

Homework 1, Monte Carlo Simulation

Erik Norlin

November, 2022

2.1a

The time to reach the equilibrium state with $E = 2k_bT$ was approximately 40000 time steps most of the runs. After this time the position probability barely fluctuates. The transition frequency shows that the particle make any transitions.

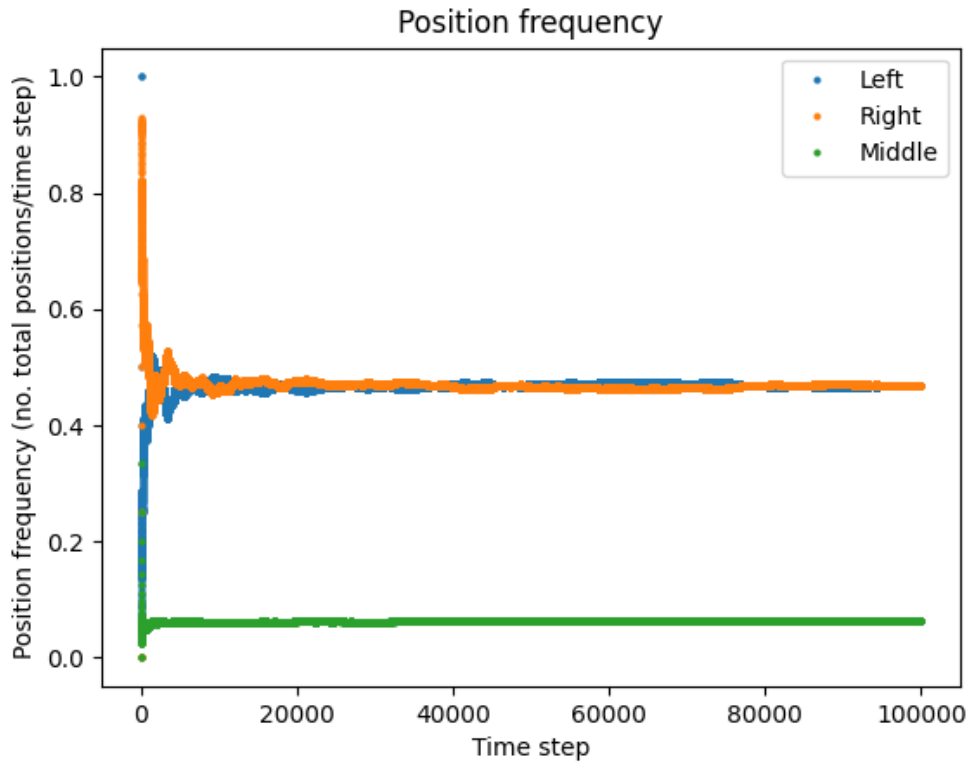


Figure 1: Position frequency (position probability) with $E = 2k_bT$.

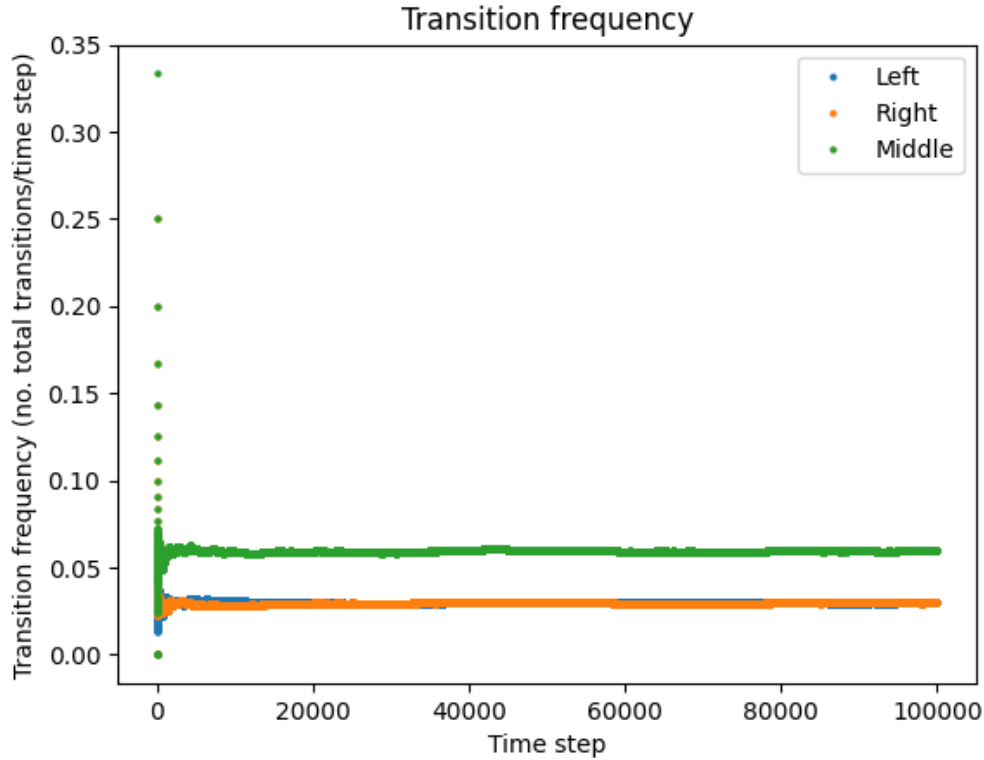


Figure 2: Transition frequency with $E = 2k_bT$.

