

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Settlement of a soil layer

1EL5000–Continuum Mechanics – Tutorial Class #4

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Outline

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Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

1 Some algebra

- Vector & tensor products
- Vector & tensor analysis

2 Material behavior

3 4.1 Settlement of a soil layer

Outline

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

1 Some algebra

- Vector & tensor products
- Vector & tensor analysis

2 Material behavior

3 4.1 Settlement of a soil layer

Some algebra

Vector & tensor products

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É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

■ Scalar product:

$$\mathbf{a}, \mathbf{b} \in \mathbb{R}^a, \quad \langle \mathbf{a}, \mathbf{b} \rangle = \sum_{j=1}^a a_j b_j = a_j b_j,$$

The last equality is [Einstein's summation convention](#).

■ Tensors and tensor product (or outer product):

$$\mathbf{A} \in \mathbb{R}^a \rightarrow \mathbb{R}^b, \quad \mathbf{A} = \mathbf{a} \otimes \mathbf{b}, \quad \mathbf{a} \in \mathbb{R}^a, \mathbf{b} \in \mathbb{R}^b.$$

■ Tensor application to vectors:

$$\mathbf{A} = \mathbf{a} \otimes \mathbf{b} \in \mathbb{R}^a \rightarrow \mathbb{R}^b, \mathbf{c} \in \mathbb{R}^b, \quad \mathbf{Ac} = \langle \mathbf{b}, \mathbf{c} \rangle \mathbf{a}.$$

■ Product of tensors \equiv composition of linear maps:

$$\mathbf{A} = \mathbf{a} \otimes \mathbf{b}, \mathbf{B} = \mathbf{c} \otimes \mathbf{d}, \quad \mathbf{AB} = \langle \mathbf{b}, \mathbf{c} \rangle \mathbf{a} \otimes \mathbf{d}.$$

Some algebra

Vector & tensor products

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Scalar product of tensors:

$$\langle \mathbf{A}, \mathbf{B} \rangle = \text{Tr}(\mathbf{AB}^T) := \mathbf{A} : \mathbf{B} = A_{jk} B_{jk}.$$

- Let $\{\mathbf{e}_j\}_{j=1}^d$ be a Cartesian basis in \mathbb{R}^d . Then:

$$a_j = \langle \mathbf{a}, \mathbf{e}_j \rangle,$$

$$\begin{aligned} A_{jk} &= \langle \mathbf{A}, \mathbf{e}_j \otimes \mathbf{e}_k \rangle = \mathbf{A} : \mathbf{e}_j \otimes \mathbf{e}_k \\ &= \langle \mathbf{A}\mathbf{e}_k, \mathbf{e}_j \rangle, \end{aligned}$$

such that:

$$\mathbf{a} = a_j \mathbf{e}_j,$$

$$\mathbf{A} = A_{jk} \mathbf{e}_j \otimes \mathbf{e}_k.$$

- Example: the identity matrix

$$\mathbf{I} = \mathbf{e}_j \otimes \mathbf{e}_j.$$

Some analysis

Vector & tensor analysis

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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Gradient of a vector function $\mathbf{a}(\mathbf{x})$, $\mathbf{x} \in \mathbb{R}^d$:

$$\mathbb{D}_{\mathbf{x}} \mathbf{a} = \frac{\partial \mathbf{a}}{\partial x_j} \otimes \mathbf{e}_j .$$

- Divergence of a vector function $\mathbf{a}(\mathbf{x})$, $\mathbf{x} \in \mathbb{R}^d$:

$$\operatorname{div}_{\mathbf{x}} \mathbf{a} = \langle \nabla_{\mathbf{x}} \mathbf{a}, \mathbf{a} \rangle = \operatorname{Tr}(\mathbb{D}_{\mathbf{x}} \mathbf{a}) = \frac{\partial a_j}{\partial x_j} .$$

- Divergence of a tensor function $\mathbf{A}(\mathbf{x})$, $\mathbf{x} \in \mathbb{R}^d$:

$$\operatorname{Div}_{\mathbf{x}} \mathbf{A} = \frac{\partial (\mathbf{A} \mathbf{e}_j)}{\partial x_j} .$$

Some analysis

Vector & tensor analysis in cylindrical coordinates

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Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Gradient of a vector function $\mathbf{a}(r, \theta, z)$:

$$\mathbb{D}_{\mathbf{x}} \mathbf{a} = \frac{\partial \mathbf{a}}{\partial r} \otimes \mathbf{e}_r + \frac{\partial \mathbf{a}}{\partial \theta} \otimes \frac{\mathbf{e}_\theta}{r} + \frac{\partial \mathbf{a}}{\partial z} \otimes \mathbf{e}_z.$$

- Divergence of a vector function $\mathbf{a}(r, \theta, z)$:

$$\operatorname{div}_{\mathbf{x}} \mathbf{a} = \left\langle \frac{\partial \mathbf{a}}{\partial r}, \mathbf{e}_r \right\rangle + \left\langle \frac{\partial \mathbf{a}}{\partial \theta}, \frac{\mathbf{e}_\theta}{r} \right\rangle + \left\langle \frac{\partial \mathbf{a}}{\partial z}, \mathbf{e}_z \right\rangle.$$

