

# 16-811 hw 4

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## 1

the surface area of a curve  $y$  rotated about the  $x$ -axis is given by the integral:

$$\int_{x_0}^{x_1} 2\pi y \sqrt{1 + y'^2} dx$$

let's take the constant  $2\pi$  out of the integral and then we have:

$$F = y \sqrt{1 + y'^2}$$

we also have:

$$F_{y'} = \frac{yy'}{\sqrt{1 + y'^2}}$$

from pg. 26 of the Calculus of Variations notes:

$$y' F_{y'} - F = c_0$$

plugging things in:

$$\frac{yy'^2}{\sqrt{1 + y'^2}} - y \sqrt{1 + y'^2} = c_0$$

simplify:

$$c_0 + y(1 + y'^2) - yy'^2 = 0$$

simplifying some more:

$$\begin{aligned} y &= -c_0 \sqrt{1 + y'^2} \\ \sqrt{\frac{y^2 - c_0^2}{c_0^2}} &= y' \\ \frac{dy}{dx} &= \frac{\sqrt{y^2 - c_0^2}}{c_0} \\ dy &= \frac{\sqrt{y^2 - c_0^2}}{c_0} dx \end{aligned}$$

$$\int \frac{c_0}{\sqrt{y^2 - c_0^2}} dy = \int dx$$

$$c_0 \cosh^{-1} \left( \frac{y}{c_0} \right) + c_1 = x$$

$$y = c_0 \cosh \left( \frac{x - c_1}{c_0} \right)$$

plug in initial conditions to solve for  $c_0$  and  $c_1$ :

$$y_0 = c_0 \cosh \left( \frac{x_0 - c_1}{c_0} \right)$$

$$y_1 = c_0 \cosh \left( \frac{x_1 - c_1}{c_0} \right)$$

Not really sure about how to solve these implicit equations for  $c_0$  and  $c_1$ . probably can't. so basically, there's not always a  $C^2$  solution. the solution forms a catenary between the points  $(x_0, y_0)$  and  $(x_1, y_1)$ . because  $y$  is constrained to be  $> 0$ , in the case where  $(x_0, y_0)$  and  $(x_1, y_1)$  causes the catenary that connects them to drop below the x-axis, then it's hard to say what the minimizing solution will be.

## 2

objective is 3D arc length, given by the integral:

$$\int_{x_0}^{x_1} \sqrt{dx^2 + dy^2 + dz^2}$$

change of variables to make life easier (change of vars allows us to weave the constraint that the points must be on a sphere into the change of variables; we get treat the variable  $R$  as a constant to constrain the points to lie on the sphere of radius  $R$ ):

$$x = R \sin v \cos u$$

$$dx = R \cos v \cos u dv - R \sin v \sin u du$$

$$y = R \sin v \sin u$$

$$dy = R \cos v \sin u dv + R \sin v \cos u du$$

$$z = R \cos v$$

$$dz = -R \sin v dv$$

plug things in:

$$\int_{\cos^{-1}(\frac{x_0}{R \sin v})}^{\cos^{-1}(\frac{x_1}{R \sin v})} \sqrt{(R \cos v \cos u dv - R \sin v \sin u du)^2 + (R \cos v \sin u dv + R \sin v \cos u du)^2 + (-R \sin v dv)^2}$$

simplify:

$$R \int_{\cos^{-1}(\frac{x_0}{R \sin v})}^{\cos^{-1}(\frac{x_1}{R \sin v})} \sqrt{dv^2 + \sin^2 v} du^2$$

$$R \int_{\cos^{-1}(\frac{x_0}{R \sin v})}^{\cos^{-1}(\frac{x_1}{R \sin v})} \sqrt{v'^2 + \sin^2 v} du$$

So we have:

$$F = \sqrt{v'^2 + \sin^2 v}$$

as there is no  $u$  in our  $F$  we will be plugging things into this ELE:

$$v' F_{v'} - F = c$$

$$F_{v'} = \frac{v'}{\sqrt{v'^2 + \sin^2 v}}$$

plug things in:

$$\frac{v'^2}{\sqrt{v'^2 + \sin^2 v}} - \sqrt{v'^2 + \sin^2 v} = c$$

simplify:

$$v'^2 - (v'^2 + \sin^2 v) = c \sqrt{v'^2 + \sin^2 v}$$

$$\sin^4 v = c^2 (v'^2 + \sin^2 v)$$

$$v' = \sqrt{\frac{\sin^4 v - c^2 \sin^2 v}{c^2}}$$

$$\frac{c}{\sqrt{\sin^4 v - c^2 \sin^2 v}} dv = du$$

integrate both sides using the helpfully provided identity:

$$u = \int \frac{c}{\sqrt{\sin^4 v - c^2 \sin^2 v}} dv = -\sin^{-1} \left( \frac{\cot v}{\sqrt{\frac{1}{c^2} - 1}} \right) + c_1$$

substitute  $c_0 = \frac{1}{\sqrt{\frac{1}{c^2} - 1}}$ :

$$c_0 \cot v = \sin(c_1 - u)$$

using the trig identity  $\sin(a - b) = \sin a \cos b - \cos a \sin b$ :

$$c_0 \cot v = \sin c_1 \cos u - \cos c_1 \sin u$$

$$c_0 \cos v = \sin c_1 \sin v \cos u - \cos c_1 \sin v \sin u$$

changing variables back to  $x$ ,  $y$ , and  $z$ :

$$\frac{c_0 z}{R} = \frac{x}{R} \sin c_1 - \frac{y}{R} \cos c_1$$

$$c_0 z = x \sin c_1 - y \cos c_1$$

$$x \sin c_1 - y \cos c_1 - c_0 z = 0$$

this is a plane that passes through the origin. by the way we parametrized the sphere, its center is at the origin. thus, this plane intersects the sphere along a great circle, which is where our minimizing curve must lie on.

### 3

The objective to minimize is the same as that used during lecture, so we will copy it here:

$$\frac{1}{\sqrt{-2g}} \int_{x_0}^{x_1} \frac{\sqrt{1+y'^2}}{\sqrt{y_0-y}} dx$$

$$F = \frac{\sqrt{1+y'^2}}{\sqrt{y_0-y}}$$

$$F_{y'} = \frac{y'}{\sqrt{(y_0-y)(1+y'^2)}}$$

We have the equation that was given during lecture:

$$F_{y'} \Big|_{x_1} = \frac{g_y F}{g_x + y' g_y}$$

plug things in (for  $x = x_1$ :

$$\frac{y'}{\sqrt{(y_0-y)(1+y'^2)}} = \frac{g_y}{g_x + y' g_y} \frac{\sqrt{1+y'^2}}{\sqrt{y_0-y}}$$

simplify:

$$y' = \frac{g_y}{g_x + y' g_y} (1 + y'^2)$$

$$g_y (1 + y'^2) = y' (g_x + y' g_y)$$

$$g_y = g_x y'$$

$$y' = \frac{g_y}{g_x}$$

this means that the slope of the curve  $y$  is orthogonal to the curve  $g(x, y)$  at  $x = x_1$ . (this is bc the gradient of  $g(x, y) = \langle g_x, g_y \rangle$  is orthogonal to the curve; the slope of the tangent line written in terms of  $x$  is  $\frac{g_y}{g_x}$ , and this is orthogonal to the curve.)

## 4

### 4a

our Lagrangian equation is:

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i} = \tau_i$$

we will substitute in the given  $q$ s and  $L$  for our given manipulator into the above equilibrium equation to obtain the relationship between torque and the state of the system.

$$L = T - V$$

we use  $v_{21}$  to refer to the upper  $m_2$ 's velocity and  $v_{22}$  for the lower  $m_2$ 's velocity, as they have the same mass but different velocities:

$$T = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_{21}^2 + \frac{1}{2}m_2v_{22}^2$$

There is no gravity, so:

$$V = 0$$

putting it together:

$$L = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_{21}^2 + \frac{1}{2}m_2v_{22}^2$$

here are our equations for  $v$  in terms of the time derivatives of  $\ell$  and  $\theta$ :

$$v_1 = \ell_1 \dot{\theta}_1$$

$$v_1^2 = \ell_1^2 \dot{\theta}_1^2$$

here, we use  $x_{21}$  and  $y_{21}$  to refer to the  $x$ - and  $y$ -position of the upper mass, as measured from the manipulator's point of contact with the ground, and analogously for  $x_{22}$  and  $y_{22}$ . Similar to as we did class, we will use  $c_i$  as stand-in for  $\cos \theta_i$ ,  $c_{12}$  for  $\cos(\theta_1 + \theta_2)$ ,  $s_i$  for  $\sin \theta_i$ , and  $s_{12}$  for  $\sin(\theta_1 + \theta_2)$ .

