

HYDROGEN-LIKE SCHRÖDINGER OPERATORS AT LOW ENERGIES

ETHAN SUSSMAN

ABSTRACT. Consider a Schrödinger operator on an asymptotically Euclidean manifold X of dimension at least two, and suppose that the potential is of attractive Coulomb-like type. Using Vasy’s second 2nd-microlocal approach, “the Lagrangian approach,” we analyze – uniformly, all the way down to $E = 0$ – the output of the limiting resolvent $R(E \pm i0) = \lim_{\epsilon \rightarrow 0^+} R(E \pm i\epsilon)$. The Coulomb potential causes the output of the low-energy resolvent to possess oscillatory asymptotics which differ substantially from the sorts of asymptotics observed in the short-range case by Guillarmou, Hassell, Sikora, and (more recently) Hintz and Vasy. Specifically, the compound asymptotics at low energy and large spatial scales are more delicate, and the resolvent output is smooth all the way down to $E = 0$. In fact, we will construct a compactification of $(0, 1]_E \times X$ on which the resolvent output is given by a specified (and relatively complicated) function that oscillates as $r \rightarrow \infty$ times something polyhomogeneous. As a corollary, we get complete and compatible asymptotic expansions for solutions to the scattering problem as functions of both position and energy, with a transitional regime. So, in summary, we develop the low-energy scattering theory of attractive Coulombic potentials in the time-independent formalism, in the process studying delicate behavior at low energy and large scales.

CONTENTS

1. Introduction	1
2. The leC-calculus	20
3. The conjugated perspective	39
4. The situation at zero energy	49
5. Symbolic estimates	52
6. Proof of main theorem	78
Appendix A. The model ODE	100
Acknowledgements	106
References	106

1. INTRODUCTION

In this work we are concerned with scattering theory in the presence of Coulomb-like potentials. The long-range nature of the Coulomb potential modifies the large-scale asymptotics of the solutions of the relevant PDEs [Yaf82][Mel94; Mel95][Vas21a], the most well-studied of which being the time-independent Schrödinger equation, the time-dependent Schrödinger equation, and the wave equation. It has been known for a long while that the standard definitions of the S-matrix fail for the exact Coulomb potential, in both the time-independent and time-dependent settings — see e.g. [Dol64][RS79, §XI.9]. The culprit is the behavior at large- r .

The reason is easiest to explain in the time-independent (i.e. spectral side) setting, in which case the relevant PDE is the Schrödinger–Helmholtz equation on \mathbb{R}^n . First, we recall the short-range case. When the potential is short-range, the relevant ansatz for an incoming spherical wave is $e^{i\sigma r}$,

Date: December 24th, 2024 (Last Updated; typos fixed), April 18th, 2022 (Preprint).

2020 Mathematics Subject Classification. Primary 35P25; Secondary 58J40, 58J47, 58J50.

where r is the distance from the origin and $E = \sigma^2$ is the energy of the wave. That is, σ is the frequency. Scattering theory involves the study of solutions u of the PDE of the form

$$u = e^{i\sigma r} A + e^{-i\sigma r} B, \quad (1)$$

where A, B are some sufficiently nice functions. Specifically, we want A, B to be non-oscillatory. Roughly, they should be smooth functions of $1/(1+r^2)^{1/2}$ and the angular coordinates $\theta \in \mathbb{S}^{n-1}$, times some polynomial weight (which is not important in this introductory discussion). Thus, a function such as eq. (1) is the sum of an incoming spherical wave and an outgoing spherical wave. The “incoming profile” is the leading term $A_0(\theta)$ in the large- r asymptotic expansion of A . Usually, it is possible to prove that there exists a unique solution u of the form eq. (1) for specified A_0 . The all-important S-matrix is then the function $A_0 \mapsto B_0$ that sends the incoming profile to the outgoing profile B_0 , the leading term in the large- r asymptotic expansion of B_0 . Likewise, the *Poisson map* is the map $A_0 \mapsto u$ which takes the incoming profile and produces the unique solution to the PDE with that incoming data. This is how the story goes in the short-range case. One can find proofs of these results in the more general setting of scattering on asymptotically conic manifolds in [Mel94; Mel95], the perspective in which matches our own.

In the long range case, solutions no longer look like eq. (1). Instead, solutions look like

$$u = r^{iZ/2\sigma} e^{i\sigma r} A + r^{-iZ/2\sigma} e^{-i\sigma r} B, \quad (2)$$

for non-oscillatory A, B , where Z is the strength of the Coulomb potential. In atomic physics, where the equation is modeling an electron orbiting a positively charged nucleus, Z is the atomic number, which is traditionally denoted with the letter “ Z ”. So, $e^{i\sigma r}$ is no longer a sufficient ansatz for an incoming spherical wave — rather $r^{iZ/2\sigma} e^{i\sigma r}$ is. The outgoing oscillations are no longer at frequency σ . Rather, the phase receives a logarithmic correction $Z \log(r)/2\sigma$, known in the literature as the *Coulomb phase*. This is the reason why the traditional definitions of the S-matrix fail for Coulomb-like potentials: $r^{iZ/2\sigma} A$ is not non-oscillatory as $r \rightarrow \infty$. The story is similar in the time-dependent setting; the usual definitions of the S-matrix (or related objects like the wave/Møller operators) receive corrections from the Coulomb phase. The resulting objects are known as the Dollard-modified operators, after Dollard’s original work on the subject [Dol64]. With some abuse of terminology, we can refer to eq. (2) as the Dollard-modified ansatz as well.

So far, this has all been for positive energy/frequency, $\sigma > 0$. The behavior near $\sigma = 0$ is of interest in part because, in the passage from the time-independent picture to the time-dependent picture, it is necessary to take a Fourier transform in σ or $E = \sigma^2$, and $\sigma = 0$ is included in the domain of integration. For example, for the wave equation and the time-dependent Schrödinger equation with a short-range potential, the dominant contributions to the long-time behavior in spatially compact regions comes from low-energy behavior. Physically, this makes sense — high-energy wavepackets leave spatially compact regions rapidly, whereas low-energy wavepackets can linger and therefore contribute to the long-time asymptotics. So, if we want to understand the asymptotics in spacetime of solutions of these sorts of PDE, it is necessary to understand the low-energy limit of scattering theory.

Unfortunately, now we see a defect in the Dollard ansatz: the Coulomb phase diverges as $\sigma \rightarrow 0^+$. Thus, while the ansatz eq. (2) is sufficient for understanding the Schrödinger–Helmholtz equation at positive energy, if we want to understand the corresponding time-dependent PDE then something better is needed. Dollard managed to make do because the multiplier $r^{\pm iZ/2\sigma}$, while highly oscillatory as $\sigma \rightarrow 0^+$, is a bounded operator on L^2 . However, the extraction of long-time asymptotics is a much more delicate matter than the existence of a suitable S-matrix as an operator $L^2 \rightarrow L^2$. It is for this reason that we find it necessary to go beyond Dollard’s ansatz. Note that these difficulties do not arise in the short range case: $e^{i\sigma r}$ is a well-behaved ansatz all the way down to and including $\sigma = 0$. Interesting phenomena still occur in this regime — see e.g. [GH08; GH09; GHS13][Vas21b;

Vas21c][Hin22] for works in the same amount of generality as that pursued here — but these have to do with the asymptotics of the coefficients A, B in eq. (1), not the oscillatory prefactor.

