

# Learning Dynamics

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## Abstract

The work presents a fundamental theory of artificial learning that aims to explain existing optimization algorithms and support the development of new ones. It is based on a thermodynamic framework, using geometric and Hamiltonian formalisms to derive the evolution equations of learning models. These equations provide a physical grounding for the learning process and show that learning can be understood as a thermodynamic process in which models evolve according to the information they perceive.

## Introduction

This work aims to develop a fundamental theory of artificial learning that explains existing optimization algorithms and facilitates their improvement, as well as the creation of new ones. To this end, only three assumptions were made about a model capable of learning:

- It can be described by a finite set of parameters.
- It is differentiable with respect to its parameters.
- It makes no assumptions about the information it ingests.

The theory is constructed based on the axiomatic thermodynamic framework proposed by Callen [1]. Assuming a quasi-static regime, symplectic geometry is used to describe the geometry of the model's phase space [2], and a Hamiltonian formalism [3] is used to derive the evolution equations of the model.

Finally, we utilize these equations to re-derive established optimization algorithms, thereby validating the theory and offering a physical grounding for the learning process. This shows that algorithms like momentum-based Stochastic Gradient Descent [4] or regularization techniques like Weight Decay [5] emerge naturally as consequences of the proposed equations of learning.

## Model

The connection between the concepts of information and entropy [6] suggests that learning can be modeled as a thermodynamic process, in which its participants, known as models, evolve based on the information they perceive.

A model is a simplified representation of a system, defined by a set of parameters that determine its behavior. A specific choice of a set of parameters defines what we will call a configuration.

We say that a model is differentiable if its configurations are points on a smooth differential manifold such that its learning can be defined in terms of a curve parametrized over it [7].

On the other hand, we will only focus on models that are domain-agnostic, meaning they do not incorporate assumptions about the information they will ingest. In this way, the dynamics of their learning can be designed solely in terms of their configurations.

## Equilibrium

Let's consider a differentiable model whose parameters reside in a smooth manifold  $S$ . For each dimension of the manifold, we can associate a parameter so that each configuration  $s \in S$  can be described in terms of coordinates.

$$(U/c, w^1, \dots, w^d) = (w^0, \mathbf{w}) \in \mathbb{R}^{d+1} \quad (1)$$

Where  $U$  represents the internal energy,  $w^1, \dots, w^d$  are extensive parameters of the model known as weights, and  $c$  is a constant such that  $w^0$  is dimensionless.

We say that the model is in an equilibrium state [1] if an entropy function  $S$  can be defined over it, with units of information, and monotonically increasing with respect to energy, that is:

$$\frac{\partial S}{\partial U} > 0 \quad (2)$$

By differentiating  $S$ , we can see how it changes under infinitesimal displacements of the configuration:

$$dS = \frac{\partial S}{\partial U} dU + \sum_j \frac{\partial S}{\partial w^j} dw^j. \quad (3)$$

The rates of change of entropy along the directions of energy and weights give rise to the conjugate variables:

$$\beta = \frac{\partial S}{\partial U} \quad Y_j = \frac{\partial S}{\partial w^j} \quad j = 1, \dots, d \quad (4)$$

These variables are known as the intensive parameters of the model. We will refer to the  $\mathbf{Y}$  intensive parameters as *entropic moments*. We can also identify the temperature  $T$  as the reciprocal of the parameter  $\beta$  conjugate to the energy, that is:

$$T \equiv \frac{1}{\beta} > 0 \quad (5)$$

The entropy function is local, that is, it is only defined for each equilibrium state. Therefore, if one seeks to describe the states of the model over the entire state space, it is necessary to resort to the phase space [3] defined over  $S$ .

The phase space is a construction over the state space that assigns to each point its cotangent space; that is, if  $(w_0, \mathbf{w})$  are coordinates of the state space, then  $(w_0, \mathbf{w}, Y^0, \mathbf{Y})$  are coordinates of the phase space. Let us now consider the 1-form living in the phase space  $\Omega$  given by:

$$\omega = \beta dU + \sum_j Y_j dw^j \in \Omega \quad (6)$$

This form generalizes the notion of the differential of entropy, such that the model is in an equilibrium state if there exists an entropy function  $S$  such that:

$$\omega = dS. \quad (7)$$

Expanding the exterior derivative of the differential 1-form  $\omega$ , we obtain the differential 2-form:

$$d\omega = d\beta \wedge dU + \sum_j dY_j \wedge dw^j \quad (8)$$

The latter is known as the *symplectic form* [2] and allows the phase space  $\Omega$  to be endowed with a Hamiltonian geometric structure.

