real}] \\ x_o(t) = \mathcal{O}d\{x(t)\} & [x(t) \text{ real}] \end{cases}$	$j\mathfrak{Gm}\{a_k\}$

Periodic Signal

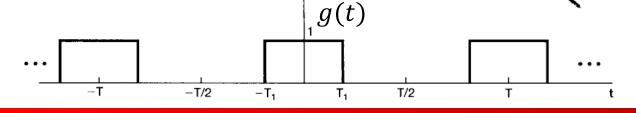
Section

**Fourier Series Coefficients** 

Parseval's Relation for Periodic Signals

$$\frac{1}{T}\int_{T}|x(t)|^{2}dt = \sum_{k=-\infty}^{+\infty}|a_{k}|^{2}$$

## Properties of col ...

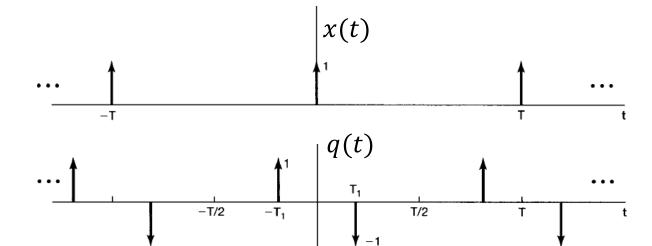




#### **□** Solution

• Let  $x(t) = \sum_{k=-\infty}^{\infty} \delta(t - kT)$ 

$$a_k = \frac{1}{T} \int_{-T/2}^{T/2} \delta(t) e^{-jk\omega_0 t} dt = \frac{1}{T}$$



• Let 
$$q(t) = x(t + T_1) - x(t - T_1)$$

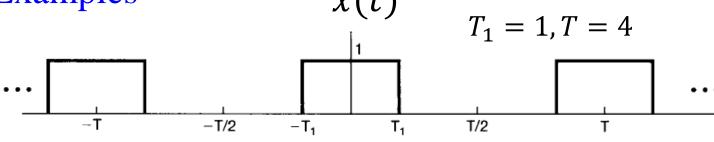
$$b_k = e^{jk\omega_0 T_1} a_k - e^{-jk\omega_0 T_1} a_k = \frac{1}{T} \left( e^{jk\omega_0 T_1} - e^{-jk\omega_0 T_1} \right) = \frac{2j\sin(k\omega_0 T_1)}{T}$$

• 
$$g(t) = \int_{-\infty}^{t} q(\tau) d\tau$$

$$\therefore c_k = \frac{b_k}{jk\omega_0} = \frac{2j\sin(k\omega_0 T_1)}{jk\omega_0 T} = \frac{\sin(k\omega_0 T_1)}{k\pi}, k \neq 0$$







$$a_k = \frac{\sin(k\omega_0 T_1)}{k\pi}, k \neq 0$$

$$= \frac{\sin(k\pi/2)}{k\pi}, k \neq 0$$

$$g(t) = x(t-1) - 1/2$$

FS coefficients of g(t)?

#### **□** Solution

$$x(t-1) \stackrel{\mathcal{FS}}{\leftrightarrow} e^{-jk\omega_0 t_0} a_k = e^{-jk\pi/2} a_k, k \neq 0$$

$$-1/2 \stackrel{\mathcal{FS}}{\leftrightarrow} \begin{cases} 0, k \neq 0 \\ -\frac{1}{2}, k = 0 \end{cases} \quad \therefore x(t-1) - 1/2 \stackrel{\mathcal{FS}}{\leftrightarrow} \begin{cases} e^{-jk\pi/2} a_k, k \neq 0 \\ a_0 - \frac{1}{2}, k = 0 \end{cases}$$



#### ■ Examples

Given a signal x(t) with the following facts, determine x(t)

- 1. x(t) is real;
- 2. x(t) is periodic with T=4 and FS coefficients  $a_k = 0$  for  $|\mathbf{k}| > 1$ ;
- 3. A signal with FS coefficients  $b_k = e^{-j\pi k/2}a_{-k}$  is odd;
- 4.  $\frac{1}{4} \int_4 |x(t)|^2 dt = \frac{1}{2}$ .

