

# Homework 7

## Erik Worth

1.

a)

```
j0[x_] := Sin[x] / x
y0[x_] := Cos[x] / x

jn[x_, n_] := Module[{u = j0[x]},
  Do[u = Simplify[(1/x)*D[u, x]], {n}]; Simplify[((-1)^n) (x^n) u]]
yn[x_, n_] := Module[{u = y0[x]},
  Do[u = Simplify[(1/x)*D[u, x]], {n}]; Simplify[((-1)^n) (-x^n) u]]
```

b)

```
jn[x, 0]
Sin[x]
-----
x

j1[x_] = jn[x, 1]
-x Cos[x] + Sin[x]
-----
x^2

jn[x, 10]
- 1
x^11
(55 x (11 904 165 - 1 670 760 x^2 + 51 597 x^4 - 468 x^6 + x^8) Cos[x] +
(- 654 729 075 + 310 134 825 x^2 - 18 918 900 x^4 + 315 315 x^6 - 1485 x^8 + x^10) Sin[x])

yn[x, 0]
Cos[x]
-----
x

y1[x_] = yn[x, 1]
-Cos[x] + x Sin[x]
-----
x^2

yn[x, 10]
1
x^11
((- 654 729 075 + 310 134 825 x^2 - 18 918 900 x^4 + 315 315 x^6 - 1485 x^8 + x^10) Cos[x] -
55 x (11 904 165 - 1 670 760 x^2 + 51 597 x^4 - 468 x^6 + x^8) Sin[x])
```

c)

```

Series[jn[x, 10], {x, 0, 14}]


$$\frac{x^{10}}{13\ 749\ 310\ 575} - \frac{x^{12}}{632\ 468\ 286\ 450} + \frac{x^{14}}{63\ 246\ 828\ 645\ 000} + O[x]^{15}$$


Series[yn[x, 10], {x, 0, -7}]


$$-\frac{654\ 729\ 075}{x^{11}} - \frac{34\ 459\ 425}{2\ x^9} - \frac{2\ 027\ 025}{8\ x^7} + \frac{1}{O[x]^6}$$

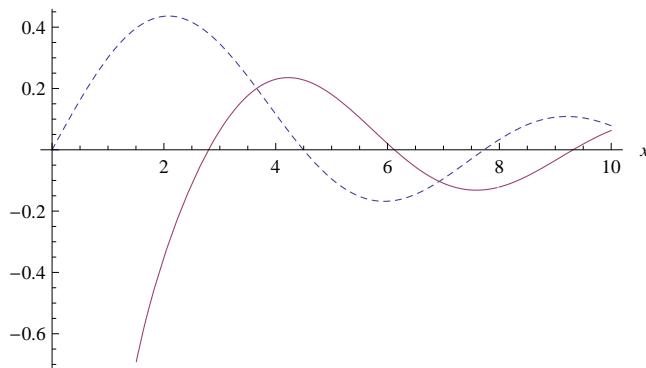

```

d)

```

Plot[{j1[x], y1[x]}, {x, 0, 10},
PlotStyle -> {Dashing[{0.01, 0.01}], Dashing[{1, 0}]}, AxesLabel -> Automatic]

```



The spherical Bessel functions with  $n = 0$   $j_0$  and  $y_0$  are defined as functions of  $x$ . The general spherical Bessel functions ( $j_n$  and  $y_n$ ) are defined as functions of  $n$  and  $x$  which iteratively apply Raleigh's formula to an internal variable  $u$ , which is initially defined to be  $j_0$  or  $y_0$ , respectively. In part b) the 0th, 1st and 10th spherical bessel functions are produced by  $j_n$  and  $y_n$ . In part d)  $j_n[x, 1]$  and  $y_n[x, 1]$  are plotted on the same graph.  $j_1$  is the dashed line;  $y_1$  is the solid line.

2.

a)

```

ser = Series[Tanh[x], {x, 0, 20}]


$$x - \frac{x^3}{3} + \frac{2x^5}{15} - \frac{17x^7}{315} + \frac{62x^9}{2835} - \frac{1382x^{11}}{155\ 925} + \frac{21\ 844x^{13}}{6\ 081\ 075} -$$


$$\frac{929\ 569x^{15}}{638\ 512\ 875} + \frac{6\ 404\ 582x^{17}}{10\ 854\ 718\ 875} - \frac{443\ 861\ 162x^{19}}{1\ 856\ 156\ 927\ 625} + O[x]^{21}$$


```

b)

```

InverseSeries[ser]


$$x + \frac{x^3}{3} + \frac{x^5}{5} + \frac{x^7}{7} + \frac{x^9}{9} + \frac{x^{11}}{11} + \frac{x^{13}}{13} + \frac{x^{15}}{15} + \frac{x^{17}}{17} + \frac{x^{19}}{19} + O[x]^{21}$$


```

The series ser is created as the series representation of  $\tanh(x)$  about  $x = 0$  up to order 20. InverseSeries then acts on ser to produce the series for  $\tan(x)$  about 0.

