or as this (find out which),

$$\delta_{\nu} = (4.26)$$

$$\left(\frac{3}{10}a - \frac{1}{5}K\right)\delta_{cdm,0}a + \cos\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\left(c_1 - \frac{4}{5}\delta_{cdm,0}K^2\operatorname{Ci}\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\right) \tag{4.27}$$

$$-2\sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\left(c_2 + \frac{4}{5}\delta_{cdm,0}K^2\operatorname{Si}\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\right) \tag{4.28}$$

$$\delta_{\nu} \simeq c_1 \cos\left(\frac{2\sqrt{K}}{\sqrt{a}}\right) - \left(\frac{2}{5}\delta_{cdm,0}K^2\pi + c_2\right) \sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)$$
 (4.29)

$$\delta_{\nu} \simeq c_1 \cos\left(\frac{2\sqrt{K}}{\sqrt{a}}\right) - \left(\frac{2}{5}\delta_{0,2}K^2\pi + c_2\right) \sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)$$
 (4.30)

Ansatz:For early time

$$\delta_{\nu} \simeq A \cos \left(\frac{2\sqrt{K}}{\sqrt{a}} + \Phi \right)$$
 (4.31)

$$= A\cos\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\cos\left(\Phi\right) - A\sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right)\sin\left(\Phi\right) \tag{4.32}$$

$$\Phi = 0 \tag{4.33}$$

$$c_1 = A\cos\left(\Phi\right) \quad c_2 = A\sin\left(\Phi\right) - \frac{2}{5}\delta_{cdm,0}K^2\pi$$
 (4.34)

$$c_1 = A \quad c_2 = -\frac{2}{5} \delta_{cdm,0} K^2 \pi \tag{4.35}$$

$$\dot{\delta_{\nu}} = \frac{2\sqrt{K}}{2\sqrt{a^3}} \sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right) = \sqrt{\frac{K}{a^3}} \sin\left(\frac{2\sqrt{K}}{\sqrt{a}}\right) \tag{4.36}$$

$$A = \sqrt{\delta_{\nu}^2 + \left(\frac{\dot{\delta_{\nu}}}{\frac{\sqrt{K}}{\sqrt{a^3}}}\right)^2} = \sqrt{\delta_{\nu}^2 + \left(\frac{\dot{\delta_{\nu}}\sqrt{a^3}}{\sqrt{K}}\right)^2} \bigg|_{a=a_0}$$

$$(4.37)$$

4.2 second order

Dette er nok forkert. (det er forresten i conformal time, tror jeg).

$$\ddot{\delta}^{(2)} + \mathcal{H}\dot{\delta}^{(2)} + 2\mathcal{H}\theta^{(2)} + 2k^2c_s(a)^2\left(\delta^{(1)}\delta^{(1)} + \frac{1}{2}\delta^{(2)} + (\partial_j\delta^{(1)})\partial_j\nabla^{-2}\delta^{(1)}\right)$$
(4.38)

$$= 2\theta^{(1)}\theta^{(1)} + 2(\partial_i\partial_j\nabla^{-2}\theta^{(1)})^2 + 4(\partial_j\nabla^{-2}\theta^{(1)})(\partial_j\theta^{(1)})$$
(4.39)

$$+3H_0^2 \frac{\Omega_m}{a} \left(\delta^{(1)} \delta_{cdm}^{(1)} + \frac{1}{2} \delta_{cdm}^{(2)} + (\partial_j \delta^{(1)}) (\partial_j \nabla^{-2} \delta_{cdm}^{(1)}) \right)$$
(4.40)

nyt forsøg

4.2. SECOND ORDER 17

først har vi vores startligninger

$$\dot{\delta}^{(1)} = -\theta^{(1)} \quad \dot{\theta}^{(1)} = -\mathcal{H}\theta^{(1)} - \frac{3}{2} \frac{H_0^2 \Omega_M}{a} \delta_M^{(1)} - c_s^2 \nabla^2 \delta^{(1)}$$

$$\dot{\delta}_{\nu}^{(2)} = -2\partial_j \nabla^{-2} \theta^{(1)} \partial_j \delta^{(1)} - 2\delta^{(1)} \theta^{(1)} - \theta^{(2)}$$