#### □ Solution

- From 2,  $x(t) = a_0 + a_1 e^{j\pi t/2} + a_{-1} e^{-j\pi t/2}$
- $b_k = e^{-j\pi k/2}a_{-k}$  corresponds to the signal x(-t+1), which is real and odd
- $\frac{1}{4} \int_4 |x(t)|^2 dt = \frac{1}{4} \int_4 |x(-t+1)|^2 dt = \sum_{k=-\infty}^{\infty} |b_k|^2 = |b_0|^2 + |b_1|^2 + |b_{-1}|^2 = \frac{1}{2}$
- x(-t+1) is real and odd  $\Rightarrow b_k = -b_{-k} \Rightarrow b_0 = 0, b_1 = -b_{-1} = \frac{j}{2}$  or  $-\frac{j}{2}$
- $a_0 = 0$ ,  $a_1 = -1/2$ ,  $a_{-1} = 1/2$



# Chapter 3: Fourier Series Representation of Periodic signals

- ☐ The response of LTI systems to complex exponentials
- **□** Fourier series representation of continuous periodic signals
- **□** Convergence of the Fourier series
- **□** Properties of continuous-time Fourier series
- **□** Fourier series representation of discrete –time periodic signals
- Properties of discrete FS
- **□** Fourier series and LTI systems



## Fourier series representation of D-T periodic signals

#### Linear combination of harmonically related complex exponentials

☐ Harmonically related complex exponentials

$$\emptyset_k[n] = e^{jk(2\pi/N)n}, k = 0, \pm 1, \pm 2, \dots$$

- Fundamental frequency  $|k|(\frac{2\pi}{N})$
- Only N distinct signals in  $\emptyset_k[n]$ , since  $\emptyset_k[n] = \emptyset_{k+rN}[n]$
- $\square$  Linear combination of  $\emptyset_k[n]$  is also periodic

$$x[n] = \sum_{k=\langle N \rangle} a_k \emptyset_k[n] = \sum_{k=\langle N \rangle} a_k e^{jk\omega_0 n} = \sum_{k=\langle N \rangle} a_k e^{jk(2\pi/N)n}$$

 $\square \sum_{k=\langle N \rangle} a_k e^{jk(2\pi/N)n}$ : Discrete-Time Fourier Series;  $a_k$ : Fourier Series coefficients

## Fourier series representation of D-T periodic signals

#### Determine the Fourier Series Representation

$$\sum_{n=\langle N \rangle} x[n] e^{-jr(2\pi/N)n} = \sum_{n=\langle N \rangle} \sum_{k=\langle N \rangle} a_k e^{jk(2\pi/N)n} e^{-jr(2\pi/N)n}$$

$$= \begin{cases} N, k = r \\ 0, k \neq r \end{cases} = N\delta[k-r]$$

$$= \sum_{k=\langle N \rangle} a_k \sum_{n=\langle N \rangle} e^{j(k-r)(2\pi/N)n} = Na_r$$

$$\therefore a_k = \frac{1}{N} \sum_{n = \langle N \rangle} x[n] e^{-jk(2\pi/N)n}$$



#### Fourier series representation of D-T periodic signals

#### Determine the Fourier Series Representation

Discrete Fourier series pair

$$a_k = \frac{1}{N} \sum_{n=\langle N \rangle} x[n] e^{-jk(2\pi/N)n}$$
 Analysis equation;  $a_k$ : Fourier Series coefficients or spectral coefficients

$$x[n] = \sum_{k=\langle N \rangle}^{n} a_k e^{jk(2\pi/N)n}$$

Synthesis equation; Fourier Series (Finite)

$$\Box \ a_k \text{ is periodic} \quad x[n] = \sum_{k = \langle N \rangle} a_k \emptyset_k[n] = a_0 \emptyset_0[n] + a_1 \emptyset_1[n] + \dots + a_{N-1} \emptyset_{N-1}[n]$$

$$= a_1 \emptyset_1[n] + a_2 \emptyset_2[n] + \dots + a_N \emptyset_N[n]$$

$$= a_2 \emptyset_2[n] + a_3 \emptyset_3[n] + \dots + a_{N+1} \emptyset_{N+1}[n]$$

$$\vdots \ a_k = a_{k+rN}$$

#### Fourier series representation of D-T periodic signals

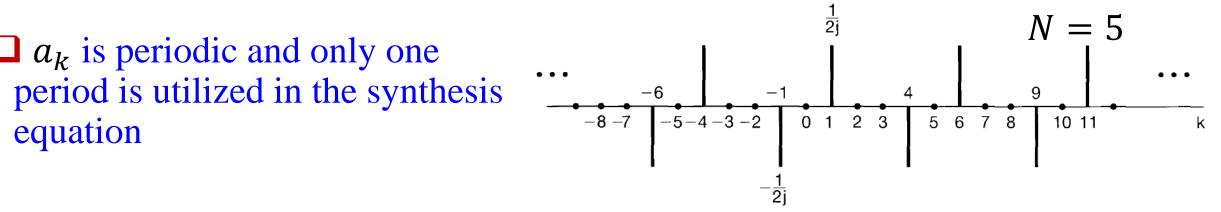
#### Determine the Fourier Series Representation