## Evolution

For each point in the phase space, entropy is defined only for equilibrium states. This means that, to remain within the scope of a thermodynamic description, the system's evolution must be slow enough to preserve the quasi-static approximation, then the learning curve can be viewed as a succession of equilibrium states what allows us to define canonical pairs over the entire manifold through Poisson brackets:

$$\{U, \beta\} = 1 \quad \{w^i, Y_j\} = \delta_j^i \quad \text{with } \delta_j^i = \begin{cases} 1 & i = j \\ 0 & i \neq j \end{cases} \quad (9)$$

Then, we can recover an analogue to Hamilton's equations [8] for thermodynamic parameters to describe the evolution of the model parameters without constraints:

$$-h \frac{dw^i}{dt} = k\{w^i, H\} = k \frac{\partial H}{\partial Y_i} \quad - \frac{dY_i}{dt} = k\{Y_i, H\} = -k \frac{\partial H}{\partial w^i} \quad (10)$$

$$-\frac{dU}{dt} = k\{U, H\} = k \frac{\partial H}{\partial \beta} \quad -h \frac{d\beta}{dt} = k\{\beta, H\} = -k \frac{\partial H}{\partial U} \quad (11)$$

where  $h$  is the unit of action and  $k$  is the unit of information, introduced to maintain consistent units.

The problem with this formulation is that it leads to dynamics in which the model evolves in closed orbits. To address this, we introduce a coupling of the intensive parameters with the temperature of the form:

$$Y_i = \beta X_i \quad i = 1, \dots, d. \quad (12)$$

We will refer to the  $\mathbf{X}$  parameters as *energy moments*. This coupling is not arbitrary; rather, it arises directly from the energy representation of thermodynamics:

$$dU = T dS - \sum_j X_j dw^j \quad (13)$$

The coupling deforms the symplectic structure that describes the geometry of the phase space. By substituting the coupling into the 2-form  $d\omega$ , we obtain:

$$d\omega = d\beta \wedge \left( dU + \sum_j X_j dw^j \right) + \beta \sum_j dX_j \wedge dw^j \quad (14)$$

Which remains a non-degenerate symplectic form for  $\beta > 0$ , a condition already imposed. The new non-negative Poisson brackets yield:

$$\{U, \beta\} = 1 \quad \{w^i, X_j\} = \delta_j^i \quad \{U, X_i\} = -\frac{1}{\beta} X_i \quad (15)$$

Their respective equations of motion are given by:

$$-h \frac{dw^i}{dt} = \frac{k}{\beta} \frac{\partial H}{\partial X_i} \quad -h \frac{dX_i}{dt} = -\frac{k}{\beta} \frac{\partial H}{\partial w^i} + \frac{kX_i}{\beta} \frac{\partial H}{\partial U} \quad (16)$$

$$-h \frac{dU}{dt} = k \frac{\partial H}{\partial \beta} - \frac{k}{\beta} \sum_j X_j \frac{\partial H}{\partial X_j} \quad -h \frac{d\beta}{dt} = -k \frac{\partial H}{\partial U} \quad (17)$$

We will refer to these as the *Hermosis equations of learning*. While they can be rigorously derived from the new symplectic form, a more streamlined derivation based on the properties of Poisson brackets is provided in the Appendix.

## Integration

The presented four equations allow us to describe the learning process of a model. Since they must be numerically integrated to perform optimization, we seek their integral form. By substituting the last equation into the second one and rearranging terms, we obtain the evolution equation for the momenta  $X_i$  as:

$$h\beta \frac{dX_i}{dt} + h \frac{d\beta}{dt} X_i = h \frac{d}{dt}(\beta X_i) = k \frac{\partial H}{\partial w^i} \quad (18)$$

Integrating over the interval  $[t - \tau, t]$ , we obtain an update rule for the momenta:

$$\beta(t)\mathbf{X}(t) = \beta(t - \tau)\mathbf{X}(t - \tau) - \int_{t-\tau}^t \mathbf{F}(t') dt' \quad (19)$$

