

Homework 5

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To complete this homework, I worked with **Andrew Binder** and **Aren Martinian**

1 Problem 1

From lecture, we have the following relationships for impedances in circuit elements:

Inductors: $Z_L = i\omega L$

Resistors: $Z_R = R$

Capacitors: $Z_C = \frac{i}{\omega C}$

We also have a relationship between the impedances of circuit elements in series and parallel:

Series: $Z_{eq} = Z_1 + Z_2$

Parallel: $Z_{eq} = \frac{Z_1 Z_2}{Z_1 + Z_2}$

Combining these two we get:

$$\begin{cases} Z_{top} &= i\omega L + R \\ Z_{bottom} &= \frac{i}{\omega C} \end{cases}$$

And since both the top and bottom are connected in parallel we have:

$$\begin{aligned} Z_{eq}(\omega) &= \frac{\frac{i}{\omega C}(i\omega L + R)}{i\omega L + R + \frac{i}{\omega C}} \\ &= \frac{iR - \omega L}{i(\omega^2 LC + 1) + R\omega C} \end{aligned}$$

Analyzing the limits of this function as ω goes to infinity, we see that constant terms don't really matter, so we look only at the terms that depend on ω . Eliminating the non ω -dependent terms, we get:

$$\begin{aligned} Z_{eq}(\omega) &= \frac{-\omega L}{\omega(i\omega LC + RC)} \\ &= \frac{-L}{i\omega LC + RC} \\ &= \frac{-L}{i\omega LC} \\ &= \frac{i}{\omega C} \end{aligned}$$

This expression for $Z_{eq}(\omega)$ is the same as the impedance of a capacitor. This makes sense intuitively, since at very high frequencies, the electrons that move around the circuit cannot move, and thus the impedance should drop to zero, just like a capacitor. We could have also deduced this fact by taking the $\lim_{\omega \rightarrow \infty} Z_{eq}(\omega)$ and showing that this limit also goes to zero. On the other hand, we notice that:

$$\lim_{\omega \rightarrow 0} Z_{eq}(\omega) = \frac{iR}{i} = R$$

This limit also makes sense intuitively because an AC power supply that has zero frequency is essentially a steady-state DC current. Since the current is steady-state, then the back \mathcal{E} from the inductor is zero (as a result of $\dot{I} = 0$), leading to the impedance from the inductor to be zero. Simultaneously, since the current is steady-state, then the capacitor acts like a break in the circuit, and thus there is no current going through that section of wire, also leading to a zero impedance. As a result, the only element that has the ability to impede the current is the resistor, which has an impedance of R , as we expected.

The approximation $Z_{eq}(\omega) \approx R + iX_L(\omega)$ is valid whenever the capacitance is very small compared to R and L . This will cause the impedance through the capacitor to be larger than that through the resistor and inductor, causing the circuit to approximately be an RL circuit. Thus, we aim to solve:

$$\begin{aligned} Z_c &\ll Z_R + Z_L \\ \frac{i}{\omega C} &\ll R + i\omega L \\ C &\ll \frac{i}{\omega(R + i\omega L)} \\ C &\ll \frac{i(\omega R - i\omega^2 L)}{\omega^2 R^2 + \omega^4 L^2} \\ C &\ll \frac{\omega^2 L}{\omega^2 R^2 + \omega^4 L^2} + \frac{i\omega R}{\omega^2 R^2 + \omega^4 L^2} \end{aligned}$$

Taking the real part of this, we get:

$$C \ll \frac{L}{R^2 + L^2 \omega^2}$$

2 Problem 2

Assume that the circuit is connected to a voltage V_T . From Kirchhoff's loop rule, we have:

$$\begin{cases} V - I_1 Z_1 - M I_2 = 0 \\ V - I_2 Z_2 - M I_1 = 0 \end{cases}$$

We can subtract these two equations, and here's the algebra:

$$\begin{aligned} I_2 Z_2 - I_1 Z_1 + M I_1 - M I_2 &= 0 \\ I_1 Z_2 - (I_T - I_2) Z_1 + M(I_T - 2I_2) &= 0 & I_T = I_1 + I_2 \text{ from Kirchhoff} \\ I_2 Z_2 - I_T Z_1 + I_2 Z_1 + M I_T - 2M I_2 &= 0 \\ I_2(Z_2 + Z_1 - 2M) &= I_T(Z_1 - M) \\ I_2(Z_2 + Z_1 - 2M) &= \frac{V_T}{Z_T}(Z_1 - M) & \text{Use } V_T = I_T Z_T \\ \therefore Z_T &= \frac{V_T}{I_2} \frac{Z_1 - M}{Z_2 + Z_1 - 2M} & \frac{V_T}{I_2} = Z_2 \text{ since they are in parallel} \\ &= \frac{Z_2(Z_1 - M)}{Z_2 + Z_1 - 2M} \end{aligned}$$

Which is the equivalent impedance.