If  $\omega_0 = \frac{2\pi}{N}$ , x[n] is periodic with fundamental period of N.

$$x[n] = \sin \omega_0 n = \frac{1}{2j} e^{j(2\pi/N)n} - \frac{1}{2j} e^{j(2\pi/N)n}$$

: 
$$a_1 = \frac{1}{2j}$$
  $a_{-1} = -\frac{1}{2j}$   $a_k = 0$ , for  $k \neq \pm 1$  in one period

 $\square$   $a_k$  is periodic and only one



## Fourier series representation of D-T periodic signals

#### Determine the Fourier Series Representation

Examples 
$$x[n] = 1 + \sin(\frac{2\pi}{N})n + 3\cos(\frac{2\pi}{N})n + \cos(\frac{4\pi}{N}n + \frac{\pi}{2})$$
  
 $x[n] = 1 + \frac{1}{2j} \left[ e^{j(2\pi/N)n} - e^{-j(2\pi/N)n} \right] + \frac{3}{2} \left[ e^{j(2\pi/N)n} + e^{-j(2\pi/N)n} \right]$ 

$$+\frac{1}{2}\left(e^{j(4\pi n/N+\pi/2)}+e^{-j(4\pi n/N+\pi/2)}\right)$$

$$\therefore x[n] = 1 + \left(\frac{3}{2} + \frac{1}{2j}\right) e^{j(2\pi/N)n} + \left(\frac{3}{2} - \frac{1}{2j}\right) e^{-j(2\pi/N)n}$$

$$+\frac{1}{2}e^{j\pi/2}e^{j2(2\pi/N)n} + \frac{1}{2}e^{-j\pi/2}e^{-j2(2\pi/N)n}$$

# Fourier series representation of D-T periodic signals

#### Linear combination of harmonically related complex exponentials

$$a_k = a_{-k}^*$$
, or  $a_k^* = a_{-k}$ 

$$x[n] = \sum_{k=\langle N \rangle} a_k e^{jk(2\pi/N)n}$$

$$x^*[n] = \sum_{k=\langle N \rangle} a_k^* e^{-jk(2\pi/N)n} = \sum_{k=\langle N \rangle} a_{-k}^* e^{jk(2\pi/N)n}$$

$$x[n] = x^*[n] \implies a_k = a_{-k}^*$$

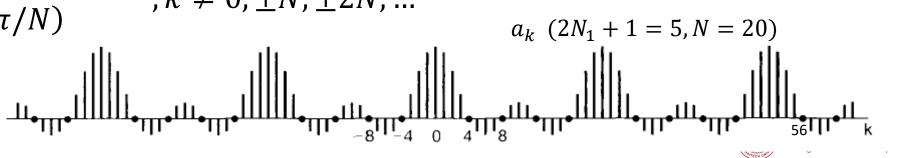


## Fourier series representation of D-T periodic signals

$$a_k = \frac{1}{N} \sum_{n=-N_1}^{N_1} x[n] e^{-jk(2\pi/N)n} = \frac{1}{N} \sum_{n=-N_1}^{N_1} e^{-jk(2\pi/N)n}$$

$$m = n + N_1 = \frac{1}{N} \sum_{m=0}^{2N_1} e^{-jk(2\pi/N)(m-N_1)} = \frac{1}{N} e^{jk(2\pi/N)N_1} \sum_{m=0}^{2N_1} e^{-jk(2\pi/N)m}$$

$$= \begin{cases} \frac{2N_1 + 1}{N}, k = 0, \pm N, \pm 2N, \dots \\ \frac{1}{N} \frac{\sin[2k\pi(N_1 + 1/2)/N]}{\sin(k\pi/N)}, k \neq 0, \pm N, \pm 2N, \dots \end{cases}$$