$$(4.41)$$

$$\dot{\delta}_{\nu}^{(2)} = -2\partial_{i}\nabla^{-2}\theta^{(1)}\partial_{i}\delta^{(1)} - 2\delta^{(1)}\theta^{(1)} - \theta^{(2)} \tag{4.42}$$

$$\dot{\theta}^{(2)} = -\mathcal{H}\theta^{(2)} - 2\left(\partial_i\partial_j\nabla^{-2}\theta^{(1)}\partial_i\partial_j\theta^{(1)}\right) - 2\partial_j\nabla^{-2}\theta^{(1)}\partial_j\theta^{(1)}$$

$$\tag{4.43}$$

$$-\frac{3}{2}\frac{H_0^2\Omega_M}{a}\delta_M^{(2)} - c_s^2\nabla^2\delta^{(2)} \tag{4.44}$$

$$\ddot{\delta}^{(2)} = -2\partial_{j}\nabla^{-2}\theta^{(1)}\partial_{j}\delta^{(1)} - 2\partial_{j}\nabla^{-2}\dot{\theta}^{(1)}\partial_{j}\dot{\delta}^{(1)} - 2\dot{\delta}^{(1)}\theta^{(1)} - 2\delta^{(1)}\dot{\theta}^{(1)} - \dot{\theta}^{(2)}$$
(4.45)

$$\ddot{\delta}^{(2)} = -2\partial_j \nabla^{-2} \theta^{(1)} \partial_j \delta^{(1)} - 2\partial_j \nabla^{-2} \dot{\theta}^{(1)} \partial_j \dot{\delta}^{(1)} - 2\dot{\delta}^{(1)} \theta^{(1)} - 2\delta^{(1)} \dot{\theta}^{(1)}$$
 (4.46)

$$+\frac{3}{2}\frac{H_0^2\Omega_M}{a}\delta_M^{(2)} + c_s^2\nabla^2\delta^{(2)} + 2\left(\partial_i\partial_j\nabla^{-2}\theta^{(1)}\partial_i\partial_j\theta^{(1)}\right) + 2\partial_j\nabla^{-2}\theta^{(1)}\partial_j\theta^{(1)} + \mathcal{H}\theta^{(2)}$$
(4.47)

$$\ddot{\delta}^{(2)} = -2\partial_j \nabla^{-2} \theta^{(1)} \partial_j \delta^{(1)} - 2\partial_j \nabla^{-2} \dot{\theta}^{(1)} \partial_j \dot{\delta}^{(1)} - 2\dot{\delta}^{(1)} \theta^{(1)} - 2\delta^{(1)} \dot{\theta}^{(1)}$$
 (4.48)

$$+\frac{3}{2}\frac{H_0^2\Omega_M}{a}\delta_M^{(2)} + c_s^2\nabla^2\delta^{(2)} + 2\left(\partial_i\partial_j\nabla^{-2}\theta^{(1)}\partial_i\partial_j\theta^{(1)}\right) + 2\partial_j\nabla^{-2}\theta^{(1)}\partial_j\theta^{(1)}$$
(4.49)

$$+\mathcal{H}(-2\partial_{j}\nabla^{-2}\theta^{(1)}\partial_{j}\delta^{(1)}-2\delta^{(1)}\theta^{(1)}-\dot{\delta}_{\nu}^{(2)}) \quad (4.50)$$

(4.51)

$$\ddot{\delta}^{(2)} = 2\partial_j \nabla^{-2} \delta^{(1)} \partial_j \delta^{(1)} + 2\partial_j \nabla^{-2} (\mathcal{H}\delta^{(1)} - \frac{3}{2} \frac{H_0^2 \Omega_M}{a} \delta_M^{(1)} - c_s^2 \nabla^2 \delta^{(1)})^{(1)} \partial_j \dot{\delta}^{(1)}$$
(4.52)

$$+2\dot{\delta}^{(1)}\dot{\delta}^{(1)} - 2\delta^{(1)}(\mathcal{H}\delta^{(1)} - \frac{3}{2}\frac{H_0^2\Omega_M}{a}\delta_M^{(1)} - c_s^2\nabla^2\delta^{(1)})$$
 (4.53)