3 Problem 3

First, we calculate Z_{out} , since the system is slightly simpler:

$$Z_{out} = \frac{R_2 \frac{i}{\omega C_2}}{R_2 + \frac{i}{\omega C_2}} = \frac{R_2 i}{R_2 \omega C_2 + i}$$

Now we calculate Z_{in} , in a similar fashion to how we calculated Z_{out} :

$$Z_{in} = \frac{R_1 i}{R_1 \omega C_1} + \frac{R_2 i}{R_2 \omega C_2 + i}$$

We also have $V_{out} = I Z_{out}$ and $V_{in} = I Z_{in}$ for both circuits. Notice that the current is the same in both loops, because of Kirchhoff's junction rules - there are paths for the current to flow in that aren't considered by both systems, so the current for both should be the same. Now, we calculate $\frac{V_{in}}{V_{out}}$ for convenience:

$$\begin{aligned} \frac{V_{in}}{V_{out}} &= \frac{Z_{in}}{Z_{out}} = \frac{\frac{R_1 i}{R_1 \omega C_1 + i} + \frac{R_2 i}{R_2 \omega C_2 + i}}{\frac{R_2 i}{R_2 \omega C_2 + i}} \\ &= \frac{\frac{R_1 i}{R_1 \omega C_1 + i}}{\frac{R_2 i}{R_2 \omega C_2 + i}} + 1 \end{aligned}$$

$$= \frac{R_1}{R_1\omega C_1 + i} \cdot \frac{R_2\omega C_2 + i}{R_2} + 1$$

Now notice that if we introduce $R_1C_1 = R_2C_2$, then $R_1\omega C_1 + i = R_2\omega C_2 + i$, since ω is the same across both circuits, so those terms cancel out. Thus, we get:

$$\frac{V_{in}}{V_{out}} = \frac{R_1}{R_2} + 1 = \frac{R_1 + R_2}{R_2}$$

If we reciprocate this so that it's in the form of the problem statement, we get:

$$\frac{V_{out}}{V_{in}} = \frac{R_2}{R_1 + R_2}$$

As desired.

This final equation for $\frac{V_{out}}{V_{in}}$ makes sense to only depend on R , since our goal from the beginning was to make this voltage divider independent of the frequency, thus meaning that our voltage divider is now compensated.

4 Problem 4

Calculating the equivalent impedances for both circuits:

$$Z_{left} = \frac{Z_R Z_L}{Z_R + Z_L} = \frac{iR^2\omega L + R\omega^2 L^2}{R^2 + \omega^2 L}$$

$$Z_{right} = Z_R + Z_C = R + \frac{i}{\omega C}$$

Thus, for them to be equal, we set $Z_{left} = Z_{right}$:

$$R + \frac{i}{\omega C} = \frac{iR^2\omega L + R\omega^2 L^2}{R^2 + \omega^2 L}$$

$$R + \frac{i}{\omega C} = \frac{R\omega^2 L^2}{R^2 + \omega^2 L} + \frac{iR^2\omega L}{R^2 + \omega^2 L}$$

To find equivalence, we need to solve the equivalences:

$$\begin{cases} R &= \frac{R\omega^2 L^2}{R^2 + \omega^2 L} \\ \frac{i}{\omega C} &= \frac{iR^2\omega L}{R^2 + \omega^2 L} \end{cases}$$

The only way that these values are equivalent to one another is if both sides equal zero (i.e. there is no other real, simultaneous solution for R , L and C that satisfies both equations). Taking the case for equating the

real values, we have $R = 0$, or $L = \infty$ (while $R \neq 0$). In the case that $R = 0$, then we get $C = \infty$, and in the case that $L = \infty$, then we get $C = \frac{1}{\omega}$.

In both these cases, we took the value of some element in the circuit to be infinity. This implies that there are no real solutions for R, L, C that causes the equivalent impedances to be exactly equal. However, we can often get a good approximate equivalence for them by setting the relative sizes of R, L and C to be as described.