## Fourier series representatio signals

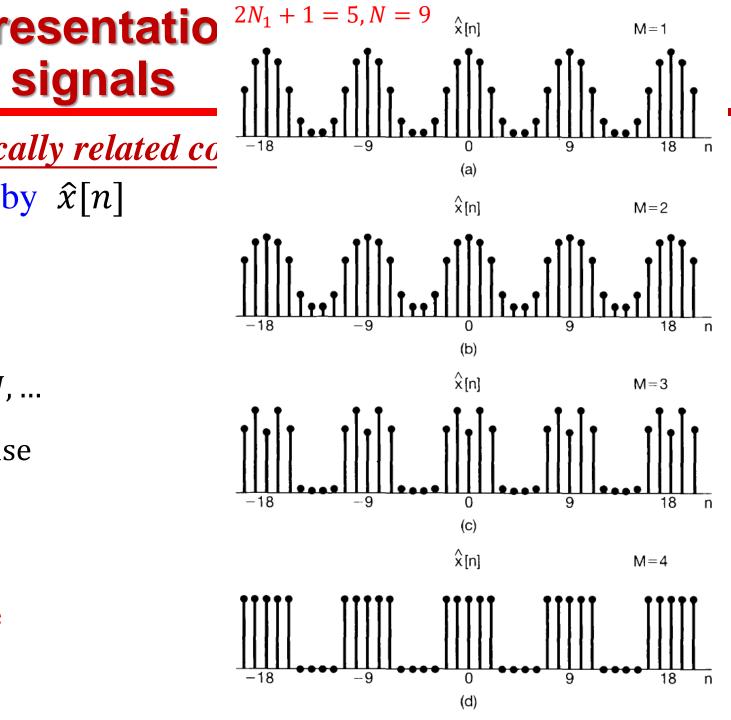
#### Linear combination of harmonically related co

 $\square$  Approximate a discrete square by  $\hat{x}[n]$ 

$$\hat{x}[n] = \sum_{k=-M}^{M} a_k e^{jk(2\pi/N)n}$$

With 
$$a_k = \begin{cases} \frac{2N_1+1}{N}, k = 0, \pm N, \pm 2N, \dots \\ \frac{1}{N} \frac{\sin[2k\pi(N_1+1/2)/N]}{\sin(k\pi/N)}, \text{ else} \end{cases}$$

- $\square$  For M=4,  $\widehat{x}[n] = x[n]$
- ☐ No convergence issues for the discrete—time Fourier series!



# Chapter 3: Fourier Series Representation of Periodic signals

- ☐ The response of LTI systems to complex exponentials
- **□** Fourier series representation of continuous periodic signals
- **□** Convergence of the Fourier series
- **□** Properties of continuous-time Fourier series
- **□** Fourier series representation of discrete –time periodic signals
- **□** Properties of discrete FS
- **□** Fourier series and LTI systems



## **Properties**

г -	$\mathcal{FS}$	1_

■ Multiplication

$$x[n]y[n] \xrightarrow{\mathcal{FS}} \sum_{l=\langle N \rangle} a_l b_{k-l}$$

☐ First difference

$$x[n] - x[n-1] \xrightarrow{\mathcal{FS}} (1 - e^{-jk(2\pi/N)}) a_{k_{\text{onjugate Symmetry for}}}$$

☐ Parseval's relation

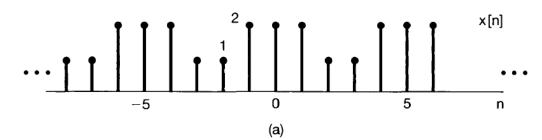
$$\frac{1}{N} \sum_{n=\langle N \rangle} |x[n]|^2 = \sum_{k=\langle N \rangle} |a_k|^2$$