### 3.

a)

```
Assuming[{T < θ, T > 0, θ > 0},
{9 ((T/θ)^3) Integrate[(x^4) Exp[x] / (Exp[x] - 1)^2, {x, 0, Infinity}]}]

$$\left\{ \frac{12 \pi^4 T^3}{5 \theta^3} \right\}$$

```

b)

```
9 ((T/θ)^3) Integrate[(x^4) / (x)^2, {x, 0, θ/T}]

$$3$$

```

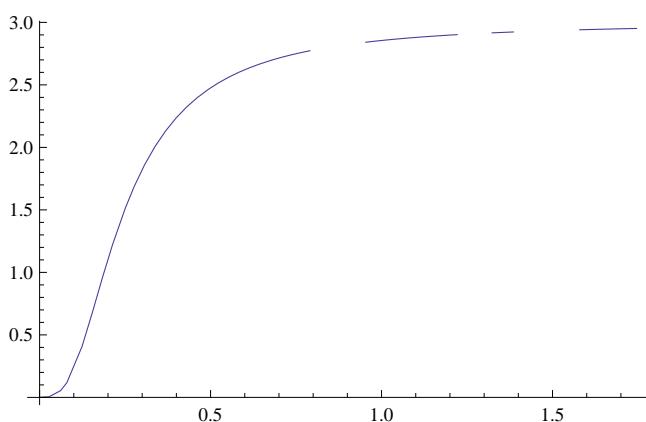
c)

```
Clear[x]
f[x_] = Assuming[{x > 0},
  9 ((1/x)^3) Integrate[(x^4) Exp[x] / (Exp[x] - 1)^2, {x, 0, x}] // Simplify

$$\frac{1}{x^3} 9 \left( -\frac{4 \pi^4}{15} + 4 i \pi x^3 + \frac{e^x x^4}{1 - e^x} + 4 x^3 \text{Log}[-1 + e^x] + 12 x (x \text{PolyLog}[2, e^x] - 2 \text{PolyLog}[3, e^x]) + 24 \text{PolyLog}[4, e^x] \right)$$

```

```
y = 1/x;
Plot[f[y], {x, 0, 100}, PlotRange → Automatic]
```

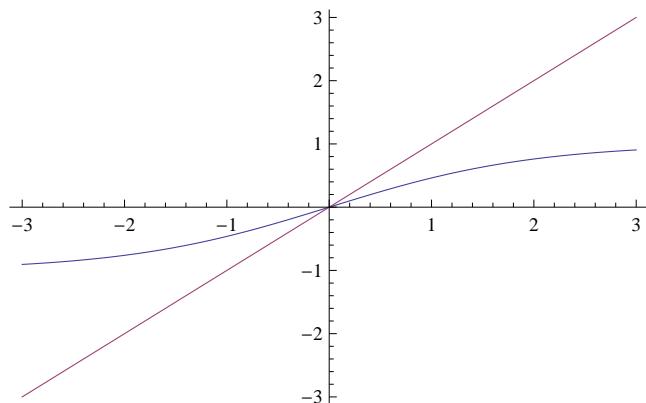


In part a), since  $T \ll \theta$ , the upper limit of integration is taken to be infinity. `Integrate[]` is used to analytically integrate the function, which is helped along by assuming that  $T < \theta$ , and both  $T$  and  $\theta$  are positive. In part b), since  $T \gg \theta$ ,  $x$  tends to be very small, so  $\text{Exp}[x]$  tends towards 1 and  $\text{Exp}[x] - 1$  can be approximated as  $x$ . `Integrate[]` is then used to do the simplified integral. In part c), the integral is performed without any approximations and plotted against  $T/\theta$ . The function appears to approach 3 asymptotically as  $T/\theta$  gets larger, which is consistent with the result from part b).

#### 4.

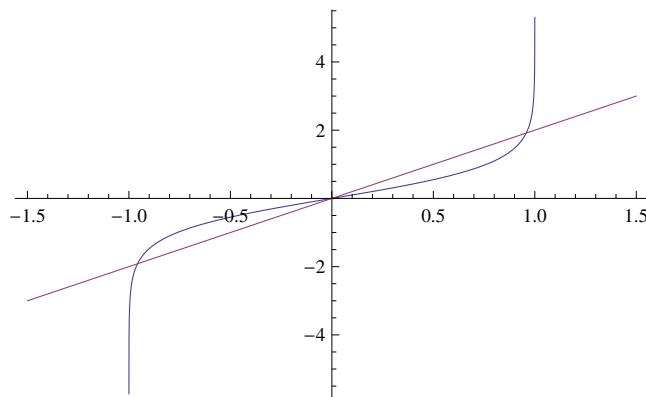
a)

```
Plot[{Tanh[0.5 u], u}, {u, -3, 3}]
```



b)

```
Plot[{ArcTanh[u], 2 u}, {u, -1.5, 1.5}]
```

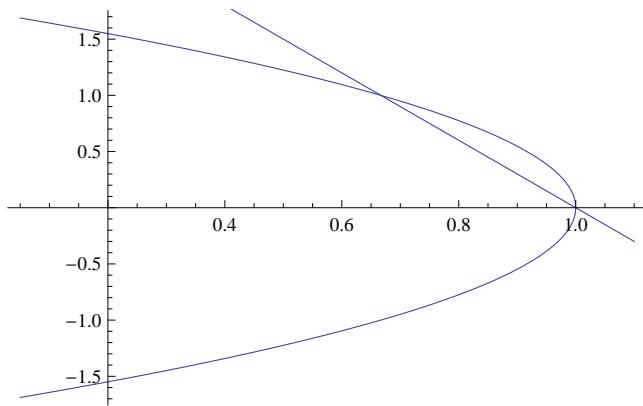


c) see attached paper

d-e)

```
g1 = Plot[m /. Solve[m^2 - 3(1 - T) == 0, m], {T, 0.05, 1.1}];  
g2 = Plot[3(1 - T), {T, 0.05, 1.1}];
```

```
Show[g1, g2]
```



In part a),  $T < J$ , and so  $m$  is multiplied by some factor between 0 and 1. I chose  $T/J$  to be 0.5 arbitrarily and plotted  $\text{Tanh}[0.5m]$  and clearly the only solution is at  $m = 0$ . For  $T > J$  the argument of  $\text{Tanh}[]$  must be greater than 1. I arbitrarily chose the ratio to be 2 and plotted  $\text{Tanh}[2m]$  and there are two non-zero solutions in the vicinity of  $\pm 1$ . In parts d) and e)  $m$  and the high temperature approximation for  $m^2$  are plotted against  $T/T_c$  on the same graph.