One way to see that eq. (2) fails in the $\sigma \rightarrow 0^+$ limit is to consider exactly what happens exactly at $\sigma = 0$. At $\sigma = 0$, it turns out that the solutions of the PDE are still oscillatory, at least if $Z > 0$, which is the case we treat here. Specifically, the solutions look like

$$u = e^{2i\sqrt{Z}r} A + e^{2i\sqrt{Z}r} B \quad (3)$$

for non-oscillatory A, B — we will essentially prove a precise version of this in §4. This fact depends crucially on the presence of a Coulomb potential and on its sign. In the short-range case ($Z = 0$), the PDE is essentially Laplace’s equation, for which solutions look like polynomials, and in the $Z < 0$ case, a generic solution of the PDE is exponentially growing as $r \rightarrow \infty$ (though the physically relevant solutions are typically exponentially decaying), as can be seen by replacing $Z \rightarrow -Z$ in eq. (3). The \sqrt{r} -type oscillatory behavior seen in eq. (3) is specific to the case of attractive Coulombic potentials. In the Dollard ansatz eq. (2), there is no hint of this sort of oscillatory behavior. This is why Dollard’s ansatz must fail in the $\sigma \rightarrow 0^+$ limit. Somehow, the phases σr and $Z \log(r)/2\sigma$, one of which goes to 0 and the other of which blows up as $\sigma \rightarrow 0^+$, have to give way to the intermediate $2\sqrt{Z}r$.

The problem before us is therefore this: to come up with an alternative to the Dollard ansatz

$$r^{\pm iZ/2\sigma} e^{\pm i\sigma r} \quad (4)$$

that correctly describes the behavior in the $\sigma \rightarrow 0^+$ limit, and which reduces to eq. (3) exactly at $\sigma = 0$. The answer we present is¹

$$\exp \left[\pm i \left(r \sqrt{\sigma^2 + \frac{Z}{r}} + \frac{Z}{\sigma} \operatorname{arcsinh} \left(\sigma \sqrt{\frac{r}{Z}} \right) \right) \right]. \quad (5)$$

(The apparent singularity of the $\operatorname{arcsinh}$ term at $\sigma = 0$ is removable because $\operatorname{arcsinh}$ is odd. Thus, eq. (5) is better behaved than eq. (4) as $\sigma \rightarrow 0^+$.) It is an elementary exercise to show that, for each individual $\sigma > 0$, this reduces to the Dollard spherical wave modulo smooth factors which can be absorbed into A, B , and likewise that, for $\sigma = 0$, this reduces to the correct $\exp(\pm 2i\sqrt{Z}r)$, again modulo nice factors. For $\sigma > 0$, the $r\sqrt{\sigma^2 + Z/r}$ contributes the main term, while $\operatorname{arcsinh}$ term gives the Coulomb phase $Z \log(r)/2\sigma$. On the other hand, when $\sigma = 0$ exactly, both terms contribute equally, each giving one factor of $\sqrt{Z}r$. So, the ansatz eq. (5) successfully interpolates between the positive energy and low energy behavior while still approximately solving the relevant PDE, and it is the essentially unique (meaning modulo continuous factors) ansatz that does so. This is made precise by Corollary 1.3 below.

For the reader wondering where the complicated formula eq. (5) comes from, we explain in the next subsection how the low-energy problem has a semiclassical structure at large-radii; the ansatz above is precisely the Liouville–Green (a.k.a. WKB or JWKB) approximation that results. The essential uniqueness of eq. (5) comes from the essential uniqueness of the Liouville–Green approximation. Usually, it is the high-energy regime in which one encounters semiclassical phenomena. The fact that the low-energy, large-radius hydrogen atom has a semiclassical structure is not as well-known as it perhaps should be.

Our goal in this manuscript is to develop the low-energy scattering theory of attractive Coulomb-like potentials along these lines, taking eq. (5) as the ansatz describing an incoming spherical wave. Our focus is on describing the behavior of the limiting resolvent $R(\sigma^2 \pm i0)$. The reason for this focus is that the S-matrix and Poisson operators can both be constructed in terms of the two limiting resolvents, as can the spectral projector via one of Stone’s theorems. Thus, the limiting resolvents

¹Actually, this is only correct for potential scattering. When the metric is variable, the ansatz receives another correction coming from one particular subleading term in the metric. We call this term “ a_{00} ” below.

can be taken as the fundamental objects in time-independent scattering theory. This is why our main theorem, Theorem 1.2, only talks about the resolvents. However, the reader more interested in scattering theory than spectral theory might prefer to think of Corollary 1.3 as the main result.

We say nothing about the time-dependent formalism except for a few brief sketches. While time-independent scattering theory is important in its own right, the most convincing motivation comes from understanding the corresponding PDE in spacetime. Passing to the time-dependent formalism is, in some sense, “just” a matter of taking the Fourier transform. However:

- whether one takes the Fourier transform in σ or E depends on whether one wants to study the wave/Klein–Gordon equation or the time-dependent Schrödinger equation, and different asymptotic forms result,
- in order to control the Fourier transform, one needs also the high energy, $\sigma \rightarrow \infty$, asymptotics; this is actually already understood, but it is orthogonal to the low-energy phenomena that we are focused on,
- producing full asymptotic expansions on the spacetime side can involve a nontrivial amount of work, even once full asymptotic expansions are known on the spectral side. This work should not be difficult — it boils down to studying the asymptotics of some simple oscillatory integrals — but it also should not be trivial.

Thus, we have chosen to only deal with time-independent scattering theory here. The time-dependent equations are well worth studying, of course, but it is our view that, once the spectral side is understood, understanding the time-dependent side is more or less straightforward. In the short-range case, this is well established, and the long-range case should not be different in this regard. We intend to tackle the time-dependent Schrödinger equation for ionized hydrogen in future work.

This manuscript has been written in the language of the Melrose school of geometric microlocal analysis. This language is an incisive tool for studying multifaceted asymptotic expansions like those we encounter below. It provides a precise notion of what it means to have complete understanding of the asymptotics. Unfortunately, this language can also make for difficult reading. Let us therefore point out that, in §A, we consider the case of an exact Coulomb potential. Then, the solutions of the differential equation are special functions, namely Whittaker functions. In that appendix, we apply our main results to make concrete statements about the asymptotics of the Whittaker functions in a poorly understood asymptotic regime. There, we reproduce and refine classical results from the literature of special functions. Hopefully, that appendix provides an accessible introduction to the problem studied here. The results are illustrated with numerics, unlike elsewhere in this paper.

The next subsection is a technical introduction to the paper, for general asymptotically conic manifolds, this being the natural amount of generality in which our results hold. One of its purposes is to record the main results of the paper. We will repeat there some of the discussion here, but with much more detail. Our main result is Theorem 1.2, which expresses, in precise terms, the fact that the resolvent output admits full asymptotic expansions in terms of elementary functions on a particular manifold-with-corners compactifying the low-energy regime — see Figure 1. To make this more concrete, we present in Corollary 1.4, Corollary 1.5, Corollary 1.6 the content of this result vis-a-vis explicit asymptotic expansions. Roughly, our result that the resolvent output $R(\sigma^2 \pm i0)f$ is, for Schwartz f , of the form

$$\exp \left[\pm i \left(r \sqrt{\sigma^2 + Zr} + \frac{Z}{\sigma} \operatorname{arcsinh} \left(\sigma \sqrt{\frac{r}{Z}} \right) \right) \right] A \quad (6)$$

for A which is almost a smooth function of σ^2 , $1/r$, and the angular variables θ . There are two reasons why the “almost” in the previous sentence is necessary:

- The functions are not smooth, but instead polyhomogeneous. Polyhomogeneity is the natural generalization of smoothness to allow any generalized Taylor series, which may include

logarithmic terms. For example, a polyhomogeneous function on $[0, 1)_x$ is a function which admits well-behaved asymptotic expansions in functions of the form $x^j(\log x)^k$.

