

Trajectory Optimization for an Inertial Cube

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Abstract—Inertially-actuated robots pose interesting problems for underactuated control. Here we present a trajectory optimization method to control a 2d cube with a torque-limited flywheel mounted inside. Inspired by earlier work done with the Cubli [2] and M-Blocks [3], we use contact-implicit trajectory optimization to generate trajectories for the inertial cube. We use one actuator (the flywheel) to control the 8 states of the robot over time. Using Sparse Nonlinear OPTimizer (SNOPT) in Drake, the inertial cube can be described as a floating body with contact implicit movements via linear complementarity constraints. This leads to elegant optimal trajectory control for the inertial cube.

Index Terms—trajectory optimization, underactuated, inertia, contact-implicit

I. INTRODUCTION

The inertial cube is particularly interesting both due to its extreme underactuation and its simplicity. After seeing the potential of M-Blocks [3], it's clear that inertially-actuated cubes have potential to be used in a variety of applications. The possibilities of movements are great, as seen in this video <https://www.youtube.com/watch?v=Nns0qzd8Noo>. The cubli also presents very nice control in stabilization to resist movement https://www.youtube.com/watch?v=n_6p-1J551Y. By combining prior work with ideas from optimal control and trajectory optimization, the inertial cube has potential for much more complex behavior. With elegant algorithms, we may be able to help robots explore extraterrestrial space [6] or create configurable robots like the nanobots in the movie *Big Hero 6*. We are interested in applying our best underactuated control algorithms to the system.

II. DEFINING THE MODEL

In this section, we describe the inertial cube in floating-body coordinates with a few simplifying assumptions. The dynamics work out nicely when described in this way. We can then proceed to optimize over the state trajectory, input torque, and external contact forces over time.

A. State Space Explanation

Fig. 1 is a diagram depicting the states of the cube. The cube is defined in floating-coordinates, meaning there is no notion of the ground in the state description. However, there are 4 contact vectors, 1 for each corner. These are described in the figure.

Following from the cube diagram, the full states can be described in the following way.

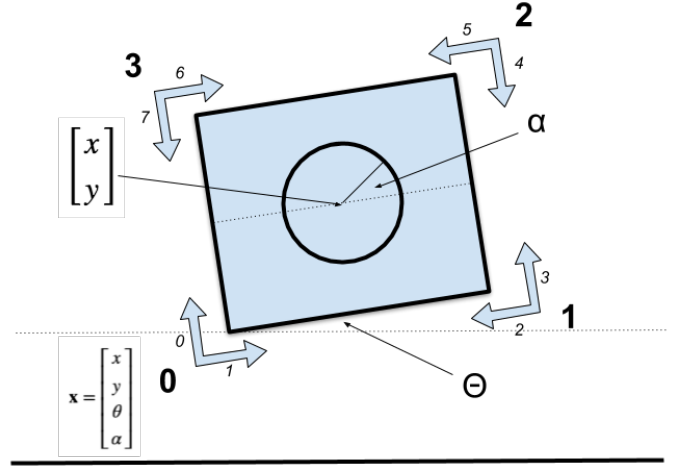


Fig. 1. The states of the floating inertial cube. The derivatives of these states are also included in the total state vector \mathbf{x} but are not shown here for convenience.

$$\mathbf{x} = [x \quad y \quad \theta \quad \alpha \quad \dot{x} \quad \dot{y} \quad \dot{\theta} \quad \dot{\alpha}]^T \quad (1)$$

$$\dot{\mathbf{x}} = [\dot{x} \quad \dot{y} \quad \dot{\theta} \quad \dot{\alpha} \quad \ddot{x} \quad \ddot{y} \quad \ddot{\theta} \quad \ddot{\alpha}]^T \quad (2)$$

These state vector \mathbf{x} includes both the positions, angles, and their velocities. This is standard state-space notation. Following from Fig. 1, the dynamics are the following. These dynamics are informed by prior work done at Chalmers University of Technology [4] but modified for this state-space description and free-body coordinates.

$$\ddot{x} = (f_1 - f_2 + f_6 - f_5) \cos \theta - (f_0 + f_3 - f_4 - f_7) \sin \theta \quad (3)$$

$$\ddot{y} = (f_1 - f_2 + f_6 - f_5) \sin \theta + (f_0 + f_3 - f_4 - f_7) \cos \theta - g \quad (4)$$

$$\ddot{\theta} = \frac{-u + b_w \dot{\alpha} - b_c \dot{\theta}}{I_c} + \frac{1}{2} \left(\sum_{n \in \{1,3,5,7\}} f_n - \sum_{n \in \{0,2,4,6\}} f_n \right) \quad (5)$$

$$\ddot{\alpha} = \frac{u(I_c + I_w) + b_c I_w \dot{\theta} - b_w \frac{I_c + I_w}{2} \dot{\alpha}}{I_w I_c} \quad (6)$$

Here we explain the notation used in the dynamics and the few simplifying assumptions made.

- f_n is the force acting on the corner based on Fig. 1.

- u is the torque on the interial wheel.
- The cube is unit sized, meaning its dimensions are 1 for each edge length.
- b_w and b_c are the friction coefficients for the wheel and cube respectively.
- I_w and I_c are the moments of inertia for the wheel and cube respectively.
- The mass of the cube and wheel are each 0.5 kg.