## 5.

```
Clear["Global`*"]

g = 9.81;
k = 0.004;

xmax[vv_] := (k = 0.004; v = vv;
  tfinal[theta_] := (sol = NDSolve [{x''[t] == -k x'[t] Sqrt[y'[t]^2 + x'[t]^2],
    y''[t] == -k y'[t] Sqrt[y'[t]^2 + x'[t]^2] - g, x[0] == 0, y[0] == 0,
    x'[0] == v Cos[theta], y'[0] == v Sin[theta]}, {x, y}, {t, 0, 100}];
  yy[t_] = y[t] /. sol[[1]]; xx[t_] = x[t] /. sol[[1]];
  t /. FindRoot[yy[t], {t, 1, 40}, MaxIterations → 500]);
  xfinal[theta_?NumericQ] := xx[tfinal[theta]];
  FindMaximum[xfinal[theta], {theta, 0.1, 1.3}] [[1]])

xmax2[vv_] := (k = 0; v = vv;
  tfinal[theta_] := (sol = NDSolve [{x''[t] == -k x'[t] Sqrt[y'[t]^2 + x'[t]^2],
    y''[t] == -k y'[t] Sqrt[y'[t]^2 + x'[t]^2] - g, x[0] == 0, y[0] == 0,
    x'[0] == v Cos[theta], y'[0] == v Sin[theta]}, {x, y}, {t, 0, 100}];
  yy[t_] = y[t] /. sol[[1]]; xx[t_] = x[t] /. sol[[1]];
  t /. FindRoot[yy[t], {t, 1, 40}, MaxIterations → 500]);
  xfinal[theta_?NumericQ] := xx[tfinal[theta]];
  FindMaximum[xfinal[theta], {theta, 0.1, 1.3}] [[1]])
```

a)

```
xmax[50]
```

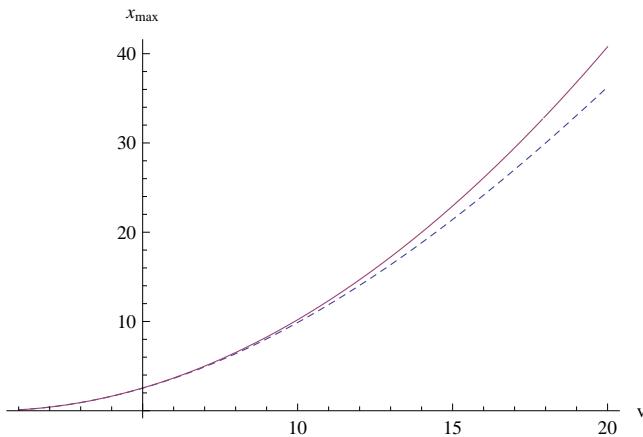
150.21

```
xmax2[50]
```

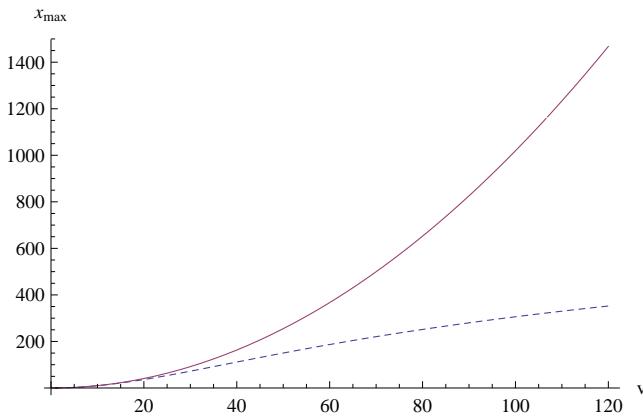
```
254.842
```

b)

```
Plot[{xmax[v], xmax2[v]}, {v, 1, 20}, AxesLabel -> {"v", "xmax"}, PlotStyle -> {Dashing[{0.01, 0.01}], Dashing[{1, 0}]}]
```



```
Plot[{xmax[v], xmax2[v]}, {v, 1, 120}, AxesLabel -> {"v", "xmax"}, PlotStyle -> {Dashing[{0.01, 0.01}], Dashing[{1, 0}]}]
```



Functions xmax and xmax2 are defined as functions of initial velocity vv. Xmax has a friction coefficient  $k = 0.004$ , while xmax2 does not have friction. Both functions define the function tfinal, which is a function of the launch angle  $\theta$ . It solves the equations of motion and produces the total time the ball spends in the air. Xfinal finds the x coordinate at tfinal. FindMaximum varies  $\theta$

from 0.1 to 1.3 and selects the largest value of  $x_{final}$ , producing the longest range possible.

We can see that the difference between the simulations with friction and without friction produce very different ranges; the frictionless launch goes over 100 meters farther.