- Divergence of a tensor function $\mathbf{A}(r, \theta, z)$:

$$\operatorname{Div}_{\mathbf{x}} \mathbf{A} = \frac{\partial \mathbf{A}}{\partial r} \mathbf{e}_r + \frac{\partial \mathbf{A}}{\partial \theta} \frac{\mathbf{e}_\theta}{r} + \frac{\partial \mathbf{A}}{\partial z} \mathbf{e}_z.$$

Outline

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

1 Some algebra

- Vector & tensor products
- Vector & tensor analysis

2 Material behavior

3 4.1 Settlement of a soil layer

Recap

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Some
algebra

Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Local equilibrium equation:

$$\mathbf{Div}_x \boldsymbol{\sigma} + \mathbf{f}_v = \varrho \ddot{\mathbf{u}} .$$

- Small strains assumption:

$$\boldsymbol{\epsilon} = \frac{1}{2} (\mathbb{D}_x \mathbf{u} + \mathbb{D}_x \mathbf{u}^\top) .$$

- Closure is missing: $\boldsymbol{\sigma} = f(\boldsymbol{\epsilon})$ or $\boldsymbol{\epsilon} = g(\boldsymbol{\sigma})$, the material constitutive equation.

Recap

Traction test

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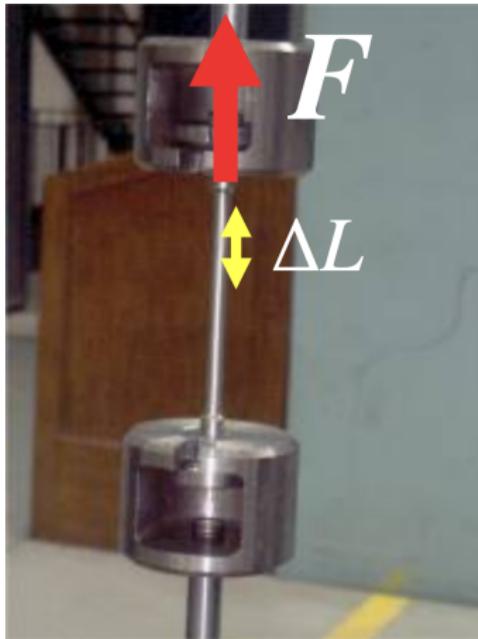
Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer



Measurements along e :

- $\sigma_{ee} = \frac{F}{S};$
- $\varepsilon_{ee} = \frac{\Delta L}{L}.$

Recap

Traction test

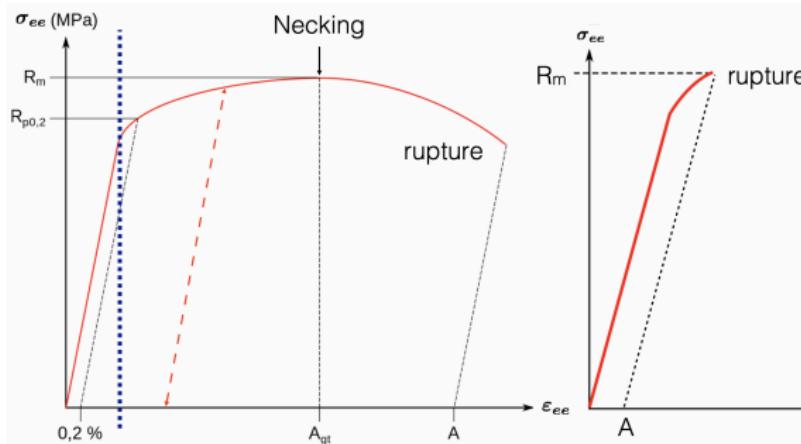
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Some algebra
Vector & tensor products
Vector & tensor analysis

Material behavior

4.1
Settlement of a soil layer



- Ductile vs. fragile materials;
- At small strains $F = K\Delta L$ or $\sigma_{ee} = k\varepsilon_{ee}$;
- Generalization: $\boldsymbol{\sigma} = \mathbf{C}\boldsymbol{\varepsilon}$ where \mathbf{C} is the fourth-order (C_{jklm}) elasticity tensor.

Elasticity tensor

Symmetries

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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Fourth-order tensor: linear application between second-order tensors,

$$\boldsymbol{\sigma}(\boldsymbol{x}, t) = \mathbf{C}(\boldsymbol{x}) \boldsymbol{\varepsilon}(\boldsymbol{x}, t),$$
$$\sigma_{jk}(\boldsymbol{x}, t) = C_{jklm}(\boldsymbol{x}) \varepsilon_{lm}(\boldsymbol{x}, t), \quad \forall \boldsymbol{x} \in \Omega, \forall t \in \mathbb{R}.$$