In this paper, the static values are the following. We chose not to spend much time on choosing accurate values. The purpose of this project is to use contact-implicit trajectory optimization, which is described in the next section.

- $b_w = b_c = 0.5$
- $I_w = I_c = 0.5$
- $g = 9.81$

III. OPTIMIZATION FORMULATION

In this section, we explain the necessary nonlinear program formulation to solve for a cube trajectory through space. We'll start by explaining this in the context of a swing up of the cube. This motion involves contact with the ground, which is why we are choosing to use contact-implicit trajectory optimization to avoid dealing with many modes of the system [1].

A. The Swing-Up Formulation

The swing-up is performed by getting the cube to stand on exactly one corner in a stable position. By using the contact-implicit trajectory optimization, this is possible. We formulate the optimization problem with direct transcription and a complementarity constraint to keep the cube above the ground.

$$\begin{aligned}
 & \text{find} \\
 & x[0:N], u[0:N], f[0:N] \\
 & \text{subject to} \\
 & x[n+1] = x[n] + f(x[n], u[n], f[n])dt, \\
 & n \in [0, N-1], \\
 & -u_{max} < u < u_{max}, \\
 & 0 < f[n][i] < f_{max}, i \in [0, 7], \\
 & c[n][:] \cdot \phi[:] = 0.0, \\
 & x[0] = x_{initial}, \\
 & x[N] = x_{final},
 \end{aligned}$$

Here is a description of the values used in the optimization.

- $f(x[n], u[n], f[n])$ is \dot{x} .
- dt is the timestep.
- N is the number of knot points, which will implicitly define the timestep.
- $0 < f[n][:] < f_{max}$ to make sure contact forces only push away from the ground
- $c[n][i]$ is an intermediate value describing if there is a force on the i^{th} corner. $c[n][0] = f[n][0] * f[n][1]$
 $c[n][1] = f[n][2] * f[n][3]$
 $c[n][2] = f[n][4] * f[n][5]$
 $c[n][3] = f[n][6] * f[n][7]$
- $\phi[i]$ is the distance from the i^{th} corner to the ground

B. Swing-Up Results

The computed trajectory values are shown in Fig. 2, 3, and 4. This computed result is formatted without a cost function. Rather, we are looking for a solution that satisfied all of our constraints. The trajectory in these figures start on the ground and stands on corner 0. See this video for live graphical depiction of the states over time <https://www.youtube.com/watch?v=EH4mXZSPz5o>.

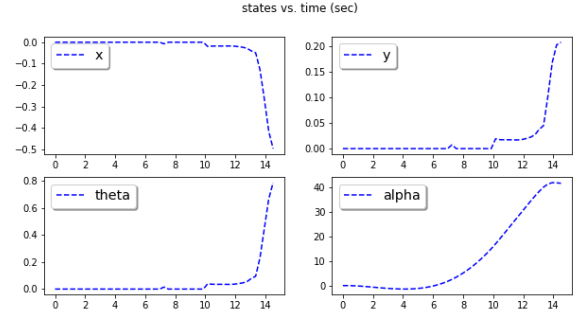


Fig. 2. Swing up trajectory states.

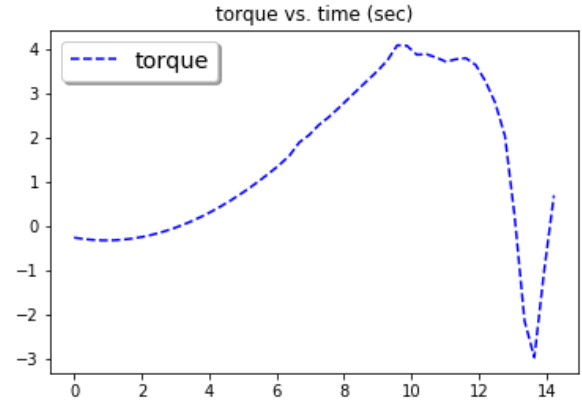


Fig. 3. Swing up trajectory torque input.

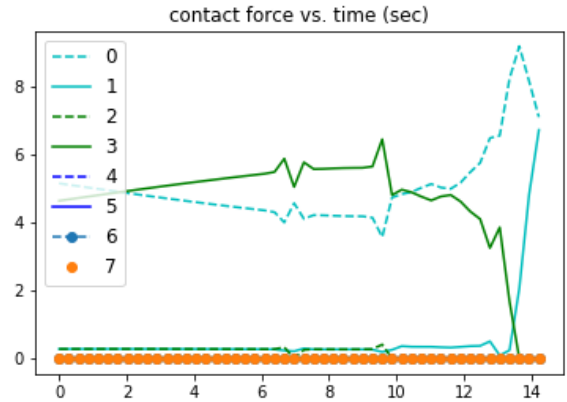


Fig. 4. Swing up trajectory ground forces.

After achieving the swing up in this way, one could most definitely control the position using LQR [1]. However, this

was not attempted in our trajectory optimization work.

C. Experiments in Higher Dimensional State Space

Explain how 3D seemed to help in some cases even though some states don't enter dynamics. We did a few experiments are reading some work done on nonlinear optimization in high dimensions [5].

D. Stable Walking Motion

Explain how to achieve stable walking motion.

RESULTS

In this section we discuss overall results and trends discovered throughout the more specific examples covered thus far.

FUTURE WORK

Future work section.

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