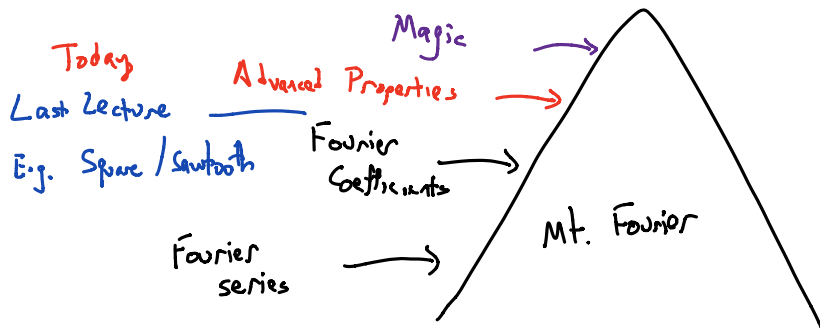


Ethan CEZ  
Jason EB1  
Timmy CBZ



# EE102

## Lecture 11

# EE102 Announcements

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- Syllabus link is [tiny.cc/ucla102](https://tiny.cc/ucla102)
- HW4 due Friday
- Student Feedback (not hitting response rate, but good faith bonus applied).  
Feedback is extremely helpful.

**Slide Credits:** This lecture adapted from Prof. Jonathan Kao (UCLA) and recursive credits apply to Prof. Cabric (UCLA), Prof. Yu (Carnegie Mellon), Prof. Pauly (Stanford), and Prof. Fragouli (UCLA)

# Review from Last Time

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- **Fourier Series:** Learned that periodic signal  $\Rightarrow$  broken down into sines/cosines, or more formally, complex exponentials.
- **Fourier Coefficients:** The coefficients of these sines/cosines can be solved for.
- **Previous Learning Milestones**
  - Given a signal, you should be able to compute its fourier coefficients (we did this for square/sawtooth)
- Today's goals are two-fold
  - Advanced Fourier Properties
  - Does Fourier apply for periodicity alone?

# Summary of Fourier Series

---

These are the main mathematical results of this lecture, written here for convenience.

If  $f(t)$  is a well-behaved periodic signal with period  $T_0$ , then  $f(t)$  can be written as a Fourier series

Fourier  
Series

$$f(t) = \sum_{k=-\infty}^{\infty} c_k e^{jk\omega_0 t}$$

where  $\omega_0 = \frac{2\pi}{T_0}$  and

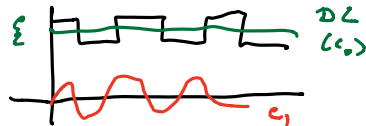
Fourier Series  
Coefficients

$$c_k = \frac{1}{T_0} \int_{\tau}^{\tau+T_0} f(t) e^{-jk\omega_0 t} dt$$

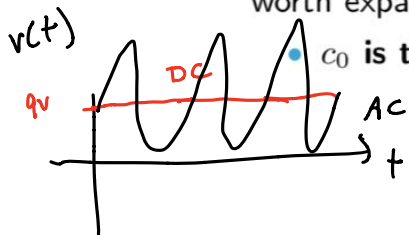
for all integers  $k$ . The  $c_k$  are called the *Fourier coefficients* of  $f(t)$ .

Here,  $f(t)$  is the *weighted average* of complex exponentials (which are simply complex sines and cosines).

# Fourier Series Properties



There are interesting symmetries and properties of the Fourier series that are worth expanding upon.

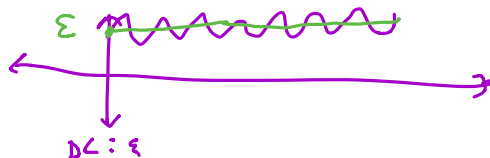
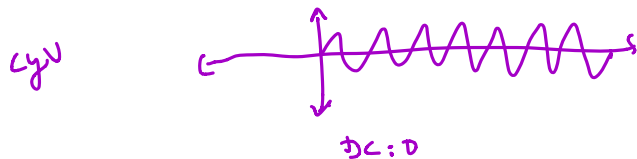


$c_0$  is the average of the signal. Note that for  $k = 0$ , we have that

$$c_0 = \frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) dt$$

Same expression  
as average of a function

Thus,  $c_0$  is exactly the time-averaged mean of the signal and corresponds to a constant value (i.e., it has no sinusoidal component). For this reason, it is sometimes called the “DC component.” DC stands for direct current in circuits, and refers to non-alternating (sinusoidal) currents. The DC component is the average value taken on by a signal.