- Minor symmetries from the symmetry of $\boldsymbol{\varepsilon}$ and $\boldsymbol{\sigma}$:

$$C_{jklm}(\boldsymbol{x}) = C_{kjl m}(\boldsymbol{x}),$$

$$C_{jklm}(\boldsymbol{x}) = C_{kjml}(\boldsymbol{x}).$$

- Major symmetry from thermodynamics first principle:

$$C_{jklm}(\boldsymbol{x}) = C_{lmjk}(\boldsymbol{x}).$$

- # coefficients: 81 $\xrightarrow{\text{minor symmetries}} 36 \xrightarrow{\text{major symmetry}} 21.$

Elasticity tensor

Isotropy

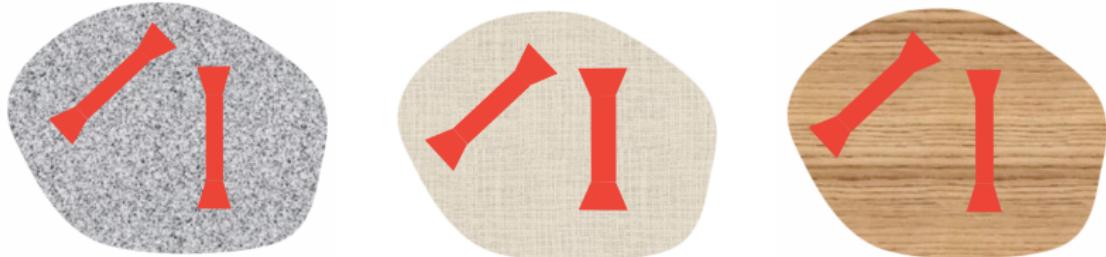
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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer



- Isotropy: the constitutive equation is the same whatever the orientation of the sample is,

$$\boldsymbol{\sigma} = \mathbf{C}\boldsymbol{\epsilon} \text{ (sample #1)}, \quad \boldsymbol{\sigma}^* = \mathbf{C}\boldsymbol{\epsilon}^* \text{ (sample #2)}.$$

- Letting \mathbf{R} be an arbitrary rotation ($\mathbf{R}\mathbf{R}^\top = \mathbf{I}$):

$$\boldsymbol{\epsilon}^* = \mathbf{R}\boldsymbol{\epsilon}\mathbf{R}^\top, \quad \boldsymbol{\sigma}^* = \mathbf{R}\boldsymbol{\sigma}\mathbf{R}^\top,$$

then:

$$\boldsymbol{\sigma}^* = \mathbf{R}\boldsymbol{\sigma}\mathbf{R}^\top = \mathbf{C}(\mathbf{R}\boldsymbol{\epsilon}\mathbf{R}^\top) \Rightarrow \mathbf{R}(\mathbf{C}\boldsymbol{\epsilon})\mathbf{R}^\top = \mathbf{C}(\mathbf{R}\boldsymbol{\epsilon}\mathbf{R}^\top).$$

Elasticity tensor

Isotropy

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Some
algebra

Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

■ Rivlin-Ericksen theorem:

$$\mathbf{R}(\mathbf{C}\boldsymbol{\varepsilon})\mathbf{R}^T = \mathbf{C}(\mathbf{R}\boldsymbol{\varepsilon}\mathbf{R}^T) \Rightarrow \mathbf{C}\boldsymbol{\varepsilon} = \alpha_0 \mathbf{I} + \alpha_1 \boldsymbol{\varepsilon} + \alpha_2 \boldsymbol{\varepsilon}^2,$$

where $\alpha_m(\text{Tr}(\boldsymbol{\varepsilon}), \text{Tr}(\boldsymbol{\varepsilon}^2), \text{Tr}(\boldsymbol{\varepsilon}^3))$, $m = 1, 2, 3$.



Ronald Rivlin [1915–2005]



Jerald Ericksen [1925–]

Elasticity tensor

Isotropy

- To the leading order in $O(\boldsymbol{\epsilon})$:

$$\mathbf{C}\boldsymbol{\epsilon} = \lambda \operatorname{Tr}(\boldsymbol{\epsilon})\mathbf{I} + 2\mu\boldsymbol{\epsilon}$$

where λ, μ are **Lamé's modulii**.

- This relationship can be inverted:

$$\begin{aligned}\boldsymbol{\epsilon} &= -\frac{\lambda}{2\mu(3\lambda+2\mu)} \operatorname{Tr}(\boldsymbol{\sigma})\mathbf{I} + \frac{1}{2\mu}\boldsymbol{\sigma} \\ &= \alpha \operatorname{Tr}(\boldsymbol{\sigma})\mathbf{I} + \beta\boldsymbol{\sigma}\end{aligned}$$

where α, β are obtained from measurements.

- More generally $\boldsymbol{\epsilon} = \mathbf{S}\boldsymbol{\sigma}$, where \mathbf{S} is the compliance (fourth-order) tensor.