Table 1: Probability distribution of positions and transition frequency for $E = 2k_bT$.

| | <i>Left</i> | <i>Middle</i> | <i>Right</i> |
|-------------|-------------|---------------|--------------|
| Positions | 0.46774 | 0.06329 | 0.46897 |
| Transitions | 0.02992 | 0.05975 | 0.02983 |

2.1b

The results show that the lower the temperature is in relation to the energy, the particle tend to transition to another state with lower probability. In contrast, for higher temperature, the particle tend to transition with higher probability. This can be seen through out the figure series. The theory supports this as well. For very high temperatures $e^{-E/k_bT} \rightarrow 1$, and for very low temperatures $e^{-E/k_bT} \rightarrow 0$. This means that the probability for transitioning to another state increases as the temperature increases.

The time to reach equilibrium varied for the different temperatures. The time to reach equilibrium for extremely low temperature was instantaneous because of the very high probability to never change state. For very high temperature, the probabilities of changing state at any state becomes almost equivalent. With approximately equal probabilities the equilibrium is reached almost instantly. In between very high and very low temperatures it takes longer to reach equilibrium, the total time can vary from run to run and depends on the temperature. For some temperatures it can take up to 60,000 iterations to reach equilibrium.

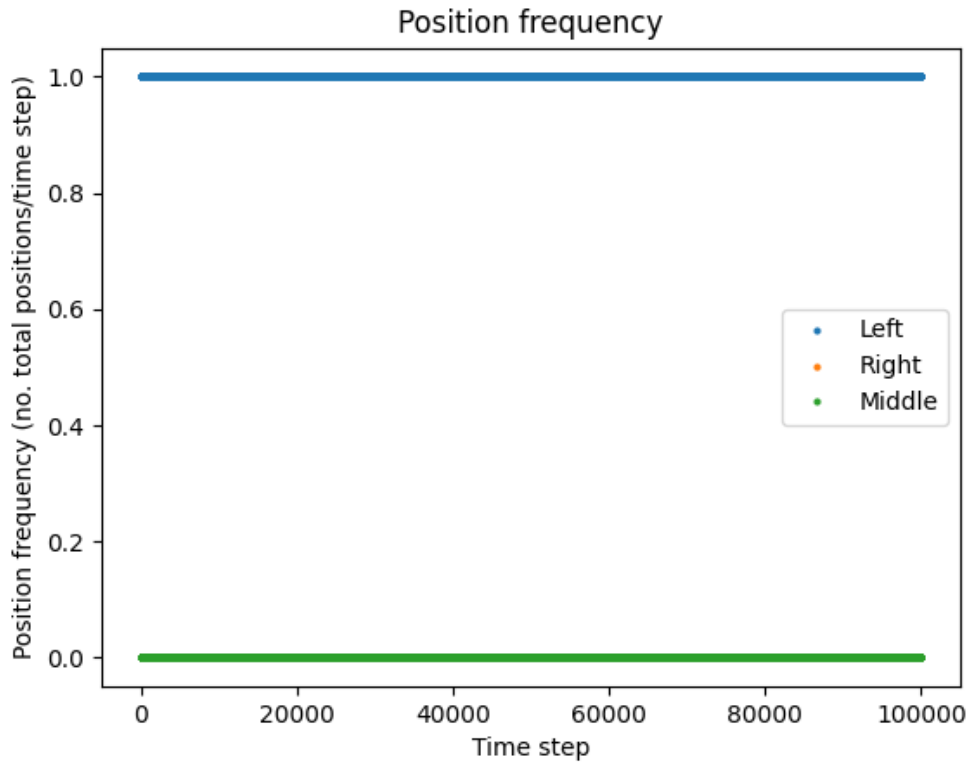


Figure 3: Position frequency (position probability) with $E = 10, T = 0.1$.