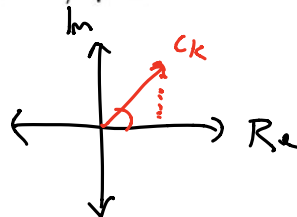


# Fourier Series Properties

- **Complex representation.** In general, the  $c_k$  may be complex, and so they can be expressed in their real / imaginary form or in magnitude / phase form. i.e.,

$$\begin{aligned} c_k &= \Re(c_k) + j\Im(c_k) \\ &= |c_k| e^{j\angle c_k} \end{aligned}$$

Amplitude Spectrum      Phase Spectrum



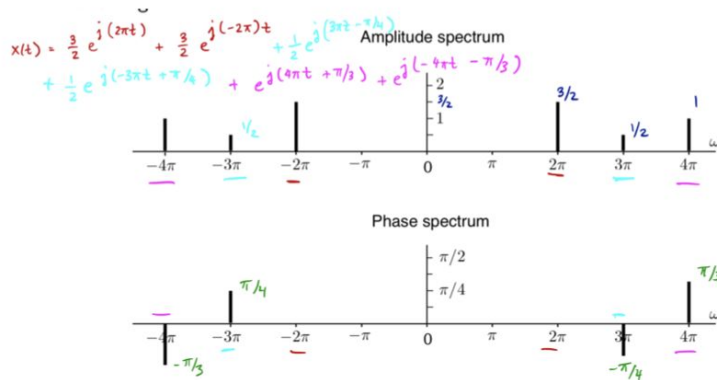
$c_k$  is  
complex often

$a + jb$

# Fourier Series Properties

- **Complex representation.** In general, the  $c_k$  may be complex, and so they can be expressed in their real / imaginary form or in magnitude / phase form. i.e.,

$$\begin{aligned} c_k &= \Re(c_k) + j\Im(c_k) \\ &= |c_k| e^{j\angle c_k} \end{aligned}$$



// Amplitude Spectrum.  
It tells you how much magnitude of frequency  $\omega_0$  is in the signal

// Phase Spectrum: Mixture of cosines and sines at that frequency.

$$\# = \int f(t) g(t) dt \quad \rightarrow \quad \sum f_i g_i = \langle f, g \rangle$$

# Fourier Series Properties

## Fourier symmetry

We can apply Euler's formula to re-write the Fourier coefficients, and reveal some symmetries:

$$c_k = a - jb$$

F.C. //

Apply Euler's

Algebra

$$\begin{aligned} c_k &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) e^{-j \frac{2\pi k t}{T_0}} dt \\ &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) \left[ \cos\left(\frac{2\pi k}{T_0} t\right) - j \sin\left(\frac{2\pi k}{T_0} t\right) \right] dt \\ &= \boxed{\frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) \cos\left(\frac{2\pi k}{T_0} t\right) dt} - \boxed{j \frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) \sin\left(\frac{2\pi k}{T_0} t\right) dt} \end{aligned}$$

If  $f(t)$  is real, then so are:

$$\Re(c_k) = \frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) \cos\left(\frac{2\pi k}{T_0} t\right) dt$$

$$\Im(c_k) = -\frac{1}{T_0} \int_{t_0}^{t_0+T_0} f(t) \sin\left(\frac{2\pi k}{T_0} t\right) dt$$



