

Pressure variation in a fluid at rest:

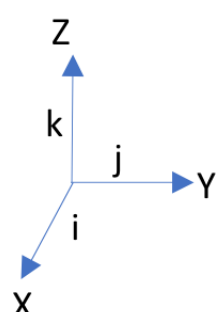
**Concept:** For a fluid at rest the pressure at a point depends only on the weight of fluid vertically above that point ("Vertical" = parallel to gravitational vector) and any horizontal plane is an isobar (constant pressure plane).

**Aim :** To show that

(1) for liquids  $P_1 = P_{\text{surface}} + \rho g h$

(2) for gases (isothermal)  $P_1 = P_0 e^{\left[ \frac{-g(z_1 - z_0)}{RT_0} \right]}$

(3) for real atmosphere  $P_1 = P_0 \left( 1 - \frac{\beta z_1}{T_0} \right)^{g/\beta}$



$\beta = 0.0065 \text{ K/m}$  (Kelvin/m)

$T_0$  = absolute temperature at sea level  
 $P_0$  = absolute pressure at sea level  
 $R = 287 \text{ J/Kg-K}$

(Σ) → top side of the liquid..

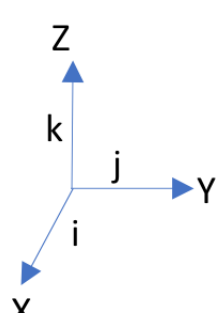
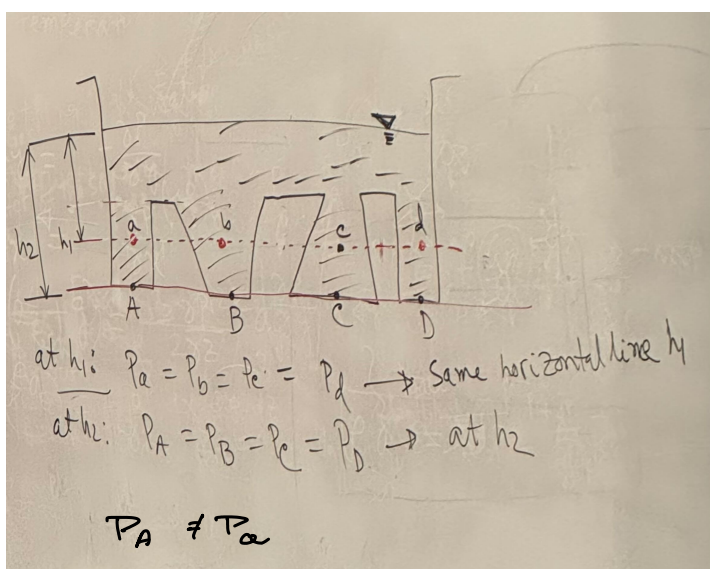
Pressure variation in a fluid at rest:

We have seen that

$\rho a_x = \frac{\partial P}{\partial x} = 0$

$\rho a_y = \frac{\partial P}{\partial y} = 0$

$\rho a_z = -\frac{\partial P}{\partial z} - \rho g = 0$



For a fluid at rest:

$a_x = a_y = a_z = 0$

①  $\frac{\partial P}{\partial x} = 0$       ②  $\frac{\partial P}{\partial y} = 0$

x-y plane  
acts as an isobar

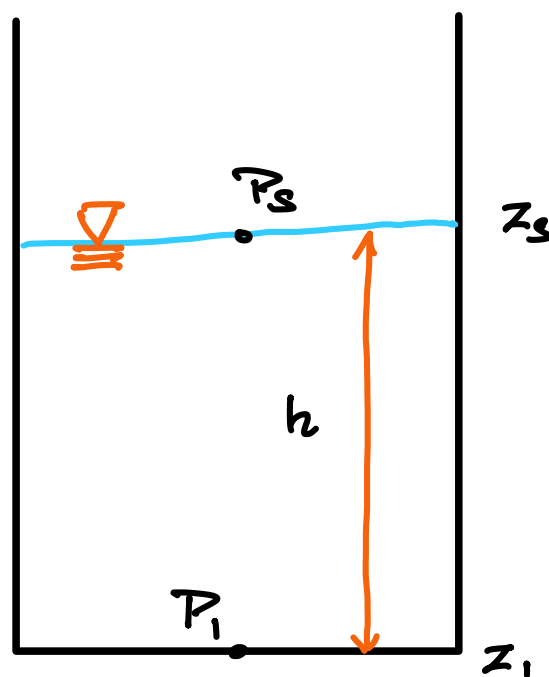
③  $-\frac{\partial P}{\partial z} - \rho g = 0$

$\frac{dP}{dz} = -\rho g$

Pressure varies only in z direction (vertically)

(1) Liquids: Hydrostatic distribution: (assume  $\rho g = \gamma$  = constant)