Elasticity tensor

Isotropy

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Some
algebra

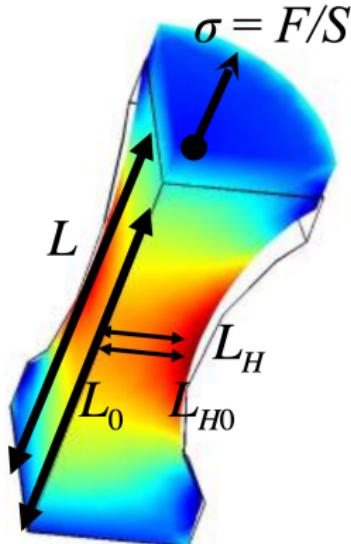
Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Traction test $\sigma = \frac{F}{S} e \otimes e$:



Young's modulus E :

$$E = \frac{\sigma_{ee}}{\varepsilon_{ee}} = \frac{\frac{F}{S}}{\frac{\Delta L}{L_0}} ;$$

Poisson's coefficient ν :

$$\nu = -\frac{\varepsilon_{hh}}{\varepsilon_{ee}} = -\frac{\frac{\Delta L_H}{L_{H0}}}{\frac{\Delta L}{L_0}} .$$

- E is in Pa (GPa), and ν is dimensionless.

Elasticity tensor

Isotropy

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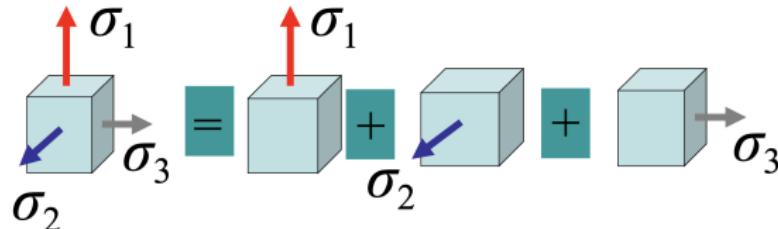
Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer



- σ and ϵ have the same principal directions:

- Along e.g. the principal direction #1:

$$\varepsilon_1 = \frac{\sigma_1}{E} - \nu \frac{\sigma_2}{E} - \nu \frac{\sigma_3}{E} .$$

- In the basis $(\mathbf{i}_1, \mathbf{i}_2, \mathbf{i}_3)$ of the principal directions:

$$\begin{aligned}\boldsymbol{\epsilon} &= \left(\frac{\sigma_1}{E} - \nu \frac{\sigma_2 + \sigma_3}{E} \right) \mathbf{i}_1 \otimes \mathbf{i}_1 + \left(\frac{\sigma_2}{E} - \nu \frac{\sigma_1 + \sigma_3}{E} \right) \mathbf{i}_2 \otimes \mathbf{i}_2 + \left(\frac{\sigma_3}{E} - \nu \frac{\sigma_1 + \sigma_2}{E} \right) \mathbf{i}_3 \otimes \mathbf{i}_3 \\ &= \frac{1+\nu}{E} \boldsymbol{\sigma} - \frac{\nu}{E} \text{Tr}(\boldsymbol{\sigma}) \mathbf{I} .\end{aligned}$$

Elasticity tensor

Isotropy

- Recap: linear elastic isotropy has 2 coefficients,

$$\boldsymbol{\sigma} = \mathbf{C}\boldsymbol{\varepsilon} = \lambda \operatorname{Tr}(\boldsymbol{\varepsilon})\mathbf{I} + 2\mu\boldsymbol{\varepsilon},$$

$$\boldsymbol{\varepsilon} = \mathbf{S}\boldsymbol{\sigma} = \frac{1+\nu}{E}\boldsymbol{\sigma} - \frac{\nu}{E}\operatorname{Tr}(\boldsymbol{\sigma})\mathbf{I}.$$

- $\lambda(\mathbf{x}), \mu(\mathbf{x})$ are Lamé's modulii, $E(\mathbf{x})$ is Young's modulus, $\nu(\mathbf{x})$ is Poisson's ratio:

$$\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}, \quad \mu = \frac{E}{2(1+\nu)}.$$

- $E > 0, \mu > 0, 3\lambda + 2\mu > 0$, and $-1 < \nu < 0.5$.

Thermoelasticity

Isotropic case

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Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

- Thermal strains in the isotropic case with temperature gradient ΔT :

$$\boldsymbol{\epsilon}_{\text{th}} = \alpha \Delta T \mathbf{I}$$

where α is the coefficient of linear thermal expansion.

- Total strain tensor:

$$\begin{aligned}\boldsymbol{\epsilon} &= \boldsymbol{\epsilon}_{\text{elas}} + \boldsymbol{\epsilon}_{\text{th}} \\ &= \mathbf{S} \boldsymbol{\sigma} + \alpha \Delta T \mathbf{I}.\end{aligned}$$

- Linear thermoelastic constitutive equation:

$$\boxed{\boldsymbol{\sigma} = \mathbf{C}(\boldsymbol{\epsilon} - \alpha \Delta T \mathbf{I})}.$$

Outline

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É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

1 Some algebra

- Vector & tensor products
- Vector & tensor analysis

2 Material behavior

3 4.1 Settlement of a soil layer

Settlement of a soil layer

Setup

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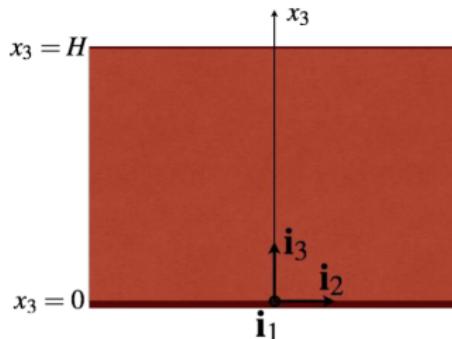
Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer



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$$\mathbf{u}(\mathbf{x}) = u_3(x_3)\mathbf{i}_3, \quad \mathbf{x} = (x_1, x_2, x_3) \in \Omega = \mathbb{R}^2 \times]0, H[.$$

Settlement of a soil layer

Solution

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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #0: $u_3(x_3)\mathbf{i}_3$?