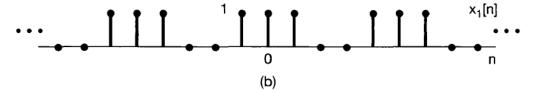
#### Periodic Signal Fourier Series Coefficients **Property** x[n]) Periodic with period N and $a_k$ ) Periodic with y[n] | fundamental frequency $\omega_0 = 2\pi/N$ $b_k \mid \text{period } N$ Linearity Ax[n] + By[n] $Aa_k + Bb_k$ $a_{\nu}e^{-jk(2\pi/\hat{N})n_0}$ Time Shifting $x[n-n_0]$ $e^{jM(2\pi/N)n}x[n]$ Frequency Shifting Conjugation $x^*[n]$ $a_{-k}^*$ Time Reversal x[-n] $a_{-k}$ $x_{(m)}[n] = \begin{cases} x[n/m], & \text{if } n \text{ is a multiple of } m \\ 0, & \text{if } n \text{ is a multiple of } m \end{cases}$ $\frac{1}{m}a_k$ (viewed as periodic) with period mNTime Scaling if n is not a multiple of m(periodic with period mN) $\sum_{r=\langle N\rangle} x[r]y[n-r]$ Periodic Convolution $Na_kb_k$ $\sum_{l=\langle N\rangle}a_lb_{k-l}$ Multiplication x[n]y[n] $(1 - e^{-jk(2\pi/N)})a_t$ First Difference x[n] - x[n-1]Running Sum $a_k = a_{-k}^*$ $\Re\{a_k\} = \Re\{a_{-k}\}$ $\mathfrak{I}m\{a_k\} = -\mathfrak{I}m\{a_{-k}\}$ x[n] real Real Signals $|a_k| = |a_{-k}|$ Real and Even Signals x[n] real and even $a_{\nu}$ real and even Real and Odd Signals x[n] real and odd $a_k$ purely imaginary and odd Even-Odd Decomposition $x_{\rho}[n] = \mathcal{E}\nu\{x[n]\}$ $\Re\{a_k\}$ [x[n] real] of Real Signals $x_o[n] = Od\{x[n]\}$ [x[n] real] $j\mathfrak{I}m\{a_k\}$ Parseval's Relation for Periodic Signals

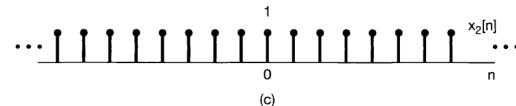
 $\frac{1}{N} \sum_{n=(N)} |x[n]|^2 = \sum_{k=(N)} |a_k|^2$ 

### Properties of discrete-time FS

 $\square$   $x_1[n]$  is a square wave with N=5 and  $N_1 = 1$ 







$$\square$$
 For  $x_2[n]$ 

$$c_{k} = \begin{cases} 1, k = \pm N, \pm 2N, \dots \\ 0, & \text{else} \end{cases}$$

$$\therefore a_{k} = b_{k} + c_{k} = \begin{cases} \frac{8}{5}, & k = \pm 5, \pm 10, \dots \\ \frac{1}{5} \frac{\sin(3k\pi/5)}{\sin(k\pi/5)}, & \text{else} \end{cases}$$

## Properties of discrete-time FS

#### Examples

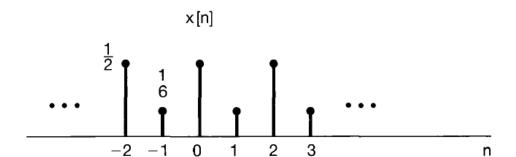
Suppose we are given the following facts about a sequence x[n]:

- 1. x[n] is periodic with period N = 6.
- **2.**  $\sum_{n=0}^{5} x[n] = 2$ .
- 3.  $\sum_{n=2}^{7} (-1)^n x[n] = 1$ .
- 4. x[n] has the minimum power per period among the set of signals satisfying the preceding three conditions.



- $\sum_{n=0}^{5} x[n] = 2 \Longrightarrow a_0 = \frac{1}{N} \sum_{n=\langle N \rangle} x[n] e^{-j0(2\pi/N)n} = 1/3.$
- $\sum_{n=2}^{7} (-1)^n x[n] = 1 \Longrightarrow \sum_{n=\langle N \rangle} x[n] e^{-j3(2\pi/N)n} = 1 \Longrightarrow a_3 = 1/6$
- from 4,  $a_1 = a_2 = a_4 = a_5 = 0$
- $x[n] = a_0 e^{-j0(2\pi/N)n} + a_3 e^{-j3(2\pi/N)n} = \frac{1}{3} + \frac{1}{6} e^{-j\pi n} = \frac{1}{3} + \frac{1}{6} (-1)^n$





# Chapter 3: Fourier Series Representation of Periodic signals

- ☐ The response of LTI systems to complex exponentials
- **□** Fourier series representation of continuous periodic signals
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- **□** Fourier series and LTI systems



Recall
$$e^{st} \longrightarrow \text{LTI} \longrightarrow H(s)e^{st} \qquad H(s) = \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau$$

$$z^{n} \longrightarrow \text{LTI} \longrightarrow H[z]z^{n} \qquad H[z] = \sum_{k=-\infty}^{\infty} h[k]z^{-k}$$

 $\square$  System functions: H(s) and H[z]

For periodic signal, CT Fourier Series (Ch3)

s pure imagery  $e^{st} o e^{j\omega t}$ For aperiodic signal, CT Fourier Transform (Ch4)

s complex number  $e^{st}$ Laplase Transform (Ch9)