We have seen that



$\frac{dP}{dz} = -\rho g$  || integrate

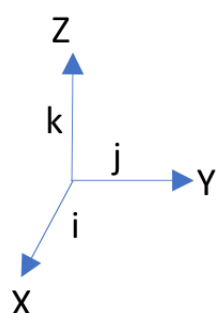
$\int_{P_1}^{P_s} dP = \int_{z_1}^{z_s} -\rho g dz$

$P_s - P_1 = -\rho g (z_s - z_1)$

$P_1 = P_s + \rho g (z_s - z_1)$

$P_1 = P_s + \rho g h$

$P \propto h$

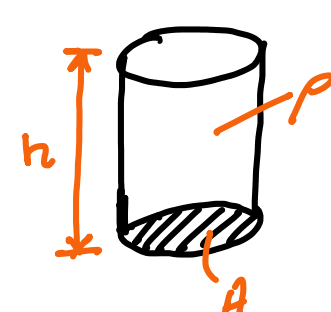


∴ Pressure *increases linearly* with depth in a liquid

(1) Liquids: Hydrostatic distribution: continued

What is the term  $\rho g h$ ? no idea

Consider liquid column of height  $h$  and small base area  $A$ :



Weight of column =  $m \cdot g$

$= \rho V g$

$= \rho A h g$

Pressure at base =  $\frac{N}{A}$

$= \frac{\rho A h g}{A}$

$= \rho h g$

∴ We can express any **pressure  $P$  (Pa, N/m<sup>2</sup>)** in terms of a height  $h$  (m) of the fluid of **density  $\rho$  (kg/m<sup>3</sup>)** i.e.,

(2) Compressible medium (isothermal gas):

$T = T_0 = \text{constant}$

- Important when height variation are large (>100 m)
- A very good approximation for the real atmosphere for heights <500 m

Assume isothermal gas ∴  $T = T_0$

$\frac{dP}{dz} = -\rho g$        $\rho = \frac{P}{RT_0}$

$\frac{dP}{dz} = -\frac{P g}{RT_0}$

$\int_{P_0}^{P_1} \frac{dP}{P} = -\frac{g}{RT_0} \int_{z_0}^{z_1} dz$

$\ln(P) \Big|_{P_0}^{P_1} = -\frac{g}{RT_0} z \Big|_{z_0}^{z_1}$

$\ln\left(\frac{P_1}{P_0}\right) = -\frac{g}{RT_0} h$

$\frac{P_1}{P_0} = e^{-\left(\frac{g h}{RT_0}\right)}$

(3) Real (Standard) atmosphere:

- Idealized or "average" representation of earth's atmosphere
- Good for design purposes
- Has linear decrease of  $T$  with  $z$  in lower 11 km (troposphere)

∴  $T = T_0 - \beta z$

$\frac{dP}{dz} = -\rho g$        $\rho = \frac{P}{RT}$

$\frac{dP}{dz} = -\frac{P g}{RT}$

$\frac{dP}{dz} = -\frac{P g}{R(T_0 - \beta z)}$

$\int_{P_0}^{P_1} \frac{dP}{P} = -\frac{g}{R} \int_{z_0}^{z_1} \frac{1}{T_0 - \beta z} dz$

$\ln\left(\frac{P_1}{P_0}\right) = \frac{g}{R \beta} \left[ \ln(T_0 - \beta z_1) - \ln(T_0 - \beta z_0) \right]$

$= \frac{g}{R \beta} \left[ \ln\left(\frac{T_0 - \beta z_1}{T_0 - \beta z_0}\right) \right]$

$\frac{P_1}{P_0} = \left( \frac{T_0 - \beta z_1}{T_0 - \beta z_0} \right)^{g/\beta}$

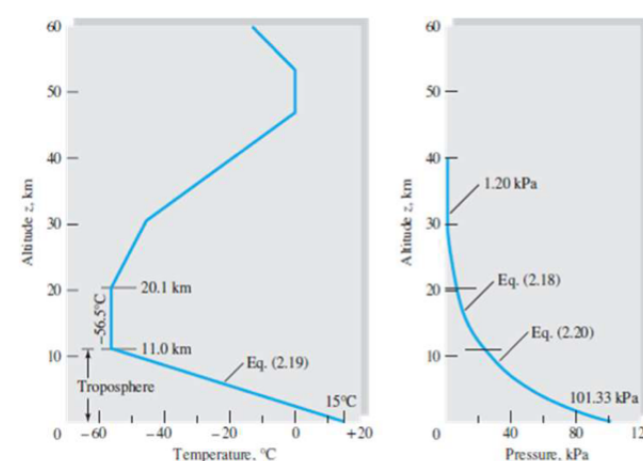


Figure 2.7 (FM White): Temperature and pressure distribution in Standard Atmosphere.