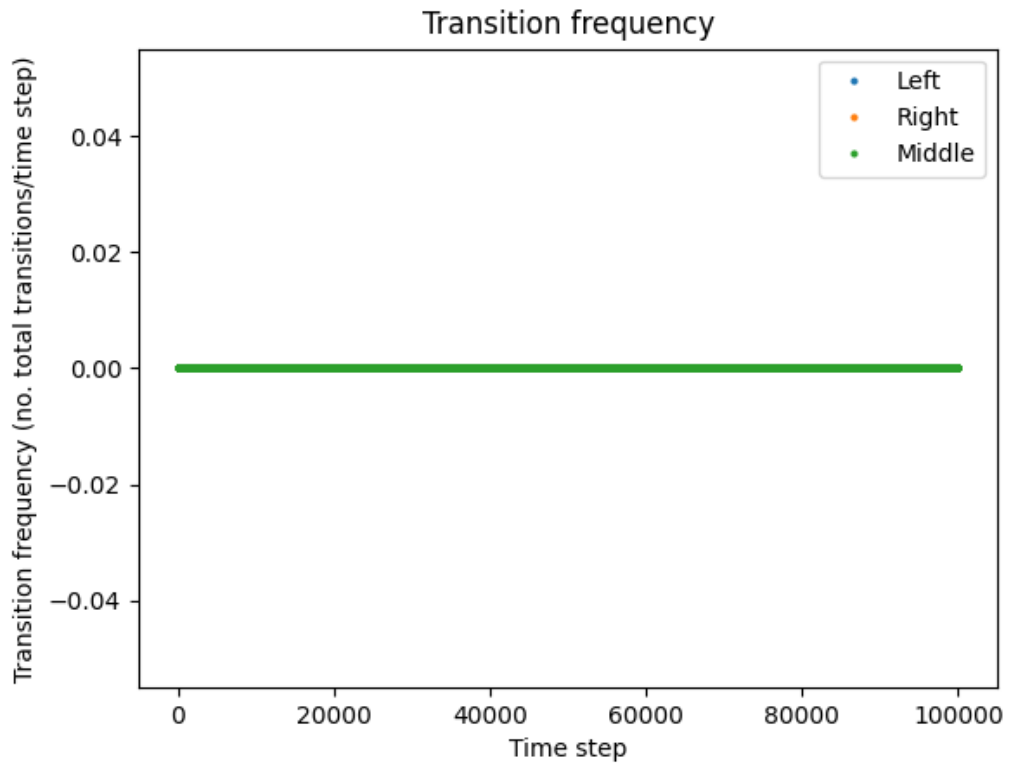


Figure 4: Transition frequency with $E = 10, T = 0.1$.

Table 2: Probability distribution of positions and transition frequency for $E = 10$, $T = 0.1$.

| | <i>Left</i> | <i>Middle</i> | <i>Right</i> |
|-------------|-------------|---------------|--------------|
| Positions | 1 | 0 | 0 |
| Transitions | 0 | 0 | 0 |

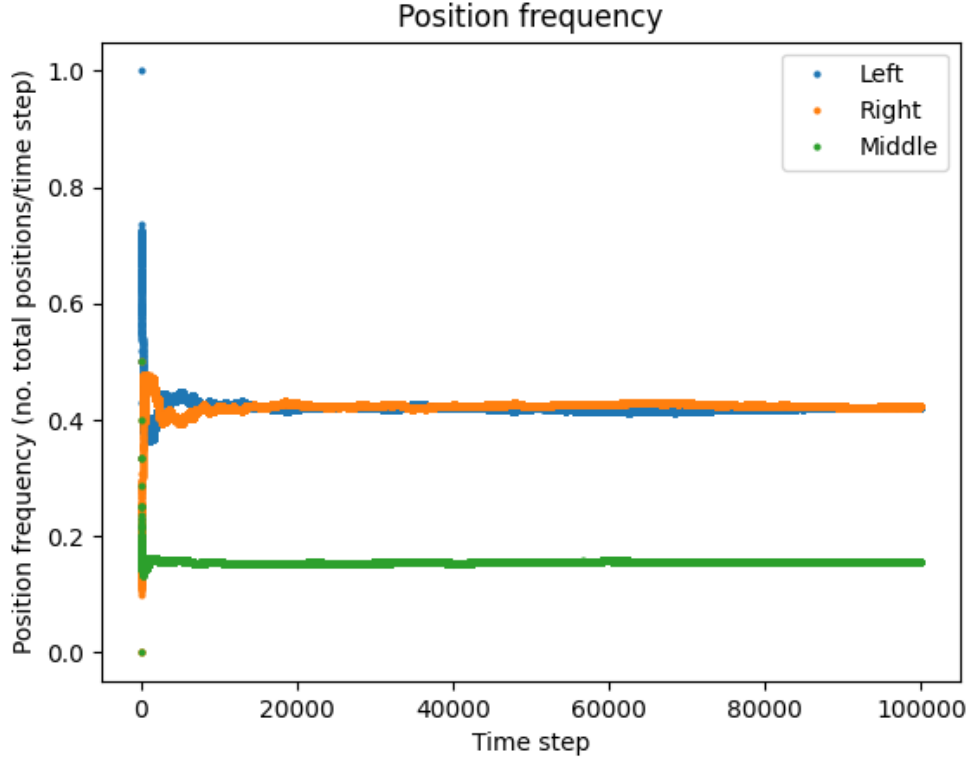


Figure 5: Position frequency (position probability) with $E = 10$, $T = 10$.

