

# Chapter 6

## Background estimation

A reliable and trustworthy estimation of the expected SM background rates in the signal regions is crucial for exercising the statistical machinery laid out in chapter 3 and making conclusive statistical statements about the SUSY scenarios studied. The background estimation approaches used herein either rely on semi-data-driven techniques, or on MC-only estimations. As estimating backgrounds only from MC simulation is sometimes problematic due to e.g. mis-modelings in the phase space targeted not being appropriately covered by the uncertainties, a (semi-)data-driven approach is often favoured. In the following, the major backgrounds  $t\bar{t}$ , single top and  $W + \text{jets}$  are estimated using a semi-data-driven approach, while the expected rates from the remaining, smaller backgrounds rely purely on MC simulations and are normalised to their theoretical cross section.

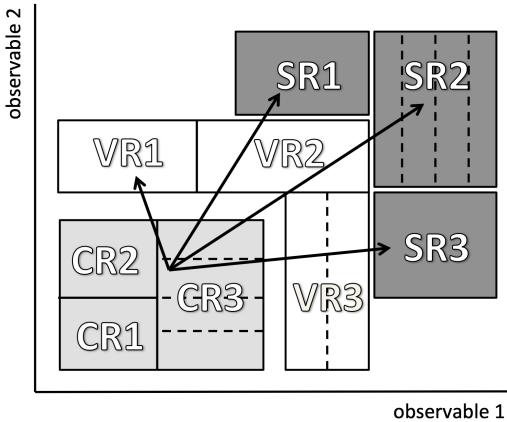
### 6.1 General strategy

#### 6.1.1 Transfer factor approach

Estimating background contributions in SRs in a semi-data-driven approach usually involves the introduction of so-called control regions (CRs), used to control dominant background processes by comparing their expected event rates to data. The CRs are designed to be enriched in events of a given background process (or type) while being approximately free of signal contamination. If  $N_p^{\text{MC}}(\text{SR})$  and  $N_p^{\text{MC}}(\text{CR})$  are the expected rates for a given background process  $p$  obtained from MC simulation in a given SR and CR, respectively, then the transfer factor  $N_p^{\text{MC}}(\text{SR})/N_p^{\text{MC}}(\text{CR})$  allows to convert the number of observed background events in the CRs,  $N_p^{\text{obs.}}(\text{CR})$ , into a background estimate in the SRs,  $N_p^{\text{est.}}(\text{SR})$ , through

$$N_p^{\text{est.}}(\text{SR}) = N_p^{\text{obs.}}(\text{CR}) \frac{N_p^{\text{MC}}(\text{SR})}{N_p^{\text{MC}}(\text{CR})} = \mu_p N_p^{\text{MC}}(\text{SR}). \quad (6.1)$$

Here,  $\mu_p$  is the process-specific normalisation factor introduced in section 3.1. An important benefit of this approach is that the impact of systematic uncertainties on the estimated background rates can be evaluated on the transfer factors, that are ratios of MC estimates. As such, systematic uncertainties can cancel in the extrapolation to the SRs. The uncertainty on



**Figure 6.1:** Schematic view of an analysis strategy including multiple control, validation and signal regions with one or multiple bins each. Extrapolations from the control regions into the signal regions can be verified in the validation regions lying in the phase space extrapolated over. All regions are designed to be statistically independent. Figure adapted from Ref. [154].

the background estimate in the SRs is then a combination of statistical uncertainties in the CRs and remaining uncertainties affecting the extrapolation [154]. For this reason, CRs are usually deliberately chosen to have large statistics, effectively reducing the uncertainties on the extrapolation to the SRs.

As indicated in eq. (6.1), the transfer factor approach is formally equivalent to using the process-specific normalisation factors from section 3.1, effectively *normalising* the number of total background events expected from MC simulation to the number of observed events in each control region. In the profile likelihood fits used in the following, implemented using HISTFITTER [154], the normalisation factors  $\mu_p$  are fit to data instead of the background processes as expected from MC simulation. Multiple disjoint CRs are used to simultaneously normalise multiple background processes to data in a combined fit. In order not to have an underdetermined minimisation problem, at least the same number of CRs as normalisation factors need to be used. Two different profile likelihood fit configurations are used in the following; the first configuration being a so-called *background-only* fit configuration, assuming no signal contribution and typically only including the CRs. The second configuration is a so-called *model-dependent* fit configuration with nominal signal contribution using all CRs as well as SRs.

In order to verify the quality of the extrapolation from the CRs to the SRs, so-called validation regions (VRs) are defined. VRs do not participate in the actual fit of the model parameters to data, but serve as intermediate regions to verify the extrapolation. For this reason, VRs are typically placed in the region between the CRs and SRs that is extrapolated over. A schematic view of an analysis strategy using all three types of regions is shown in fig. 6.1. All three types of regions can have more than one bin and are separated using suitable observables that are extrapolated over. In order to be able to use information from all control and signal regions in a single profile likelihood fit, all regions necessarily need to be statistically independent.

### 6.1.2 Analysis blinding

An important concept in the design phase of searches for new physics is the idea of *blinding* regions of interest [238], meaning that measured data are not looked at in these regions. This avoids issues of *experimenter's bias*, i.e. unintended influences on the design of the analysis

based on the observed data. If data were already known when designing the signal regions (and therefore the outcome of the analysis would be known to some extent), experimenter's bias could for example occur during the selection of the final signal region definitions.

During the design of a search for SUSY, signal regions are generally kept blinded until the complete analysis strategy is fixed. Once the SRs have been designed, the next step is to develop a suitable background estimation strategy, often involving the introduction of CRs with negligible signal contamination. This is then often followed by the design of VRs that can be unblinded once the CRs are fixed. The SRs are only unblinded after the extrapolation of the background estimate (obtained using a background-only fit) from the CRs has been verified in the VRs, allowing to either quantify potential excesses in data or set limits on model parameters.

### 6.1.3 Data versus Monte Carlo plots

In this chapter, all plots comparing data versus MC are so-called *pre-fit* plots, meaning that no background-only fit has been run in order to determine the normalisation factors and total systematic uncertainties for the background estimate. Instead, the contributions from the dominant backgrounds  $t\bar{t}$ ,  $W + \text{jets}$  and single top are normalised simultaneously in the control regions by solving the system of  $i$  equations

$$n_{\text{data}}^{\text{CR}_i} = \mu_{t\bar{t}} B_{t\bar{t}}^{\text{CR}_i} + \mu_W B_W^{\text{CR}_i} + \mu_{\text{ST}} B_{\text{ST}}^{\text{CR}_i} + B_{\text{other}}^{\text{CR}_i}, \quad (6.2)$$

where  $i$  runs over the list of CRs introduced in section 6.2 and  $\mu_{t\bar{t}}$ ,  $\mu_W$  and  $\mu_{\text{ST}}$  are the normalisation factors of the  $t\bar{t}$ ,  $W + \text{jets}$  and single top backgrounds, respectively, that are to be determined.  $B_{t\bar{t}}^{\text{CR}_i}$ ,  $B_W^{\text{CR}_i}$ ,  $B_{\text{ST}}^{\text{CR}_i}$  and  $B_{\text{other}}^{\text{CR}_i}$  are the background rates expected from MC simulation in the  $i$ -th CR. The normalisation factors obtained are 0.96 for  $t\bar{t}$ , 1.24 for  $W + \text{jets}$  and 0.73 for single top. As will be shown in section 8.1, the normalisation factors obtained using the full statistical procedure will be close to these values.

Additionally, the uncertainty bands on the background estimate will include only MC statistical uncertainty as well as experimental uncertainties. The variations of the experimental uncertainties are normalised to the nominal background estimate in the case of  $t\bar{t}$ ,  $W + \text{jets}$  and single top, such that only the shapes of the dominant backgrounds are affected. For the remaining minor backgrounds, the experimental uncertainties can affect both normalisation and shape. All experimental uncertainties are assumed to be fully correlated over all processes and bins, allowing them to be summed in quadrature. Finally, the uncertainty bars on the data points are obtained by assuming data to be Poisson distributed and correspond to the 68% confidence interval.

## 6.2 Control regions

The contributions from  $t\bar{t}$ ,  $W + \text{jets}$  production and single top processes are normalised to data in dedicated control regions. Other processes like  $Z + \text{jets}$ , diboson and multiboson,  $t\bar{t} + V$ ,  $t\bar{t} + h$  and  $V + h$  are estimated directly from MC simulation and normalised to their theoretical cross sections. All CRs are designed to be kinematically as close as possible to the respective

**Table 6.1:** Overview of the CR and VR definitions. With the exception of  $m_{\ell b_1}$ , which is not used in the definitions of the CRs and VRs, all regions share the same selection as the SRs on the remaining kinematic observables not listed here.

CR	TR-LM	TR-MM	TR-HM	WR	STR
$m_{b\bar{b}}$ [GeV]		<100 or >140		$\in [50, 80]$	>195
$m_T$ [GeV]	$\in [100, 160]$	$\in [160, 240]$	>240	$\in [50, 100]$	>100
$m_{CT}$ [GeV]		<180		>180	>180
VR	VR-onLM	VR-onMM	VR-onHM	VR-offLM	VR-offMM
$m_{b\bar{b}}$ [GeV]		$\in [100, 140]$		$\in [50, 80] \cup [160, 195]$	$\in [50, 80] \cup [160, 195]$
$m_T$ [GeV]	$\in [100, 160]$	$\in [160, 240]$	>240	$\in [100, 160]$	$\in [160, 240]$
$m_{CT}$ [GeV]		<180			>180
					VR-offHM
					>240

SRs, such that the normalisation factors derived in the CRs are also valid in the SRs. The CRs are mutually exclusive and made orthogonal to the SRs through their requirements on  $m_T$ ,  $m_{CT}$  and  $m_{b\bar{b}}$ . Apart from the requirements on these three observables, as well as the requirement on  $m_{\ell b_1}$  (removed altogether in the CRs), the CRs share the same set of cuts as the SRs. Figure 6.4(a) illustrates the configuration of all CRs, especially highlighting the fact that all CRs are located in sideband regions off the  $m_{b\bar{b}}$  window, significantly reducing signal contamination. Table 6.1 summarises the kinematic requirements separating the CRs from other regions of interest in the analysis. The pre-fit distributions of all CRs in representative observables are shown in fig. 6.2.