In part b) the functions  $x_{max}$  and  $x_{max2}$  are plotted against initial velocity on the same graph. The frictionless case is the solid line, the case with friction is dashed. One graph is a low initial velocity case, where  $v$  is varied between 1 and 20 and we see that friction makes little difference; the maximum ranges only start to differ significantly as  $v$  gets larger than 15 m/s. In the second graph,  $v$  is varied between 1 and 120 and a huge difference in maximum range is evident. The frictionless case is increasing without bound, while the case with friction is plateauing. At  $v = 120$  m/s the difference between the two maximum ranges is around 1000 meters.

## 6.

a)

```

Clear["Global`*"]
f[x_] := 4 λ x (1 - x);

λ = 0.2; g1 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.2"];

λ = 0.6; g2 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.6"];

λ = 0.8; g3 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.8"];

λ = 0.84; g4 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.84"];

λ = 0.90; g5 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.90"];

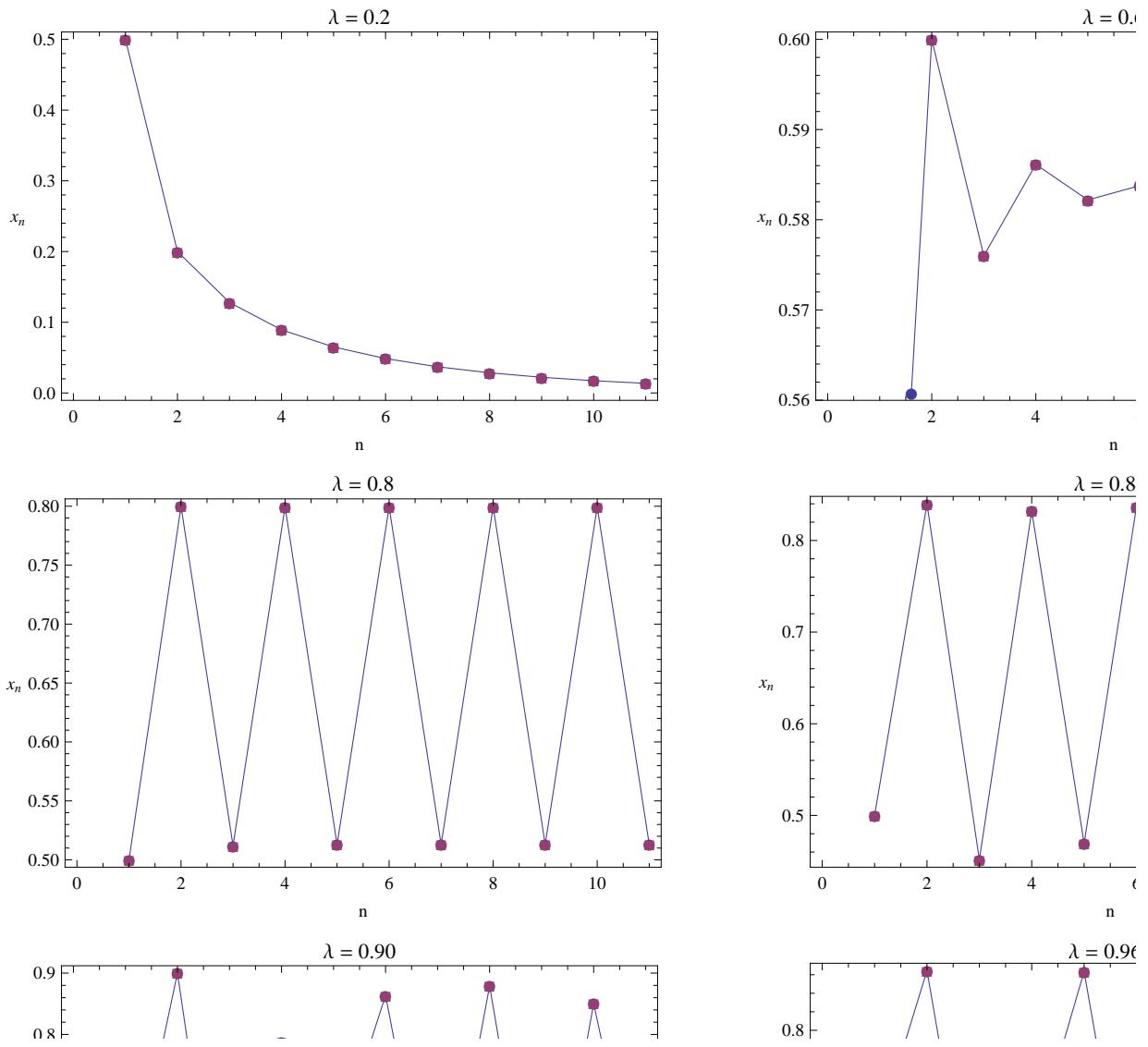
λ = 0.961; g6 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.961"];

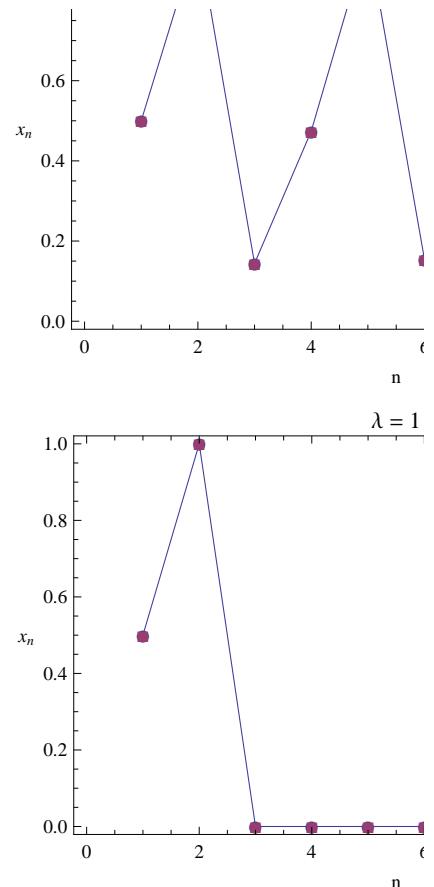
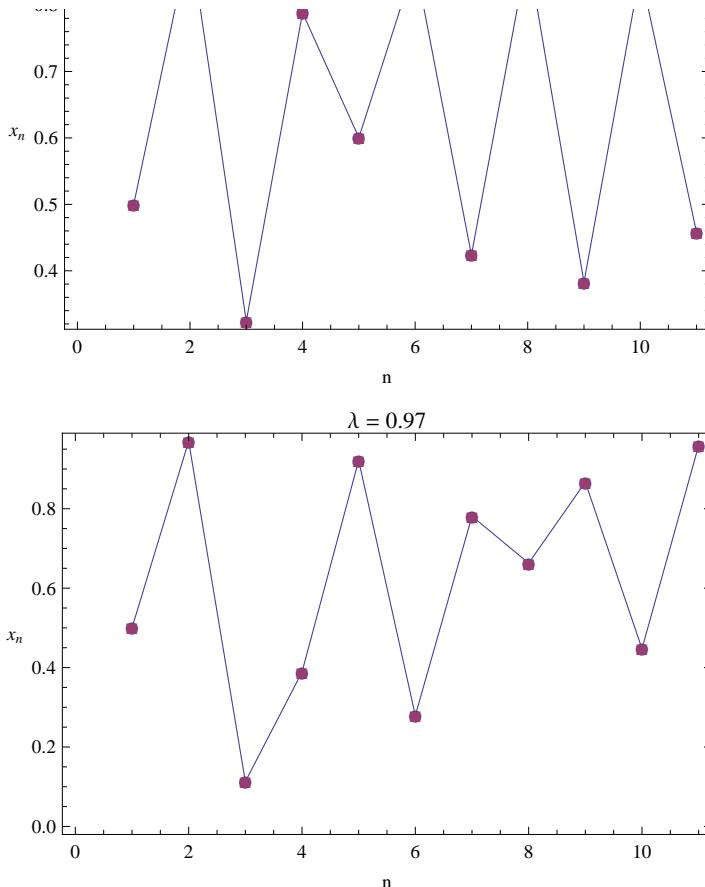
λ = 0.97; g7 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 0.97"];

λ = 1;
g8 = ListPlot [{NestList[f, 0.5, 10], NestList[f, 0.5, 10]} ,
 Joined → {True, False}, Axes → False, Frame → True, FrameLabel → {"n", "xn"} ,
 PlotMarkers → Automatic, RotateLabel → False, PlotLabel → "λ = 1"];

Show[GraphicsGrid[{{g1, g2}, {g3, g4}, {g5, g6}, {g7, g8}}], Spacings → Scaled[0.05]]

