

Section 6: Ionization

AE435

Spring 2018

1 Equilibrium Ionization

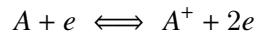
Here we derive the ionization fraction for an equilibrium gas. This results in what is called the Saha Equation.

Contents

1	Equilibrium Ionization	2
1.1	Detailed Balancing	3
1.2	Partition Functions	4
1.3	Saha Equation	8
1.4	Conditions for Equilibrium	11
1.4.1	Deviations from Equilibrium	13

1.1 Detailed Balancing

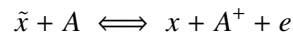
Assume **detailed balancing**, where every process is balanced by the reverse process:



An atom collision with an electron can knock off an electron on the atom creating an ion.

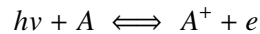
Examples of the sort of reactions involved in detailed balancing:

1. Inelastic collision with a 3-body recombination.



(A high-energy particle collides with an atom, knocking off an electron turning it into an ion)

2. Photoionization with radiative recombination



(A photon-particle collision ionizes the particle by knocking out an electron)

The overall reaction can be described by an equilibrium constant, a function of pressure and temperature:

$$K_n(P, T) = \frac{n_+ n_e}{n_A} \quad (2)$$

We can show from statistical mechanics that at equilibrium...

Boltzmann Distribution

$$\frac{N_i}{N} = \frac{n_i}{n} = \frac{e^{-\varepsilon_i/kT}}{\sum_i e^{-\varepsilon_i/kT}} \quad (3)$$

Where

N is the total population within the control volume

N_i is the population within the CV at a given state, i

ε_i is the energy state i

k is Boltzmann constant.

The Boltzmann Distribution is a probability distribution or frequency distribution of particles in a system over various possible states at equilibrium.

We will use this when moving onto partition functions.

1.2 Partition Functions

We can define...

We can define a **Total Partition Function** for a given species that is the product of a series of individual partial partition functions for all the different energy modes that could exist,

$$F = f^t f^e f^r f^v \quad (4)$$

Where for each energy mode of a given species:

f^t is the Translational Partial Partition Function

f^e is the Electronic Partial Partition Function

f^r is the Rotational Partial Partition Function

f^v is the Vibrational Partial Partition Function

We have a partial partition function.

A **Partial Partition Function** is given by

$$f = \sum_i \exp\left(-\frac{\varepsilon_i}{kT}\right) \quad (5)$$

This tells us how the energy is distributed across any available energy states. There are distributions amongst every form of energy within a system. As an example, the translational energy states of an atom can be different in one area of the system compared to the other. Jahn defines the partial partition function as the sum of accessible energy states, appropriately weighted by degeneracies and by the Boltzmann factors in the relative energies of those states. The Boltzmann factor is that exponential term we see in the Equation 5

Other ways of writing the partial partition function include:

$$f = \sum_i \exp\left(-\frac{\varepsilon_i - \varepsilon_o}{kT}\right) \quad (6)$$

where ε_o is the ground-state energy. (Note: this is a more general form as it indicates that energy is relative $\varepsilon_o = 0$ for atoms and $\varepsilon_o = \varepsilon_i$ for the ionization energy.

$$f = \sum_i g_i \exp\left(-\frac{\varepsilon_i - \varepsilon_o}{kT}\right) \quad (7)$$

includes the degeneracy g_i , making it useful for electronic states.

The partition function is a link between classical thermo and the microscopic behavior of gases (statistical kinetic theory). It is fairly easy to show that the total energy for N particles is:

$$E = NkT^2 \frac{\partial}{\partial T} \ln f \quad (8)$$

For instance...