Where  $\mathbf{F}$  represents a generalized force, whose components are given by:

$$F_i = -\frac{k}{h} \frac{\partial H}{\partial w^i} \quad (20)$$

On the other hand, integrating the first equation over the same interval, we obtain an update rule for the weights:

$$\mathbf{w}(t) = \mathbf{w}(t - \tau) - \int_{t-\tau}^t \mathbf{v}(t') dt' \quad (21)$$

where  $\mathbf{v}$  denotes the learning velocity, with components:

$$v^i = \frac{k}{h\beta} \frac{\partial H}{\partial X_i} \quad (22)$$

This last equation tells us something important, the  $\beta$  parameter determines the system's inertia throughout its evolution. Large values of  $\beta$  slow down learning, while small values accelerate it.

## Application

Let us now examine the connection between the proposed dynamics and current algorithms used in machine learning.

In practice, a model is trained by minimizing a loss function  $L$ , which measures the distance between a model's current state and an expected state. Drawing from classical mechanics, we propose a potential analogous to the gravitational potential:

$$V = \frac{\beta c^2}{k} L \quad (23)$$

where  $c^2$ , with units of square energy, is introduced to ensure that  $L$  remains dimensionless. The choice of this potential is not arbitrary, it is based on the interpretation of the term  $\beta/k$  as a thermal mass that amplifies the importance of the distance within the potential energy. Furthermore, we propose a kinetic energy in terms of a mass tensor  $M_{ij}$  of the form:

$$K = \frac{1}{2} \sum_{ij} M^{ij} Y_i Y_j = \frac{\beta}{2k} \sum_{ij} g^{ij} X_i X_j = \frac{\beta}{2k} \mathbf{X}^2 \quad (24)$$

Here  $g_{ij}$  is a dimensionless metric tensor, which we assume for now to be constant. In this way, the Hamiltonian is defined as:

$$H = \frac{\beta}{2k} \mathbf{X}^2 + \frac{\beta c^2}{k} L + E(U, \beta) \quad (25)$$

Where  $E$  is a function of the temperature, enabling us to specify an arbitrary thermal profile. Under this Hamiltonian, the velocity components are given by:

$$v^i = \frac{k}{h\beta} \frac{\partial H}{\partial X_i} = \frac{1}{h} X^i \quad X^i = \sum_j g^{ij} X_j \quad (26)$$

Applying an Euler discretization, the weight update rule can be approximated as:

$$\mathbf{w}(t) \approx \mathbf{w}(t - \tau) - \frac{\tau}{h} \mathbf{X}(t) \quad (27)$$

On the other hand, the generalized force driving the learning process has the following components:

$$F_i = -\frac{\beta c^2}{h} \frac{\partial L}{\partial w^i} \quad (28)$$

By evaluating the components and performing the integration, we arrive at an expression for the impulse as a function of the loss gradient.

$$\mathbf{I}(t) = - \int_{t-\tau}^t \mathbf{F}(t') dt' = \frac{c^2}{h} \nabla L \int_{t-\tau}^t \beta(t') dt' \quad (29)$$

Substituting into the momenta update rule:

$$\mathbf{X}(t) = \frac{\beta(t - \tau)}{\beta(t)} \mathbf{X}(t - \tau) + \frac{c^2}{h} \left( \frac{1}{\beta(t)} \int_{t-\tau}^t \beta(t') dt' \right) \nabla L \quad (30)$$

This tells us that the momenta are updated in terms of time averages of the thermal mass.

## Constant temperature

Assuming a constant temperature, that is, a constant thermal mass:

$$\beta(t) = \beta$$

We obtain a constant impulse that depends on the step size:

$$\mathbf{I}(t) = \frac{\beta c^2}{h} \int_{t-\tau}^t \nabla L dt' = \frac{\beta c^2 \tau}{h} \nabla L \quad (31)$$

And the update rule will be a historical accumulation of the loss gradient:

$$\mathbf{X}(t) = \mathbf{X}(t - \tau) + \frac{\tau c^2}{h} \nabla L. \quad (32)$$

Under suitable reparameterization, we recover the stochastic gradient descent or SGD update rule [4]:

$$\begin{aligned} \mathbf{X}(t) &= \mathbf{X}(t - \tau) + \zeta \nabla L, \\ \mathbf{w}(t) &\approx \mathbf{w}(t - \tau) - \eta \mathbf{X}(t) \end{aligned} \quad (33)$$

Where  $\eta$  is the learning rate and  $\zeta$  controls the influence of the loss gradient on momentum.

