6 Amperes Law

From Equation 80, we arrived that

$$\nabla \times \vec{H} = \vec{j}$$

For a current density \vec{j} not due to gradient in magnetization (but rather due to "free" current). By Stokes's theorem (Equation 73)

$$\oint_C \vec{H} \cdot d\vec{l} = \int_S \vec{j} \cdot \hat{n} \, dA$$

But this FAILS in the general case. Recall back from Equation 81 that this was the result for static cases only. Equation 81 must be modified to be completely general. In other words, we need to include the "displacement" current?

Consider a current J(t) charging a capacitor (Jahn Fig 2-3):

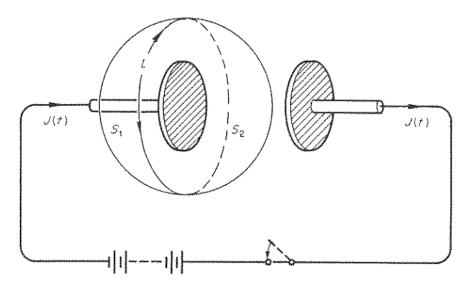


Fig. 2-3 Hypothetical experiment to illustrate Maxwell displacement current.

In this, we allow one plate of a capacitor in a simple charging circuit to be enclosed by a control surface S, around which is a closed line element l, which divides the surface into two sections, S_1 and S_2 . S_1 is the input end while S_2 is the output end of the electrode. When the switch is closed, a current J(t) flows around the circuit to charge the two capacitor plates.

Applying the previously-derived version of Ampere's law (Equation 80 and Equation 81) to contour C and surface S_1 we get:

$$\oint_C \vec{H} \cdot d\vec{l} = \int_{S_1} \vec{j} \cdot \hat{n} \, dA = J(t)$$

But l is also a boundary for S_2 , and no current flows through that surface. Hence, we should infer the contradiction:

$$\oint_C \vec{H} \cdot d\vec{l} = \int_{S_2} \vec{j} \cdot \hat{n} \, dA = 0$$

Clearly this needs modification for transient situations since

$$\int_{S_2} \vec{j} \cdot \hat{n} \, dA - \int_{S_1} \vec{j} \cdot \hat{n} \, dA \neq 0$$

The result of which means that in order to sum to 0, $\frac{\partial \rho}{\partial t} = 0$ which is not ideal. The reason why requires us to consider current continuity:

$$\frac{\partial \rho_e}{\partial t} + \nabla \cdot \vec{j} = 0$$

If

$$\frac{\partial \rho_e}{\partial t} = 0 \quad \text{(Static)} \to \nabla \cdot \overrightarrow{j} = 0$$

But we are missing a time-varying ρ_e

The net charge within the control volume is changing with time and therefore is NOT equal to zero because charge is accumulating on the capacitor plate. The capacitor is being charged.

Now we see why Equation 80, $\nabla \times \vec{H} = \vec{j}$, is inadequate, it doesn't satisfy current continuity.

We can fix discrepancy by adding a term. Since (Equation 35)

$$\nabla \cdot \vec{D} = \rho_e$$

Then

$$\frac{\partial}{\partial t}\nabla\cdot\overrightarrow{D}=\frac{\partial\rho_{e}}{\partial t}$$

Now, we can write current continuity (Equation 55) as

$$\frac{\partial}{\partial t}(\nabla \cdot \vec{D}) + \nabla \cdot \vec{j} = 0$$

Or

$$\nabla \cdot (\overrightarrow{j} + \frac{\partial \overrightarrow{D}}{\partial t}) = 0 \tag{95}$$

The correct form of the magnetic intensity ("H"-Field) equation then becomes the:

Generalized Form of Ampere's Law

$$\nabla \times \overrightarrow{H} = \overrightarrow{j} + \frac{\partial \overrightarrow{D}}{\partial t} \tag{96}$$

Where

$$\frac{\partial \vec{D}}{\partial t}$$
 = Displacement Current

$$\vec{j} + \frac{\partial \vec{D}}{\partial t} = \text{Generalized Current Density}$$

Previously we assumed (in Equation 81) that fields were independent of time or slowly varying, so $\frac{\partial \vec{D}}{\partial t} << \vec{j}$ for Equation 80 and Equation 81.

To summarize and tie things back to Fig. 2-3 above, Maxwell resolved this inconsistency by arguing that in the process of charging the capacitor in this way, an electric field is caused to increase between the plates and that this time rate of change of \vec{E} is equivalent to a real current density for the purposes of induction of an "H"-field. However, to be dimensionally consistent, it is the time derivative of \vec{D} , the source field, not \vec{E} , that is needed.

7 Maxwell's Equations - Summary

We now have all the background information required to write Maxwell's Equations:

$$abla imes \overrightarrow{H} = \overrightarrow{j} + rac{\partial \overrightarrow{D}}{\partial t}$$
 Loi d'Ampère

The curl of the H-field is due to a time-varying displacement current and the applied current.

Meaning:

The magnetic field induced around a closed loop is proportional to the electric current plus displacement current (rate of change of electric field) that the loop encloses.

$$abla imes ec{E} = rac{\partial ec{B}}{\partial t}$$
 Loi de Faraday

The curl of the E-field is due to a time-varying B-field *Meaning:*

The voltage induced in a closed loop is proportional to the rate of change of the magnetic flux that the loop encloses.

$$\nabla \cdot \overrightarrow{D} = \rho_f$$
 Loi de Gauss

The divergence of the displacement current is due to the free charge *Meaning:*

The electric flux leaving a volume is proportional to the charge inside.

$$\nabla \cdot \vec{B} = 0$$
 Loi Monopole

The divergence of the B-field is 0

Meaning:

Magnetic Monopoles do not exist; the total magnetic

flux through a closed surface is zero.

When dealing with material bodies, we need constitutive relations, meaning we need to know how those materials respond to Electric and Magnetic fields:

$$\vec{D} = \epsilon(\vec{E}) \vec{E}$$
 Equation 39

$$\vec{H} = \mu(\vec{B}) \vec{B}$$
 Equation 83

Finally, we need a force equation to model the forces on charged particles, Lorentz Force Equation (Equation 66):

$$\vec{F} = q \left(\vec{E} + \vec{v} \times \vec{B} \right)$$