For periodic signal, DT Fourier Series (Ch3)

For aperiodic signal, DT Fourier Transform (Ch5)

z complex number  $z^n$ Z-Transform (Ch10)

 $\square$  Frequency response for CT system:  $H(j\omega)$ 

$$H(s) = \int_{-\infty}^{\infty} h(\tau)e^{-s\tau}d\tau \quad \stackrel{s=j\omega}{\Longrightarrow} \quad H(j\omega) = \int_{-\infty}^{\infty} h(\tau)e^{-j\omega\tau}d\tau$$

$$e^{j\omega t}$$
 LTI  $H(j\omega)e^{j\omega t}$ 

$$x(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \longrightarrow \underbrace{\text{LTI}} y(t) = \sum_{k=-\infty}^{\infty} a_k H(j\omega) e^{jk\omega_0 t}$$

$$f(t) = \sum_{k=-\infty}^{\infty} a_k e^{jk\omega_0 t} \longrightarrow \underbrace{\text{LTI}} y(t) = \sum_{k=-\infty}^{\infty} a_k H(j\omega) e^{jk\omega_0 t}$$

$$b_k = a_k H(j\omega)$$



☐ Frequency response for CT system: example

$$x(t) = \sum_{k=-3}^{3} a_k e^{jk2\pi t}$$
  $(a_0 = 1, a_1 = a_{-1} = \frac{1}{4}, a_2 = a_{-2} = \frac{1}{2}, a_3 = a_{-3} = \frac{1}{3})$  is the input of a LTI system with  $h(t) = e^{-t}u(t)$ , determine  $y(t)$ 

Solution

$$y(t) = \sum_{k=-\infty}^{\infty} a_k H(j\omega) e^{jk2\pi t} \quad H(j\omega) = \int_0^{\infty} e^{-\tau} e^{-j\omega\tau} d\tau = \frac{1}{1+j\omega}$$

$$b_k = a_k H(j\omega) = a_k \frac{1}{1+jk2\pi} \qquad b_0 = 1 \cdot 1 = 1 \qquad b_1 = \frac{1}{4} \frac{1}{1+j2\pi} \quad b_{-1} = \frac{1}{4} \frac{1}{1-j2\pi}$$

$$b_2 = \frac{1}{2} \frac{1}{1 + j4\pi} \qquad b_{-2} = \frac{1}{2} \frac{1}{1 - j4\pi} \qquad b_3 = \frac{1}{3} \frac{1}{1 + j6\pi} \qquad b_{-3} = \frac{1}{3} \frac{1}{1 - j6\pi}$$



 $\Box$  Frequency response DT system:  $H(e^{j\omega})$ 

$$H[z] = \sum_{n = -\infty}^{\infty} h[k]z^{-n} \qquad \stackrel{z = e^{j\omega}}{\Longrightarrow} \qquad H(e^{j\omega}) = \sum_{n = -\infty}^{\infty} h[n]e^{-j\omega n}$$

$$e^{j\omega n} \longrightarrow \left\{ \qquad \text{LTI} \qquad \right\} \longrightarrow H(e^{j\omega})e^{j\omega n}$$

$$\longrightarrow \boxed{\text{LTI}} \qquad y[n] = \sum_{k=\langle N \rangle} a_k H(e^{j\omega}) e^{jk(2\pi/N)n}$$

$$b_k = a_k H(e^{j\omega})$$



☐ Frequency response DT system: example

$$h[n] = \alpha^n u[n], |\alpha| < 1$$

$$x[n] = \cos \frac{2\pi n}{N} \longrightarrow \boxed{\text{LTI}} \qquad y[n]?$$

**□** Solution

$$x[n] = \frac{1}{2}e^{j(2\pi/N)n} + \frac{1}{2}e^{-j(2\pi/N)n}$$

$$H(e^{j\omega}) = \sum_{n=-\infty}^{\infty} h[n]e^{-j\omega n} = \sum_{n=0}^{\infty} \alpha^n e^{-j\omega n} = \frac{1}{1 - \alpha e^{-j\omega}}$$

$$x[n] = \frac{1}{2} \left( \frac{1}{1 - \alpha e^{-j2\pi/N}} \right) e^{j(2\pi/N)n} + \frac{1}{2} \left( \frac{1}{1 - \alpha e^{j2\pi/N}} \right) e^{-j(2\pi/N)n}$$