# Fourier Series Properties

## Fourier symmetry (cont.)

Therefore, for  $f(t)$  real, and using the fact that  $\cos(k)$  is even and  $\sin(k)$  is odd, we have the following symmetries:

$$|c_k| = \sqrt{(\operatorname{Re}\{c_k\})^2 + (\operatorname{Im}\{c_k\})^2}$$

$$\Re(c_k) = \Re(c_{-k})$$

$$\Im(c_k) = -\Im(c_{-k})$$

$$c_k^* = c_{-k}$$

$$|c_k| = |c_{-k}|$$

$$\angle c_k = -\angle c_k^*$$

$$\angle c_k = \operatorname{atan} \left[ \frac{\operatorname{Im}\{c_k\}}{\operatorname{Re}\{c_k\}} \right]$$

$$\text{Even: } f(t) = f(-t)$$

$$\text{Odd: } f(t) = -f(-t)$$

$$c_k = \operatorname{Re}\{c_k\} + j \operatorname{Im}\{c_k\}$$

$$c_k^* = \operatorname{Re}\{c_k\} - j \operatorname{Im}\{c_k\}$$

$$= \operatorname{Re}\{c_{-k}\} + j \operatorname{Im}\{c_{-k}\}$$

$c_{-k}$

cyv: Prove that  $c_k^* = c_{-k}$

# CYU

$$\tau \triangleq -t$$
$$d\tau = -dt$$

What is the relationship between the Fourier series coefficients  $c_k$  and  $c_{-k}$  if  $x(t)$  is even?

Refresher: 
$$c_k = \frac{1}{T_0} \int_{t_0}^{t_0+T_0} x(t) e^{-jk\omega_0 t} dt$$

$$c_k = \frac{1}{T_0} \int_0^{T_0} x(t) e^{-jk\omega_0 t} dt$$

$$c_{-k} = \frac{1}{T_0} \int_0^{T_0} x(t) e^{jk\omega_0 t} dt$$

$$= -\frac{1}{T_0} \int_0^{-T_0} x(-\tau) e^{-jk\omega_0 \tau} d\tau$$

$$= \frac{1}{T_0} \int_{-T_0}^0 x(-\tau) e^{-jk\omega_0 \tau} d\tau$$

$$= \frac{1}{T_0} \int_{-T_0}^0 x(\tau) e^{-jk\omega_0 \tau} d\tau$$

Extra Practice: CYU if  $x(t)$  is odd.

# Fourier Series Properties

- Combining facts, we have that if  $f(t)$  is even and real, then  $c_k = c_{-k}$  and  $c_{-k} = c_k^*$ , and so  $c_k = c_k^*$ . This means that the  $c_k$  must be real.

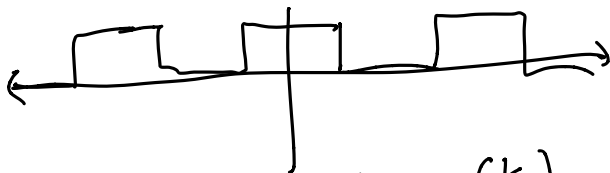
$$f(t) \text{ even and real} \implies c_k \text{ real}$$

- If  $f(t)$  is odd and real, then  $c_k = -c_{-k}$  and because  $c_{-k} = c_k^*$ , then  $c_k = -c_k^*$ . This means the  $c_k$  must be imaginary.

$$f(t) \text{ odd and real} \implies c_k \text{ imaginary}$$

① Math Exercises

② Practical Value



$$c_k \text{ is real: } \frac{1}{2} \sin\left(\frac{k}{2}\right)$$

Signal  $x(t)$

# Parseval's Theorem

Parseval's  
Thm

$$\frac{1}{T_0} \int_{t_0}^{t_0+T_0} |x(t)|^2 dt = \sum_{k=-\infty}^{\infty} |c_k|^2$$