- The A in eq. (6) is not polyhomogeneous in σ and $1/r$. Rather, it is polyhomogeneous on the manifold-with-corners in Figure 1. Concretely, this will mean polyhomogeneity in σ and $\hat{x} = 1/(\sigma^2 r)$ as $\hat{x} \rightarrow 0^+$ and polyhomogeneity in $x = 1/r$ and $\hat{E} = \sigma^2 r$ as $\hat{E} \rightarrow 0^+$, as well as in the angular variables θ .

For the reader not familiar with the use of blowups for this sort of asymptotic analysis, it is worth keeping in mind the simplest example: on the two-dimensional quadrant $[0, \infty)^2 \subset \mathbb{R}_{x,y}^2$, the polar angle $\theta = \arctan(y/x)$ is not polyhomogeneous. Rather, it is polyhomogeneous on the blowup

$$[[0, \infty)^2; \text{origin}] \cong [0, \pi/2]_\theta \times [0, \infty)_r. \quad (7)$$

Thus, in this example, a blowup just means using polar coordinates, a coordinate system which is singular with respect to the smooth structure on the original manifold $[0, \infty)^2$.

For another example, consider $x/(x+y)$. When $y = 0$, this is identically 1. When $x = 0$, this is identically 0. The limits $\lim_{x \rightarrow 0^+}, \lim_{y \rightarrow 0^+}$ do not commute:

$$\lim_{x \rightarrow 0^+} \lim_{y \rightarrow 0^+} \frac{x}{x+y} \neq \lim_{y \rightarrow 0^+} \lim_{x \rightarrow 0^+} \frac{x}{x+y}. \quad (8)$$

This is exactly the sort of behavior that polyhomogeneity excludes. The function $x/(x+y)$ is not polyhomogeneous on the quadrant $[0, \infty)^2$, but rather on the blowup $[[0, \infty)^2; \text{origin}]$. Indeed,

$$\frac{x}{x+y} = \frac{\cos(\theta)}{\cos(\theta) + \sin(\theta)} \quad (9)$$

is a perfectly smooth function of $\theta \in [0, \pi/2]$. We will make use of more complicated manifolds-with-corners below, but the one $X_{\text{res}}^{\text{sp}}$ we use in Theorem 1.2 to describe the resolvents is just $[[0, \infty)^2; \text{origin}]$ times \mathbb{S}^{n-1} , where the Cartesian coordinates x, y are to be replaced by $\sigma^2, 1/r$. (The definition of $X_{\text{res}}^{\text{sp}}$ involves an additional change of smooth structure, but we can ignore that for now.)

Let us emphasize that we do not allow our potential to be singular at $r = 0$ (except in the appendix). Thus, the exact Coulomb potential $1/r$ is excluded. This is for reasons of simplicity. It is not expected to be difficult to extend the results below to handle such singularities. The low-energy asymptotics near them is straightforward since, as far as the $r \rightarrow 0$ limit is concerned, the differential equation does not degenerate at low-energy. Consequently, unlike in the $r \rightarrow \infty$ case where the $\sigma \rightarrow 0^+$ limit requires a transitional asymptotic regime, no such complexities are present near a singularity. Unfortunately, developing this requires a detailed discussion of regular singular PDE beyond the elementary results required here. The added burden was felt to be not worth it. The ODE case is simple enough, so the discussion in §A includes a singularity at $r = 0$.

1.1. Technical introduction. In [Mel94], Melrose introduced the programme of understanding the limiting absorption principle on asymptotically conic manifolds, denoted X below, from the microlocal point of view, initially for the Laplacian and then for Schrödinger operators more generally [Mel94, §16]. The case of fixed energy $E > 0$ was dealt with first by Melrose [Mel94], then by Hassell & Vasy [HV99] and – using a more modern approach – Vasy [Vas21a] again, while uniform estimates in the high energy (a.k.a. “semiclassical,” $E \rightarrow \infty$) limit have been established by Vasy & Zworski [VZ00] and Vasy [Vas21a, §5]. More recently, the low energy $E \rightarrow 0^+$ behavior has been understood to a highly satisfactory degree in the works of Guillarmou, Hassell, and Sikora [GHS13] and Vasy [Vas21b][Vas21c] — see also Hintz [Hin22, §2]. (Complementarily, Guillarmou and Hassell [GH08][GH09] consider the $E \rightarrow 0^-$ limit of the resolvent kernel.) While the previous results applied to all Schrödinger operators, sometimes under standard dynamical assumptions, the low energy results proven so far in this level of generality apply only to Schrödinger operators without Coulomb-like (a.k.a. long range) terms. Indeed, a Coulomb-like potential has a serious

effect on the asymptotics of formal Schrödinger eigenfunctions in the low energy limit — this is true in the context of Euclidean potential scattering, and the general case inherits that complexity.

In this paper, we study the case of an *attractive* Coulomb-like potential using the framework of [Vas21a][Vas21c], Vasy’s 2nd second microlocal approach (the “Lagrangian approach”). This is in contrast to his 1st second microlocal approach to the low energy limit, pursued in [Vas21b], which utilized variable order Sobolev spaces (as previously used in e.g. [Vas18, Proposition 5.28] for the $E > 0$ case). The use of the term “Lagrangian” belies the fact that the framework of Lagrangian (more properly *Legendrian*) distributions does not appear explicitly in this approach — in fact this is precisely the point: instead of “module regularity [HMV04, §6][GR+20] for a Lagrangian submanifold” — which implies extra regularity *outside* of that submanifold via some elliptic estimates — we need only consider b-Sobolev regularity. We use a conjugation to move what would otherwise have been the energy-dependent Legendrian submanifold for which module regularity (in this case essentially the Sommerfeld radiation condition) is established to the zero section of the scattering cotangent bundle (which can be blown up to the fibers of the b-cotangent bundle over the boundary). The upshot is that the Sommerfeld radiation condition, in one of its forms, is replaced by a condition regarding b-Sobolev regularity. This is convenient for getting estimates which are uniform in E because, unlike the relevant notion of module regularity prior to conjugation (see e.g. [HMV04][GR+20]), in which the relevant test module depends on E , the relevant notion of b-regularity (as stated in [Vas21a, Theorem 1.1]) does not depend so much on E . (The sharpest form of [Vas21a, Theorem 1.1] is phrased using sc,b-Sobolev regularity, which is a form of module regularity — the key point here is that the relevant test module is E -independent.)

We will apply similar considerations to the study of the $E \rightarrow 0^+$ limit of the limiting resolvents “ $P((E \pm i0)^{1/2})^{-1} = (P(0) - E \mp i0)^{-1}$ ” of the Schrödinger operator

$$P(0) = \Delta_g - Zx + V, \tag{10}$$

where V is short-range and $Z > 0$, so the total potential $W = -Zx + V$ is of attractive Coulomb-like type. Here $x \in C^\infty(X; \mathbb{R}^{\geq 0})$ is a boundary defining function (bdf), e.g. $1/\langle r \rangle$ when X is asymptotically Euclidean, r denoting the Euclidean radial coordinate. (Note: we follow the convention of parametrizing the spectral family $\{P(0) - E\}_{E \geq 0}$ of $P(0)$ in terms of $\sigma = E^{1/2}$, so we write $P(\sigma) = P(0) - \sigma^2$.) Singular versions of the operator eq. (10) first appeared in Schrödinger’s model of atomic hydrogen — or more generally hydrogenic ions — hence the operators we consider could also be called “hydrogen-like.” In eq. (10), $\Delta_g \geq 0$ is the positive semidefinite Laplacian associated with an asymptotically conic metric. The conjugated perspective complicates the family of operators under consideration (see §3) — more so than in the case $Z = 0$, when the total potential is short-range — but it greatly facilitates the derivation of low energy asymptotics. See below for a heuristic discussion (and §6 for details). It should be noted that although we work with general asymptotically conic manifolds, our results are new even on exact Euclidean space. The existing literature on low energy asymptotics in the presence of a Coulomb-like term is quite sparse — the only previous treatments the author is aware of are [Yaf82][Nak94][FS04][DS09b][DS09a][Bou11][Ski13]. In comparison to these earlier works, we require more of our potential, but the payoff is a complete understanding of asymptotics at spatial infinity.