The **Translational Partial Partition Function** is given on quantum mechanics basis, as Equation 9 for atoms of mass M_A :

$$f_A^t = \frac{(2\pi M_A kT)^{\frac{3}{2}}}{h^3} = KT^{\frac{3}{2}} \quad (9)$$

Using Equation 9, we can take the logarithmic derivative to find...

$$\frac{\partial}{\partial T} \ln f_A^t = \frac{1}{f_A^t} \frac{\partial}{\partial T} f_A^t \quad (10)$$

Such that

$$\frac{\partial}{\partial T} \ln f_A^t = \frac{1}{KT^{\frac{3}{2}}} \frac{3}{2} KT^{\frac{1}{2}} = \frac{3}{2} \frac{1}{T} \quad (11)$$

The total translational energy is then: (substituting equation 11 into equation 8)

$$E_{tr} = NkT^2 \frac{\partial}{\partial T} \ln f_A^t = NkT^2 \frac{3}{2} \frac{1}{T} = \frac{3}{2} NkT \quad (12)$$

which is the result we had before in Chapter 4.1.2. when we discussed Kinetic Theory and Internal Translational Energy.

We can also find that...

electronic partial partition function

$$f^e = \sum_{j=0}^n g_j \exp\left(-\frac{\varepsilon_j}{kT}\right) \quad (13)$$

Where

g_i degeneracy for state j
 $j = 0$ represents the ground state
 $j = n$ is the highest bound state

Note that this series converges rapidly, so we only to take the first few terms at low to moderate temperatures. In hot plasmas, may need to take many terms for an accurate solution. The reason why at low moderate temperatures we can simply take the first two terms is because the energy of each atomic state increases significantly with every orbital.

Now lets consider a species with negligible rotational and vibrational modes meaning we are only considering atoms or a monatomic species. We can define:

Total Atomic Partition Function

$$F_A = f_A^t f_A^e \quad (14)$$

Substituting in the partial partition functions,

$$F_A = \frac{(2\pi M_A kT)^{\frac{3}{2}}}{h^3} \sum_{j=0}^n g_j \exp\left(-\frac{\varepsilon_j}{kT}\right) \quad (15)$$

Likewise,

Total Ionic Partition Function

$$F_+ = f_+^t f_+^e \quad (16)$$

Substituting in the partial partition functions,

$$F_+ = \frac{(2\pi M_+ kT)^{\frac{3}{2}}}{h^3} \sum_{j=0}^n g_j \exp\left(-\frac{\bar{\varepsilon}_j^+}{kT}\right) \quad (17)$$

We need to make sure that we have consistent assumptions about the energy reference state. Doesn't matter what it is, but it has to be the same for all.

- For atoms, we use the ground state, $\varepsilon_j = 0$
- For ions, we need to add the ionization energy to reference back to atomic ground, thus...

Ionic Ground State

$$\varepsilon_j^+ = \bar{\varepsilon}_j^+ - \varepsilon_i \quad (18)$$

Where the following terms express an electronic energy level in the ion:

ε_j^+ = is referenced to the ionic ground state

$\bar{\varepsilon}_j^+$ = is referenced to the atomic ground state

ε_i = is the ionization energy (referenced to the atomic ground state)

Finally, we can define a partition function for free electrons:

Total Electron Partition Function

$$F_e = f_e^+ g_e \quad (19)$$

where $g_e = 2$ is the spin degeneracy (the only internal degree of freedom).

Substituting in the partial partition functions,

$$F_e = \frac{2 (2\pi M_e kT)^{\frac{3}{2}}}{h^3} \quad (20)$$

Electrons can't have electronic energy if they are free since electronic energy comes from an electric potential. There is also no rotational or vibrational energy. Free electrons only have translational energy.

So far, we have only looked at monatomic species. Now we are going to combine all of these together using Equation 2 in order to derive the SAHA equation.

1.3 Saha Equation

The equilibrium constant for ionization as seen from Equation 2 is

$$K_n = \frac{n_+ n_e}{n_A} \quad (21)$$

which can be expressed as the ratio of total partition functions,

$$K_n = \frac{F_+ F_e}{F_A} \quad (22)$$

Substituting in these partial partition functions, we get

$$K_n = \frac{2 (2\pi M_e kT)^{\frac{3}{2}}}{h^3} \left(\frac{f_+^e}{f_A^e} \right) \exp \left(- \frac{\varepsilon_i}{kT} \right) \quad (23)$$

Question: Why don't we have translational partition functions for atoms or ions?

Answer: The terms canceled because their masses are about the same. The only difference between the two is in the mass. Understanding that the masses are very similar, the equations are identical and cancel out.