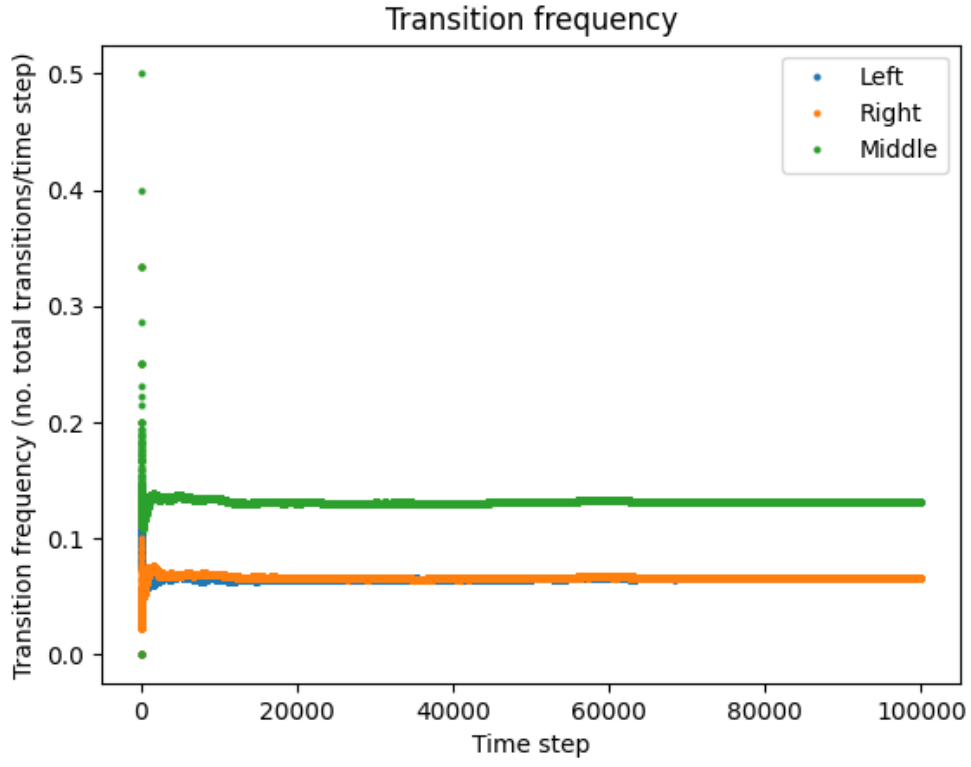


Figure 6: Transition frequency with $E = 10, T = 10$.

Table 3: Probability distribution of positions and transition frequency for $E = 10, T = 10$.

| | <i>Left</i> | <i>Middle</i> | <i>Right</i> |
|-------------|-------------|---------------|--------------|
| Positions | 0.42121 | 0.15641 | 0.42238 |
| Transitions | 0.06553 | 0.13177 | 0.06624 |

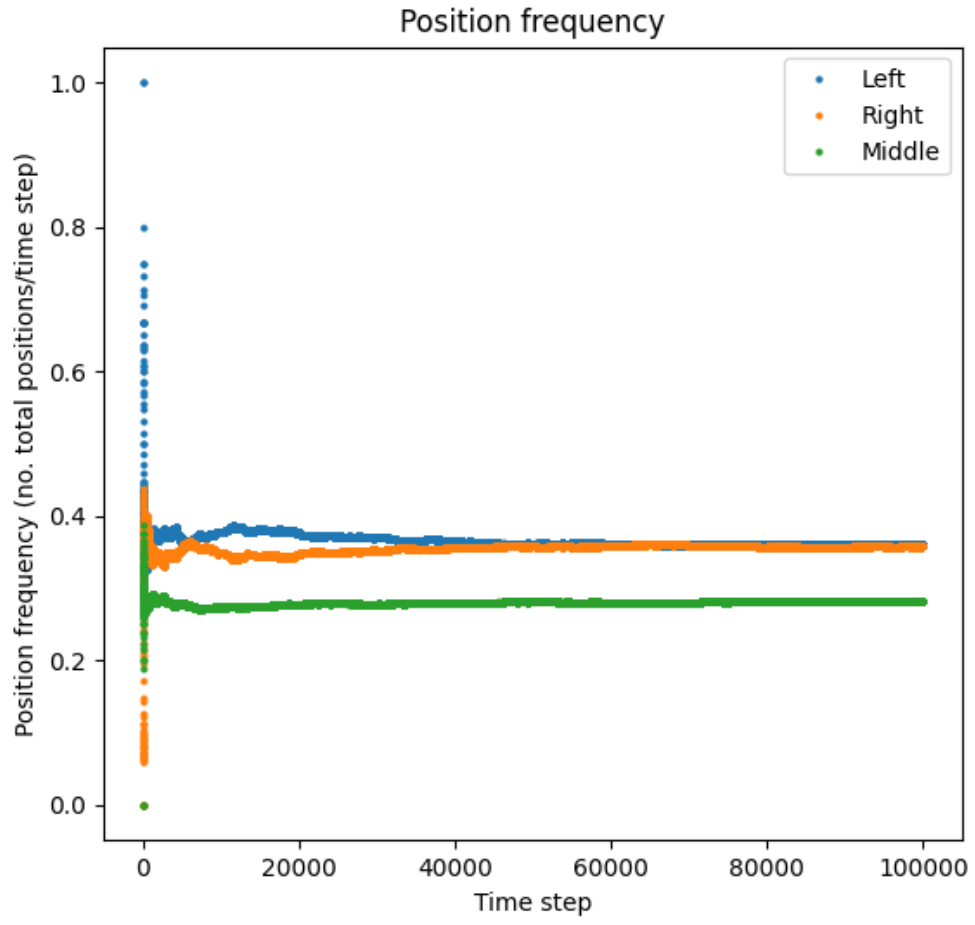


Figure 7: Position frequency (position probability) with $E = 10, T = 40$.