### Control regions for $t\bar{t}$

As events from  $t\bar{t}$  processes constitute the dominant SM background in the majority of the SRs, it is necessary to have a precise and reliable estimate of their contributions. Three CRs are defined for  $t\bar{t}$ , following the same binning in  $m_T$ , and thus called TR-LM, TR-MM and TR-HM in the following. A good purity of  $t\bar{t}$  processes as well as the necessary high MC statistics are achieved by inverting the requirement on  $m_{CT}$ , selecting events below the kinematic endpoint for  $t\bar{t}$  processes. The achieved pre-fit  $t\bar{t}$  purities are 79.6% in TR-LM, 85.9% in TR-MM and 84.1% in TR-HM. The remaining contributions stem mostly from single top and  $W +$  jets processes and vary between 8.6%–14.1% and 1.8%–4.3%, respectively, depending on the SR.

For a trustworthy estimate of the contributions from  $t\bar{t}$  processes, it is important that the control regions associated to each signal region exhibit approximately the same composition of  $t\bar{t}$  decay modes. The decay mode most relevant to the  $1\ell$  search at relatively low and moderate values of  $m_T$  is the semi-leptonic decay ( $\ell\nu qq$ ), where one of the  $W$  bosons decays leptonically, while the other one undergoes a hadronic decay. The semi-leptonic decay mode exhibits the well-known kinematic endpoint in  $m_T$  and thus quickly loses importance at high transverse mass values. Events involving a hadronic decay of a  $\tau$ -lepton originating from  $W \rightarrow \tau_{\text{had}}\nu$  in one of the two branches and a leptonic  $W$  boson decay in the other branch ( $\ell\nu\tau_{\text{had}}\nu$ ), are the dominant decay mode in selections with high values of  $m_T$ . Due to the additional neutrino in such events, the  $\ell\nu\tau_{\text{had}}\nu$  decay mode does not exhibit the same kinematic endpoint as the semi-leptonic one. Finally, di-leptonic decays ( $\ell\nu\ell\nu$ ) and events with a leptonically decaying  $\tau$ -lepton ( $\ell\nu\tau_\ell\nu$ ), where one of the two leptons is not reconstructed, play a sub-dominant but non-negligible role in all regions. Other  $t\bar{t}$  decay modes are negligible in all analysis selections.

In the low-mass regions with moderate values in  $m_T$  not far above its kinematic endpoint, 80% (40%) of  $t\bar{t}$  events involve the semi-leptonic decay mode in the control region (signal region). The sub-dominant decay mode in these regions involves the  $\ell\nu\tau_{\text{had}}\nu$  decay mode, with a contribution of 25% and 10% in TR-LM and SR-LM, respectively. Di-leptonic and  $\ell\nu\tau_\ell\nu$  decay modes each contribute about 15% of all events in TR-LM and about 3% in SR-LM. Overall, the composition in the low-mass regions is hence not exactly the same in the control and signal regions, but the agreement is still considered to be acceptable. With about 45% (36%) and 30% (35%), the largest contributions in TR-MM (SR-MM) originate from  $\ell\nu\tau_{\text{had}}\nu$  decays and di-leptonic events, respectively. Events with a  $\ell\nu\tau_\ell\nu$  decay contribute to about 10% (15%) in TR-MM (SR-MM). In the high-mass control and signal regions with a high requirement on  $m_T$ , the majority (about 50%) of events involve the  $\ell\nu\tau_{\text{had}}\nu$  decay mode, while the di-leptonic and  $\ell\nu\tau_\ell\nu$  decay modes contribute with about 30% and 20%, respectively. Overall, the compositions of the different  $t\bar{t}$  decay modes in each control region are thus similar to the contributions in the respective signal region, meaning that the proportions of  $t\bar{t}$  processes constrained through the log-likelihood fit in the CRs are the same as those to be estimated in the SRs.

Signal contamination in the  $t\bar{t}$  CRs is avoided by inverting the requirement on  $m_{b\bar{b}}$ , i.e. placing the  $t\bar{t}$  CRs in the  $m_{b\bar{b}}$  sideband. The maximum signal contamination over the entire signal grid is 0.8%, 1.1% and 1.9% for TR-LM, TR-MM and TR-HM, respectively, and thus negligible. Figures 6.3(a) to 6.3(c) show the signal contamination in the  $t\bar{t}$  CRs over the full signal grid.

### Control region for $W + \text{jets}$

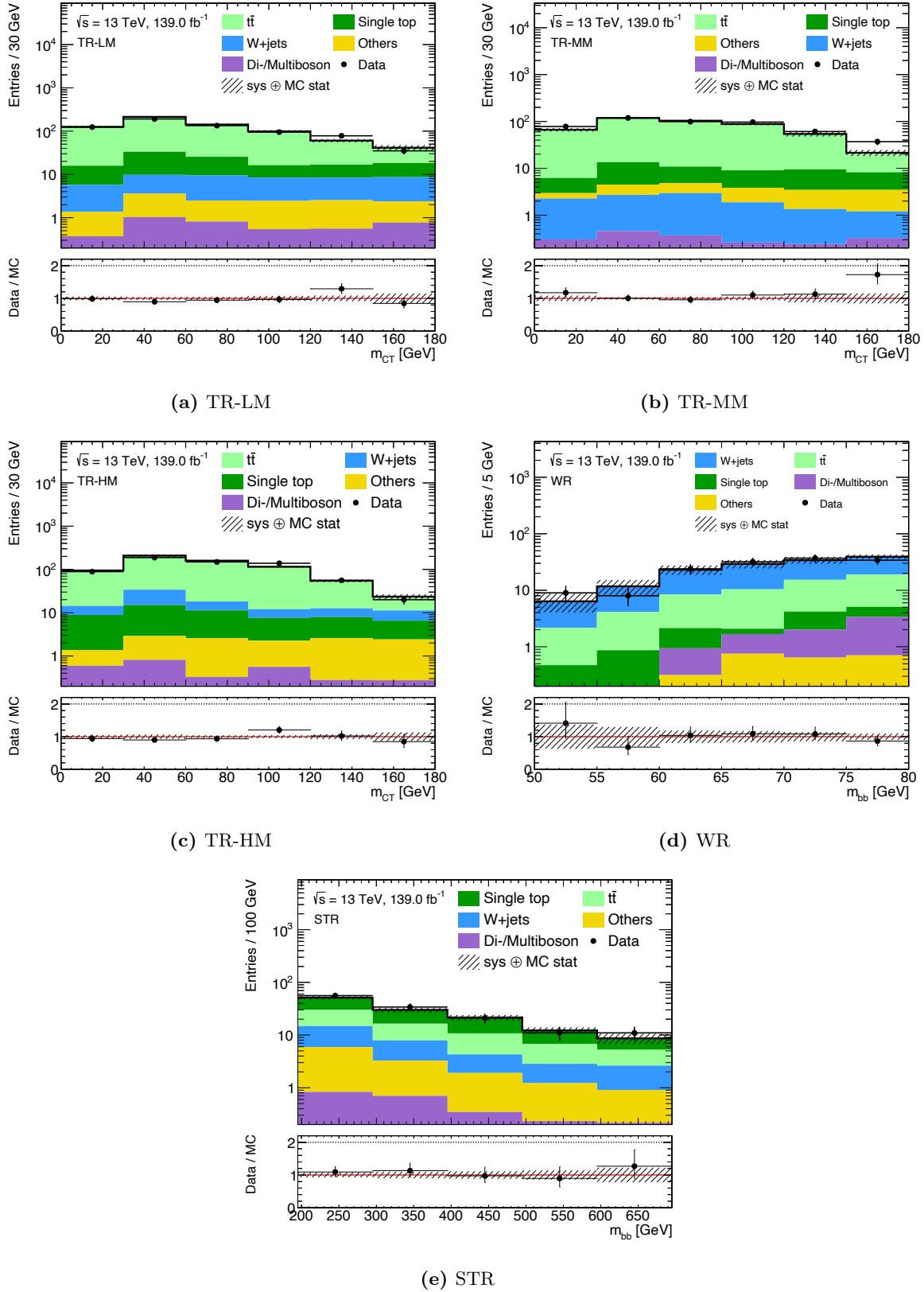
Events from  $W + \text{jets}$  production represent the second largest contribution of SM background processes in most SRs. A single  $W + \text{jets}$  control region, called WR in the following, is defined by replacing the signal region requirements on  $m_T$  and  $m_{b\bar{b}}$  with  $50 \text{ GeV} < m_T < 100 \text{ GeV}$  and  $50 \text{ GeV} < m_{b\bar{b}} < 80 \text{ GeV}$ , respectively. No bins in  $m_{\text{CT}}$  or  $m_T$  are defined for WR, as the composition of  $W + \text{jets}$  is approximately constant in all regions.