```





```

lya[l_, xinit_, n_, ndrop_] := (λ = l; xlist = Drop[NestList[f, xinit, n], ndrop + 1];
Apply[Plus, Log[Abs[f'[xlist]]]] / Length[xlist])

l1 = lyा[0.2, 0.7, 50000, 20]
-0.223144

l2 = lyा[0.6, 0.7, 50000, 20]
-0.916291

l3 = lyा[0.8, 0.7, 50000, 20]
-0.91629

l4 = lyा[0.84, 0.7, 50000, 20]
-0.281411

l5 = lyा[0.90, 0.7, 50000, 20]
0.181605

l6 = lyा[0.961, 0.7, 50000, 20]
-0.280341

l7 = lyा[0.97, 0.7, 50000, 20]
0.463027

```

```
18 = lya[1, 0.7, 50000, 20]
0.693132

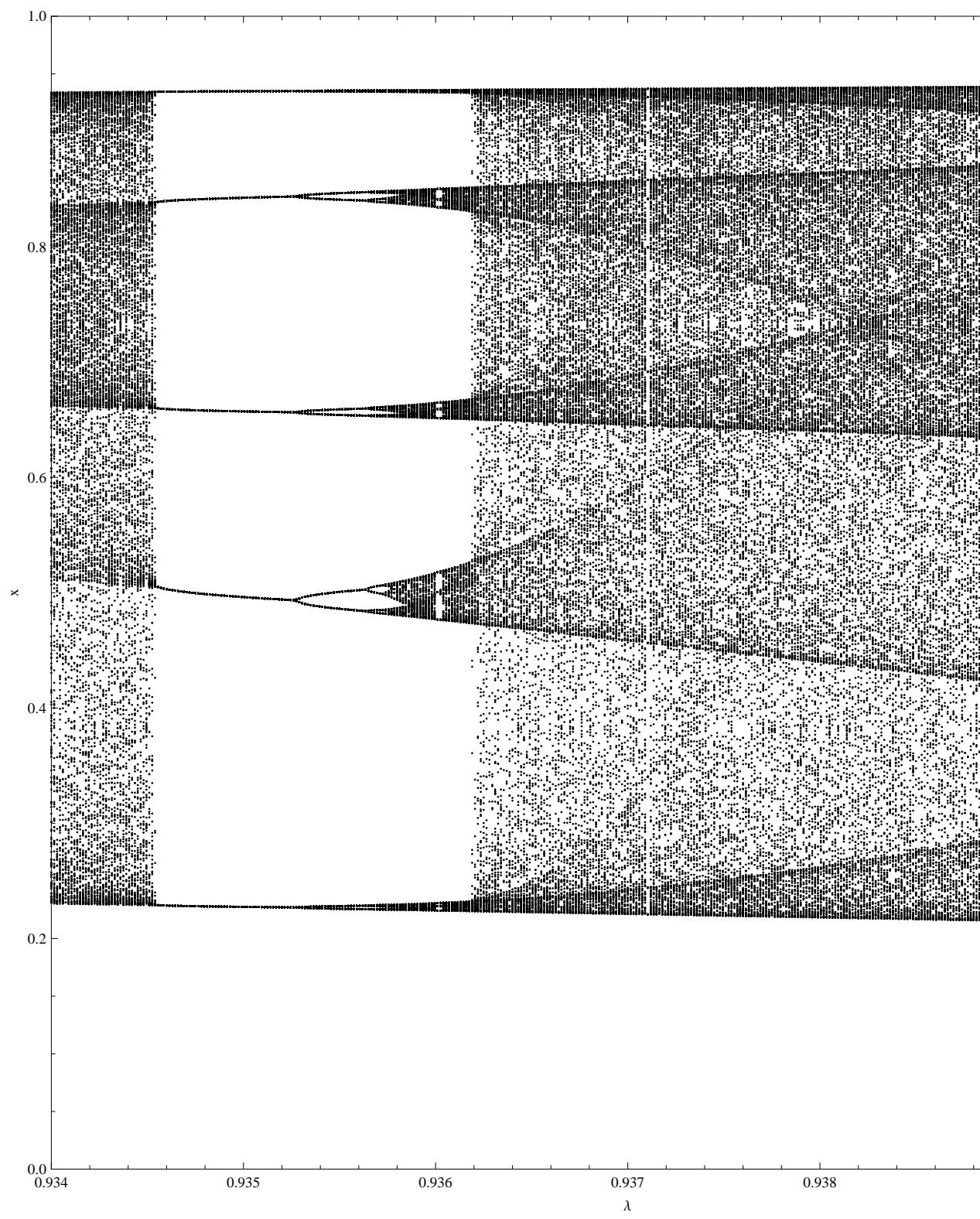
c)

iterate [m_, n_] := Drop[NestList[f, 0.5, n], m]

drawpt[y_] := Point[{λ, y}]

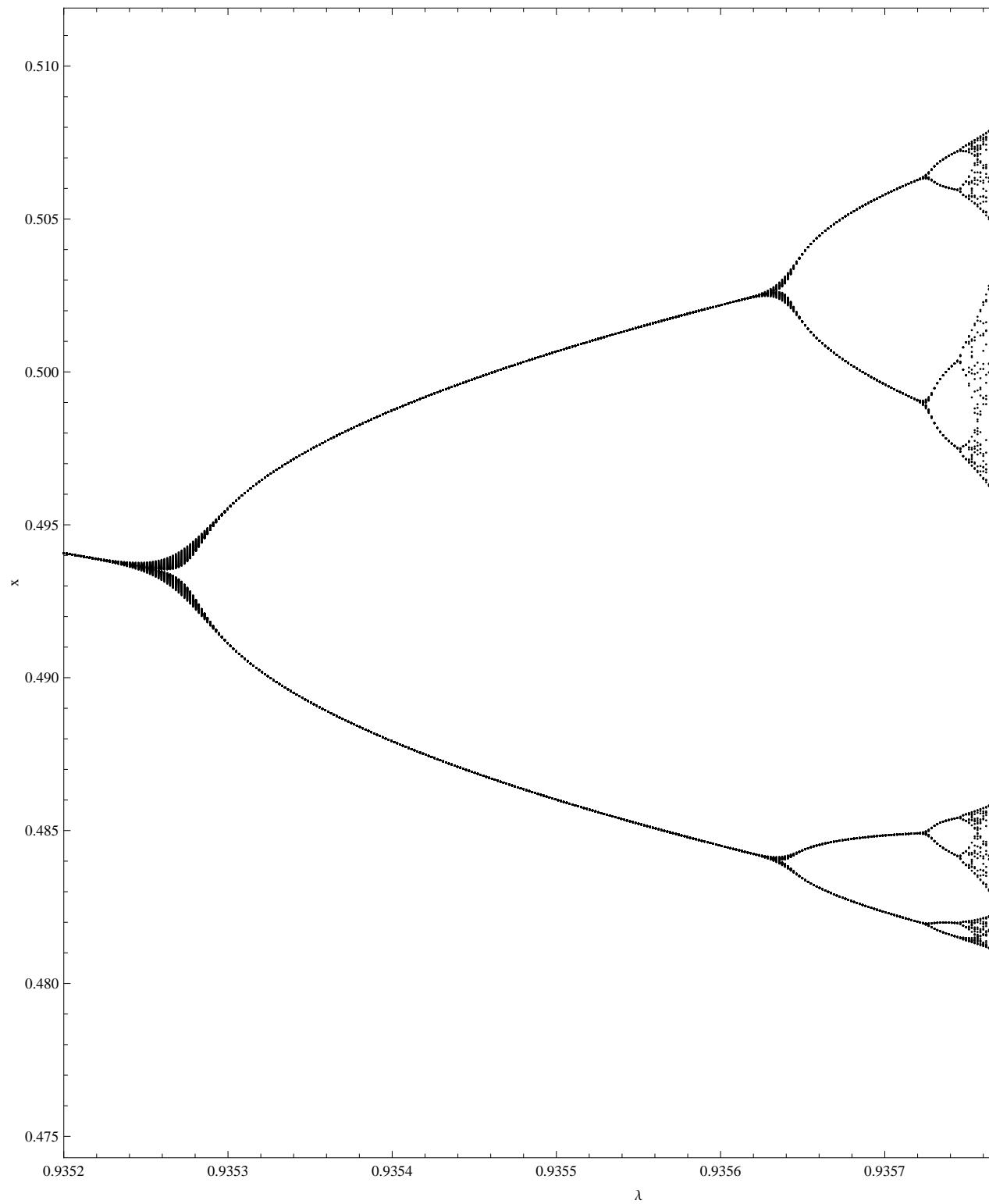
graph[λmin_, λmax_, nλ_, mdrop_, n_] := Graphics[{PointSize[0.001],
  Table[Map[drawpt, iterate[mdrop, n]], {λ, λmin, λmax, (λmax - λmin) / nλ}]]}
```

```
Show[graph[0.934, 0.94, 400, 300, 700], Axes → False, Frame → True,  
FrameLabel → {"λ", "x"}, PlotRange → {{0.934, 0.94}, {0, 1}}, AspectRatio → 1]
```