LHS RHS

Measure in Time  $\longleftrightarrow$  Measure in Frequency

// cycl: Prove Parseval's Thm

Start LHS  $\rightarrow$  work it algebraically  $\rightarrow$  RHS

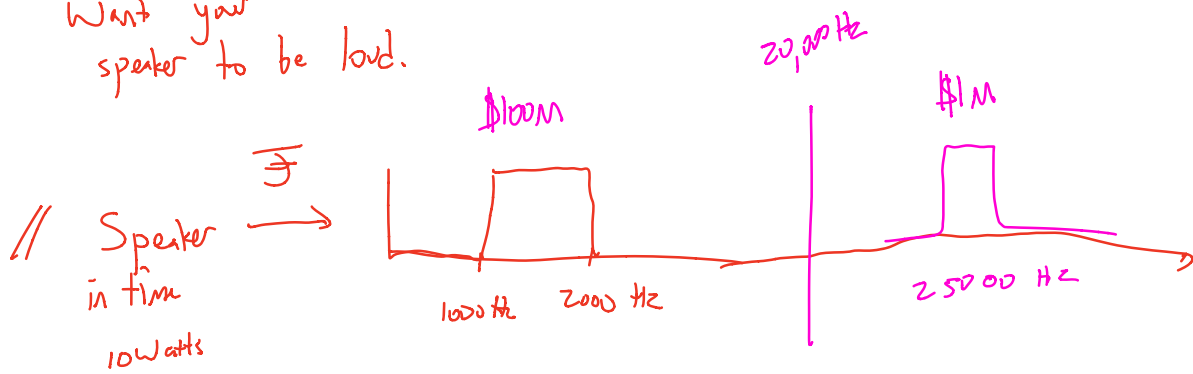
$$\begin{aligned} \frac{1}{T_0} \int_{t_0}^{t_0+T_0} |x(t)|^2 dt &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} x(t) \cdot x^*(t) dt \\ &= \frac{1}{T_0} \int_{t_0}^{t_0+T_0} \left( \sum_k c_k e^{jk\omega_0 t} \right) \left( \sum_n c_n^* e^{-jn\omega_0 t} \right) dt \\ &= \frac{1}{T_0} \sum_k \sum_n c_k c_n^* \int_{t_0}^{t_0+T_0} e^{j(k-n)\omega_0 t} dt \\ &= \frac{1}{T_0} \sum_k c_k c_k^* T_0 \longrightarrow \sum_k |c_k|^2 \end{aligned}$$

Integration trick of complex exp  
 $\begin{cases} T_0 & \text{if } k=n \\ 0 & \text{if } k \neq n \end{cases}$

## Practical Use of Parseval's Thm.

Making a speaker: Sending info here:  
1000 - 2000 Hz

Want your  
speaker to be loud.



Customer A

# What about non-periodic signals?

---

This is all well and good, but most signals we care about are not periodic.

## Motivation

Last lecture, we learned about the Fourier series, which can model (almost) any **periodic** or **time-limited** function as a sum of complex exponentials.

But the Fourier series is limited because it requires the signals be **periodic** or **time-limited**.

The Fourier transform allows us to calculate the spectrum of aperiodic signals.

# How do we go to aperiodic signals?

---

## Intuition of going from Fourier series to Fourier transform

Extending Fourier series to the Fourier transform is fairly intuitive.

The idea is the following.

- We can calculate the Fourier series of a periodic or time-limited signal, over some interval of length  $T_0$ .
- A signal that is not periodic can be viewed as a periodic signal, where  $T_0$  is infinite. As  $T_0$  is infinite, it never repeats.
- But the point is that we can replace our Fourier series calculation as, instead of being over a finite period,  $T_0$ , being over all time, from  $t = -\infty$  to  $\infty$ .

# Arriving at the Fourier Transform

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## Intuition (cont.)

Mathematically, we can calculate the Fourier series of  $f(t)$  over the interval  $[-T/2, T/2)$  via:

$$f(t) = \sum_{k=-\infty}^{\infty} c_k e^{jk\omega_0 t}$$

with

$$c_k = \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-jk\omega_0 t} dt$$

where  $\omega_0 = 2\pi/T$ .

In the Fourier transform, we're now going to let  $T \rightarrow \infty$ .