Applying a low requirement on  $m_T$  allows to predominantly select events below the kinematic endpoint of the transverse mass of the  $W$  boson, resulting in a high statistics control region with a pre-fit  $W + \text{jets}$  purity of roughly 52.5%. The sub-dominant background component of WR is  $t\bar{t}$  with 35.2%. Minor contributions of 7.0% and 4.2% originate from single top and diboson processes, respectively. The composition of  $W + \text{jets}$  events in WR and all signal regions is found to be dominated by  $W$  boson production in association with two real  $b$ -jets. Minor contributions originate from processes with mis-tagged  $c$ -jets or light-flavour jets.

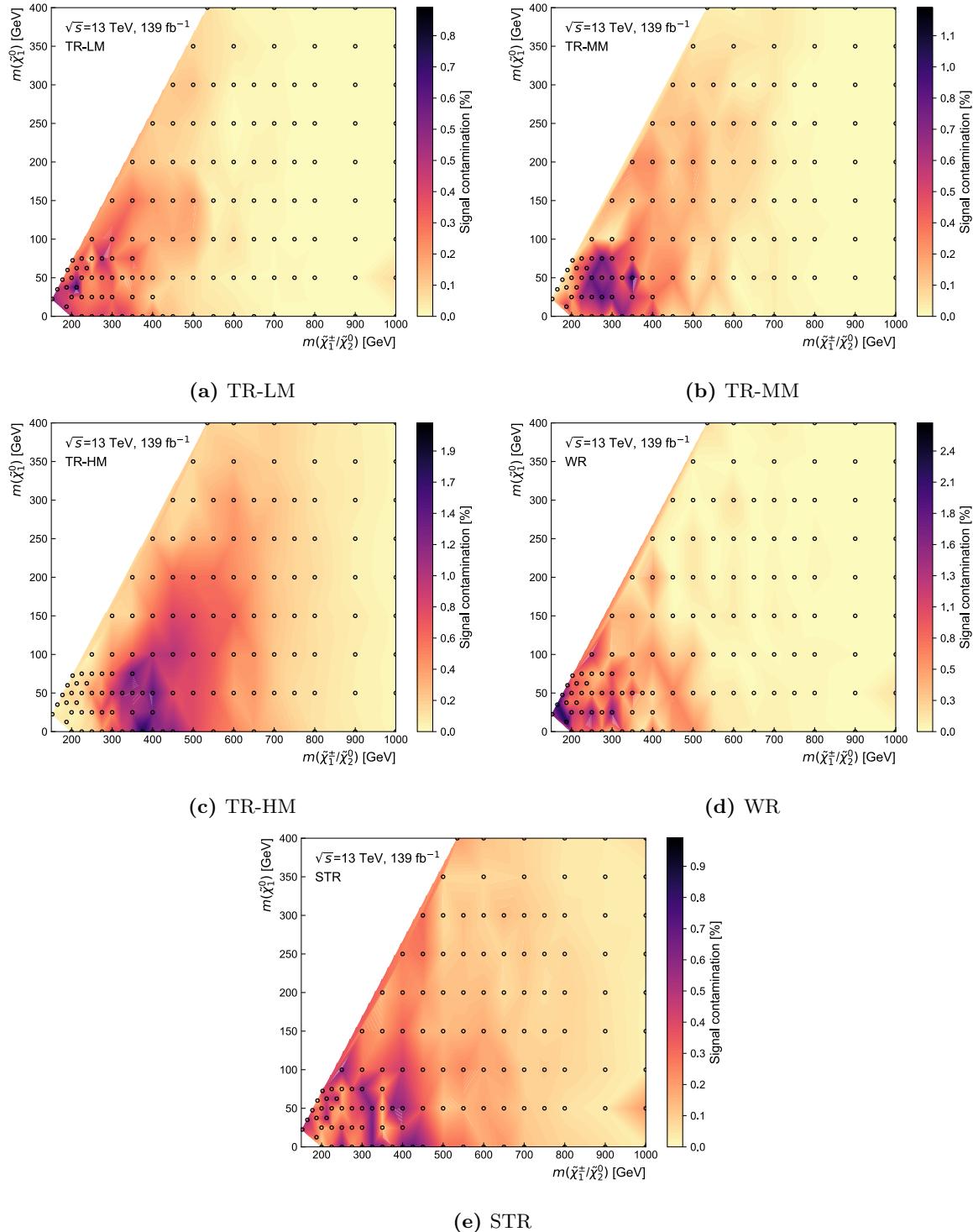
As was the case for the  $t\bar{t}$  control regions, placing WR off the Higgs mass peak allows to achieve a tolerable maximum signal contamination of only 2.4% without affecting the composition of processes in the  $W + \text{jets}$  background too much. Most signal points have significantly less than 1% signal contamination in WR, as can be seen in fig. 6.3(d).

### Control region for single top

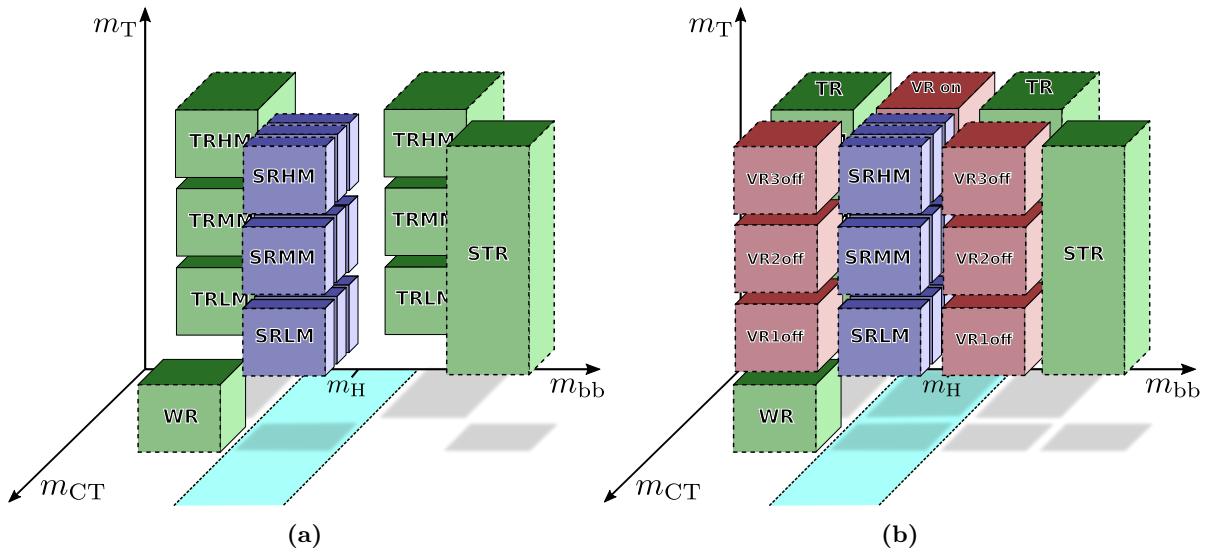
Single top processes result in significant background contributions in some SRs, necessitating a proper semi-data-driven estimation. A single top CR (STR) is defined starting from the SRs by



**Figure 6.2:** Exemplary pre-fit distributions for each control region. As laid out in the beginning of this chapter, the shaded region includes MC statistical uncertainty as well as experimental uncertainties, added in quadrature. A good agreement between MC expectation and data is observed in all CRs.



**Figure 6.3:** Signal contamination (shown on the  $z$ -axis) for all CRs throughout the signal grid. The space between the signal points (indicated by the black circles) is interpolated using Delaunay triangles.



**Figure 6.4:** Configuration of (a) the CRs placed around the SRs off the  $m_{b\bar{b}}$  window as well as (b) the validation regions in the phase space between the CRs and SRs. The VRs are arranged such that each of the extrapolations can be validated separately for SR-LM, SR-MM and SR-HM.

replacing the Higgs mass window cut on  $m_{b\bar{b}}$  with  $m_{b\bar{b}} > 195$  GeV and removing the bins in  $m_{CT}$ .

The sideband approach achieves again a low maximum signal contamination of roughly 0.8%. The signal contamination across the entire signal grid is shown in fig. 6.3(e). The pre-fit purity of the single top processes in STR is 51.7% and sub-dominant contributions arise from  $t\bar{t}$  processes (29%),  $W + \text{jets}$  (10%) and  $t\bar{t} + V$  (6%) production.

### 6.3 Validation regions

Two sets of VRs regions are introduced in order to verify the extrapolations over the different distributions. The selections defining all VRs are summarised in table 6.1. The first set, called VR-on is situated on the Higgs boson mass peak but with the  $m_{CT}$  requirement inverted to  $m_{CT} < 180$  GeV. This allows the VR-on regions to validate the extrapolation over  $m_{b\bar{b}}$ , performed when extrapolating the background estimate from the control regions into the signal regions. Three disjunct VR-on regions are introduced, with  $m_T$  requirements matching those of the SRs, such that the extrapolations can be validated separately for each signal region. The three VR-on regions are aptly named VR-onLM, VR-onMM and VR-onHM. A similar composition of  $t\bar{t}$  decay modes as in the control and signal regions is observed in the VR-on regions, necessary for a trustworthy validation of the  $t\bar{t}$  estimate. A maximum signal contamination of about 5%–14% is achieved, depending on the requirement in  $m_T$ . As can be seen from fig. A.9, most signal points have a signal contamination well below 5% for all VR-on regions.

