

MATH134: Worksheet 11

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Professor Ebru Bekyel

Eric Xia

Problem 1

(a) Keeping in mind that the velocity is $v(t) = x'(t) = \frac{dx}{dt}$, acceleration is $a(t) = \frac{dv}{dt} = x''(t)$ and using the equation $x'' + \omega^2 x = 0$ together with the chain rule show that

$$\frac{dv}{dx} = -\frac{1}{v} \cdot \omega^2 x$$

(b) Now use separation of variables together with the initial conditions $x(0) = x_0$ and $v(0) = 0$ to get an equation for v in terms of x . (What is the sign of v ?)

(c) From $\frac{dx}{dt} = v$, it follows that $dt = \frac{dx}{v}$. Compute the integral

$$\int_{x=x_0}^{x=0} dt$$

(with x being your variable of integration) to show that the time it takes the object to reach the origin is

$$\frac{\pi}{2\omega}$$

which you can see is independent of the starting position x_0 !

(a)

By definition,

$$\begin{aligned} x' &= \frac{dx}{dt} \\ \frac{1}{x'} &= \frac{dt}{dx} \\ \frac{1}{v} &= \frac{dt}{dx} \end{aligned}$$

Since $x'' = \frac{dv}{dt}$ we can rewrite the given equation in Leibniz notation and plug in $\frac{dt}{dx} = \frac{1}{v}$:

$$\begin{aligned} \frac{dv}{dt} + \omega^2 x &= 0 \\ \frac{dv}{dt} \cdot \frac{dt}{dx} &= -\omega^2 x \cdot \frac{dt}{dx} \\ \frac{dv}{dx} &= -\frac{1}{v} \omega^2 x \end{aligned}$$

(b)

$$\begin{aligned} \frac{dv}{dx} &= -\frac{1}{v} \cdot \omega^2 x \\ \int v dv &= -\omega^2 \int x dx \\ \frac{v^2}{2} + C_1 &= -\frac{\omega^2}{2} x^2 + C_2 \\ C_1 \text{ and } C_2 \text{ are arbitrary, let } C &= C_2 - C_1 \\ \frac{v^2}{2} &= -\frac{\omega^2}{2} x^2 + C \\ v^2 &= -\omega^2 x^2 + 2C \end{aligned}$$

Since C is arbitrary, redefine C from $2C$, substituting in the initial-value $(x_0, 0)$. This initial value is since $x(t)$ and $v(t)$ are functions of t which we do not consider in the new function $v(x)$.

$$\begin{aligned} v^2 &= -\omega^2 x^2 + C \\ 0 &= -\omega^2 x_0^2 + C \\ C &= \omega^2 x_0^2 \end{aligned}$$

Plugging this value in for C ,

$$v^2 = -\omega^2 x^2 + \omega^2 x_0^2$$

Since v should be negative,

$$\begin{aligned} v(x) &= -\sqrt{\omega^2 x_0^2 - \omega^2 x^2} \\ &= -\omega \sqrt{x_0^2 - x^2} \end{aligned}$$

(c)

$$\begin{aligned}
\int_{x=x_0}^{x=0} dt &= \int_{x=x_0}^{x=0} \frac{dx}{v} \\
&= \int_{x_0}^0 \frac{dx}{-\omega \sqrt{x_0^2 - x^2}} \\
&= -\frac{1}{\omega} \int_{x_0}^0 \frac{dx}{\sqrt{x_0^2 - x^2}} \\
&= -\frac{1}{\omega} \int_{x_0}^0 \frac{dx}{x_0 \sqrt{1 - \left(\frac{x}{x_0}\right)^2}} \\
u = \frac{x}{x_0} &\implies du = \frac{dx}{x_0} \\
&= -\frac{x_0}{\omega x_0} \int_{x_0}^0 \frac{du}{\sqrt{1 - u^2}} \\
&= -\frac{1}{\omega} \arcsin u \Big|_1^0 \\
&= \frac{1}{\omega} \cdot \frac{\pi}{2} \\
&= \frac{\pi}{2\omega}
\end{aligned}$$

Problem 2

Now we solve the initial value problem $x'' + \omega^2 x = 0$, $x(0) = x_0$, $x'(0) = 0$.

