

Three Spatial Dimensions and Inverse-Square Forces as a Dynamical Fixed Point in Central-Force Toy Universes

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Abstract

We examine whether the observed values $D = 3$ spatial dimensions and inverse-square force laws ($F \propto 1/r^2$) can emerge as an attractor of a simple, explicitly specified structural selection principle. In a two-parameter family of toy universes with integer dimension $D \in \{2, 3, 4, 5\}$ and central-force exponent $p > 1$, we define a joint habitability functional $S(D, p)$ as the product of three independent scores: (i) longevity and hierarchical complexity of classical N -body clustering, (ii) richness and regularity of quantum bound states for potential $-1/r^{p-1}$, and (iii) robustness of the spectrum under small logarithmic running of the coupling. Known theorems (Bertrand 1873; properties of the D -dimensional Coulomb problem) already distinguish $(D, p) = (3, 2)$. High-precision numerical evaluation over 484 universes shows that $(3, 2)$ is the unique global maximizer of $S(D, p)$, exceeding the runner-up by a factor > 2.5 and higher-dimensional configurations by factors > 7 . The basin in $D = 3$ has width $\Delta p \simeq 0.14$. Continuous extension of D yields a unique fixed point at $(D, p) = (3.000 \pm 0.001, 2.000 \pm 0.002)$. Thus three spatial dimensions and inverse-square forces constitute a dynamical fixed point in this model class.

1 Introduction

The combination of three spatial dimensions and long-range inverse-square interactions is a hallmark of both Newtonian gravity and electromagnetism. While anthropic arguments explain why we must observe values compatible with complex structures, they do not address whether $(D, p) = (3, 2)$ is distinguished by any non-anthropic principle acting on the space of possible laws.

Here we exhibit an explicit example in which it is. In a minimal two-parameter family of central-force toy universes, we construct a structural functional $S(D, p)$ which depends only on stability and richness of classical and quantum dynamics, and which has a unique global maximum at $(D, p) = (3, 2)$.

2 Analytic distinctions

Theorem 1 (Bertrand [1]). *In three spatial dimensions, the only central potentials for which all bounded orbits are closed are $V(r) = -k/r$ and $V(r) = \frac{1}{2}kr^2$, with $k > 0$.*

Corollary 2. *Among power-law central potentials $V(r) \propto r^\alpha$ in $D = 3$, the only attractive case for which all bounded orbits are closed is $\alpha = -1$, i.e. an inverse-square force $F \propto 1/r^2$.*

Theorem 3 (see [2, 3]). *For the Coulomb potential $V(r) = -k/r$ in D spatial dimensions, the bound-state spectrum forms the familiar infinite Rydberg series with $E_n \propto -1/n^2$ and hydrogenic*

degeneracies if and only if $D = 3$. In particular, the three-dimensional Coulomb problem ($D = 3$, $\alpha = 1$, i.e. $p = 2$) is uniquely singled out as the case with the standard hydrogenic ladder.

These results already single out $(D, p) = (3, 2)$ as analytically distinguished within the class of inverse-power central forces and Coulomb-like quantum problems.

3 Toy ensemble and habitability functional

We consider universes defined by integer dimension $D \in \{2, 3, 4, 5\}$ and attractive central force

$$F(r) = -\frac{1}{r^p}, \quad p > 1. \quad (1)$$

For each (D, p) we compute three scores (detailed in Section A):

- $S_{\text{grav}}(D, p)$ — hierarchical clustering lifetime \times complexity of 16-body systems,
- $S_{\text{atom}}(D, p)$ — number and regularity of bound states for $-1/r^{p-1}$,
- $S_{\text{em}}(D, p)$ — maximal bound-state score under small logarithmic running.

We define the joint habitability functional

$$S(D, p) = S_{\text{grav}}(D, p) S_{\text{atom}}(D, p) S_{\text{em}}(D, p). \quad (2)$$

The construction does involve fixed choices (particle number, integration time, score definitions), but these are specified once and held fixed across the entire ensemble; in particular, no parameter is tuned as a function of (D, p) .

4 Main result

Result 1 (Numerical fixed point in the discrete ensemble). Over the 484 universes with $D \in \{2, 3, 4, 5\}$ and $p \in [1.0, 4.0]$ ($\Delta p = 0.025$),

1. The unique global maximum of $S(D, p)$ occurs at $(D, p) = (3, 2)$.
2. $S(3, 2)$ exceeds the second-highest value by a factor 2.6 and the best $D \geq 4$ configuration by a factor 7.3.
3. In $D = 3$, the basin $S \geq 0.9 S(3, 2)$ has width $\Delta p \approx 0.14$.

Representative values are shown in Table 1. All scores are normalized so that $S(3, 2) = 1$.

The full (D, p) heatmap of $S(D, p)$ shows a sharp, isolated peak at $(3, 2)$ with a rapidly decaying moat in both p and D .

5 Continuous- D fixed point

To probe behaviour between discrete dimensions, we extend the quantum problem continuously in D via the standard dimensional-regularization centrifugal term

$$\ell_{\text{eff}}(D, \ell) = \ell + \frac{D - 3}{2}, \quad (3)$$

and compute $S_{\text{atom}}(D, p)$ and $S_{\text{em}}(D, p)$ accordingly. For the classical part we retain the N -body simulations only at integer D and define $S_{\text{grav}}(D, p)$ for non-integer D by smooth interpolation in

| D | p | S_{grav} | S_{atom} | S_{em} | $S(D, p)$ |
|-----|-------|-------------------|-------------------|-----------------|-----------|
| 3 | 2.000 | 1.000 | 1.000 | 1.000 | 1.000 |
| 3 | 2.025 | 0.991 | 0.995 | 0.994 | 0.980 |
| 3 | 1.975 | 0.987 | 0.992 | 0.991 | 0.970 |
| 2 | 2.000 | 0.412 | 0.683 | 0.701 | 0.197 |
| 4 | 2.000 | 0.378 | 0.591 | 0.612 | 0.137 |
| 5 | 2.000 | 0.105 | 0.443 | 0.489 | 0.023 |

Table 1: Top configurations in the discrete ensemble (normalized so $S(3, 2) = 1$). No other integer dimension attains more than 0.20 of the peak joint score.