- Symmetry of geometry **AND** loads.
- Translational invariance in $(x_1, x_2) \in \mathbb{R}^2$:

$$\mathbf{u}(x_1 + \Delta_1, x_2 + \Delta_2, x_3) = \mathbf{u}(x_1, x_2, x_3), \quad \forall (\Delta_1, \Delta_2) \in \mathbb{R}^2, \forall \mathbf{x} \in \Omega,$$

thus $\mathbf{u}(\mathbf{x}) = \mathbf{u}(x_3)$.

- Mirror symmetries:

$$\langle \mathbf{u}(x_3), \mathbf{i}_1 \rangle = -\langle \mathbf{u}(x_3), \mathbf{i}_1 \rangle = 0, \quad \forall x_3 \in]0, H[,$$

$$\langle \mathbf{u}(x_3), \mathbf{i}_2 \rangle = -\langle \mathbf{u}(x_3), \mathbf{i}_2 \rangle = 0, \quad \forall x_3 \in]0, H[,$$

thus $\mathbf{u}(x_3) = u_3(x_3)\mathbf{i}_3$.

Settlement of a soil layer

Solution

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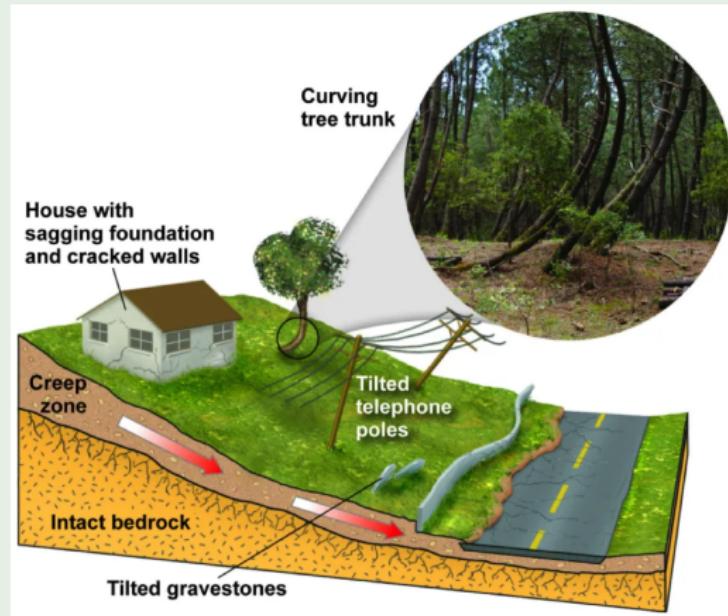
É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #0: $u_3(x_3)\mathbf{i}_3$?



Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #1: Differential equation satisfied by $x_3 \mapsto u_3(x_3)$?

- Local equilibrium equation $\text{Div}_x \sigma + f_v = \varrho a$, where $a = \mathbf{0}$ because "we can neglect the effects of inertia," and $f_v = -\varrho g i_3$.

Settlement of a soil layer

Solution

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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #1: Differential equation satisfied by $x_3 \mapsto u_3(x_3)$?

- Local equilibrium equation $\text{Div}_x \boldsymbol{\sigma} - \varrho g \mathbf{i}_3 = \mathbf{0}$;
- Constitutive relation $\boldsymbol{\sigma} = \lambda(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu \boldsymbol{\epsilon}$, because soil's "behavior is linear elastic and isotropic" and "we can adopt the framework of the infinitesimal deformation hypothesis."

Settlement of a soil layer

Solution

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É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #1: Differential equation satisfied by $x_3 \mapsto u_3(x_3)$?

- Local equilibrium equation $\text{Div}_x \boldsymbol{\sigma} - \varrho g \mathbf{i}_3 = \mathbf{0}$;
- Constitutive relation $\boldsymbol{\sigma} = \lambda(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu \boldsymbol{\epsilon}$;
- Compute $\boldsymbol{\epsilon}$:

$$\begin{aligned}\boldsymbol{\epsilon} &= \frac{\partial \mathbf{u}}{\partial x_j} \otimes_s \mathbf{i}_j \\ &= u'_3(x_3) \mathbf{i}_3 \otimes \mathbf{i}_3,\end{aligned}$$

where $u'_3(x_3) = \frac{du_3}{dx_3}$.