## Harmonic potential

By adding a harmonic potential to the Hamiltonian:

$$H = \frac{\beta}{2k} \mathbf{X}^2 + \frac{\lambda}{2} \mathbf{w}^2 + V \quad (34)$$

The generalized force now includes a term consistent with that used in gradient descent with weight decay [5], which proposes adding a term  $\lambda \mathbf{w}$  to the loss gradient. This is because an extra term is now added to the force:

$$F'_i = \lambda w_i \quad (35)$$

Further physical explanations of the underlying mechanism, as well as possible corrections, are deferred to future work.

## Exponential cooling

Assuming that the system temperature decreases to zero exponentially, that is, an exponentially growing thermal mass:

$$\beta(t) = \beta e^{\gamma t} \quad (36)$$

We obtain an impulse given by:

$$\mathbf{I}(t) = \frac{c^2}{h} \frac{1 - e^{-\gamma t}}{\gamma} \beta e^{\gamma t} \nabla L \quad (37)$$

Then, momentum will be an exponential moving average [9] of the loss gradient:

$$\mathbf{X}(t) = e^{-\gamma \tau} \mathbf{X}(t - \tau) + \frac{1 - e^{-\gamma \tau}}{\gamma} \frac{c^2}{h} \nabla L \quad (38)$$

The exponential moving average is biased toward zero at early times due to its initialization, therefore, a bias correction factor should be applied [10]:

$$\hat{\mathbf{X}}(t) = \frac{\mathbf{X}(t)}{1 - e^{-\gamma t}} \quad (39)$$

Under reparameterization, this recovers gradient descent with momentum and friction [11]:

$$\begin{aligned} \mathbf{X}(t) &= \mu \mathbf{X}(t - \tau) + (1 - \mu) \zeta \nabla L \\ \mathbf{w}(t) &= \mathbf{w}(t - \tau) - \eta \mathbf{X}(t) \end{aligned} \quad (40)$$

Most implementations of SGD with momentum ignore the  $\gamma$  parameter and skip the bias correction.

## Relativistic Hamiltonian

Lastly, we adopt a classical Hamiltonian incorporating relativistic kinetic energy, given by:

$$H = \sqrt{\mathbf{X}^2 + \frac{k^2}{\beta^2}} + V + E(U, \beta) \quad (41)$$

This Hamiltonian is a generalization of the previous one, since for small momenta  $|\mathbf{X}| \ll k/\beta$ , the kinetic energy can be approximated as:

$$\sqrt{\mathbf{X}^2 + \frac{k^2}{\beta^2}} \approx \frac{k}{\beta} + \frac{\beta}{2k} \mathbf{X}^2 - \frac{\beta^3 (\mathbf{X}^2)^2}{8k^3} + \dots \quad (42)$$

The first term,  $kT$ , corresponds to a rest energy, while higher-order terms are typically discarded. However, recent work on physics-inspired optimizers [12] has shown that retaining higher-order terms in the series can enhance optimization. Under a relativistic regime, the learning velocity will then be:

$$v^i = \frac{k}{h\beta} \frac{\partial H}{\partial X_i} = \frac{k}{h\beta} \frac{X^i}{\sqrt{\mathbf{X}^2 + k^2/\beta^2}} = \frac{1}{h} \frac{X^i}{\sqrt{\beta^2 \mathbf{X}^2/k^2 + 1}} \quad (43)$$

If we consider  $\beta(t) = \beta$  to be constant, we recover the relativistic gradient descent [13] which proposes weight updates of the form:

$$\mathbf{w}(t) = \mathbf{w}(t - \tau) - \eta \frac{\mathbf{X}}{\sqrt{\mathbf{X}^2 + k^2/\beta^2}} \quad (44)$$

However, considering  $\beta(t) = \beta e^{\gamma t}$ , we observe that  $\mathbf{v} \rightarrow 0$  rapidly and the system ceases to learn. This is not an error in the theory, rather, it is due to the fact that we are working under a classical approximation in which the potential is considered decoupled from the metric tensor.

The work of Guskov and Vanchurin on covariant gradient descent [14] suggests that by embedding the potential into the metric tensor, recovers adaptive momentum-based optimizers such as Adam [10] and, as a specific case, the RMSProp optimizer [15]. Verifying this correspondence explicitly within the current theoretical framework remains as future work.