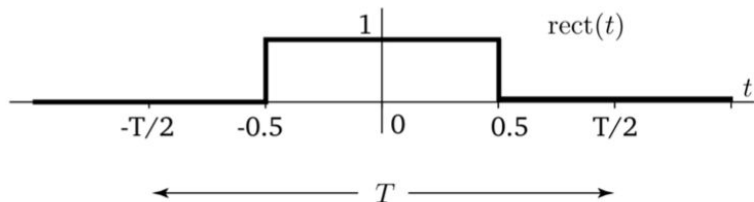


# Example: rect

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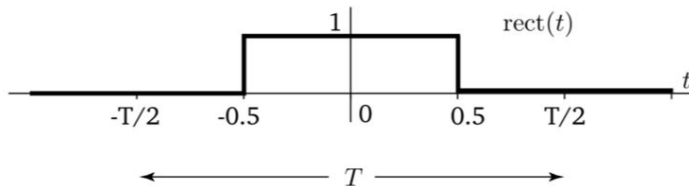
## The Fourier series of rect as $T \rightarrow \infty$

Consider  $x(t) = \text{rect}(t)$ . Further, let's define a period  $T$  over which, if we made a periodic extension of the rect, it would repeat every  $T$ . (Again, we're going to set  $T \rightarrow \infty$  eventually so that it doesn't repeat.)



# Example: rect

---



The Fourier series of this signal is related to the one we did last lecture, but we'll do it again for the sake of completeness:

$$\begin{aligned} c_k &= \frac{1}{T} \int_{-T/2}^{T/2} \text{rect}(t) e^{-jk \frac{2\pi}{T} t} dt \\ &= \frac{1}{T} \int_{-1/2}^{1/2} e^{-jk \frac{2\pi}{T} t} dt \\ &= \frac{1}{T} \left. \frac{T e^{-jk \frac{2\pi}{T} t}}{-jk 2\pi} \right|_{-1/2}^{1/2} \end{aligned}$$

# Example: rect

---

**The Fourier series of rect as  $T \rightarrow \infty$**

Continuing ...

$$\begin{aligned}c_k &= \left. \frac{e^{-jk\frac{2\pi}{T}t}}{-jk2\pi} \right|_{-1/2}^{1/2} \\&= \frac{-j\sin(\pi k/T) - j\sin(\pi k/T)}{-j2\pi k} \\&= \frac{\sin(\pi k/T)}{\pi k} \\&= \frac{1}{T} \frac{\sin(\pi k/T)}{\pi k/T} \\&= \frac{1}{T} \operatorname{sinc}\left(\frac{k}{T}\right)\end{aligned}$$

Therefore, the Fourier series of  $\operatorname{rect}(t)$  with a periodic extension every  $T$  is:

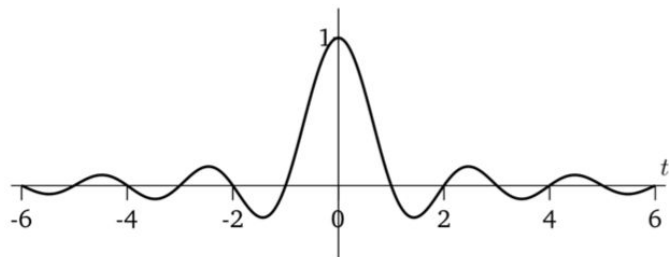
$$\frac{1}{T} \sum_{k=-\infty}^{\infty} \operatorname{sinc}(k/T) e^{jk\omega_0 t}$$

for  $\omega_0 = 2\pi/T$ .

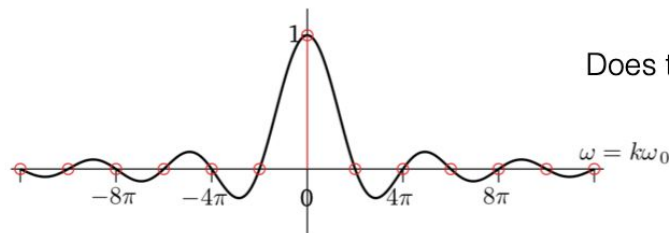
# Example: rect

## The Fourier series of rect as $T \rightarrow \infty$

Let's now look at what the coefficients look like for varying values of  $T$ . First, for a refresher, let's recall what the sinc function looks like.



Now, let's set  $T = 1$  ( $\omega_0 = 2\pi$ ) and calculate each of the Fourier coefficients.