We comment briefly on our focus on the case $E > 0$ and $Z > 0$. Considering the spectral family of a not-necessarily-attractive Coulomb-like Schrödinger operator at energy $E \in \mathbb{R} \setminus \{0\}$ and with attractivity strength $Z \in \mathbb{R} \setminus \{0\}$ (the atomic number for the case of an electron orbiting an atomic nucleus, using appropriately natural units): out of the four cases (I) $E > 0, Z > 0$ (attractive, scattering near zero energy), (II) $E < 0, Z > 0$ (attractive, ellipticity near zero energy), (III) $E > 0, Z < 0$ (repulsive, scattering near zero energy), (IV) $E < 0, Z < 0$ (repulsive, ellipticity near zero energy), the first and fourth are the most tractable, as evidenced by the state of the literature on similar problems in exact Euclidean space — see [Yaf82][Nak94] for work on cases (I) and (III) and [FS04][Ski13] for case (I). In the more difficult case (II), for example, one can encounter an

infinite sequence of bound states, as in the hydrogen atom — see [Ivr19, Theorem 11.6.7]. Case (III), which has been studied partially in [Yaf82], is expected to be intermediate between (II) and (I) in terms of difficulty. While our focus is on (I), the pseudodifferential technology developed yields an easy treatment of case (IV).

The jumping off point for our analysis is the proof of a symbolic estimate, Theorem 5.12, of $u \in \mathcal{S}'(X)$ in terms of $f = P(E^{1/2})u$ that is uniform down to $E = 0$. The estimate is structurally similar to the combined radial point and propagation estimates proven for $E > 0$ in [Mel94, §8]. We will formulate the estimate in a “second microlocal” framework akin to that in [Vas21a], as this dovetails with the conjugated perspective, but a similar estimate can be articulated using function spaces analogous to the somewhat more standard *variable order* sc-Sobolev spaces [Vas13][Vas18][GR+20]. To illustrate the idea, we rewrite the operator $P(\sigma)$ in terms of the coordinate $\hat{x} = x/\sigma^2$, which is appropriate for homogenizing the spectral parameter and Coulomb potential. To the relevant order, $\sigma^{-2}P(\sigma)$ is given in terms of $\hat{x}, \sigma > 0$, the latter of which we suggestively rename ‘ h ,’ by

$$\hat{P}(h) = -h^2(\hat{x}^2\partial_{\hat{x}})^2 + h^2\hat{x}^2\Delta_{\partial X} - 1 - Z\hat{x}, \quad (11)$$

which we consider as a 1-parameter family of operators on the exact cone $[0, \infty)_{\hat{x}} \times \partial X$. Note that $\sigma = h$ appears on the right-hand side of eq. (11) as an effective *semiclassical* parameter (somewhat surprisingly, since semiclassical problems typically arise in the study of the high energy regime rather than the low energy regime of interest here), and so \hat{P} can be studied as a semiclassical operator on an exact cone, an approach that seems to have first been undertaken by Nakamura [Nak94]. This conic problem differs from the conic problem arising in [Wan06][GH08][GHS13][Vas21c], which has no semiclassical parameter. The qualitative features of the semiclassical family defined by eq. (11) are as follows:

- At the “large” end of the cone, $\{\hat{x} = 0\}$, $\hat{P}(h)$ is of real principal type for each $h > 0$. This holds regardless of the sign of Z and reflects the dynamical structure of $P(\sigma)$ for $\sigma > 0$.
- At $h = 0$, one has real principal type propagation from $\hat{x} = 0$ to $\hat{x} = \infty$ (at finite frequencies).

The second of these is of secondary importance below.

The situation changes completely if $Z < 0$ or if $E < 0$. Then, eq. (11) is to be replaced by

$$\hat{P}(h) = -h^2(\hat{x}^2\partial_{\hat{x}})^2 + h^2\hat{x}^2\Delta_{\partial X} \pm 1 \pm Z\hat{x}, \quad (12)$$

where the first sign is that of $-E$ and the second sign is that of $-Z$, each possible pair of signs corresponding to one of the four cases (I), (II), (III), (IV) above. If $E < 0$, then eq. (12) is elliptic at the large end of the cone, and if $Z < 0$ then eq. (12) is semiclassically elliptic near the small end of the cone. Even for the case when X is one-dimensional, one subtlety of the cases (II) $E < 0, Z > 0$ and (III) $E > 0, Z < 0$ is understanding precisely the transition from ellipticity to nonellipticity that occurs at $h = 0$ and $\hat{x} = 1/|Z|$. This subtlety arises, for example, in physicists’ treatment of the WKB approximation, where Airy functions are used to patch quasimodes in the classically allowed and classically forbidden regions — see [SN17] for a standard treatment at a physicist’s level of rigor. In case (II), the full partial differential operator eq. (12) has an additional subtlety: its Hamiltonian flow has closed loops corresponding to classical Keplerian orbits. See [Ivr19, §11.6] for a discussion of the effects of this on eigenvalue counting.

Zooming back from the boundary, we will analyze $P(\sigma)$ on a mwc (“mwc” standing for “manifold-with-corners,” in the sense of Melrose’s school [Mel92]) which we will denote $X_{\text{res}}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \partial X; \frac{1}{2}]$, depicted in Figure 1. This is the result of modifying the smooth structure of

$$X_{\text{res},0}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \partial X] \quad (13)$$

at the front face of the blow-up so that $(E + x)^{1/2}$ becomes a bdf. One can analyze $P(\sigma)$ by quantizing the Lie algebra $\mathcal{V}_{\text{sc,leC}}(X)$ of smooth vector fields on $X_{\text{res}}^{\text{sp}}$ that

- (I) are tangent to the level sets of E ,

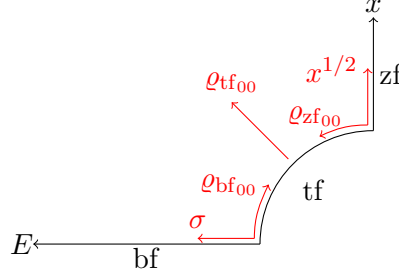


FIGURE 1. The mwc $X_{\text{res}}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \partial X; 1/2]$, with bdfs $\varrho_{\text{bf}00} = Zx/(\sigma^2 + Zx)$, $\varrho_{\text{tf}00} = (\sigma^2 + Zx)^{1/2}$, and $\varrho_{\text{zf}00} = \sigma^2/(\sigma^2 + Zx)$ in terms of $E^{1/2} = \sigma$. This is $X_{\text{res},0}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \partial X]$ with the smooth structure at the front face of the blow-up modified. The interior of the mwc is to the upper-left of the drawn boundary. (Degrees of freedom associated with ∂X omitted from the diagram.)

(II) lie in $\varrho_{\text{bf}00}\varrho_{\text{tf}00}C^\infty(X_{\text{res}}^{\text{sp}}) \otimes \mathcal{V}_b(X) \subset \mathcal{V}_E(X_{\text{res}}^{\text{sp}})$,

where $\varrho_{\text{bf}00}, \varrho_{\text{tf}00}$ are as in Figure 1 and $\mathcal{V}_E(X_{\text{res}}^{\text{sp}})$ is the Lie algebra of smooth vector fields on $X_{\text{res}}^{\text{sp}}$. The quotient algebra $\mathcal{V}_{\text{sc}, \text{leC}}(X)/\varrho_{\text{bf}00}\varrho_{\text{tf}00}\mathcal{V}_{\text{sc}, \text{leC}}(X)$ is commutative, so the corresponding ΨDO calculus (which is closely related to the leC-calculus discussed below) is under symbolic control. This allows us to prove a half-Fredholm estimate involving variable order “sc,leC”-Sobolev spaces that is uniform down to zero energy. The second-microlocal estimate Theorem 5.12 is a sharper version of this.