The second set of VRs is located on both sides off the Higgs boson mass peak at same values in  $m_{CT}$  than the SRs. This set of *off-peak* VRs, called VR-off, is used to validate the extrapolation over the  $m_{b\bar{b}}$  distribution performed in the case of STR. Additionally, the VR-off regions validate

the extrapolation over the  $m_{\text{CT}}$  performed in the  $t\bar{t}$  control regions. Similar to the on-peak validation regions, the VR-off regions are split into  $m_{\text{T}}$  bins matching the signal regions, allowing a validation of the background estimate in their respective signal region. The bins in VR-off are called VR-offLM, VR-offMM and VR-offHM. The maximum signal contamination in the VR-off regions is found to be about 7%–13%, depending on the requirement on  $m_{\text{T}}$ . Most signal points, however, reveal a signal contamination in the VR-off regions of less than 3% (cf. fig. A.9).



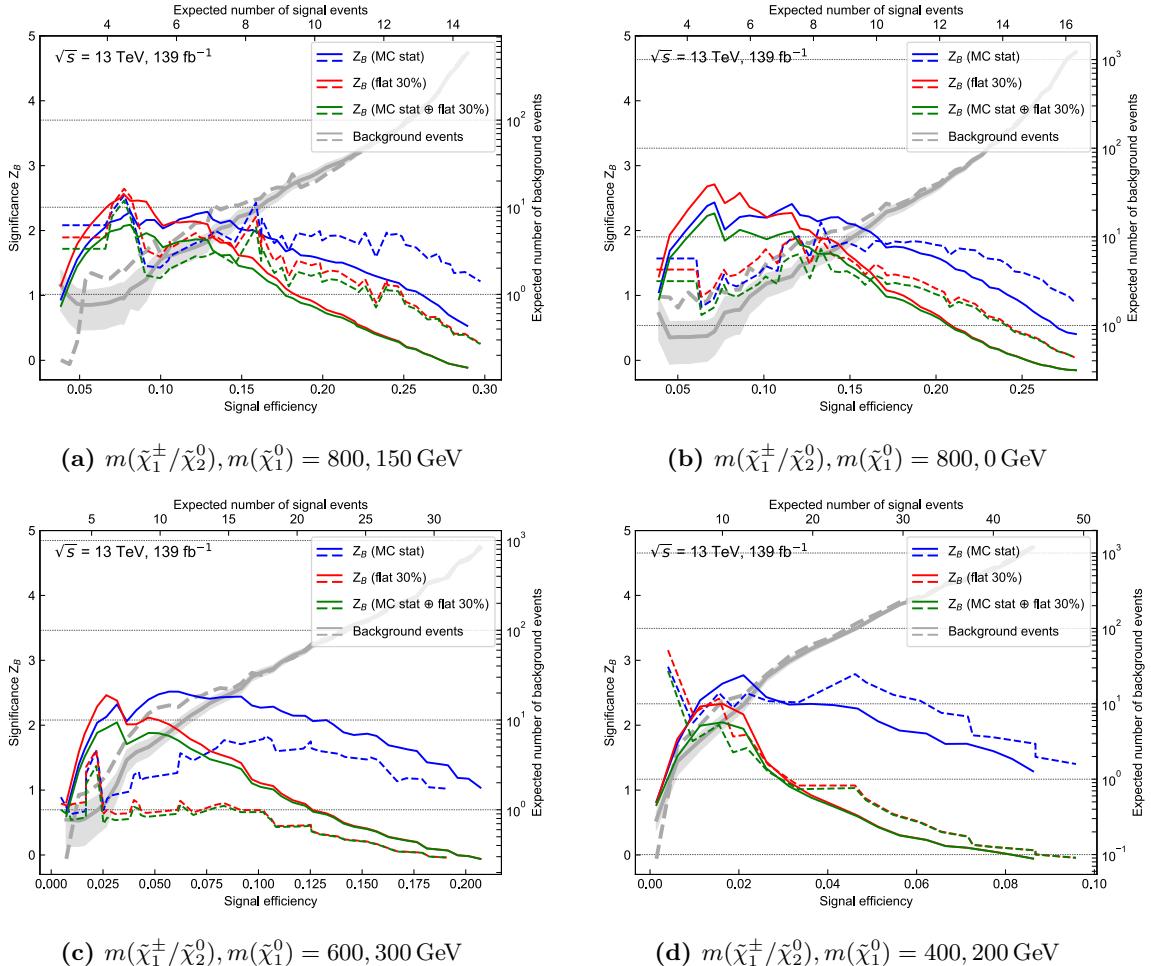
# Appendix A

## A.1 Additional information on signal region optimisation

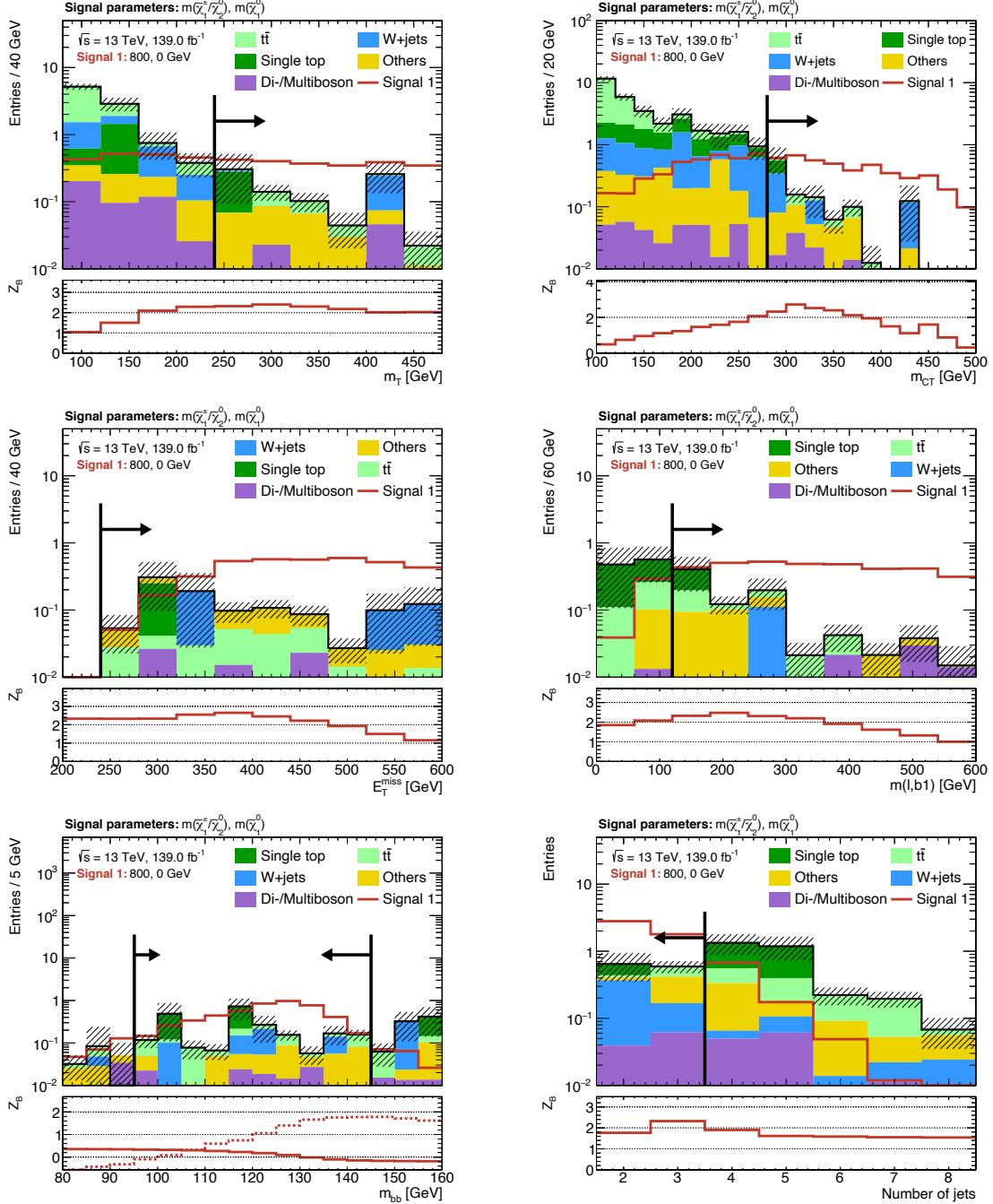
The following figures provide additional information on the signal region optimisation performed in chapter 5. Figure A.1 shows the results of the  $N$ -dimensional cut scan for the remaining benchmark signal points considered. As before, three different uncertainty configurations are used for computing the significance  $Z_B$ , and all values are computed for the two statistically independent subsets used during the  $N$ -dimensional scan. This approach allows to gauge the impact of statistical fluctuations on the cut combinations tested.

By choosing a well-performing cut combination for each benchmark point, the optimised selections in figs. A.2 to A.7 are found after a round of  $N - 1$  plots. As discussed in section 5.2.2 the optimal cut combinations for each benchmark signal point are consolidated into multiple signal regions designed to be sensitive to different kinematic regions of the model parameter space.