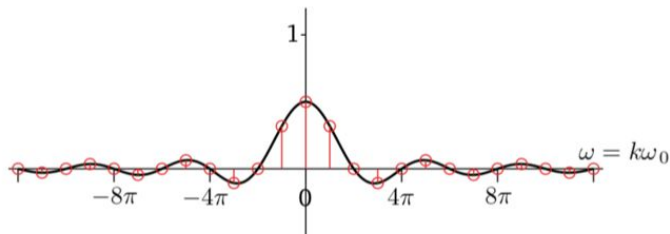


Does this make sense?

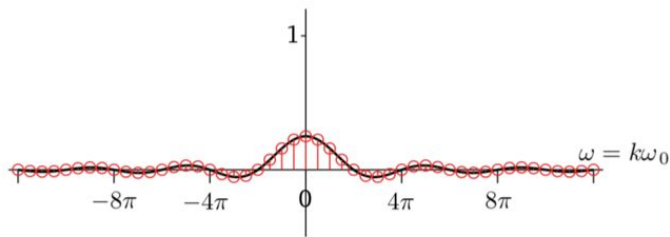
# Example: rect

## The Fourier series of rect as $T \rightarrow \infty$

Set  $T = 2$  ( $\omega_0 = \pi$ ) and calculate each of the Fourier coefficients.



Set  $T = 4$  ( $\omega_0 = \pi/2$ ) and calculate each of the Fourier coefficients.



What if we replace  $k\omega_0$  with  $\omega$ ?

# Example: rect

---

## The Fourier series of rect as $T \rightarrow \infty$

The trend we see is that as we set  $T$  larger, we more densely sample the sinc function. This gives us reason to believe that the spectrum of the rect signal, when  $T \rightarrow \infty$ , is a sinc function.

Let's formalize this intuition with math.

---

## Arriving at the Fourier transform

The Fourier series of  $f(t)$  on an interval  $[-T/2, T/2)$  is given by:

$$f(t) = \sum_{k=-\infty}^{\infty} c_k e^{jk\omega_0 t}$$

with

$$c_k = \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-jk\omega_0 t} dt$$

## Arriving at the Fourier transform

The Fourier series of  $f(t)$  on an interval  $[-T/2, T/2)$  is given by:

$$f(t) = \sum_{k=-\infty}^{\infty} c_k e^{jk\omega_0 t}$$

with

$$c_k = \frac{1}{T} \int_{-T/2}^{T/2} f(t) e^{-jk\omega_0 t} dt$$

We define the truncated Fourier transform as:

$$F_T(j\omega) = \int_{-T/2}^{T/2} f(t) e^{-j\omega t} dt$$

so that

$$c_k = \frac{1}{T} F_T(jk\omega_0)$$

Then,

$$f_T(t) = \sum_{k=-\infty}^{\infty} \frac{1}{T} F_T(jk\omega_0) e^{jk\omega_0 t}$$

Remember, we're going to replace  $k\omega_0$  with  $\omega$ .



## Arriving at the Fourier transform (cont.)

Now, let's set  $T \rightarrow \infty$ . If we do this, then  $\omega_0 = 2\pi/T$  will approach 0. So suppose instead that we define a continuous variable,

$$\omega = \frac{2\pi k}{T}$$

which means that  $k$  increases with  $T$ , so that  $\omega = k\omega_0$  is fixed.

The Fourier transform is the limit of the truncated Fourier transform.

$$\begin{aligned} F(j\omega) &= \lim_{T \rightarrow \infty} F_T(j\omega) \\ &= \lim_{T \rightarrow \infty} \int_{-T/2}^{T/2} f(t) e^{-j\omega t} dt \\ &= \int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt \end{aligned}$$

This is the Fourier transform, which takes you from the time domain,  $f(t)$ , to the frequency domain,  $F(j\omega)$ .

# The Fourier Transform

---

The Fourier transform is:

$$F(j\omega) = \int_{-\infty}^{\infty} f(t)e^{-j\omega t} dt$$

## CYU (hw/exam problem difficulty)

---

Show that  $f(t) = \cos(\omega_0 t)$  is not an eigenfunction of an LTI system.