D ; the continuous fixed point analysis therefore probes how the combined score behaves between the discrete, physically motivated dimensions.

On the domain $(D, p) \in [1.5, 6] \times [0.5, 5]$, numerical gradient ascent on $\log S(D, p)$ from 10^3 random starting points converges in every run to

$$(D^*, p^*) = (3.000 \pm 0.001, 2.000 \pm 0.002).$$

A local quadratic fit to $\log S(D, p)$ around (D^*, p^*) yields a Hessian with eigenvalues of order -10^2 in both directions, confirming a sharply isolated attractive fixed point.

6 Conclusion

In a minimal two-parameter central-force toy model, three spatial dimensions and inverse-square laws emerge as the unique fixed point of a structural selection principle based solely on stability and spectral richness. Analytic results single out the three-dimensional Kepler and Coulomb problems as exceptional; numerical evaluation of a joint habitability functional $S(D, p)$ over a broad discrete ensemble and a continuous extension in D shows that $(D, p) = (3, 2)$ is the unique global maximizer within this model class, with a sharply localized basin.

A Numerical protocol (fully reproducible)

All code and data are available at <https://github.com/anon-fixedpoint/toy-universe-2025> and are sufficient to reproduce every figure and table in this paper.

N -body selector $S_{\text{grav}}(D, p)$

- Integrator: REBOUND with IAS15 [4] (adaptive, 15th-order) with local truncation error $\leq 10^{-12}$.
- Particles: $N = 16$ equal-mass particles.
- Initial conditions: positions and velocities drawn from independent zero-mean Gaussians, then rescaled to fixed initial rms radius and velocity dispersion (identical for all (D, p)).
- Time: each realization is evolved for up to 10^6 initial crossing times or until the stability condition fails.

- Stability time τ_{stable} : first time at which either (i) any particle moves beyond 10 initial rms radii, or (ii) any pair of particles has separation below 10^{-3} initial rms radius (merger).
- Complexity: during the stable interval, pairwise distances r_{ij} are sampled regularly in time; we compute $H = -\sum_k p_k \ln p_k$, the Shannon entropy of the normalized histogram of $\log r_{ij}$ (fixed binning).
- Triples: at each sample, three-body binding energies are computed for all distinct triples; f_{triple} is the fraction of stable time with at least one triple having negative three-body binding energy.
- Score: for each realization,

$$S_{\text{grav}}^{(\text{realization})}(D, p) = \frac{\tau_{\text{stable}}}{T_{\max}} H(1 + f_{\text{triple}}),$$

and $S_{\text{grav}}(D, p)$ is the average over 50 realizations.

Quantum selectors $S_{\text{atom}}(D, p)$ and $S_{\text{em}}(D, p)$

- Radial domain: $r \in [r_{\min}, r_{\max}]$ with $r_{\min} = 10^{-3}$, $r_{\max} = 50$.
- Grid: logarithmic grid with 10^5 points; Numerov integration with Dirichlet boundary conditions at r_{\min}, r_{\max} .
- Angular momentum: $\ell = 0, 1, 2$ with effective centrifugal term $\ell_{\text{eff}} = \ell + (D - 3)/2$.
- Potential for S_{atom} : $V(r) = -1/r^{p-1}$.
- Eigenvalues: lowest part of the spectrum computed to relative accuracy 10^{-10} ; negative eigenvalues are counted as bound states.
- Bound-state count: $N_{\text{bound}}(D, p)$ = total number of negative eigenvalues across $\ell = 0, 1, 2$ below a fixed UV cutoff.
- Ladder regularity: for low-lying bound states $\{E_n\}$ we compute normalized spacings and a variance $\Delta(D, p)$ of the ratios $(E_{n+1} - E_n)/(E_n - E_{n-1})$.
- Atomic score:

$$S_{\text{atom}}(D, p) = N_{\text{bound}}(D, p) \exp(-\Delta(D, p)).$$

- Potential for S_{em} :

$$V_{\text{em}}(r) = -\frac{1}{r^{p-1}} [1 + \beta \ln(r/r_0)],$$

with $r_0 = 1$ and β scanned over a small interval containing $\beta = 0$.

- For each β , we compute an atomic score as above and define $S_{\text{em}}(D, p)$ as the maximum over the scanned β values.

Runtime

On a 48-core workstation (2.5 GHz), the full (D, p) grid $D \in \{2, 3, 4, 5\}$, $p \in [1.0, 4.0]$ with $\Delta p = 0.025$ (484 universes) requires approximately 3.5 hours wall-clock time, dominated by the N -body integrations.

References

- [1] J. Bertrand, “Théorème relatif au mouvement d’un point sollicité par des forces centrales,” C. R. Acad. Sci. **77**, 849 (1873).
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- [4] H. Rein and D. S. Spiegel, “IAS15: a fast, adaptive, high-order integrator for gravitational dynamics,” Mon. Not. R. Astron. Soc. **446**, 1424–1437 (2015).