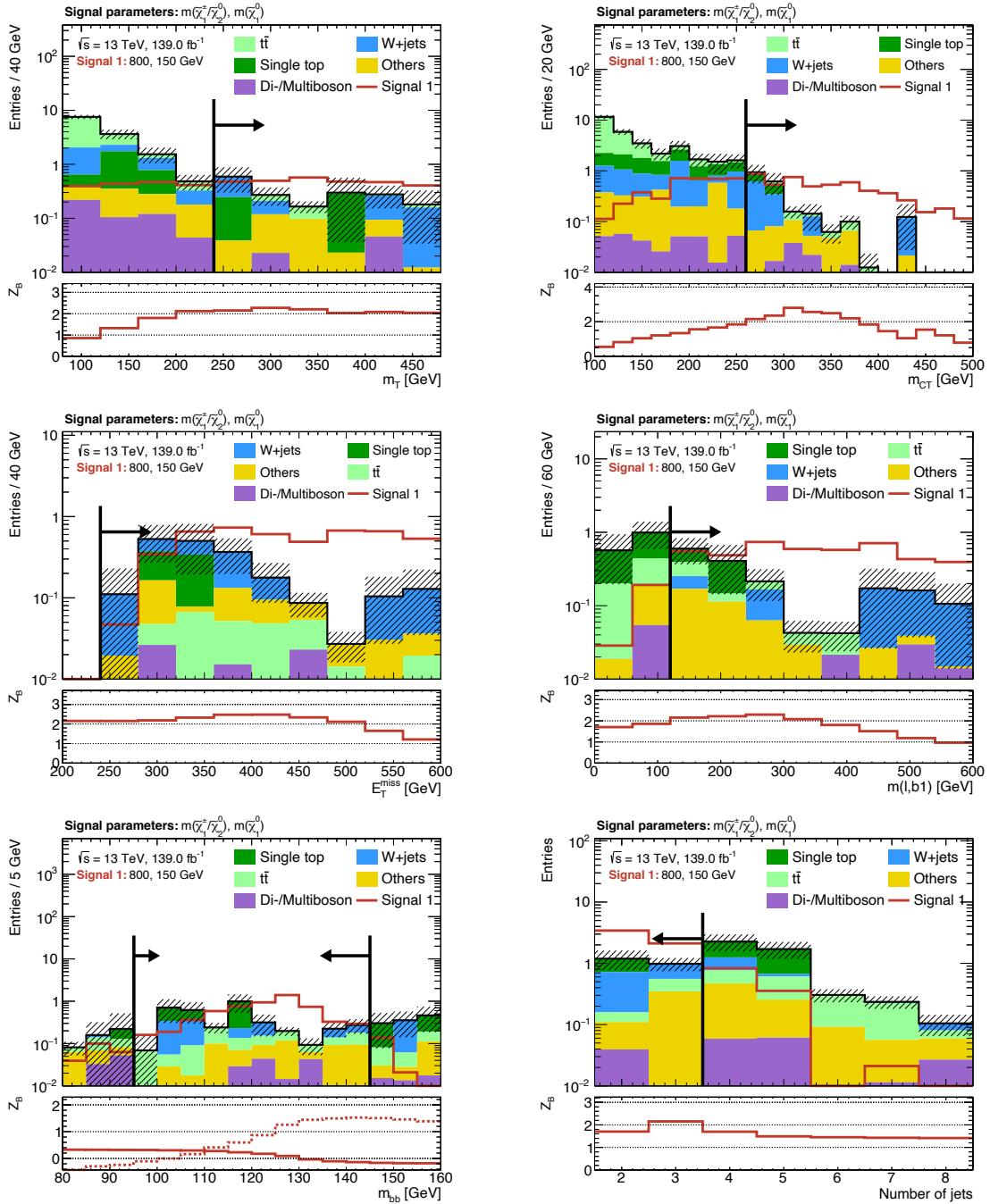
Figure A.8 shows additional information on some of the investigated simplified shape-fit configurations. Two-dimensional shape-fits in  $(m_T, m_{b\bar{b}})$ ,  $(m_T, E_T^{\text{miss}})$  and  $(m_T, m_{\text{CT}})$  with  $3 \times 3$  bins each are compared. The configuration using bins in  $m_T$  and  $m_{\text{CT}}$  results in the best expected  $\text{CL}_s$  values throughout the entire signal grid. Adding a requirement on high values of  $m_{\ell b_1}$  to SR-HM further increases the expected sensitivity. Overall, the expected sensitivity achieved through introduction of the two-dimensional shape fit significantly exceeds the sensitivity of the previous analysis iteration (using a one-dimensional shape-fit in  $m_T$ , scaled to  $139 \text{ fb}^{-1}$ ).



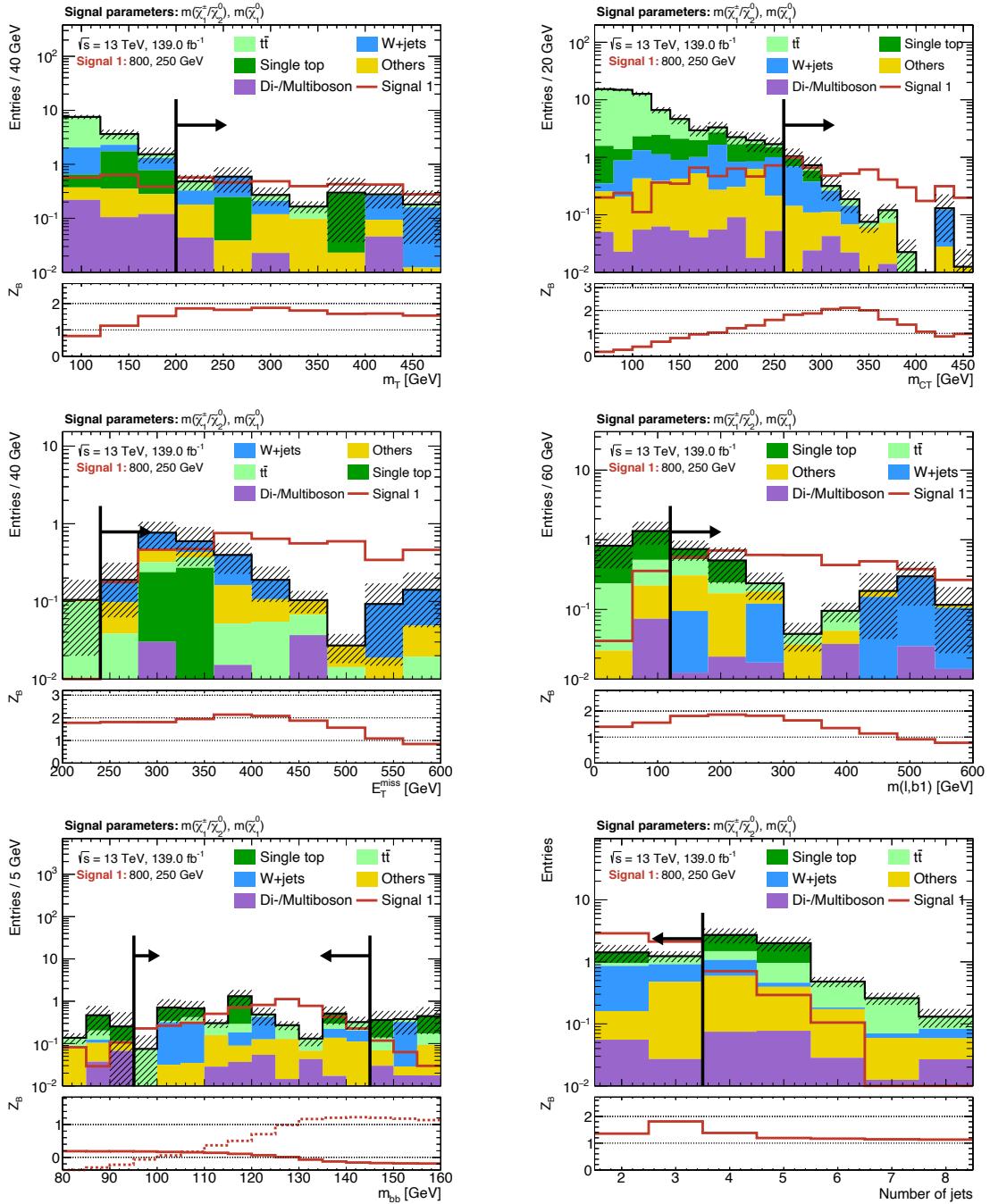
**Figure A.1:** Results of the  $N$ -dimensional cut scan for the remaining four benchmark points. The binomial discovery significance  $Z_B$  is plotted against the signal efficiency for varying uncertainty configurations. Additionally, the expected SM background rates are shown, including statistical uncertainty for one of the two statistically independent samples (shaded area). The solid and dashed lines represent the two statistically independent subset that the MC samples are split into.