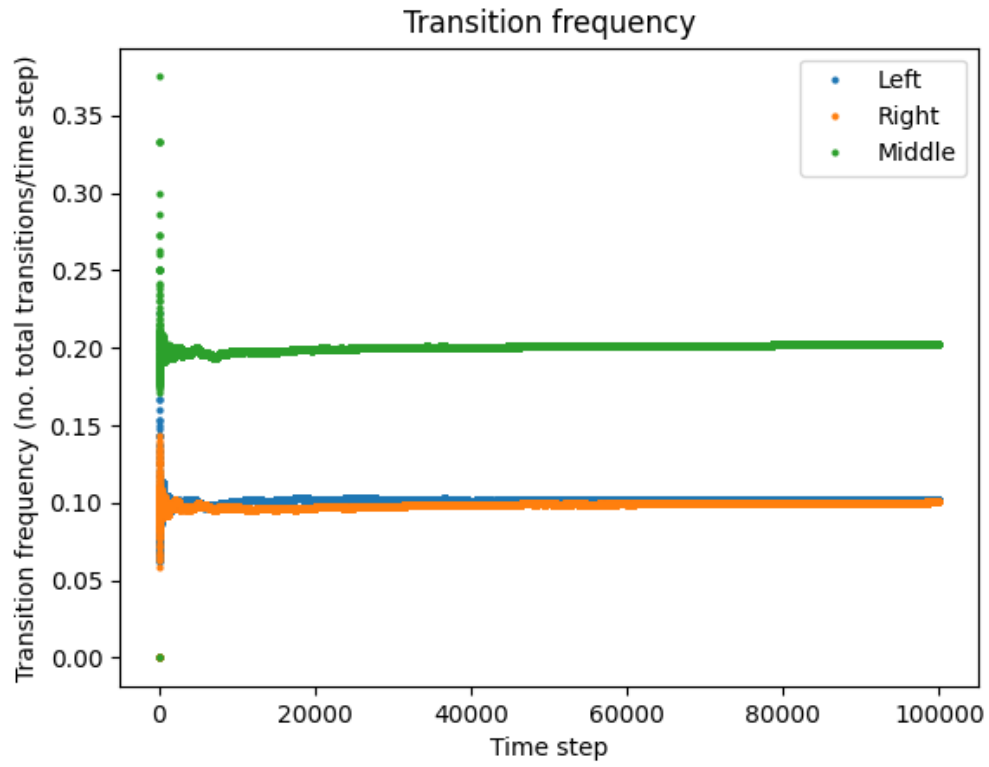


Figure 8: Transition frequency with $E = 10, T = 40$.

Table 4: Probability distribution of positions and transition frequency for $E = 40, T = 40$.

| | <i>Left</i> | <i>Middle</i> | <i>Right</i> |
|-------------|-------------|---------------|--------------|
| Positions | 0.3605 | 0.28202 | 0.35748 |
| Transitions | 0.10219 | 0.20264 | 0.10045 |

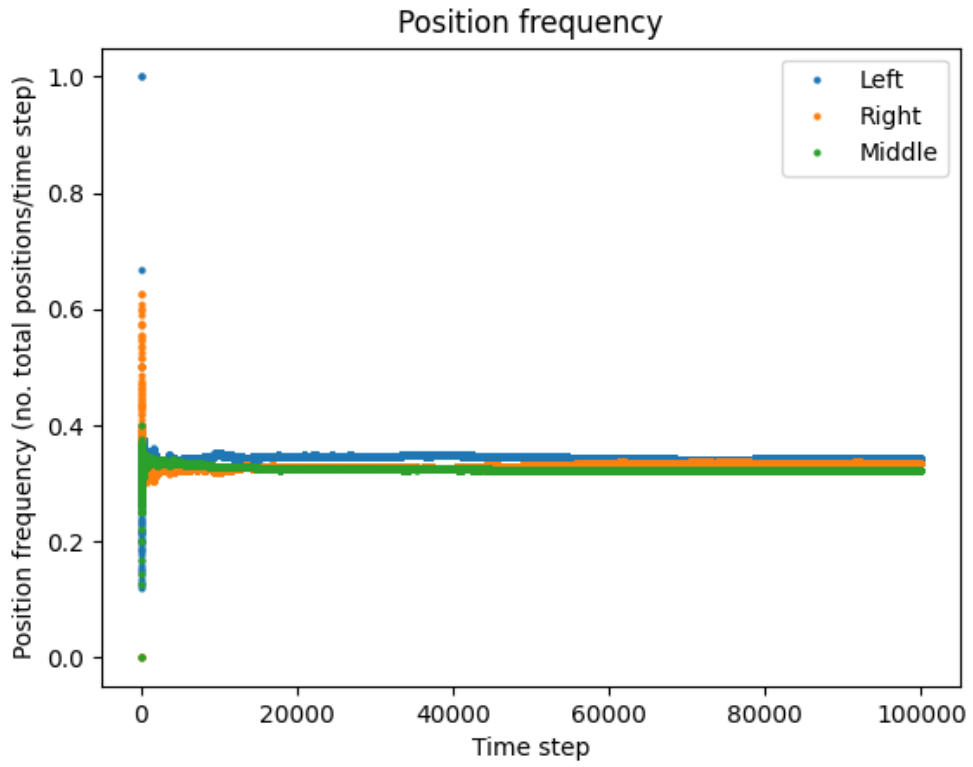


Figure 9: Position frequency (position probability) with $E = 10, T = 200$.