$$x_{21} = \ell_1 c_1 + \ell_2 c_{12}$$

$$y_{21} = \ell_1 s_1 + \ell_2 s_{12}$$

$$\dot{x}_{21} = -\ell_1 s_1 \dot{\theta}_1 - \ell_2 s_{12}(\dot{\theta}_1 + \dot{\theta}_2)$$

$$\dot{y}_{21} = \ell_1 c_1 \dot{\theta}_1 + \ell_2 c_{12}(\dot{\theta}_1 + \dot{\theta}_2)$$

$$v_{21} = \dot{x}_{21}^2 + \dot{y}_{21}^2$$

$$v_{21}^2 = \ell_1^2 \dot{\theta}_1^2 + \ell_2^2 (\dot{\theta}_1 + \dot{\theta}_2)^2 + 2\ell_1 \ell_2 c_2 (\dot{\theta}_1 + \dot{\theta}_2) \dot{\theta}_1$$

$$x_{22} = \ell_1 c_1 - \ell_2 c_{12}$$

$$\begin{aligned}
y_{22} &= \ell_1 s_1 - \ell_2 s_{12} \\
\dot{x}_{22} &= -\ell_1 s_1 \dot{\theta}_1 + \ell_2 s_{12} (\dot{\theta}_1 + \dot{\theta}_2) \\
\dot{y}_{22} &= \ell_1 c_1 \dot{\theta}_1 - \ell_2 c_{12} (\dot{\theta}_1 + \dot{\theta}_2) \\
v_{22} &= \dot{x}_{22}^2 + \dot{y}_{22}^2 \\
v_{22}^2 &= \ell_1^2 \dot{\theta}_1^2 + \ell_2^2 (\dot{\theta}_1 + \dot{\theta}_2)^2 - 2\ell_1 \ell_2 c_2 (\dot{\theta}_1 + \dot{\theta}_2) \dot{\theta}_1
\end{aligned}$$

so plugging in  $v_1^2$ ,  $v_{22}^2$ , and  $v_2^2$  into our equation for  $L$ , we have:

$$L = \frac{1}{2} m_1 \ell_1^2 \dot{\theta}_1^2 + \ell_1^2 \dot{\theta}_1^2 + \ell_2^2 (\dot{\theta}_1 + \dot{\theta}_2)^2$$

taking the necessary derivatives:

$$\frac{\partial L}{\partial \theta_1} = 0$$

$$\frac{\partial L}{\partial \theta_2} = 0$$

$$\begin{aligned}
\frac{\partial L}{\partial \dot{\theta}_1} &= m_1 \ell_1^2 \dot{\theta}_1 + 2m_2 \ell_1^2 \dot{\theta}_1 + 2m_2 \ell_2^2 (\dot{\theta}_1 + \dot{\theta}_2) \\
\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_1} &= m_1 \ell_1^2 \ddot{\theta}_1 + 2m_2 \ell_1^2 \ddot{\theta}_1 + 2m_2 \ell_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) \\
\frac{\partial L}{\partial \dot{\theta}_2} &= 2m_2 \ell_2^2 (\dot{\theta}_1 + \dot{\theta}_2) \\
\frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}_2} &= 2m_2 \ell_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2)
\end{aligned}$$

finally, plugging things into the Lagrangian equation ( $\tau_i = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} - \frac{\partial L}{\partial q_i}$ ):

$$\begin{aligned}
\tau_1 &= m_1 \ell_1^2 \ddot{\theta}_1 + 2m_2 \ell_1^2 \ddot{\theta}_1 + 2m_2 \ell_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2) \\
\tau_2 &= 2m_2 \ell_2^2 (\ddot{\theta}_1 + \ddot{\theta}_2)
\end{aligned}$$

In matrix form:

$$\begin{bmatrix} \tau_1 \\ \tau_2 \end{bmatrix} = \begin{bmatrix} m_1 \ell_1^2 + 2m_2 (\ell_1^2 + \ell_2^2) & 2m_2 \ell_2^2 \\ 2m_2 \ell_2^2 & 2m_2 \ell_2^2 \end{bmatrix} \begin{bmatrix} \ddot{\theta}_1 \\ \ddot{\theta}_2 \end{bmatrix}$$

#### 4b

if  $\ddot{\theta}_2 = 0$ , then:

$$\tau_1 = (m_1 \ell_1^2 + 2m_2 \ell_1^2 + 2m_2 \ell_2^2) \ddot{\theta}_1$$

the first term is the moment of inertia of  $m_1$ . the second term is the moment of inertia for the sum of the two  $m_2$ s about their center of mass (the point where  $m_1$  is, thus  $\ell_1$ ). the third term is the moment of inertia required to give the two  $m_2$ s their existing centripetal acceleration as they rotate about the point where  $m_1$  is located.