Although it is not our most general result, as the main propositions of §5, §6 require less smoothness of the coefficients of the PDE, we will prove the following main theorem. We state it using some standard or semistandard terminology, which, if not standard, is recalled in §1.2 below. The special case of Euclidean potential scattering off of an attractive asymptotically Coulomb-like potential is partially stated in Corollary 1.4, Corollary 1.5, and Corollary 1.6.

Theorem 1.1 (Main theorem, rough version). *The limiting resolvent output is of exponential-polyhomogeneous type on $X_{\text{res}}^{\text{sp}}$.* ■

More precisely:

Theorem 1.2 (Main theorem). *Suppose that (X, ι, g_0) is an exactly conic manifold of dimension $\dim X = n \geq 2$, and let $x \in C^\infty(X)$, $x : X \rightarrow [0, \infty)$, denote a compatible boundary-defining-function (bdf), so that, near ∂X ,*

$$g_0 = \frac{dx^2}{x^4} + \frac{g_{\partial X}}{x^2} \quad (14)$$

for some Riemannian metric $g_{\partial X}$ on ∂X . Let g denote a fully classical asymptotically conic metric on X , which in this context means a Riemannian metric on $X^\circ = X \setminus \partial X$ which near ∂X has the form

$$g = g_0 + a_{00} \frac{dx^2}{x^3} + \frac{\Gamma_{1, \partial X} \odot dx}{x^2} + \frac{h_{1, \partial X}}{x} + x^2 C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \quad (15)$$

for some $a_{00} \in \mathbb{R}$, $\Gamma_{1, \partial X} \in \Omega^1(\partial X)$, and $h_{1, \partial X} \in C^\infty(\partial X; \text{Sym}^2 T^* \partial X)$. Given $Z > 0$ and $V \in x^2 C^\infty(X; \mathbb{R})$, consider the Schrödinger operator

$$P(0) = \Delta_g - Zx + V : \mathcal{S}'(X) \rightarrow \mathcal{S}'(X), \quad (16)$$

where Δ_g is the (positive semidefinite) Laplacian. For each $E > 0$, let

$$R(E \pm i0) = R(E \pm i0; Z) : \mathcal{S}(X) \rightarrow \mathcal{S}'(X) \quad (17)$$

denote the limiting resolvent from above or below the spectrum — cf. Melrose [Mel94, §14]. (That is, for any $f \in \mathcal{S}(X)$, $u_{\pm} = R(E \pm i0)f$ is the unique solution to $P(0)u = Eu + f$ satisfying the weak Sommerfeld radiation condition [Mel94, §11].) Set

$$\Phi(x; \sigma) = \frac{1}{x} \sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} + \frac{1}{\sigma} (Z - \sigma^2 a_{00}) \operatorname{arcsinh} \left(\frac{\sigma}{x^{1/2}} \frac{1}{(Z - \sigma^2 a_{00})^{1/2}} \right) \quad (18)$$

for all $\sigma > 0$ such that $\sigma^2 a_{00} < Z$.

Then, for any Schwartz function $f \in \mathcal{S}(X)$, the function $u_{0,\pm}$ on $X_{\text{res}}^{\text{sp}} \cap \{Z > Ea_{00}\}$ defined by

$$u_{\pm} = e^{\pm i\Phi(x; E^{1/2})} x^{(n-1)/2} (E + Zx)^{-1/4} u_{0,\pm} \quad (19)$$

(for E such that $Z > Ea_{00}$) is polyhomogeneous on $X_{\text{res}}^{\text{sp}} \cap \{Z > Ea_{00}\}$, conormal of order zero, and smooth at $\text{zf}^{\circ} \cup \text{bf}^{\circ}$.

Moreover, $R(E = 0; Z \pm i0) : \mathcal{S}(X) \rightarrow \mathcal{S}'(X)$ is well-defined (e.g. as a strong limit of $R(E = 0; Z \pm i\epsilon)$ as $\epsilon \rightarrow 0^+$), and we can write $u_{\pm}(-; 0) = R(E = 0; Z \pm i0)f$ as

$$u_{\pm}(-; 0) = e^{\pm i\Phi(x; 0)} x^{(n-1)/2} (Zx)^{-1/4} u_{0,\pm}(-; 0) \quad (20)$$

($\Phi(-; 0)$ being defined by removing the removable singularity of eq. (18) at $\sigma = 0$), where

$$u_{0,\pm}(-; 0) \in C^{\infty}(X_{1/2}) \quad (21)$$

is the restriction of $u_{0,\pm}$ to $\text{zf} = \text{cl}\{\sigma = 0, x > 0\} \subset X_{\text{res}}^{\text{sp}}$. Thus, defining $u_{\pm}(-; E^{1/2})$ either as $R(E \pm i0)f$ for $E > 0$ or $R(E = 0; Z \pm i0)f$ for $E = 0$, the formula eq. (19) holds for all $E \geq 0$ sufficiently small. ■

For $z > 0$, by $\operatorname{arcsinh}(z)$ we mean $\log(z + (1 + z^2)^{1/2})$.

Remark 1. Theorem 1.2 applies equally well for σ -dependent forcing $f \in C^{\infty}([0, \infty)_{\sigma^2}; \mathcal{S}(X))$ as can be proven using the argument in §6. ■

Remark 2. In terms of the notation for spaces of polyhomogeneous functions used below, $u_{0,\pm} \in \mathcal{A}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}})$ for some index set² $\mathcal{E} \subset \{(k, \kappa) \in \mathbb{C} \times \mathbb{N} : \Re z \geq 0\}$ at tf. In fact, we can take

$$\mathcal{E} = \{(k, \kappa) \in \mathbb{N} \times \mathbb{N} : \kappa \leq \lfloor k/2 \rfloor\}. \quad (22)$$

See §6.3.

One can be even more precise than this: as the proof of Proposition 6.18 shows, the terms in the polyhomogeneous expansion of $u_{0,\pm}$ at tf having many logs are proportional to many powers of E . Specifically, the coefficient $a \in C^{\infty}(\text{tf})$ of $\varrho_{\text{tf}00}^k \log^{\kappa}(\varrho_{\text{tf}00})$ for $(k, \kappa) \in \mathcal{E}$ is actually in $\varrho_{\text{zf}00}^{\kappa} C^{\infty}(\text{tf})$. In order for $(k, \kappa) \in \mathcal{E}$ to hold, $2\kappa \leq k$, so any logarithm $\log^{\kappa}(\varrho_{\text{tf}00})$ in the asymptotic expansion at tf is suppressed by a factor of E^{κ} . This is why $u_{0,\pm}(-; 0)$ can be smooth on $X_{1/2} = \text{zf}$ even though $u_{0,\pm}(-; -)$ is not necessarily (claimed to be) smooth at $\text{zf} \cap \text{bf}$. The statement of the theorem above is therefore slightly nonoptimal, and we will not introduce the terminology needed to state an optimal version. ■

In a standard manner – cf. [Mel94, §15] – we can deduce from the conjunction of (1) an asymptotic summation argument, (2) Theorem 1.2 (strengthened slightly by the first remark), and (3) the other results in the body of the paper below the following:

Corollary 1.3. *Consider the setup of Theorem 1.2, and fix $E_0 > 0$ such that $E_0 a_{00} < Z$. If $\alpha \in C^{\infty}(\partial X)$, then there exist*

$$A, B \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) \quad (23)$$

²We will only consider index sets with the property that $(k, \kappa) \in \mathcal{E} \Rightarrow (k+1, \kappa) \in \mathcal{E}$.