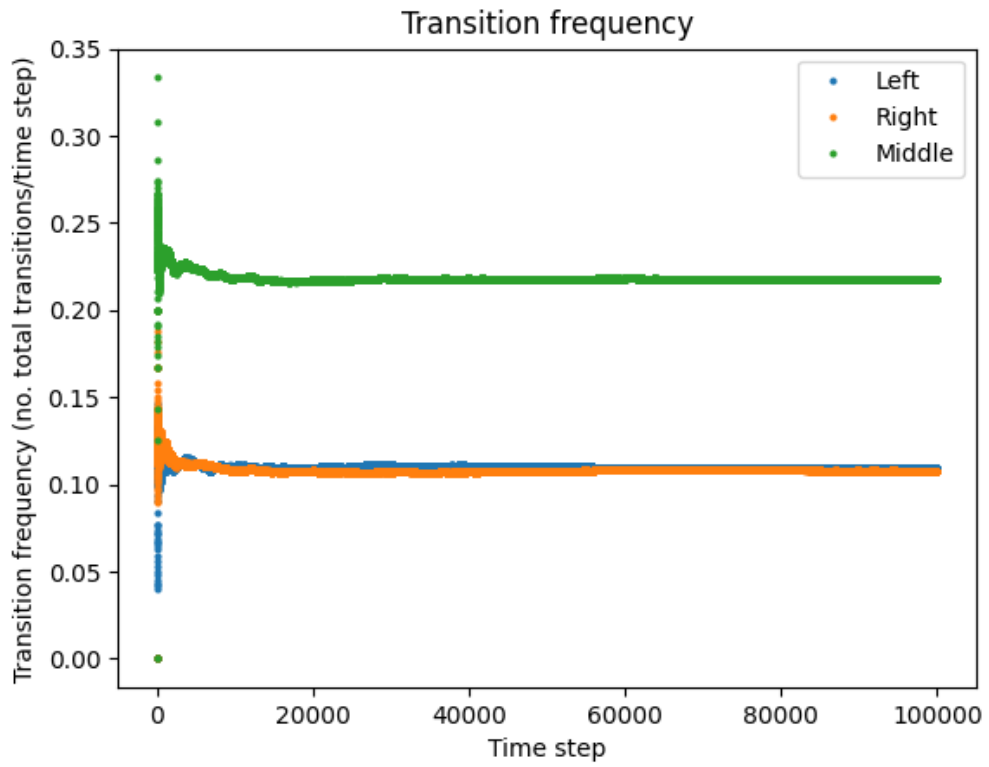


Figure 10: Transition frequency with $E = 10, T = 200$.

Table 5: Probability distribution of positions and transition frequency for $E = 10$, $T = 200$.

| | <i>Left</i> | <i>Middle</i> | <i>Right</i> |
|-------------|-------------|---------------|--------------|
| Positions | 0.34365 | 0.32168 | 0.33467 |
| Transitions | 0.10969 | 0.21727 | 0.10758 |

2.2b and c

We can see in the following figure that clusters of the same spins form at a subcritical temperature, that clusters don't form at a supercritical temperature and also that weaker defined clusters form at the critical temperature. Clusters of the same spins form in subcritical temperatures because the probability that a spin to be in a certain direction (+1 or -1) is larger than the opposite direction. Since the direction of every spin is probabilistically determined by the directions of the neighbouring spins, over time this results in clustering of the same spins. The probability of a spin being in a certain direction in supercritical temperatures become more random as the temperature increases. Instead of formed clusters, the direction of the spins results in a more stochastic nature, as one can see in the last row in the following figure. In metals, this is referred to as a *paramagnetic* property if the metal is subjected to a magnetic field without formation of clustering spins. This means that the metal react weakly to magnetic fields. The paramagnetic property can be observed in the spins in the simulation at a supercritical temperature $T = 6$ when the model is subjected to a magnetic field of $H = 1$. We can see that no clusters of the same spins are formed. However, the magnetic field still has an impact and shows a slight domination of yellow (+1) spins but there are still no clustering of the spins.