(a) Start by thinking about, i.e. guess, two familiar functions $f(t)$ and $g(t)$ which are solutions of $x'' + \omega^2 x = 0$. You can think about the simpler case $\omega = 1$ first. Then, $C_1 f(t) + C_2 g(t)$ will be the general solution to $x'' + \omega^2 x = 0$.

(b) Use your result above together with $x(0) = x_0$, $x'(0) = 0$ to show that the solution to the initial value problem is

$$x(t) = x_0 \cos(\omega t)$$

(a)

Problem 3

Now we are moving onto two dimensions with a parametric curve $(x(\theta), y(\theta))$. The object is sliding down, under (constant) gravitational force on this curve. The length of the curve from its lowest point is the arclength $s(t)$. If the arclength satisfies the initial value problem

$$s'' + \omega^2 s = 0, s(0) = s_0, s'(0) = 0$$

then the time it takes for the object to reach the bottom will be independent of s_0 , where it starts in the curve, as established above. We want to find the parametric equations of a curve whose arclength satisfies this differential equation where the acceleration is the result of a constant gravitational force.

A component of the gravitational force is normal to the curve and cancels out with the reaction force F_N from the surface of its path. The tangential component of the gravitational force F_T causes the acceleration s'' . Let θ be the angle the tangent to the curve makes with the x -axis as shown (negative in the picture).

- Compute F_T in terms of the gravitational force mg and θ .
- Use the fact that $s'' + \omega^2 s = 0$ (and $F = ma$) to relate θ and s .
- Differentiate your result above to get

$$ds = \frac{g}{\omega^2} \cos \theta$$

- Relate the slope of the curve $\frac{dy}{dx}$ to the angle θ .
- Use the definition of the arclength differential ds to show

$$ds = \frac{dx}{\cos \theta} \text{ and } ds = \frac{dy}{\sin \theta}$$

- Now use your results in (c) and (e) above to express dx and dy in the form $F(\theta)d\theta$ ready to integrate with respect to θ .
- Integrate to get x and y with θ as a parameter to get

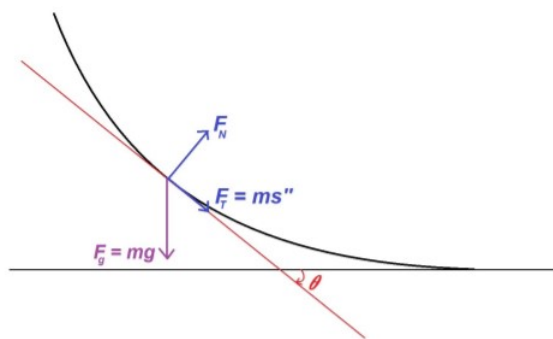
$$x = \frac{g}{4\omega^2}(2\theta - \sin 2\theta) + C_x \text{ and } y = -\frac{g}{4\omega^2} \cos 2\theta + C_y$$

where C_x and C_y are the integration constants.

(a)

$$F_T = mg \sin \theta$$

(b)



Problem 4

Now, we want to see what the shape of this parametric curve is. You can graph it before you go on (choose random integration constants) and try to make a guess.

- A cycloid is the path of a point on the circumference of a circle as the circle rolls on a surface. To get the equations for $(x(\theta), y(\theta))$, you can use the picture below.

Find equations of the coordinates of point P which initially is at the bottom of the circle of radius r when the center of the circle is at $(0, r)$. The parameter is θ , which keeps track of how much the circle has rolled, at its starts at $\theta = 0$.

- Now make the change of coordinates to your result in 3(g) including choices for the integration constants so the equations match.
- Express the time it takes the object to fall in terms of r and g . It does not depend on the initial starting position of the object!