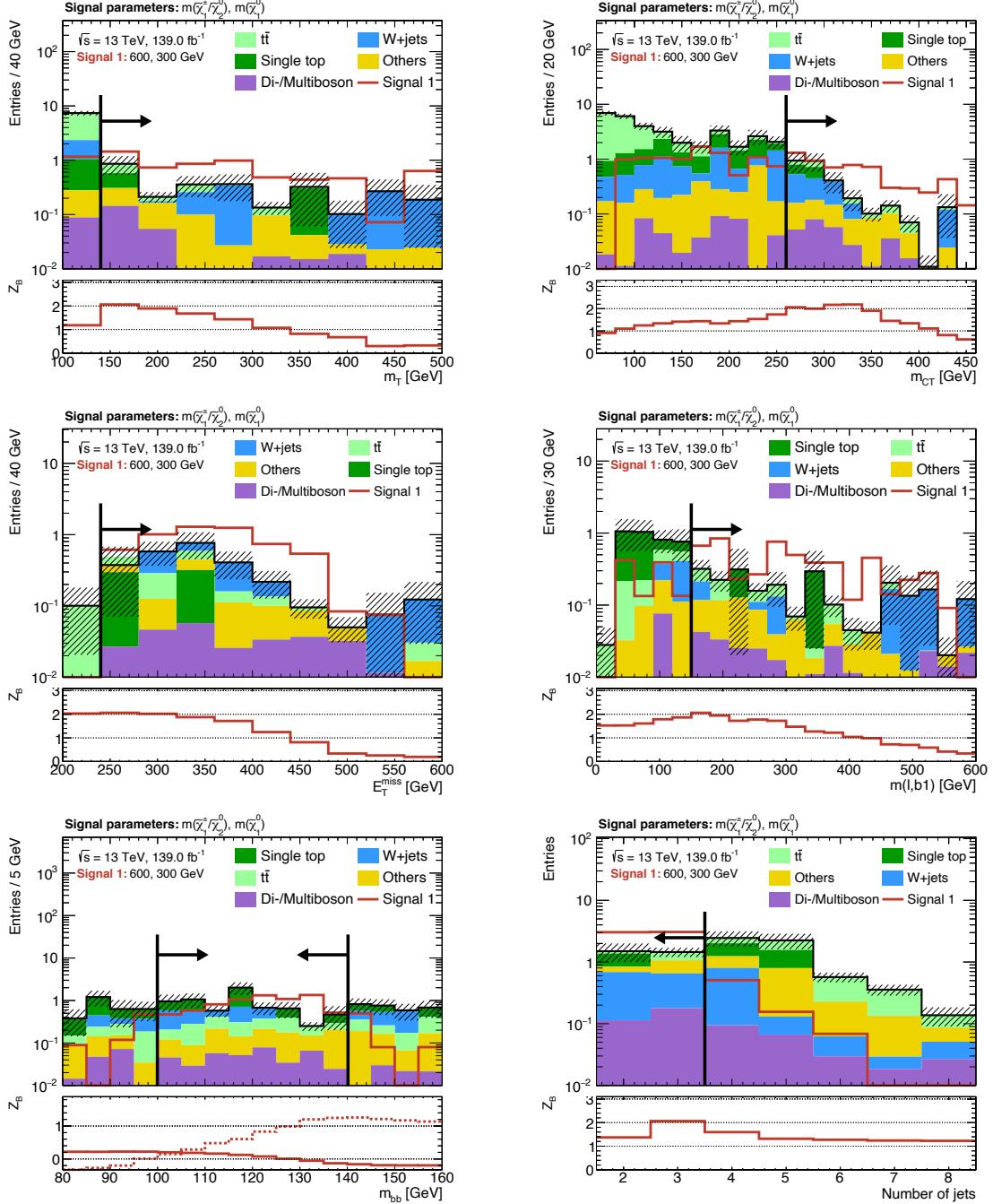
**Figure A.2:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (800 \text{ GeV}, 0 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.



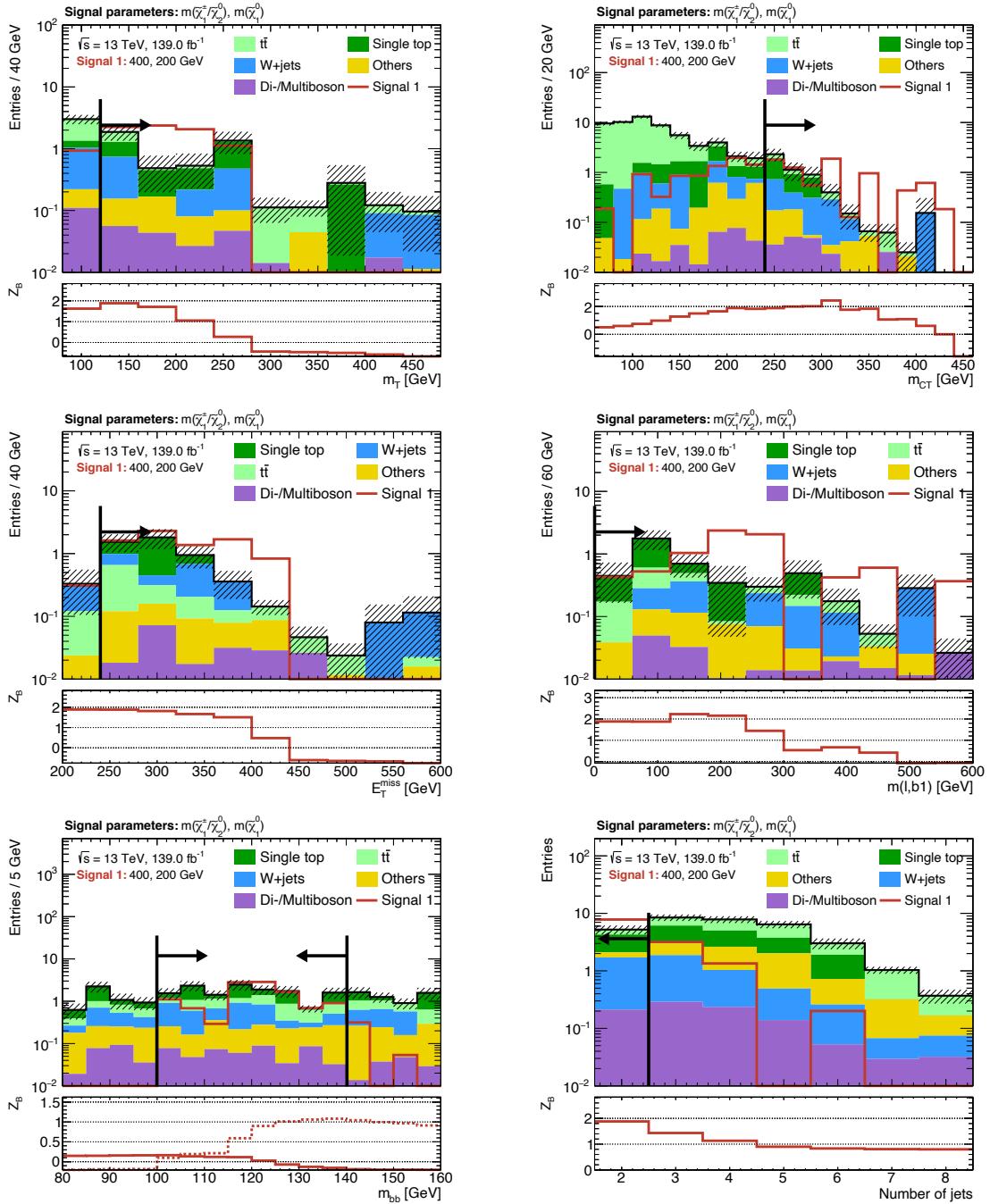
**Figure A.3:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (800 \text{ GeV}, 150 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.



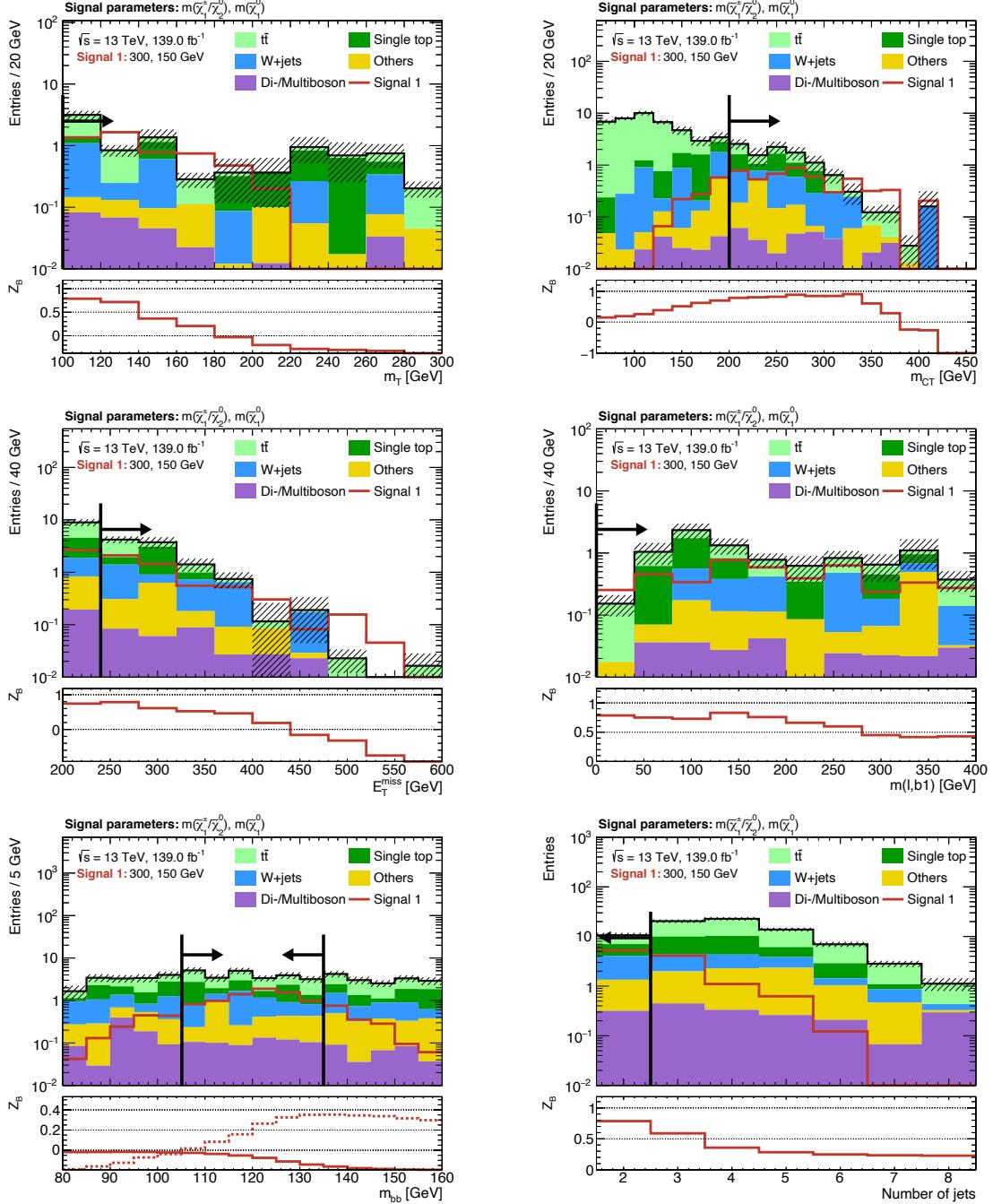
**Figure A.4:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (800 \text{ GeV}, 250 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.