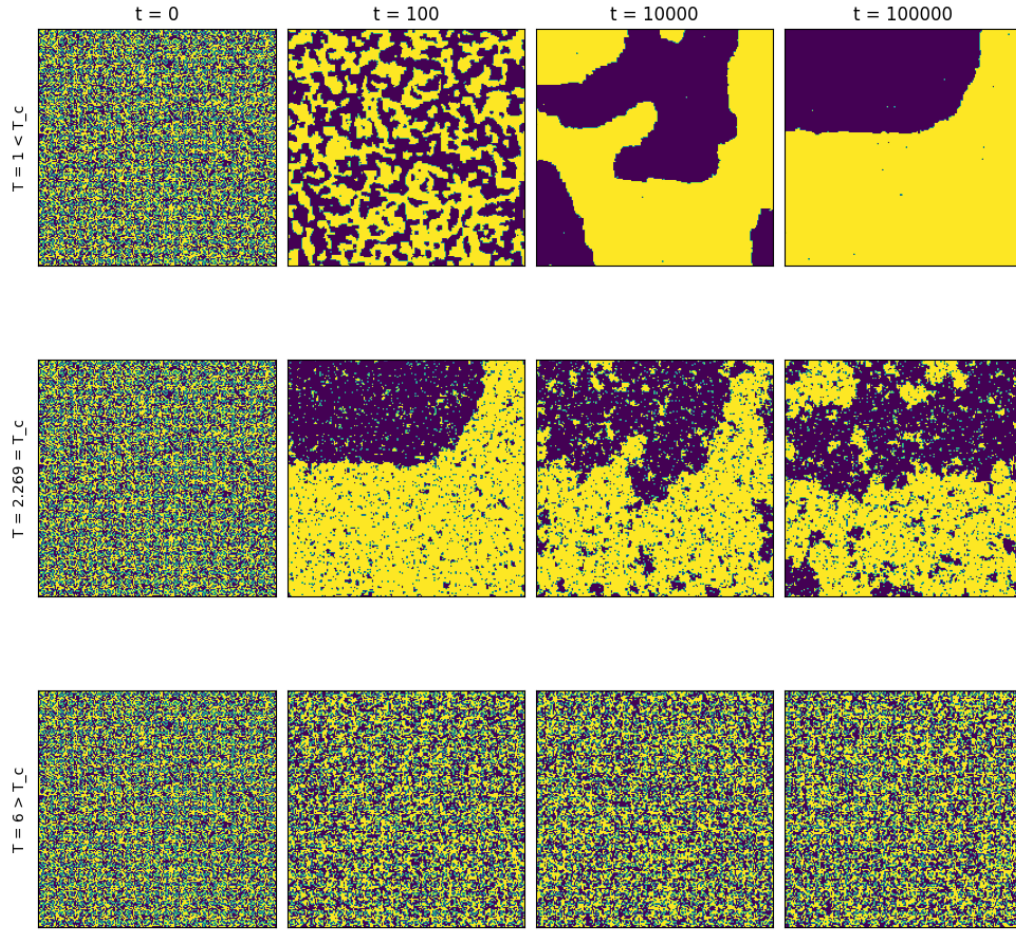


Figure 11: Spins after 100,000 iterations. Top row $T = 1$, middle row $T = 2.269$ and bottom row $T = 6$.

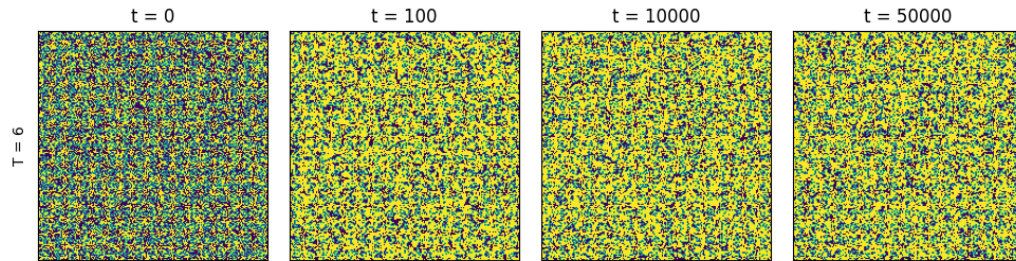


Figure 12: Paramagnetism $T = 6$, $H = 1$, 50,000 iterations.

2.2d

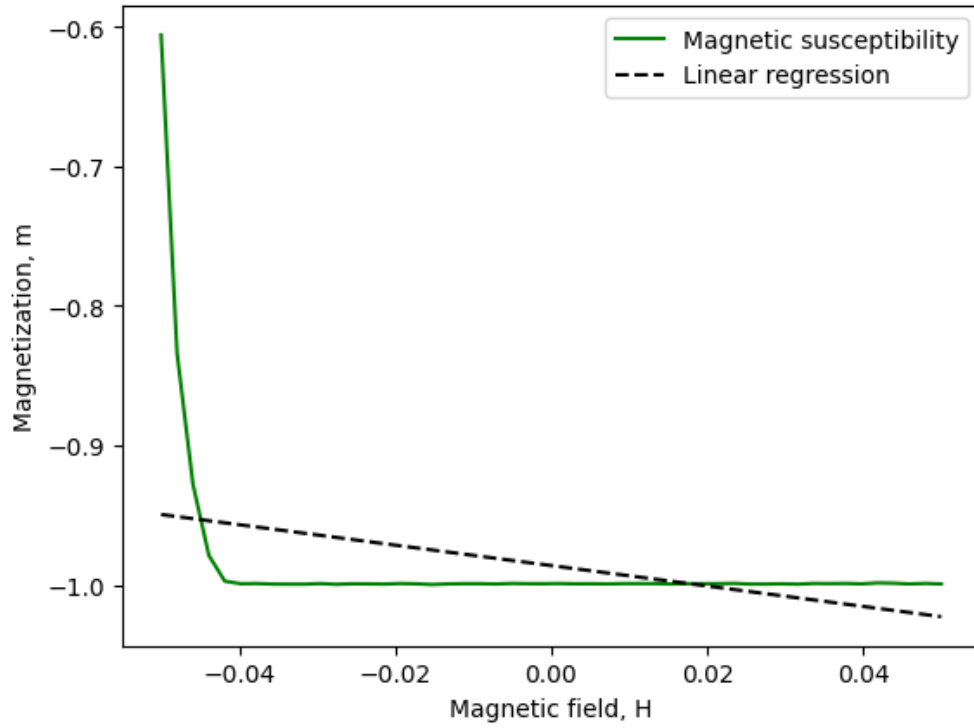


Figure 13: Spins after 1000 iterations. For $T = 1$, $X = -0.66$.