such that $A|_{\text{tf} \cup \text{bf}} = \alpha \circ \pi$ (where $\pi : \text{tf} \cup \text{bf} \rightarrow \partial X$ denotes the restriction to $\text{tf} \cup \text{bf}$ of the composition of the blowdown map $X_{\text{res}}^{\text{sp}} \rightarrow [0, \infty)_\sigma \times X$ and the projection $[0, \infty)_\sigma \times X \rightarrow X$) and

$$u = e^{-i\Phi(x; E^{1/2})} x^{(n-1)/2} (E + Zx)^{-1/4} A + e^{+i\Phi(x; E^{1/2})} (E + Zx)^{-1/4} x^{(n-1)/2} B \quad (24)$$

solves the Schrödinger–Helmholtz equation $\Delta_g u - Zxu + Vu = Eu$ for all $E \in [0, E_0]$. Moreover, u is the unique solution in $\{E \leq E_0\}$ with this property. $\blacksquare \square$

This shows that $\exp(\pm i\Phi(x; E^{1/2})) x^{(n-1)/2} (E + Zx)^{-1/4}$ serves as a notion of “incoming/outgoing spherical wave” in the presence of an attractive Coulomb-like potential that makes sense all the way down to $E = 0$. In Corollary 1.3, α serves as a notion of “incoming data.” The proof actually constructs the asymptotic expansion of A at $\text{tf} \cup \text{bf}$. A natural question, which we do not investigate here, is whether or not $B|_{\text{tf} \cup \text{bf}} = \beta \circ \text{bdn}$ for some $\beta \in C^\infty([0, \infty)_\sigma \times \partial X)$, where $\text{bdn} : X_{\text{res}}^{\text{sp}} \rightarrow [0, \infty)_\sigma \times X$ is the blowdown map. Physically, Corollary 1.3 describes the scattering of nonrelativistic electrons off of a hydrogenic nucleus, or alternatively nonrelativistic Bhaba scattering.

Compare the following corollary of Theorem 1.2 (strengthened slightly by Remark 2) with [Nak94, Cor. 1.5][FS04, Eq. 1.2]:

Corollary 1.4 (Asymptotics at zf). *For each $l \in \mathbb{R}$, let $S^l(\mathbb{R}^n) = S_{1,0}^l(\mathbb{R}^n)$ denote the space of symbols of order l . Fix $n \geq 2$. Suppose that $W \in C^\infty(\mathbb{R}^n)$ is a classical symbol of order -1 , so that there exist (unique) $W_0, W_1, W_2, W_3, \dots \in C^\infty(\mathbb{S}^{n-1})$ such that*

$$W(\mathbf{x}) - (1 - \chi(r)) \sum_{k=0}^K \frac{W_k(\mathbf{x}/r)}{r^{k+1}} \in S^{-K-2}(\mathbb{R}^n), \quad (25)$$

for each $K \in \mathbb{N}$, where $r = \|\mathbf{x}\|$ and $\chi \in C_c^\infty(\mathbb{R})$ is identically equal to one in some neighborhood of the origin. Suppose further that W is attractive and Coulomb-like, meaning that $W_0 = -Z$ for some constant $Z > 0$. Consider the Schrödinger–Helmholtz operator

$$P(\sigma) = \Delta - \sigma^2 + W \quad (26)$$

for $\sigma \geq 0$, where $\Delta = -(\partial_{x_1}^2 + \dots + \partial_{x_n}^2)$ is the positive semidefinite Laplacian. Given $f \in \mathcal{S}(\mathbb{R}^n)$ (where $\mathcal{S}(\mathbb{R}^n)$ denotes the set of Schwartz functions) let

$$u_\pm(\mathbf{x}; \sigma) = R(\sigma^2 \pm i0) f(\mathbf{x}) \quad (27)$$

denote the output of the limiting resolvent $R(\sigma^2 \pm i0) = \lim_{\epsilon \rightarrow 0^+} R(\sigma^2 \pm i\epsilon)$ applied to f for $\sigma > 0$.

Then, there exist functions $w_{\pm,0}, w_{\pm,1}, w_{\pm,2}, \dots \in C^\infty(\mathbb{R}^n)$ such that, for each $K \in \mathbb{N}$ and $\mathbf{x} \in \mathbb{R}^n$,

$$u_\pm(\mathbf{x}; \sigma) = \sum_{k=0}^K w_{\pm,k}(\mathbf{x}) \sigma^{2k} + O_{\mathbf{x},f,K}(\sigma^{2K+2}) \quad (28)$$

as $\sigma \rightarrow 0^+$, where the $O_{\mathbf{x},f,K}(\sigma^{2K+2})$ term is uniformly bounded (i.e. by $C_{\mathcal{K},f,K} \sigma^{2K+2}$) in compact subsets $\mathcal{K} \subseteq \mathbb{R}^n$ worth of $\mathbf{x} \in \mathbb{R}^n$.

In fact, we can take $w_{\pm,k} \in \langle r \rangle^k \mathcal{A}^{\mathcal{E}_k}(\text{zf})$, and then the error term in eq. (28) is an $O_{f,K}(\sigma^{2K+2})$ family of elements of $\langle r \rangle^{K+1} \mathcal{A}^{\mathcal{E}_{K+1}}(\text{zf})$, where \mathcal{E}_k is the index set

$$\mathcal{E}_k = \{(\kappa, \varkappa) \in \mathbb{N} \times \mathbb{N} : \varkappa \leq \min\{k, \lfloor \kappa/2 \rfloor\}\}. \quad (29)$$

Moreover, $w_{\pm,0}$ solves the PDE $P(0)w_{\pm,0} = f$. $\blacksquare \square$

(Of course, the $O_{\mathbf{x},f,K}(\sigma^{2K+2})$ error also depends on W , though we do not write this dependence explicitly. The same applies to the other errors in Corollary 1.5, Corollary 1.6 below.) Corollary 1.4 applies in particular to any $W \in C^\infty(\mathbb{R}^n)$ which is equal, outside of some compact set, to

$$-Z/r + \sum_{j=2}^J r^{-j} W_j(\mathbf{x}/r) + W_\infty, \quad (30)$$

for $W_j \in C^\infty(\mathbb{S}^{n-1})$, $J \in \mathbb{N}$, and $W_\infty \in \mathcal{S}(\mathbb{R}^n)$. Thus, while we impose significant restrictions on the radial behavior of the potential (and the potential's regularity) in order to get full asymptotic expansions, there are no symmetry requirements on W_1, W_2, \dots . We do, however, require that $W_0 = -Z$ is constant, so the Coulomb-like part of W is required to be spherically symmetric. We remark that Proposition 6.3 (see also Remark 5) allows us to study more general symbolic W , but then instead of full asymptotic expansions we get only partial asymptotic expansions together with symbolic estimates for the remainders. We also remark that the classical symbols of order -2 on \mathbb{R}^n are precisely those of the form $V = \langle r \rangle^{-2} U$ for U a smooth function on the radial compactification $\mathbb{B}^n = \overline{\mathbb{R}^n}$ of \mathbb{R}^n , this being diffeomorphic to a closed ball, so Corollary 1.4 applies to many potentials that are not of the form eq. (30), e.g. $W = -\langle r \rangle^{-1}$.