d)

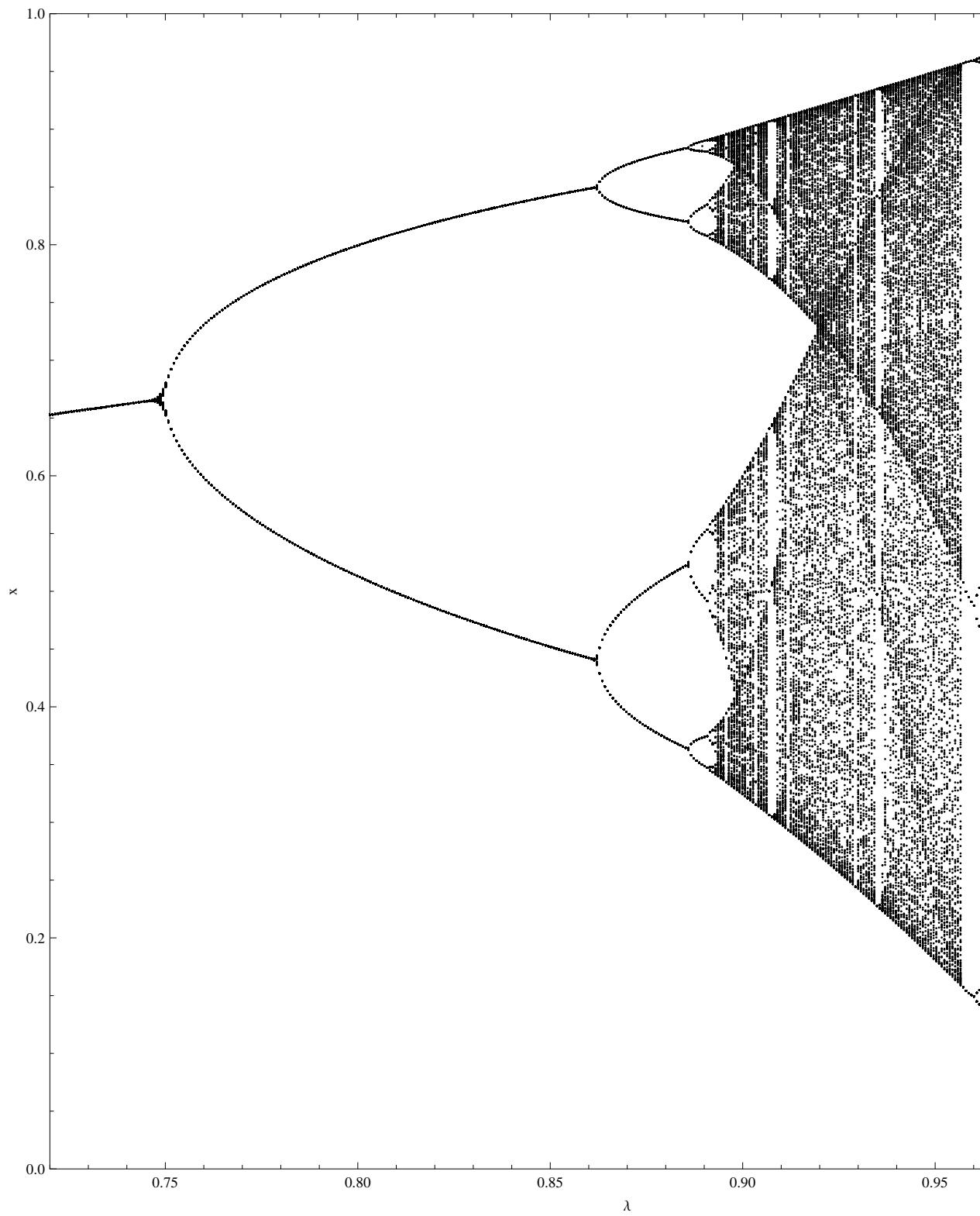
```
Show[graph[0.9352, 0.9359, 400, 300, 700], Axes → False, Frame → True,
FrameLabel → {"λ", "x"}, PlotRange → {{0.9352, 0.9359}, {0.4743, 0.5119}}, AspectRatio → 1]
```



e)

```
Clear[λ, x]

Show[graph[0.72, 1, 400, 300, 700], Axes → False, Frame → True,
FrameLabel → {"λ", "x"}, PlotRange → {{0.72, 1.02}, {0, 1}}, AspectRatio → 1]
```