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

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- Local equilibrium equation $\text{Div}_x \boldsymbol{\sigma} - \varrho g \mathbf{i}_3 = \mathbf{0}$;
- Constitutive relation $\boldsymbol{\sigma} = \lambda(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu \boldsymbol{\epsilon}$;
- $\boldsymbol{\epsilon} = u'_3(x_3) \mathbf{i}_3 \otimes \mathbf{i}_3$;
- Compute $\boldsymbol{\sigma}$:

$$\begin{aligned}\boldsymbol{\sigma} &= \lambda u'_3(x_3) \mathbf{I} + 2\mu u'_3(x_3) \mathbf{i}_3 \otimes \mathbf{i}_3 \\ &= u'_3(x_3) [\lambda (\mathbf{i}_1 \otimes \mathbf{i}_1 + \mathbf{i}_2 \otimes \mathbf{i}_2) + (\lambda + 2\mu) \mathbf{i}_3 \otimes \mathbf{i}_3].\end{aligned}$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #1: Differential equation satisfied by $x_3 \mapsto u_3(x_3)$?

- Local equilibrium equation $\mathbf{Div}_x \boldsymbol{\sigma} - \varrho g \mathbf{i}_3 = \mathbf{0}$;
- Constitutive relation $\boldsymbol{\sigma} = \lambda(\text{Tr } \boldsymbol{\epsilon})\mathbf{I} + 2\mu\boldsymbol{\epsilon}$;
- $\boldsymbol{\epsilon} = u'_3(x_3)\mathbf{i}_3 \otimes \mathbf{i}_3$;
- $\boldsymbol{\sigma} = u'_3(x_3) [\lambda(\mathbf{i}_1 \otimes \mathbf{i}_1 + \mathbf{i}_2 \otimes \mathbf{i}_2) + (\lambda + 2\mu)\mathbf{i}_3 \otimes \mathbf{i}_3]$;
- Compute $\mathbf{Div}_x \boldsymbol{\sigma}$ (remind that $(\mathbf{a} \otimes \mathbf{b})\mathbf{c} = \langle \mathbf{b}, \mathbf{c} \rangle \mathbf{a}$):

$$\begin{aligned}\mathbf{Div}_x \boldsymbol{\sigma} &= \frac{\partial \boldsymbol{\sigma}}{\partial x_j} \mathbf{i}_j \\ &= (\lambda + 2\mu) u''_3(x_3) \mathbf{i}_3 ,\end{aligned}$$

$$\text{where } u''_3(x_3) = \frac{d^2 u_3}{dx_3^2}.$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #1: Differential equation satisfied by $x_3 \mapsto u_3(x_3)$?

- Local equilibrium equation $\mathbf{Div}_x \boldsymbol{\sigma} - \varrho g \mathbf{i}_3 = \mathbf{0}$;
- Constitutive relation $\boldsymbol{\sigma} = \lambda(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu \boldsymbol{\epsilon}$;
- $\boldsymbol{\epsilon} = u'_3(x_3) \mathbf{i}_3 \otimes \mathbf{i}_3$;
- $\boldsymbol{\sigma} = u'_3(x_3) [\lambda(\mathbf{i}_1 \otimes \mathbf{i}_1 + \mathbf{i}_2 \otimes \mathbf{i}_2) + (\lambda + 2\mu) \mathbf{i}_3 \otimes \mathbf{i}_3]$;
- $\mathbf{Div}_x \boldsymbol{\sigma} = (\lambda + 2\mu) u''_3(x_3) \mathbf{i}_3$;
- From the local equilibrium equation:

$$(\lambda + 2\mu) u''_3(x_3) \mathbf{i}_3 - \varrho g \mathbf{i}_3 = \mathbf{0} \iff \boxed{u''_3(x_3) = \frac{\varrho g}{\lambda + 2\mu}}.$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #2: Boundary conditions?

- $\{x_3 \leq 0\}$: a rock mass "fixed and perfectly rigid," hence $\mathbf{u}(x_3 \leq 0) = \mathbf{0}$ yielding $u_3(0) = 0$.

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #2: Boundary conditions?

- $u_3(0) = 0$;
- $\{x_3 = H\}$: "free of forces," hence $\sigma n \mid_{x_3=H} = \mathbf{0}$ that is:

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #2: Boundary conditions?

- $u_3(0) = 0$;
- $\{x_3 = H\}$: "free of forces," hence $\sigma n|_{x_3=H} = \mathbf{0}$ that is:

$$\begin{aligned}\sigma n|_{x_3=H} &= \sigma i_3|_{x_3=H} \\ &= (\lambda + 2\mu) u'_3(H) i_3 \\ &\quad (= u'_3(H)[\lambda(i_1 \otimes i_1 + i_2 \otimes i_2) + (\lambda + 2\mu)i_3 \otimes i_3] i_3)\end{aligned}$$

Therefore $u'_3(H) = 0$.

Settlement of a soil layer

Solution

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Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #2: $x_3 \mapsto u_3(x_3)$?