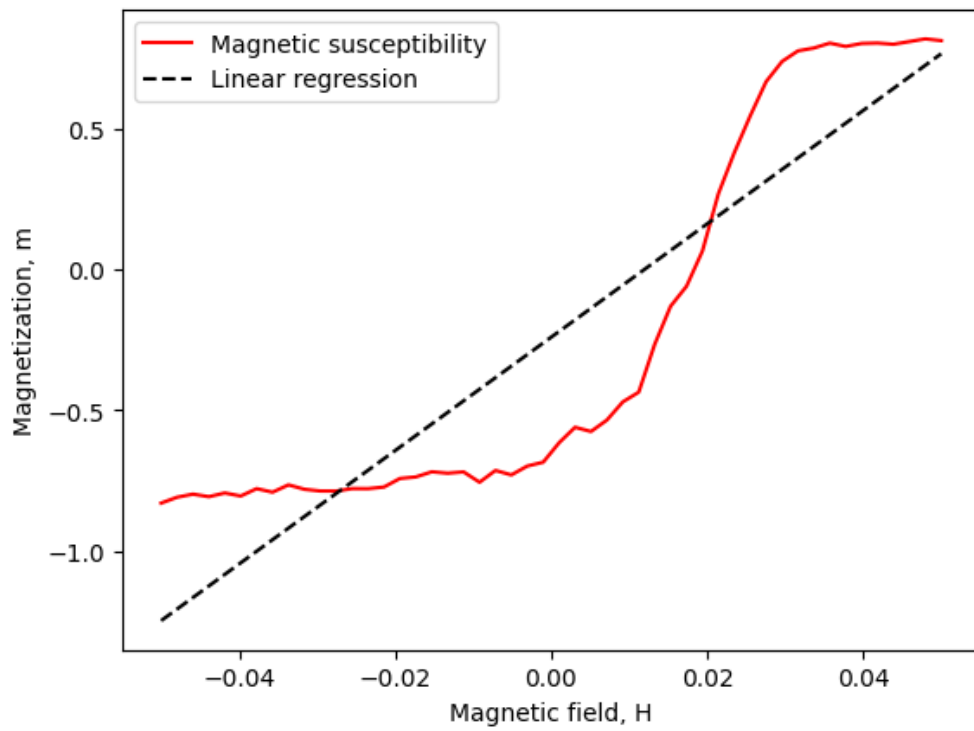


Figure 14: Spins after 1000 iterations. For $T = 2.269$, $X = 20.2$.

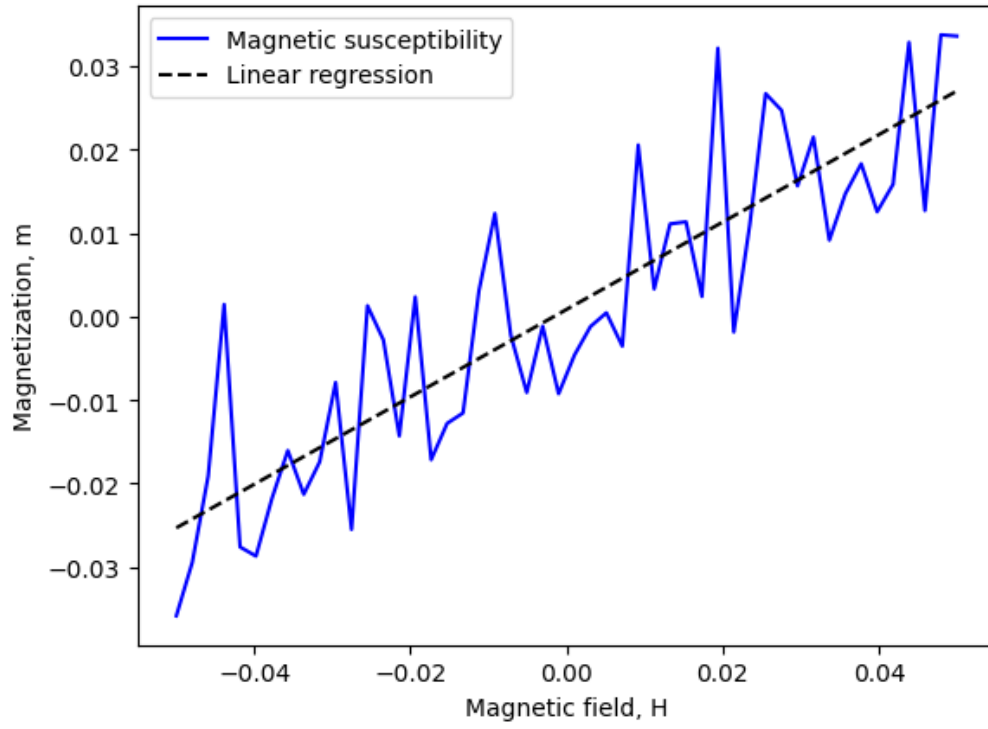


Figure 15: Spins after 1000 iterations. For $T = 5$, $X = 0.12$.