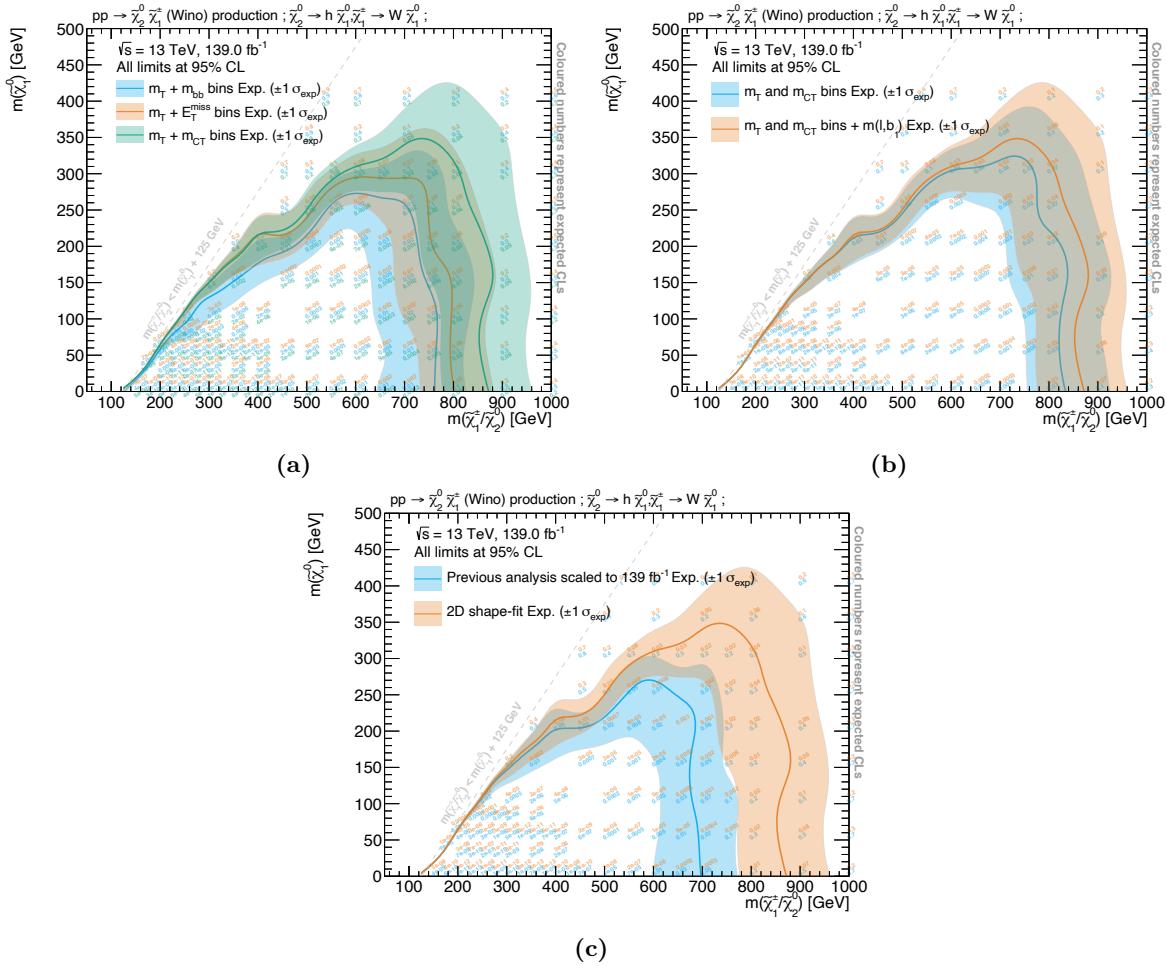
**Figure A.5:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (600 \text{ GeV}, 300 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.



**Figure A.6:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (400 \text{ GeV}, 200 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.

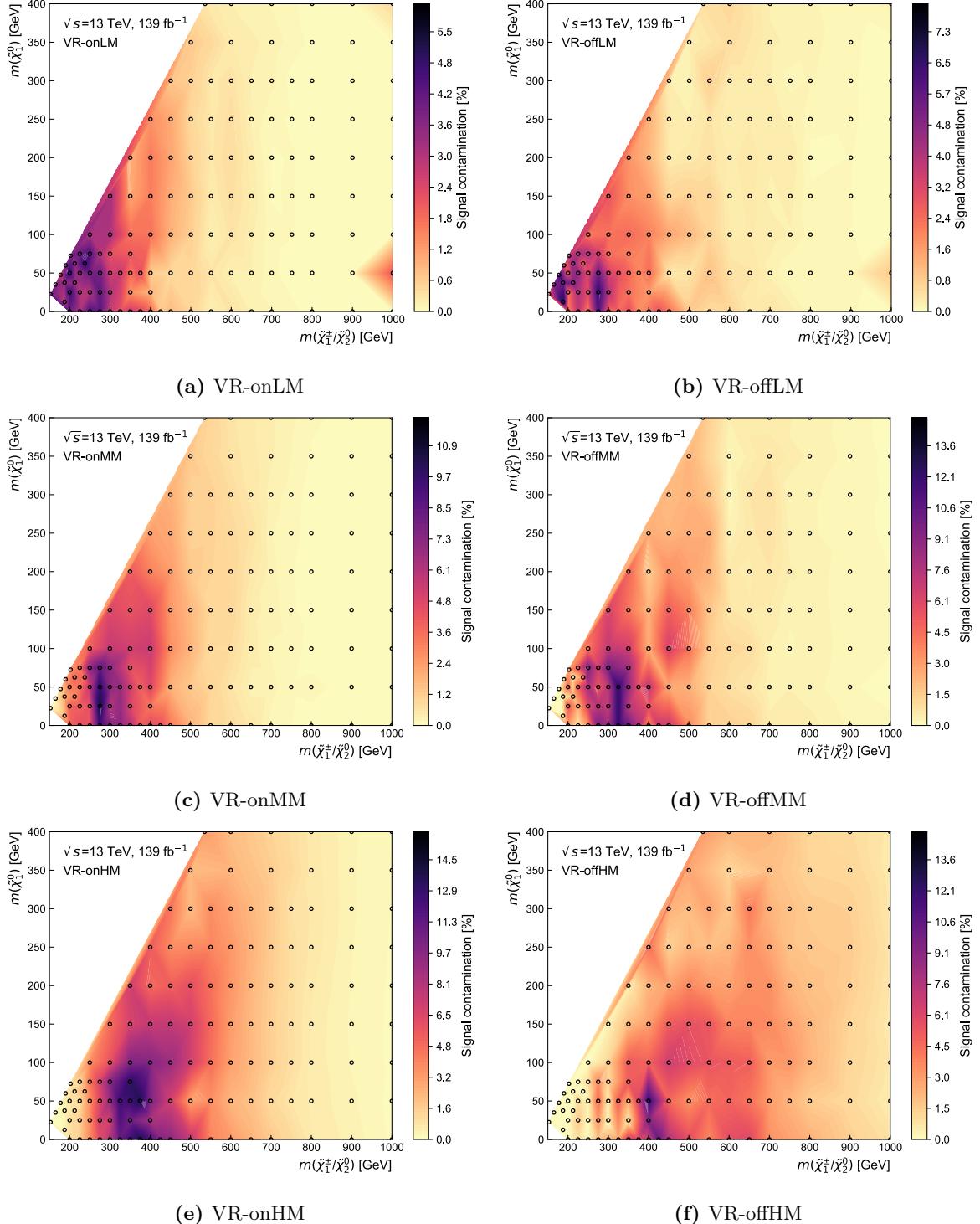


**Figure A.7:** N-1 plots for the chosen cut combination for the  $(m(\tilde{\chi}_1^\pm/\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)) = (300 \text{ GeV}, 150 \text{ GeV})$  signal point. The shaded region includes MC statistical uncertainty as well as 30% systematic uncertainty (added in quadrature) on the background. The significance is computed using the binomial discovery significance using the uncertainty on the background.



**Figure A.8:** Comparison of different shape-fit configurations. Fig. (a) compares three different two-dimensional shape-fit configurations using  $3 \times 3$  bins in  $(m_T, E_T^{\text{miss}})$ ,  $(m_T, m_{b\bar{b}})$  and  $(m_T, m_{CT})$ . Fig. (b) illustrates the sensitivity increase achieved through a requirement on high  $m_{elb_1}$  values in SR-HM on top of the two-dimensional shape-fit in  $m_T$  and  $m_{CT}$ . Fig. (c) compares the two-dimensional shape-fit in  $m_T$  and  $m_{CT}$  to the previous analysis iteration signal regions scaled to  $139 \text{ fb}^{-1}$ . All shown exclusion limits are expected limits at 95% CL, using MC statistical and 30% systematic uncertainty. Background estimation in the signal regions is taken directly from MC for all SM backgrounds.

## A.2 Background estimation



**Figure A.9:** Signal contamination (shown on the  $z$ -axis) for all VRs throughout the signal grid. The space between the signal points (indicated by the black circles) is interpolated using Delaunay triangles.