$$+\frac{3}{2}\frac{H_0^2\Omega_M}{a}\delta_M^{(2)} + c_s^2\nabla^2\delta^{(2)} + 2\left(\partial_i\partial_j\nabla^{-2}\dot{\delta}^{(1)}\partial_i\partial_j\dot{\delta}^{(1)}\right) + 2\partial_j\nabla^{-2}\dot{\delta}^{(1)}\partial_j\dot{\delta}^{(1)}$$
(4.54)

$$+\mathcal{H}(2\partial_{i}\nabla^{-2}\dot{\delta}^{(1)}\partial_{i}\delta^{(1)} + 2\delta^{(1)}\dot{\delta}^{(1)} - \dot{\delta}^{(2)}) \tag{4.55}$$

bedre opskrivning

$$\ddot{\delta}^{(2)} + \mathcal{H}\dot{\delta}^{(2)} = \frac{3}{2} \frac{H_0^2}{a} \left(2\partial_j \nabla^{-2} \delta_M^{(1)} \partial_j \delta^{(1)} + 2\delta^{(1)} \delta_M^{(1)} + \delta_M^{(2)} \right) \tag{4.56}$$

$$+c_s^2 \left(2\partial_j \delta^{(1)} \partial_j \delta^{(1)} + 2\delta^{(1)} \nabla^2 \delta^{(1)} + \nabla^2 \delta^{(2)}\right)$$
 (4.57)

$$4\partial_{j}\nabla^{-2}\dot{\delta}^{(1)}\partial_{j}\dot{\delta}^{(1)} + 2\dot{\delta}^{(1)}\dot{\delta}^{(1)} + 2\partial_{i}\partial_{j}\nabla^{-2}\dot{\delta}^{(1)}\partial_{i}\partial_{j}\nabla^{-2}\dot{\delta}^{(1)}$$

$$(4.58)$$

$$\ddot{\delta}^{(2)} + \mathcal{H}\dot{\delta}^{(2)} = \frac{3}{2} \frac{H_0^2}{a} \left(2\partial_j \nabla^{-2} \delta_M^{(1)} \partial_j \delta^{(1)} + 2\delta^{(1)} \delta_M^{(1)} + \delta_M^{(2)} \right) \tag{4.59}$$

$$+4\partial_{i}\nabla^{-2}\dot{\delta}^{(1)}\partial_{i}\dot{\delta}^{(1)} + 2\dot{\delta}^{(1)}\dot{\delta}^{(1)} + 2\partial_{i}\partial_{j}\nabla^{-2}\dot{\delta}^{(1)}\partial_{i}\partial_{j}\nabla^{-2}\dot{\delta}^{(1)}$$

$$(4.60)$$

$$a = \frac{1}{4}H_0^2\tau^2 \quad \delta^{(1)} = D\tilde{\delta} \quad \delta_M = a\tilde{\delta}_M \quad \mathcal{H} = H_0/\sqrt{a} = \frac{2}{\tau}$$
 (4.61)

(4.62)

Vi indsætter og ser hvad der sker,

$$\ddot{\delta}^{(2)} + \frac{2}{\tau}\dot{\delta}^{(2)} - c_s^2 \nabla^2 \delta^{(2)} - \frac{6}{\tau^2}\delta_M^{(2)} = \frac{6}{\tau^2} \left(2\partial_j \nabla^{-2} a \tilde{\delta}_M^{(1)} \partial_j D \tilde{\delta}^{(1)} + 2D\tilde{\delta}^{(1)} a \tilde{\delta}_M^{(1)} \right) \tag{4.63}$$

$$+4\dot{D}^2\partial_j\nabla^{-2}\tilde{\delta}^{(1)}\partial_j\tilde{\delta}^{(1)} + 2\dot{D}^2\tilde{\delta}^{(1)}\tilde{\delta}^{(1)} + 2\dot{D}^2\partial_i\partial_j\nabla^{-2}\tilde{\delta}^{(1)}\partial_i\partial_j\nabla^{-2}\tilde{\delta}^{(1)}$$
(4.64)