In this equation

$$f_+^e = \sum_j^n g_j^+ \exp \left(- \frac{\varepsilon_j^+}{kT} \right) \quad (24)$$

is the electronic partial partition function for ions, and

$$f_A^e = \sum_j^n g_j \exp \left(- \frac{\varepsilon_j}{kT} \right) \quad (25)$$

is the electronic partial partition function for atoms.

This is the classic form of the Saha Equation from astrophysics (6.23).

An alternative form used often in gas dynamics is started by defining the ionization fraction

Ionization Fraction

$$\alpha = \frac{n_+}{n_A^+ n_+} = \frac{n_+}{n_o} \quad (26)$$

Where

n_+ is the ion density

n_A is the neutral density

n_o is the total density of heavy particles

If we assume quasi-neutrality meaning we have an equal number of positive and negative particles.

$$n_+ = n_e \quad (27)$$

Then we can write

$$1 + \alpha = \frac{n_a + n_+ + n_+}{n_o} = \frac{n_a + n_+ + n_e}{n_o} \quad (28)$$

Such that the pressure from all 3 species (ions, atoms, electrons) is:

$$P = (n_a + n_+ + n_e) kT = (1 + \alpha) n_o kT \quad (29)$$

Substituting this into the equilibrium constant from Equation 21 we get:

SAHA Equation

$$\frac{\alpha^2}{1 - \alpha} = \beta(P, T) \quad (30)$$

where

$$\beta(P, T) = \frac{2 (2\pi M_e)^{\frac{3}{2}} (kT)^{\frac{5}{2}}}{P h^3} \left(\frac{f_+^e}{f_A^e} \right) \exp \left(- \frac{\varepsilon_i}{kT} \right) \quad (31)$$

The amount of ionization is dependent on temperature and pressure. A higher temperature means more ionization, everything has more energy meaning that there will be higher translational energy resulting in more collisions resulting in more electrons being knocked off atoms. Higher pressure equates to lower ionization because we get more three-body recombinations.

We can show a plot for argon:

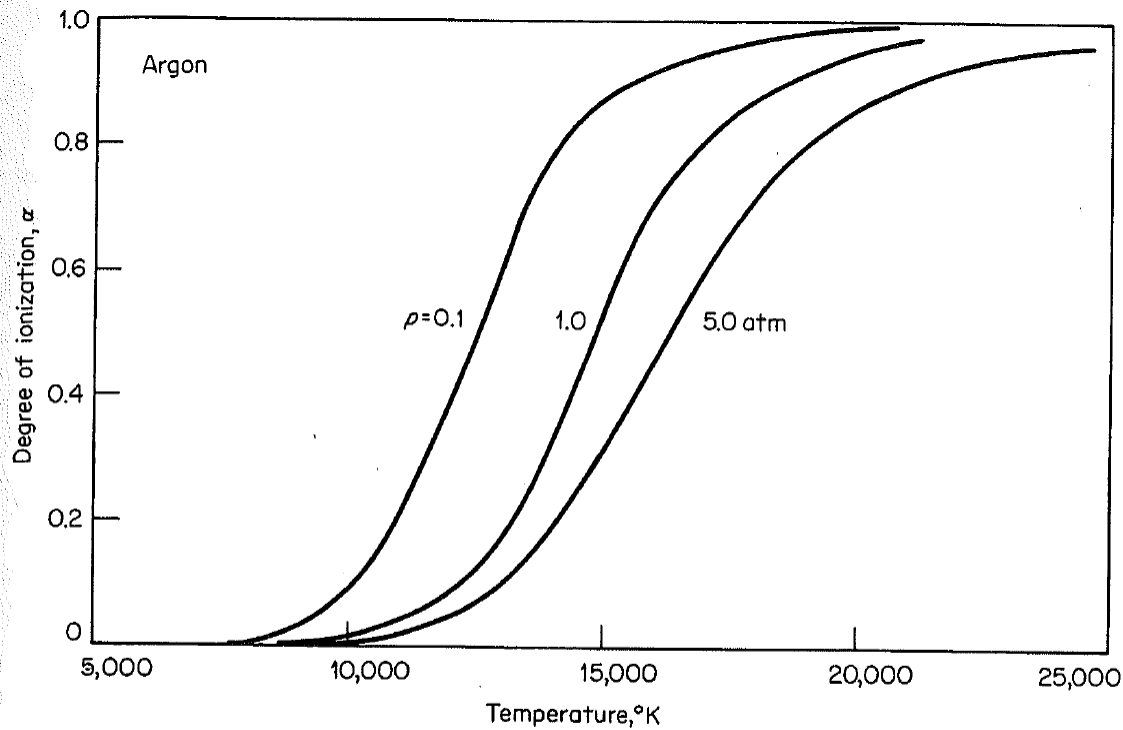


Fig. 3-4 Dependence of ionization of argon on temperature and pressure, from Saha equation. (From K. S. Drellishak, C. K. Knopp, and A. B. Cambel, *Partition Functions and Thermodynamic Properties of Argon Plasma*, *Phys. Fluids*, ser. 6, vol. 9, p. 1280, 1963.)

Note that collisional recombination drives the ionization fraction down at higher pressures.

Jahn gives a more complex analysis, required for multiple species and internal degrees of freedom. Main points:

- Vibration and rotation require their associated partition functions
- Additional reactions (dissociation, etc) can be handled the same way
- Multiple ionization (high temp, low density) adds partition function and number of reactions

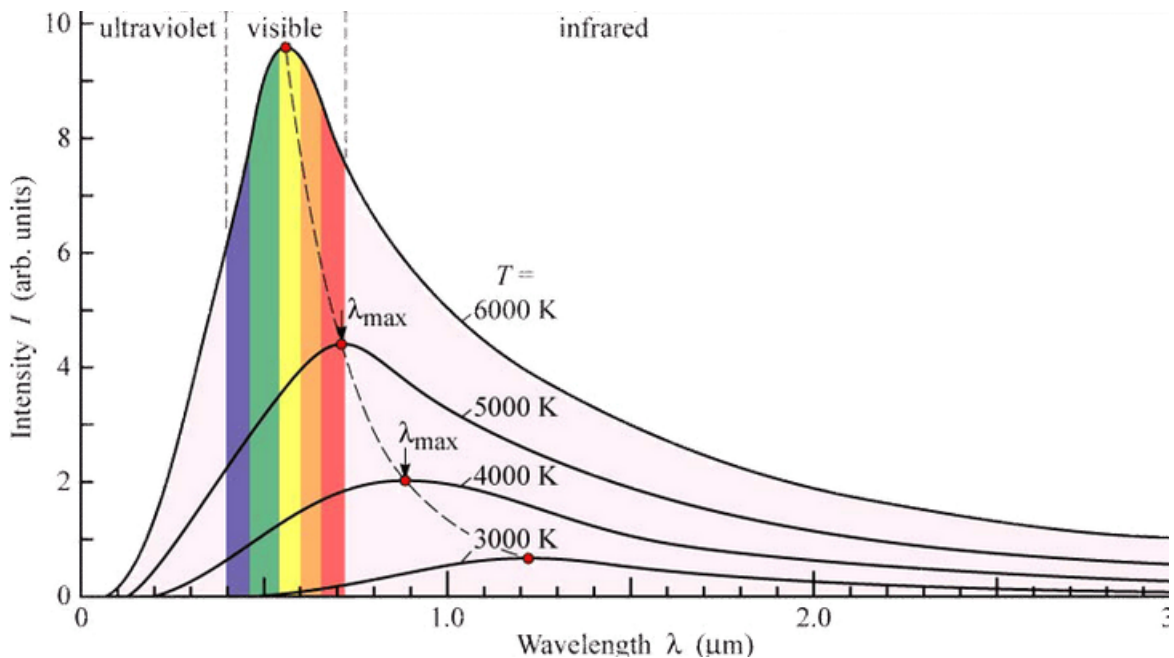
1.4 Conditions for Equilibrium

The above analysis assumes **Complete Thermodynamic Equilibrium (CTE)** meaning we assume:

1. Single temperature $T_e = T_+ = T_A$
 - All of the species we discussed had the same. The temperature of ions, atoms and electrons was the same as the temperature of the entire gas. In a lot of electric propulsion systems this is rarely the case. Maybe in the case of a magneto-thruster that operates at a higher density plasma, but it still wont be close.
2. Homogeneous plasma
 - No gradience in pressure, temperature or radiation field (I_ν)
3. Radiation field follows blackbody distribution

$$I_\nu = B_\nu(T) = \frac{2h\nu^2}{c_0^2} \frac{1}{\exp(h\nu/kt) - 1} \quad (32)$$

where I_ν is the radiation intensity (energy per unit time per unit area per unit solid angle), ν is the photon frequency and c_0 is the vacuum speed of light. Note that this requires an optically thick plasma, where all radiation is absorbed and re-radiated multiple times (as in Sun's interior).

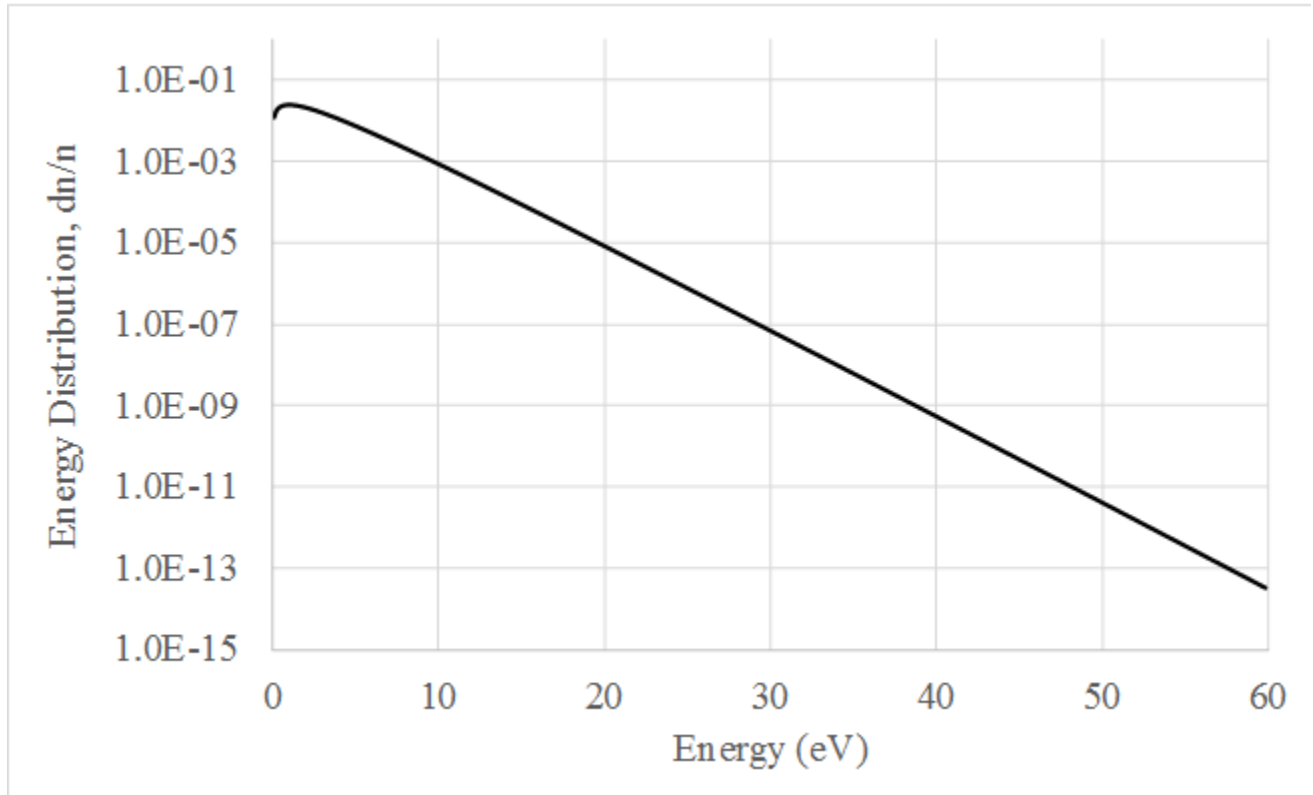


4. Maxwellian-Boltzmann Energy Distribution.

- Which follows from the Maxwellian velocity (speed) distribution from Chapter 4 Equation 30