5 Problem 5

We have $\bar{P} = \frac{1}{2}\varepsilon_0 I_0 \cos \phi$. Here, we will work with admittances instead of inductances, simply because admittances add in parallel and it's easier for us here because all three circuit elements are in parallel.

$$I_0 = |\tilde{I}| = \varepsilon_0 \sqrt{\left(\frac{1}{R}\right)^2 + \left(\omega C - \frac{1}{\omega L}\right)^2}$$

For the phase angle, we're looking at $\cos \phi = \frac{\Re \tilde{I}_0}{|\tilde{I}|}$. Computing this value we get:

$$\cos \phi = \frac{\frac{1}{R}}{\sqrt{\left(\frac{1}{R}\right)^2 + \left(\omega C - \frac{1}{\omega L}\right)^2}}$$

So calculating the power:

$$\begin{aligned} \bar{P} &= \frac{1}{2} \sqrt{\left(\frac{1}{R}\right)^2 + \left(\omega C - \frac{1}{\omega L}\right)^2} \cdot \frac{\frac{1}{R}}{\sqrt{\left(\frac{1}{R}\right)^2 + \left(\omega C - \frac{1}{\omega L}\right)^2}} \\ &= \frac{\varepsilon_0^2}{2R} \end{aligned}$$

We also have the power through the resistor as: $\bar{P}_R = \frac{1}{2} \frac{V_0^2}{R} = \bar{P}$ as desired. This result should make sense intuitively as well. The resistor is the only circuit element that dissipates energy over time, since both inductors and capacitors store energy in magnetic and electric fields respectively. As a result, the power dissipated through the circuit will have to be through the only circuit element that *can* dissipate energy, that being the resistor. Thus, we should expect the average power across the whole circuit to equal that of the resistor.

6 Problem 6

6.1 Part a

Finding the equivalent impedance:

$$\begin{aligned}
Z_{eq} &= \frac{R(R + \frac{i}{\omega C})}{R + R + \frac{i}{\omega C}} \\
&= \frac{(R^2 + \frac{iR}{\omega C})(2R - \frac{i}{\omega C})}{4R^2 + \frac{1}{\omega C}}
\end{aligned}$$

Taking the imaginary component of this, we get:

$$\begin{aligned}
\Im(z) &= \frac{\frac{iR^2}{\omega C}}{4R^2 + \frac{1}{\omega C}} \\
&= \frac{R^2}{4R^2\omega C + 1}
\end{aligned}$$

Since we choose $\omega = \frac{1}{RC}$ then we have:

$$\Im(z) = \frac{R^2}{4R + 1}$$

6.2 Part b

We have $I_0 = \frac{\varepsilon_0}{\Re(z)}$. From the previous part where we solved $Z_{eq}(\omega)$, we get:

$$\begin{aligned}
\Re(Z_{eq}) &= \frac{2R^3 + \frac{R}{\omega^2 C^2}}{4R^2 + \frac{1}{\omega C}} \\
&= \frac{2R^3\omega^2 C^2 + R}{\omega C(4R^2\omega C + 1)}
\end{aligned}$$

Again, we note down the fact that $\omega = \frac{1}{RC}$, so our equation above massively simplifies to:

$$\Re(Z_{eq}) = \frac{3R^2}{4R + 1}$$

Now, ε_0 is given in the problem, so we have:

$$I_0 = \frac{\varepsilon_0}{\Re(Z_{eq})} = \frac{\varepsilon_0(4R + 1)}{3R^2}$$

From the textbook, the phase angle can be calculated as $\phi = \tan^{-1}\left(\frac{\Im(Z)}{\Re(Z)}\right)$, so we get:

$$\phi = \tan^{-1}\left(\frac{\frac{R^2}{4R+1}}{\frac{3R^2}{4R+1}}\right) = \tan^{-1}\left(\frac{1}{3}\right) \approx 18.43^\circ$$

In regards to the current, it should make sense that we take $\Re Z_{eq}$ here because we are considering a *physical* current, and thus we should not consider its imaginary components. An interesting result of this value is the

fact that the current does not depend on the capacitor at all, but this is simply because we've set $\omega = \frac{1}{RC}$ to perfectly cancel out the capacitor.

In the general case, it should make sense that the current through the circuit should depend on all three circuit elements, and in fact, if we take $\Re(Z_{eq})$ without performing the substitution of $\omega = \frac{1}{RC}$, then this is exactly the result we get.

6.3 Part c

From the textbook (equation 8.84), we have $\bar{P} = \frac{1}{2}\varepsilon_0 I_0 \cos \phi$, so substituting our values:

$$\bar{P} = \frac{\varepsilon_0^2(4R+1)}{6R^2} \cos\left(\tan^{-1}\left(\frac{1}{3}\right)\right)$$