ATLAS SUSY Searches* - 95% CL Lower Limits									
Model	Signature	$f_{\mathcal{L}} dt$ [fb $^{-1}$ ]	Mass limit						Reference
$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{q}_1^0$	0 e, $\mu$ mono-jet	2-6 jets	$E_T^{\text{miss}}$	139	[0× Degen.]	0.43	0.71	1.9	$m(\tilde{q}^0)=400\text{ GeV}$ $m(\tilde{q})=m(\tilde{q}_1^0)=5\text{ GeV}$
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	0 e, $\mu$	2-6 jets	$E_T^{\text{miss}}$	139	[1× Bkg Degen.]	0.43	0.71	2.35	ATLAS-CONF-2019-040 1711.03301
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0 \chi_1^0$	1 e, $\mu$	2-6 jets	$E_T^{\text{miss}}$	139	[2]	Forbidden	1.15-1.95	2.2	ATLAS-CONF-2019-040 ATLAS-CONF-2020-047
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0 \chi_1^0$	ee, $\mu\mu$	2 jets	$E_T^{\text{miss}}$	36.1	[2]	1.2	1.2	2.2	ATLAS-CONF-2020-047 1805.11381
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0 Z\chi_1^0$	0 e, $\mu$	7-11 jets	$E_T^{\text{miss}}$	139	[2]	1.15	1.97	1.97	ATLAS-CONF-2020-022 1908.08457
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{g}\tilde{g}$	SS e, $\mu$	6 jets	$E_T^{\text{miss}}$	139	[2]	1.25	2.25	2.25	ATLAS-CONF-2018-041 1909.08457
$\tilde{g}\tilde{g}, \tilde{g} \rightarrow \tilde{g}\tilde{g}$	0-1 e, $\mu$	3 jets	$E_T^{\text{miss}}$	79.8	[2]	1.25	2.25	2.25	ATLAS-CONF-2018-041 1909.08457
<b>Inclusive Searches</b>	SS e, $\mu$	6 jets	$E_T^{\text{miss}}$	139	[2]	Forbidden	0.74	0.9	$m(\tilde{q}^0)=300\text{ GeV}$ $BR(\tilde{q}^0\rightarrow b\bar{b})=1$ $m(\tilde{q}^0)=300\text{ GeV}$ $BR(\tilde{q}^0\rightarrow \tilde{q}\tilde{q})=1$ $\Delta m(\tilde{q}^0)=300\text{ GeV}$
$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{q}_1^0/\tilde{b}_1^{\pm}$	0 e, $\mu$	Multiple	$E_T^{\text{miss}}$	36.1	[1]	Forbidden	0.74	0.9	ATLAS-CONF-2019-040 1708.09266-1711.03301
$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{q}_1^0$	0 e, $\mu$	2 jets	$E_T^{\text{miss}}$	139	[1]	Forbidden	0.13-0.85	0.23-1.35	1909.08457
$\tilde{b}_1\tilde{b}_1, \tilde{b}_1\rightarrow b\tilde{q}_1^0$	0 e, $\mu$	≥ 1 jet	$E_T^{\text{miss}}$	139	[1]	1.25	1.25	1.25	ATLAS-CONF-2020-031 1908.03122
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	1 e, $\mu$	3 jets	$E_T^{\text{miss}}$	139	[1]	0.44-0.59	0.44-0.59	0.44-0.59	ATLAS-CONF-2020-003 2004.14960
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	1 e, $\mu$	2 jets	$E_T^{\text{miss}}$	36.1	[1]	1.16	1.16	1.16	ATLAS-CONF-2019-017 1803.10178
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	0 e, $\mu$	2 jets	$E_T^{\text{miss}}$	36.1	[1]	0.85	0.85	0.85	1805.01849
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	0 e, $\mu$	mono-jet	$E_T^{\text{miss}}$	36.1	[1]	0.43	0.43	0.43	1711.03301
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}_2^0, \tilde{t}_2^0\rightarrow Z\tilde{h}\tilde{h}_1^0$	1-2 e, $\mu$	1-2 jets	$E_T^{\text{miss}}$	139	[1]	Forbidden	0.067-1.18	0.067-1.18	SUSY-2018-09
$\tilde{t}_1\tilde{t}_2, \tilde{t}_2\rightarrow \tilde{t}_1 + Z$	3 e, $\mu$	3 jets	$E_T^{\text{miss}}$	139	[1]	Forbidden	0.36	0.36	SUSY-2018-09
$\tilde{X}_1^{\pm 1/2}$ via $WZ$	3 e, $\mu$	1 jet	$E_T^{\text{miss}}$	139	[1]	0.64	0.64	0.64	SUSY-2018-09
$\tilde{X}_1^{\pm 1/2}$ via $WW$	2 e, $\mu$	1 jet	$E_T^{\text{miss}}$	139	[1]	0.205	0.205	0.205	SUSY-2018-09
$\tilde{X}_1^{\pm 1/2}$ via $W h$	0-1 e, $\mu$	2 jets	$E_T^{\text{miss}}$	139	[1]	0.42	0.42	0.42	SUSY-2018-09
$\tilde{X}_1^{\pm 1/2}$ via $\tilde{Z}_1/\tilde{B}$	2 e, $\mu$	2 jets	$E_T^{\text{miss}}$	139	[1]	0.74	0.74	0.74	SUSY-2018-09
<b>EW direct production</b>	0 e, $\mu$	2 jets	$E_T^{\text{miss}}$	139	[1]	0.16-0.3	0.16-0.3	0.16-0.3	ATLAS-CONF-2020-015 1911.12056
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	2 e, $\mu$	0 jets	$E_T^{\text{miss}}$	139	[1]	0.256	0.256	0.256	ATLAS-CONF-2020-015 1908.09228
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	2 e, $\mu$	1 jet	$E_T^{\text{miss}}$	139	[1]	0.13-0.23	0.13-0.23	0.13-0.23	ATLAS-CONF-2020-015 1908.09228
$\tilde{t}_1\tilde{t}_1, \tilde{t}_1\rightarrow \tilde{t}\tilde{q}_1^0$	0 e, $\mu$	0 jets	$E_T^{\text{miss}}$	139	[1]	0.55	0.55	0.55	ATLAS-CONF-2020-015 1908.09228
<b>EW direct</b>	0 e, $\mu$	4 jets	$E_T^{\text{miss}}$	139	[1]	0.55	0.55	0.55	ATLAS-CONF-2020-015 1908.09228
<b>EW direct</b>	0 e, $\mu$	1 jet	$E_T^{\text{miss}}$	36.1	[1]	0.46	0.46	0.46	ATLAS-CONF-2020-015 1908.09228
<b>Stable <math>\tilde{g}</math>-R-hadron</b>	Multiple	Multiple	$E_T^{\text{miss}}$	36.1	[1]	2.0	2.0	2.0	Pure Wine Pure Higgsino
<b>Metastable <math>\tilde{g}</math>-R-hadron, <math>\tilde{s} \rightarrow q\tilde{q}_1^0</math></b>	3 e, $\mu$	0 jets	$E_T^{\text{miss}}$	139	[1]	0.625	0.625	0.625	ATLAS-CONF-2020-015 1902.01536-1808.04095
<b>Long-lived</b>	1 $F_V$ , $m_{\tilde{q}^0}=X$	$\tilde{q}^0\rightarrow Z\ell^-\ell^+$	$E_T^{\text{miss}}$	32	[1]	1.9	1.9	1.9	Pure Wine
	$\tilde{X}_1^{\pm 1/2}/\tilde{K}_1^0 \rightarrow Z\ell^-\ell^+$	4 jets	$E_T^{\text{miss}}$	36.1	[1]	0.82	0.82	0.82	ATLAS-CONF-2020-015 1607.08079
	$\tilde{X}_1^{\pm 1/2}/\tilde{K}_1^0 \rightarrow WW/Z\ell^-\ell^+$	4-5 large-R jets	$E_T^{\text{miss}}$	36.1	[1]	1.3	1.3	1.3	1802.01536
<b>RPy</b>	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	Multiple	$E_T^{\text{miss}}$	36.1	[1]	0.05	0.05	0.05	Large $\tilde{A}_{11}^{(2)}$ $m(\tilde{q}^0)=200\text{ GeV}$ $100\text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	Multiple	$E_T^{\text{miss}}$	36.1	[1]	1.05	1.05	1.05	ATLAS-CONF-2019-003 $m(\tilde{q}^0)=200\text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	Multiple	$E_T^{\text{miss}}$	36.1	[1]	0.95	0.95	0.95	ATLAS-CONF-2020-016 $m(\tilde{q}^0)=50\text{ GeV}$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	2 jets + b	$E_T^{\text{miss}}$	36.1	[1]	0.42	0.61	0.61	ATLAS-CONF-2020-016 1710.05771
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	2 jets + b	$E_T^{\text{miss}}$	36.1	[1]	0.41-0.45	0.41-0.45	0.41-0.45	$BR(\tilde{q}^0\rightarrow \tau^+\tau^-)=20\%$
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}_1^0$	1 $\mu$	DV	136	[1]	1.0	1.0	1.0	$BR(\tilde{q}^0\rightarrow \mu^+\mu^-)=100\%$ $cos\theta_H=1$

\*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

**Figure A.10:** Summary of the ATLAS searches for SUSY. A representative selection of the available search results is shown. Results are given with respect to the nominal cross section. In some cases these additional dependencies are indicated by darker bands showing different model parameters. Figure adapted from Ref. [244].



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## **Selbstständigkeitserklärung**

Hiermit erkläre ich, die vorliegende Arbeit mit dem Titel

**Search for electroweakinos using the ATLAS detector**

Suche nach Elektroweakinos mit dem ATLAS Detektor

selbstständig verfasst zu haben und keine anderen als die in der Arbeit angegebenen Quellen und Hilfsmittel benutzt zu haben.

Eric Schanet

München, den 01. Mai 2021