$$\frac{\partial n}{n} = \frac{2\varepsilon^{\frac{1}{2}}}{\pi^{\frac{1}{2}} (kT)^{\frac{3}{2}}} \exp\left(-\frac{\varepsilon_i}{kT}\right) d\varepsilon \quad (33)$$

For $T_e = 2\text{eV}$ electrons



The plot shows the distribution of particles across energy space.

Note that the slope varies inversely with the temperature, because of $\exp(1/T)$ term in Equation 33. We use this in several diagnostics (Langmuir probes, atomic Boltzmann method in optical emission spectroscopy) to measure plasma temperature.

Also note the dip in energies less than two times the temperature. This is because the first part of Equation 33 is important whenever the energy is relatively small. It blows up after the exponential term dominates at higher temperatures.

In terms of speed,

$$\frac{dn}{n} = \left(\frac{2}{\pi}\right)^{\frac{1}{2}} \left(\frac{m}{kT}\right)^{\frac{3}{2}} c^2 \exp\left(\frac{-mc^2}{2kT}\right) dc \quad (34)$$

And, again, the mean speed for a Maxwellian is:

$$\bar{c} = \left(\frac{8kT}{\pi m}\right)^{\frac{1}{2}} \quad (35)$$

In CTE,

- All ionization follows the Saha equation
- All excitation follows the Boltzmann equation

1.4.1 Deviations from Equilibrium

1. Local thermodynamic equilibrium (LTE), where plasma is

- Sufficiently dense that collisions are dominant populating mechanism
- Optically thin, so radiation field is nonequilibrium, can't use Planck's law
- Radiative processes do not follow detailed balancing, assume we can ignore radiation processes, plasma is collision determined.

In LTE, radiative emission dominates over absorption and photoionization, depleting the number of free electrons and excited states. LTE is frequently encountered in lab plasmas.

2. Partial local thermodynamic equilibrium (PLTE), where plasma is lower density:

- Collisions are in detailed balance for upper energy levels
- Collisions are not in detailed balance for transitions to the ground state
- Often Ti ; Te

As a result, upper energy levels follow Boltzmann distribution, but the ground state is overpopulated. Again, depletes the free electron density.

3. Non-equilibrium Plasma

- Multi-thermal "equilibrium" where we have a two-temperature plasma. Typically:

$$T_e > or >> T_i \approx T_A$$

The reason for this difference in temperatures is via electron-impact collisions are an efficient method to transfer energy to gas, but a poor method of transferring momentum. So atoms get energy from electron to become ionized (takes ionization energy), but they don't get momentum, so their translational energy (temperature) doesn't change.

Fluorescent lights work by electron-impact excitation of mercury atoms. (The dominant lines are in the UV, but coatings in the tube walls downconvert to visible light).
Temperatures:

- Gas, 300 K 1/40 eV
- Electrons, 3.5 eV 35,000 K

Collisional processes are often dominated by electrons, so we can get LTE or PLTE where the relevant temperature is T_e .

- Corona "equilibrium" Not really a true equilibrium, as there is no detailed balancing. Characteristics:
 - Very low density plasmas
 - Radiative decay rates \ll collisional decay rates
 - Optically thin
 - Radiative excitation and ionization \gg collisional excitation and ionization

In this case we can balance

Energy in by collisions = energy out by radiation

For charge states Z and $Z+1$, the ionization balance is then:

$$n_e n_{Z+1} k_r = n_e n_Z k_c$$

where

k_r = is the radiative recombination rate [s^{-1}]

k_c = is the electron-impact ionization rate [s^{-1}]

Rearranging we get a simple relation for the corona model,

$$\frac{n_{Z+1}}{n_Z} = \frac{k_c}{k_r} \tag{36}$$

which you'll note is independent of n_e .