- $u_3(0) = 0$;
- $u'_3(H) = 0$;
- Therefore:

$$u'_3(x_3) = \frac{\varrho g(x_3 - H)}{\lambda + 2\mu},$$

$$u_3(x_3) = \frac{\varrho g x_3 (x_3 - 2H)}{2(\lambda + 2\mu)}.$$

Settlement of a soil layer

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Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

■ von Mises equivalent stress $\sigma_{\text{eq}} = \sqrt{3J_2(\boldsymbol{\sigma}^D)}$.

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

- von Mises equivalent stress $\sigma_{\text{eq}} = \sqrt{3J_2(\boldsymbol{\sigma}^D)}$;
- Compute $\boldsymbol{\sigma}^D = \boldsymbol{\sigma} - \frac{\text{Tr } \boldsymbol{\sigma}}{3} \mathbf{I}$:

$$\boldsymbol{\sigma} = u'_3(x_3) [\lambda(\mathbf{i}_1 \otimes \mathbf{i}_1 + \mathbf{i}_2 \otimes \mathbf{i}_2) + (\lambda + 2\mu)\mathbf{i}_3 \otimes \mathbf{i}_3]$$

$$\text{Tr } \boldsymbol{\sigma} = (3\lambda + 2\mu)u'_3(x_3)$$

$$\boldsymbol{\sigma}^D = \frac{2}{3}\mu u'_3(x_3) [-\mathbf{i}_1 \otimes \mathbf{i}_1 - \mathbf{i}_2 \otimes \mathbf{i}_2 + 2\mathbf{i}_3 \otimes \mathbf{i}_3]$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

- von Mises equivalent stress $\sigma_{\text{eq}} = \sqrt{3J_2(\boldsymbol{\sigma}^D)}$;
- $\boldsymbol{\sigma}^D = \frac{2}{3}\mu u'_3(x_3)[-i_1 \otimes i_1 - i_2 \otimes i_2 + 2i_3 \otimes i_3]$;
- Compute $J_2(\boldsymbol{\sigma}^D) = \frac{1}{2} \text{Tr}(\boldsymbol{\sigma}^{D2})$:

$$\boldsymbol{\sigma}^D = \frac{2}{3}\mu u'_3(x_3) [-i_1 \otimes i_1 - i_2 \otimes i_2 + 2i_3 \otimes i_3]$$

$$(\boldsymbol{\sigma}^D)^2 = \frac{4}{9}(\mu u'_3(x_3))^2 [i_1 \otimes i_1 + i_2 \otimes i_2 + 4i_3 \otimes i_3]$$

$$J_2(\boldsymbol{\sigma}^D) = \frac{4}{3}(\mu u'_3(x_3))^2.$$

Therefore $\sigma_{\text{eq}} = 2\mu |u'_3(x_3)|$.

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

- von Mises equivalent stress $\sigma_{\text{eq}} = \sqrt{3J_2(\boldsymbol{\sigma}^D)}$;
- $\boldsymbol{\sigma}^D = \frac{2}{3}\mu u'_3(x_3)[-i_1 \otimes i_1 - i_2 \otimes i_2 + 2i_3 \otimes i_3]$;
- $\sigma_{\text{eq}} = 2\mu |u'_3(x_3)|$;
- $\sigma_{\text{eq}} \leq \sigma_e$ provided that:

$$\frac{2\mu}{\lambda + 2\mu} \varrho g |x_3 - H| \leq \sigma_e, \quad \forall x_3 \in [0, H].$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

- $\sigma_{\text{eq}} \leq \sigma_e$ provided that:

$$\frac{2\mu}{\lambda + 2\mu} \varrho g |x_3 - H| \leq \sigma_e, \quad \forall x_3 \in [0, H].$$

- This criterion is first reached where $|x_3 - H|$ is maximum *i.e.* $x_3 = 0$;
- Conversely, the von Mises criterion is fulfilled provided that $H \leq H_{\max}$ where:

$$H_{\max} = \left[\frac{(\lambda + 2\mu)\sigma_e}{2\mu\varrho g} \right].$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #3: $\sigma_{\text{eq}} \leq \sigma_e$?

- $\sigma_{\text{eq}} \leq \sigma_e$ provided that $H \leq H_{\max}$ where:

$$H_{\max} = \frac{(\lambda + 2\mu)\sigma_e}{2\mu\varrho g}.$$

- **NB:** $\mu = \frac{E}{2(1+\nu)}$, $\lambda = \frac{\nu E}{(1+\nu)(1-2\nu)}$, thus:

$$H_{\max} = \left(\frac{1-\nu}{1-2\nu} \right) \frac{\sigma_e}{\varrho g}$$

independently of the Young's modulus E .

Settlement of a soil layer

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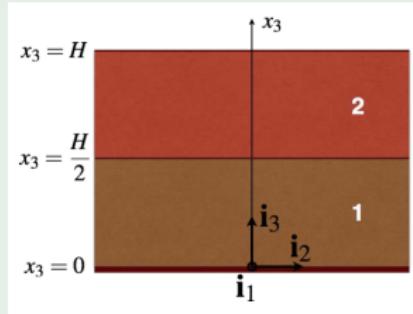
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products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #4: Two layers?