In a sense made precise by Theorem 1.2, the asymptotic expansion eq. (28) can be refined into an asymptotic expansion in powers of $\hat{E} = \sigma^2/(\sigma^2 + 1/\langle r \rangle)$ (with a complicated oscillatory prefactor) whose error terms are uniformly bounded without a loss of decay as K increases. The non-uniform expansion eq. (28) already stands in stark contrast with the situation when W is short range, where instead one only has e.g. in the case $n = 3$ (according to [Hin22, Theorem 3.1])

$$u_\pm(\mathbf{x}; \sigma) = w_{\pm,0}(\mathbf{x}) + w_{\pm,1}(\mathbf{x})\sigma + O_{\mathbf{x},f}(\sigma^2 \log \sigma). \quad (31)$$

This is sharp — the singular $\sigma^2 \log \sigma$ term is the source of the main term in Price's law.

Given the setup of Corollary 1.4, the known large- r asymptotic expansion of $u_\pm(\mathbf{x}; \sigma) = R(\sigma^2 \pm i0)f$ is [Mel94]:

- there exist functions $t_{\pm,0}, t_{\pm,1}, t_{\pm,2}, \dots \in C^\infty(\mathbb{S}^{n-1} \times \mathbb{R}^+)$ such that, for each $\sigma > 0$ and nonzero $\mathbf{x} \in \mathbb{R}^n$,

$$u_\pm(\mathbf{x}; \sigma) = r^{-(n-1)/2} e^{\pm i\sigma r} r^{\pm iZ/2\sigma} \left[\sum_{k=0}^K t_{\pm,k}(\mathbf{x}/r, \sigma) r^{-k} + O_{\sigma,f,K}(r^{-(K+1)}) \right] \quad (32)$$

for each $K \in \mathbb{N}$, where the $O_{\sigma,f,K}(r^{-(K+1)})$ term is uniformly bounded in compact subsets $\mathcal{K} \subseteq \mathbb{R}^+$ worth of σ .

See §A for the Whittaker case, where the $t_\pm(\sigma)$ are written down explicitly. Note that the existence of the expansion eq. (32) is contained in Theorem 1.2 (from the asymptotic expansion at $\text{bf}^\circ \subset \mathbb{B}_{\text{res}}^{\text{sp}}$). Note that the $r^{\pm iZ/2\sigma}$ term in eq. (32) is singular as $\sigma \rightarrow 0^+$, which renders eq. (32) unsuitable to study the low energy limit. The situation is ameliorated in the following way — according to Theorem 1.2, eq. (32) admits a repackaging that is *uniform* down to $\sigma = 0$:

Corollary 1.5 (Asymptotics at bf). *Consider the setup of Corollary 1.4. Let $\varrho = (\sigma^2 r + 1)^{-1}$, so that $r = (1 - \varrho)\varrho^{-1}\sigma^{-2}$. There exist functions $\tau_{\pm,0}, \tau_{\pm,1}, \tau_{\pm,2}, \dots \in \mathcal{A}^\mathcal{E}(\mathbb{S}^{n-1} \times [0, 1])$ such that for any $K \in \mathbb{N}$ and $\sigma > 0$,*

$$\begin{aligned} u_\pm(\theta(1 - \varrho)\varrho^{-1}\sigma^{-2}; \sigma) &= \varrho^{(n-1)/2} \sigma^{(2n-3)/2} \exp\left(\pm i\left(\frac{1 - \varrho}{\sigma\varrho}\right)\left(1 + \frac{Z\varrho}{1 - \varrho}\right)^{1/2}\right) \\ &\quad \times \left(\left(\frac{1 - \varrho}{Z\varrho}\right)^{1/2} + \left(1 + \frac{1 - \varrho}{Z\varrho}\right)^{1/2}\right)^{\pm iZ/\sigma} \left[\sum_{k=0}^K \tau_{\pm,k}(\theta, \sigma) \varrho^k + O_{f,K}(\varrho^{K+1})\right] \end{aligned} \quad (33)$$

holds for every $\theta \in \mathbb{S}^{n-1}$ as $\varrho \rightarrow 0^+$, where

$$O_{f,K}(\varrho^{K+1})|_{\sigma \leq \bar{\sigma}} \in \varrho^{K+1} L^\infty(\mathbb{R}_{\mathbf{x}}^n \setminus \mathbb{B}^n; \mathcal{A}_{\text{loc}}^\mathcal{E}[0, \bar{\sigma})_\sigma) \quad (34)$$

for every $\bar{\sigma} > 0$. ■□

Note that, for σ bounded away from zero, $\varrho \in C^\infty([0, 1]_{1/r})$ and

$$\exp\left(\pm i\frac{1}{\sigma\varrho}(1 - \varrho)\left(1 + \frac{Z\varrho}{1 - \varrho}\right)^{1/2}\right) \left(\left(\frac{1 - \varrho}{Z\varrho}\right)^{1/2} + \left(1 + \frac{1 - \varrho}{Z\varrho}\right)^{1/2}\right)^{\pm iZ/\sigma} \in e^{\pm i\sigma r} r^{\pm iZ/2\sigma} C^\infty([0, 1]_{1/r}) \quad (35)$$

uniformly, so Corollary 1.5 does imply eq. (32).

Finally, we have the following result regarding joint asymptotics as $r \rightarrow \infty$, $\sigma^2 \rightarrow 0$ together, with the product $\varsigma^2 = r\sigma^2$ fixed:

Corollary 1.6 (Asymptotics at tf). *Consider the setup of Corollary 1.4. There exist functions*

$$v_{\pm,k,\varkappa} \in \hat{E}^\varkappa C^\infty(\mathbb{S}^{n-1} \times [0, \infty)_{\hat{E}}) \quad (36)$$

such that, for every $K \in \mathbb{N}$ and $\varsigma > 0$,

$$u_{\pm}(\mathbf{x}; \varsigma/r^{1/2}) = r^{-(2n-3)/4} e^{\pm i r^{1/2} \sqrt{\varsigma^2 + Z}} \left(\frac{\varsigma}{Z^{1/2}} + \sqrt{1 + \frac{\varsigma^2}{Z}} \right)^{\pm i Z r^{1/2}/\varsigma} \left[\sum_{k=0}^K \sum_{\varkappa=0}^{\lfloor k/2 \rfloor} \frac{v_{\pm,k,\varkappa}(\mathbf{x}/r, \varsigma^2)}{r^{k/2}} \log^\varkappa(r) + O_{\varsigma,f,K} \left(\frac{\log^{\lfloor (K+1)/2 \rfloor}(r)}{r^{(K+1)/2}} \right) \right], \quad (37)$$

as $r \rightarrow \infty$, where $O_{\varsigma,f,K}(r^{-(K+1)/2} \log^{\lfloor (K+1)/2 \rfloor}(r))$ is uniformly bounded by $r^{-(K+1)/2} \log^{\lfloor (K+1)/2 \rfloor}(r)$ in compact subsets worth of $\varsigma \in [0, \infty)$.

Moreover, $\rho^{k/2} v_{\pm,k,\varkappa}(\theta, 1/\rho)$ is smooth in $\rho \in [0, \infty)$, for each $\theta \in \mathbb{S}^{n-1}$. Thus, for $\varsigma \geq 1$, $v_{\pm,k,\varkappa}(\theta, \varsigma^2) = O_{f,k,\varkappa}(\varsigma^k)$. ■□

We emphasize that these compound asymptotics are in a different regime than the one relevant to the short-range case (this being the regime of $r \rightarrow \infty$ for $r\sigma$ fixed, not $r\sigma^2$ fixed) — see Remark 6.

Statements analogous to Corollary 1.4, Corollary 1.5, and Corollary 1.6 apply to the A, B in Corollary 1.3.