## Conclusion

The proposed formalism demonstrates that a Hamiltonian framework can recover the majority of current optimization algorithms. This approach moves beyond purely heuristic updates, offering a clear physical interpretation where:

- The loss function acts as a metric distance between model's configurations.
- Temperature dictates the effective mass, weighting the importance of these distances.
- Phase space convergence is guaranteed by the evolution of internal energy.

These results lay the groundwork for a new class of physically-informed optimizers. By identifying where current classical approximations fail, such as the decoupling of the potential from the metric tensor, we open new avenues for incorporating covariant and relativistic dynamics into machine learning, potentially leading to more stable and faster convergence in complex loss landscapes.

## Appendix: Derivation of Learning Equations

The equations for the system without temperature coupling, described by the variables  $\mathbf{Y}$ , are:

$$-h \frac{dw^i}{dt} = k\{w^i, H\} = k \frac{\partial H}{\partial Y_i}, \quad -h \frac{dY_i}{dt} = k\{Y_i, H\} = -k \frac{\partial H}{\partial w^i}, \quad (45)$$

$$-\frac{dU}{dt} = k \frac{\partial H}{\partial \beta}, \quad -h \frac{d\beta}{dt} = -k \frac{\partial H}{\partial U}. \quad (46)$$

We now derive the equations with the coupling  $\mathbf{Y} = \beta \mathbf{X}$ . From the Leibniz rule for the bracket containing the coordinates of  $\mathbf{Y}$ , given a Hamiltonian  $H$ :

$$\{Y_i, H\} = \{\beta X_i, H\} = \beta \{X_i, H\} + X_i \{\beta, H\}. \quad (47)$$

Given that we know:

$$\{Y_i, H\} = -\frac{\partial H}{\partial w^i}, \quad \{\beta, H\} = -\frac{\partial H}{\partial U}, \quad (48)$$

We substitute these values to obtain the Poisson brackets of the system in terms of the momenta  $\mathbf{X}$ :

$$-\frac{\partial H}{\partial w^i} = \beta \{X_i, H\} - X_i \frac{\partial H}{\partial U} \Rightarrow \{X_i, H\} = -\frac{1}{\beta} \frac{\partial H}{\partial w^i} + \frac{X_i}{\beta} \frac{\partial H}{\partial U}. \quad (49)$$

We then recover the proposed H-equation:

$$-h \frac{dX_i}{dt} = k\{X_i, H\} = -\frac{k}{\beta} \frac{\partial H}{\partial w^i} + \frac{kX_i}{\beta} \frac{\partial H}{\partial U}. \quad (50)$$

To reconcile the evolution of the internal energy with the proposed evolution, we recall the thermodynamic identity:

$$\left( \frac{\partial H}{\partial \beta} \right)_{\mathbf{Y}} = \left( \frac{\partial H}{\partial \beta} \right)_{\mathbf{X}} + \sum_j \left( \frac{\partial H}{\partial X_j} \right)_{\beta, X_{i \neq j}} \left( \frac{\partial X_j}{\partial \beta} \right)_{\mathbf{Y}}, \quad (51)$$

Where the subscripts indicate which parameters are considered constant.

From the coupling relation:

$$X_i = \frac{Y_i}{\beta} \Rightarrow \frac{\partial X_i}{\partial \beta} = -\frac{1}{\beta^2} Y_i = -\frac{1}{\beta} X_i, \quad (52)$$

we obtain the evolution of the internal energy:

$$-h \frac{dU}{dt} = k \frac{\partial H}{\partial \beta} - \frac{1}{\beta} \sum_j \frac{\partial H}{\partial X_j} X_j. \quad (53)$$

For the weight-momentum brackets:

$$\frac{\partial H}{\partial Y_i} = \frac{\partial H}{\partial X_i} \frac{\partial X_i}{\partial Y_i} = \frac{1}{\beta} \frac{\partial H}{\partial X_i}. \quad (54)$$

Then:

$$-h \frac{dw^i}{dt} = \frac{k}{\beta} \frac{\partial H}{\partial X_i}. \quad (55)$$

Finally, the equation containing the changes of  $H$  in the direction  $\partial/\partial U$  remains identical, since the variable  $U$  is not coupled in the same way as  $\beta$ .

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