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■ Local equilibrium equation: unchanged! $\text{Div}_x \boldsymbol{\sigma} = \varrho g \mathbf{i}_3$.

Settlement of a soil layer

Solution

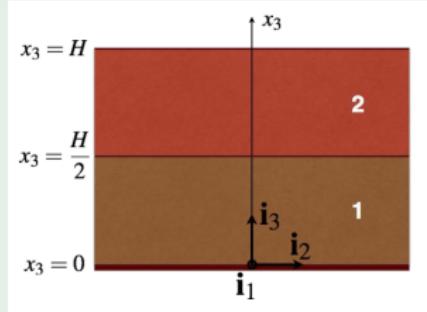
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Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis
Material
behavior

4.1
Settlement
of a soil
layer

Question #4: Two layers?



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■ Constitutive relation:

$$\boldsymbol{\sigma} = \begin{cases} = \lambda_1(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu_1 \boldsymbol{\epsilon}, & \mathbf{x} \in \Omega_1 = \mathbb{R}^2 \times \left[0, \frac{H}{2}\right], \\ = \lambda_2(\text{Tr } \boldsymbol{\epsilon}) \mathbf{I} + 2\mu_2 \boldsymbol{\epsilon}, & \mathbf{x} \in \Omega_2 = \mathbb{R}^2 \times \left[\frac{H}{2}, H\right]. \end{cases}$$

Settlement of a soil layer

Solution

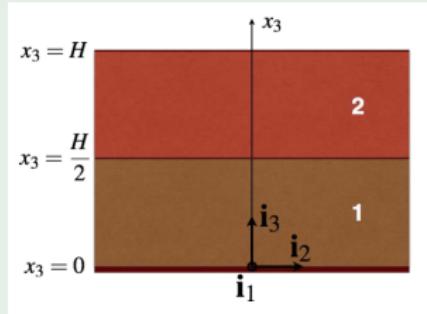
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É. Savin

Some
algebra
Vector &
tensor
products
Vector &
tensor
analysis
Material
behavior

4.1
Settlement
of a soil
layer

Question #4: Two layers?



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■ Boundary conditions:

$$\mathbf{u}(x_3) = \mathbf{0} \quad \text{on } \{x_3=0\},$$

$$\boldsymbol{\sigma}(x_3)\mathbf{i}_3 = \mathbf{0} \quad \text{on } \{x_3=H\},$$

$$\mathbf{u}(x_3)|_{\Omega_1} = \mathbf{u}(x_3)|_{\Omega_2}, \quad \boldsymbol{\sigma}(x_3)\mathbf{i}_3|_{\Omega_1} + \boldsymbol{\sigma}(x_3)(-\mathbf{i}_3)|_{\Omega_2} = \mathbf{0} \quad \text{on } \left\{x_3 = \frac{H}{2}\right\}.$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

Question #4: Two layers?

■ From Question #1:

$$u_3(x_3)|_{\Omega_1} = \frac{\varrho_1 g x_3^2}{2(\lambda_1 + 2\mu_1)} + A_1 x_3 + B_1 ,$$

$$u_3(x_3)|_{\Omega_2} = \frac{\varrho_2 g x_3^2}{2(\lambda_2 + 2\mu_2)} + A_2 x_3 + B_2 .$$

■ From boundary conditions:

$$0 = \frac{\varrho_2 g H}{\lambda_2 + 2\mu_2} + A_2 ,$$

$$\varrho_1 g \frac{H}{2} + (\lambda_1 + 2\mu_1) A_1 = \varrho_2 g \frac{H}{2} + (\lambda_2 + 2\mu_2) A_2 ,$$

$$\frac{\varrho_1 g H^2}{8(\lambda_1 + 2\mu_1)} + A_1 \frac{H}{2} + B_1 = \frac{\varrho_2 g H^2}{8(\lambda_2 + 2\mu_2)} + A_2 \frac{H}{2} + B_2 ,$$

$$B_1 = 0 .$$

Settlement of a soil layer

Solution

1EL5000/S4

É. Savin

Some
algebra

Vector &
tensor
products

Vector &
tensor
analysis

Material
behavior

4.1
Settlement
of a soil
layer

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■ From Question #1:

$$u_3(x_3)|_{\Omega_1} = \frac{\varrho_1 g x_3^2}{2(\lambda_1 + 2\mu_1)} + A_1 x_3 + B_1 ,$$

$$u_3(x_3)|_{\Omega_2} = \frac{\varrho_2 g x_3^2}{2(\lambda_2 + 2\mu_2)} + A_2 x_3 + B_2 .$$

■ From boundary conditions:

$$A_2 = -\frac{\varrho_2 g H}{\lambda_2 + 2\mu_2} ,$$

$$A_1 = -\frac{(\varrho_1 + \varrho_2)gH}{2(\lambda_1 + 2\mu_1)} ,$$

$$B_2 = \frac{gH^2}{8} \left(\frac{3\varrho_2}{\lambda_2 + 2\mu_2} - \frac{\varrho_1 + 2\varrho_2}{\lambda_1 + 2\mu_1} \right) ,$$

$$B_1 = 0 .$$