`f2[x_] = f[f[x]] ;`

```

FindRoot[{f[x] == x, f'[x] == 1}, {x, 0.2}, {λ, 0.2}]
{x → 5.97024 × 10-9, λ → 0.25}

λ[0] = 0.25
0.25

FindRoot[{f[x] == x, f'[x] == -1}, {x, 0.5}, {λ, 0.5}]
{x → 0.666667, λ → 0.75}

λ[1] = 0.75
0.75

FindRoot[{f2[x] == x, f2'[x] == -1}, {x, 0.65}, {λ, 0.75}]
{x → 0.43996, λ → 0.862372}

λ[2] = 0.8623724356957946
0.862372

f4[x_] := f2[f2[x]]

FindRoot[{f4[x] == x, f4'[x] == -1}, {x, 0.88}, {λ, 0.9}]
{x → 0.88405, λ → 0.886023}

λ[3] = 0.8860225898879807
0.886023

f8[x_] := f4[f4[x]]

FindRoot[{f8[x] == x, f8'[x] == -1}, {x, 0.891}, {λ, 0.891}]
{x → 0.890787, λ → 0.891102}

λ[4] = 0.8911018165238584
0.891102

f16[x_] := f8[f8[x]]

FindRoot[{x == f16[x], f16'[x] == -1}, {x, 0.892}, {λ, 0.892}]
{x → 0.89214, λ → 0.89219}

λ[5] = 0.8921898548851404
0.89219

Print[k, "      ", δk];
Do[Print[n, "      ", (λ[n] - λ[n - 1]) / (λ[n + 1] - λ[n])], {n, 1, 4}]

```

k	$\delta_k$
1	4.44949
2	4.75145
3	4.65625
4	4.66824

In part a) we can see fixed point behavior for  $\lambda = 0.2$  and  $0.6$ ;  $\lambda = 0.2$  is tending towards 0 while  $\lambda = 0.6$  hits a fixed point around 0.58. For  $\lambda = 0.8$  and  $0.84$  we see limit cycle behavior;  $\lambda = 0.80$  has period 2 and  $\lambda = 0.84$  appears to have been period doubled to 4.  $\lambda = 0.90$  is chaotic, and  $\lambda = 0.961$  is exhibiting limit cycle behavior with period 2.  $\lambda = 0.97$  is chaotic again, while  $\lambda = 1$  drops to a fixed point of 0 very quickly.

In part b), the Lyapunov exponents verify the graphs in part a) in most cases. The positive exponents correspond to the chaotic looking graphs,  $\lambda = 0.90$  and  $\lambda = 0.97$ . All the negative exponents correspond to graphs with non-chaotic behavior, fixed points and limit cycles. The other positive exponent corresponds to  $\lambda = 1$ , which appears in the graph to quickly converge to a fixed point of 0. I do not understand this discrepancy.

Initially the graph in part c) was plotted from  $\lambda = 0.75$  to  $\lambda = 1$ . The final graph was produced by changing the range of  $\lambda$  to zoom in on a non-chaotic region that had 5 horizontal lines, corresponding to a period 5 limit cycle.

Using the Get Coordinates function in the drop down menu for the graph in part c), new ranges for x and  $\lambda$  were plotted to zoom in on one branch in the period 5 limit cycle. Period doubling is evident again, with periods 10, 20, 40, and so on visible, then chaos appears.

In part e)  $f2[x]$  is defined as  $f[f[x]]$ ,  $f4[x] = f2[f2[x]]$  and so on. FindRoot is used to find the values of x and  $\lambda$  where  $f[x] = x$  and  $f'[x] = -1$ . The value of  $\lambda$  corresponds to the point where the previous limit cycle becomes unstable and the period doubles. The value for  $\lambda$  is stored in an array of  $\lambda$  values, and the process is repeated for  $f2[x]$  and so on. FindRoot is helped by providing guesses for its starting values for x and  $\lambda$ . These are obtained by using the Get Coordinates function on the full-scale graph produced at the beginning of part e). The first 4  $\delta_k$ s are computed and printed in a table; after 4 iterations the value for  $\delta_k$  is very close to the Feigenbaum constant. The value of  $\delta_k$  is 4.666824... and the Feigenbaum constant is 4.66920... so they are around 0.001 apart.

## 7.

```
f[x_] := λ Sin[Pi x]
```

a)

```
lya[l_, xinit_, n_, ndrop_] := (λ = l; xlist = Drop[NestList[f, xinit, n], ndrop + 1];
Apply[Plus, Log[Abs[f'[xlist]]]] / Length[xlist])
L = lya[0.9, 0.7, 50000, 20]
0.352352
```

b)

```
λ = N[9/10, 3000];
```

```
x0 = N[4 / 10, 3000];
x5000 = AccountingForm[Nest[f, x0, 5000], 5]
0.79559

Precision[x5000]
2250.98
```

In part a) the same function lya[] from problem 6 is used to compute the Lyapunov exponent for  $\lambda = 0.9$ . It is positive, so it is clear that this is a chaotic region. In part b),  $\lambda$  and  $x_0$  are defined to be 0.4 and 0.9 to 3000 digits of precision. The precision of these values was chosen because we are trying to find  $x_{5000}$  and to maintain accuracy it appears you need precision of about 0.6 times the number of iterations.  $x_{5000}$  is computed and displayed with 5 digits; its precision is computed to be 2250.