$$+c_s^2 \Big(2\partial_j D\tilde{\delta}^{(1)} \partial_j D\tilde{\delta}^{(1)} + 2D\tilde{\delta}^{(1)} \nabla^2 D\tilde{\delta}^{(1)} \Big)$$
 (4.65)

reducere

$$\ddot{\delta}^{(2)} + \frac{2}{\tau}\dot{\delta}^{(2)} - c_s^2 \nabla^2 \delta^{(2)} - \frac{6}{\tau^2} \delta_M^{(2)} = aD \frac{6}{\tau^2} \left(2\partial_j \nabla^{-2} \tilde{\delta}_M^{(1)} \partial_j \tilde{\delta}^{(1)} + 2\tilde{\delta}^{(1)} \tilde{\delta}_M^{(1)} \right) \tag{4.66}$$

$$+4\dot{D}^{2}\left(\partial_{j}\nabla^{-2}\tilde{\delta}^{(1)}\partial_{j}\tilde{\delta}^{(1)}+2\tilde{\delta}^{(1)}\tilde{\delta}^{(1)}+2\partial_{i}\partial_{j}\nabla^{-2}\tilde{\delta}^{(1)}\partial_{i}\partial_{j}\nabla^{-2}\tilde{\delta}^{(1)}\right) \tag{4.67}$$

$$+c_s^2 D^2 \left(2\partial_j \tilde{\delta}^{(1)} \partial_j \tilde{\delta}^{(1)} + 2\tilde{\delta}^{(1)} \nabla^2 \tilde{\delta}^{(1)}\right) \tag{4.68}$$

$$\ddot{\delta}^{(2)} + \frac{2}{\tau}\dot{\delta}^{(2)} - c_s^2 \nabla^2 \delta^{(2)} - \frac{6}{\tau^2}\delta_M^{(2)} = D\frac{3}{2}H_0^2 \left(2\partial_j \nabla^{-2}\tilde{\delta}_M^{(1)}\partial_j\tilde{\delta}^{(1)} + 2\tilde{\delta}^{(1)}\tilde{\delta}_M^{(1)}\right) \tag{4.69}$$

$$+4\dot{D}^2\Big(\partial_j\nabla^{-2}\tilde{\delta}^{(1)}\partial_j\tilde{\delta}^{(1)}+2\tilde{\delta}^{(1)}\tilde{\delta}^{(1)}+2\partial_i\partial_j\nabla^{-2}\tilde{\delta}^{(1)}\partial_i\partial_j\nabla^{-2}\tilde{\delta}^{(1)}\Big) \eqno(4.70)$$

$$+c_s^2 D^2 \Big(2 \partial_j \tilde{\delta}^{(1)} \partial_j \tilde{\delta}^{(1)} + 2 \tilde{\delta}^{(1)} \nabla^2 \tilde{\delta}^{(1)} \Big)$$
 (4.71)

Citat fra Steen "anden ordens termer for nuetrino bliver først relevant når neutrinoerne er stoppet med at oscillere og derfor kan vi tillade os at se bort fra lydhastigheds termer."

$$\ddot{\delta}^{(2)} + \frac{2}{\tau}\dot{\delta}^{(2)} - \frac{6}{\tau^2}\delta_M^{(2)} = D\frac{3}{2}H_0^2\Big(2\partial_j\nabla^{-2}\tilde{\delta}_M^{(1)}\partial_j\tilde{\delta}^{(1)} + 2\tilde{\delta}^{(1)}\tilde{\delta}_M^{(1)}\Big) \tag{4.72}$$

$$+4\dot{D}^2 \left(\partial_j \nabla^{-2} \tilde{\delta}^{(1)} \partial_j \tilde{\delta}^{(1)} + 2\tilde{\delta}^{(1)} \tilde{\delta}^{(1)} + 2\partial_i \partial_j \nabla^{-2} \tilde{\delta}^{(1)} \partial_i \partial_j \nabla^{-2} \tilde{\delta}^{(1)}\right) \tag{4.73}$$