A few remarks regarding Theorem 1.2:

Remark 3. The apparent singularity of u_{\pm} in Theorem 1.2 when $Z = \sigma^2 a_{00}$ is fictitious — we can write

$$u_{\pm} = e^{\pm i\sigma/x} x^{(n-1)/2 \mp iZ/2\sigma} v_{0,\pm} \quad (38)$$

for $v_{0,\pm} : (0, \infty)_{\sigma} \times X \rightarrow \mathbb{C}$ smooth. In terms of $v_{0,\pm}$, $u_{0,\pm} = e^{\mp i\Phi} (E + Zx)^{1/4} e^{\pm i\sigma/x} x^{(n-1)/2 \mp iZ/2\sigma} v_{0,\pm}$. Thus, $u_{0,\pm}$ is singular when $Z = \sigma^2 a_{00}$. The left-hand side of eq. (19) is smooth for all $\sigma > 0$, but we have written it as a product of two functions which both have singularities when $Z = \sigma^2 a_{00}$. The particular form of Φ in Theorem 1.2 is needed to get uniform estimates down to $\sigma = 0$, but it introduces a fictitious singularity at some positive σ when $a_{00} > 0$. ■

Remark 4. Some explicit bounds for $u_{0,\pm}$ in the b-Sobolev spaces

$$H_b^{m,l}(X) = \{u \in \mathcal{S}'(X) : Lu \in L^2(X, g) \text{ for all } L \in \Psi_b^{m,l}(X)\} \quad (39)$$

will be found in §6. We are indexing the b-Sobolev spaces such that $H_b^{0,0}(X)$ is equal to $L^2(X, g)$ rather than $L^2(X, g_b)$ for some b-metric g_b . This is the (somewhat nonstandard) convention followed in [Vas21a][Vas21c], and it is convenient here for the same reasons. See [Mel93][Vas18] for a pedagogical introduction to the b-calculus.

We write $X_{1/2}$ to denote the mwc canonically diffeomorphic to X over the interior, with the smooth structure at the boundary modified so that $x_{1/2} = 2^{-1/2} x^{1/2}$ becomes a bdf (the factor of $2^{-1/2}$ being a convenient choice). (Note that, in terms of the mwb X , the mwb $X_{1/2}$ is canonically defined without needing to fix a choice of bdf for X , unlike the case for what Wunsch calls X_q in [Wun99, §4]. However, since we are fixing a boundary-collar once and for all, this is not a crucial point.) The b-Sobolev spaces are convenient to work with in part because (except for indexing) they do not depend on whether we use $x_{1/2}$ or x as a bdf — that is, at the level of sets

$$H_b^{m,l}(X) = H_b^{m,2l+n/2}(X_{1/2}), \quad (40)$$

with an equivalence at the level of Banach spaces. A refined estimate of $u_{0,\pm}$ in terms of the “leC-Sobolev spaces”

$$H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) = \{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)\}_{\sigma \geq 0} \quad (41)$$

is given in Proposition 6.9. The leC-Sobolev spaces play the same role here as the sc,b,res-Sobolev spaces in [Vas21c, Theorem 1.1].³ Although it may be natural to develop *doubly*-second-microlocalized Sobolev spaces $H_{\text{sc,sc}_{1/2},\text{b}}^{m,s,\varsigma,l}(X)$, refining both $H_{\text{scb}}^{m,s,l}(X)$ and $H_{\text{scb}}^{m,\varsigma,2l}(X_{1/2})$, we will not do so here, since the leC-Sobolev spaces suffice for the proof of Theorem 1.2. ■

Remark 5. We also have analogues of Theorem 1.2 dropping the classicality assumptions regarding the metric and the potential. If g is a symbolically asymptotically conic metric in the sense below (precisely the sort of metric considered in [Vas21a]) and $V \in S^{-3/2-\delta}(X) = x^{3/2+\delta}S^0(X)$ for $\delta > 0$, the conclusion of the theorem holds except that “ $u_{0,\pm} \in \mathcal{A}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}})$ ” is replaced by the weaker

$$u_{0,\pm} \in \mathcal{A}^{0,0,0}(X_{\text{res}}^{\text{sp}}), \quad \chi u_{0,\pm} \in C^\infty([0, \infty)_E \times X) \text{ for all } \chi \in C_c^\infty(X^\circ). \quad (42)$$

Moreover, if g is classical to α_1 st order, $\alpha_1 > 1$, and V is classical to $\alpha_2 > 3/2$ order (so, for instance, $V \in x^2 C^\infty(X) + S^{-3/2-\delta_2}(X)$ is classical to $3/2 + \delta_2$ th order, and any sc-metric in the sense of [Mel94] is classical to all orders, while the metrics considered by Vasy in [Vas21a] are classical to > 1 order), then, setting $\delta_1 = \min\{\alpha_1 - 1, \alpha_2 - 1\}$ and $\delta_0 = \min\{\alpha_1 - 1, \alpha_2 - 3/2\}$, we have

$$u_{0,\pm} \in \mathcal{A}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0),\delta_1),2\delta_0-,(0,0)}(X_{\text{res}}^{\text{sp}}) \subset \mathcal{A}_{\text{loc}}^{((0,0),\delta_1),(\mathcal{E},2\delta_0-),(0,0)}(X_{\text{res}}^{\text{sp}}) \quad (43)$$

(assuming, for simplicity, that $\delta_0 \notin \mathbb{N}$) where $\mathcal{A}_{\text{loc}}^{((0,0),\delta_1),(\mathcal{E},2\delta_0-),(0,0)}(X_{\text{res}}^{\text{sp}})$ is the Fréchet space of conormal distributions on $X_{\text{res}}^{\text{sp}}$ that have partial polyhomogeneous expansions at each of bf, tf, zf, with merely conormal remainders at order δ_1 at bf and at order $2\delta_0 - \varepsilon$ at tf for every $\varepsilon > 0$, with a full expansion at zf. Elements of this space are smooth at zf in the sense of defining an element

$$C^\infty([0, \infty)_{\varrho_{\text{zf}00}}; \mathcal{A}^{(\mathcal{E},2\delta_0-)}(X_{1/2})) = C^\infty([0, \infty)_{\varrho_{\text{zf}00}}; \mathcal{A}^\mathcal{E}(X_{1/2})) + C^\infty([0, \infty)_{\varrho_{\text{zf}00}}; \mathcal{A}^{2\delta_0-}(X_{1/2})) \quad (44)$$

via restriction to a small neighborhood of $\text{zf} \subset X_{\text{res}}^{\text{sp}}$. This implies smoothness at zf° .⁴

We refer to [Mel92, Eq. 22][Hin22, Def. 2.13] for the definition of the spaces of conormal distributions with partial polyhomogeneous expansions. We follow the notational conventions in [Hin22], *mutatis mutandis*. ■

Remark 6. Note that polyhomogeneity on the mwc $X_{\text{res}}^{\text{sp}}$ with index set \mathcal{E} at tf and smoothness elsewhere is equivalent to polyhomogeneity on

$$X_{\text{res},0}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \partial X] \quad (45)$$

with index set $\{(k/2, \kappa) : (k, \kappa) \in \mathcal{E}\}$ at the front face tf of the blow-up and index set \mathbb{N} at the faces bf and zf. So, we have a Taylor series in powers of x for $\sigma > 0$ and a Taylor series in powers of $x^{1/2}$ at $\sigma = 0$, with a Taylor series in powers of $E = \sigma^2$ for fixed $x > 0$ (with some logs in the expansion at tf, except exactly at zf). We remark that the case of nonconstant a_{00} involves a more general polyhomogeneity statement on $X_{\text{res},0}^{\text{sp}}$; for fixed $E > 0$, this is discussed already in [Mel94, §14].

(The use of $(E + \text{Zx})^{1/2} = (\sigma^2 + \text{Zx})^{1/2}$ as a bdf for the front face tf of $X_{\text{res}}^{\text{sp}}$ is a convenient but otherwise arbitrary convention. The function $(E + x)^{1/2} = (\sigma^2 + x)^{1/2}$ would also work, as already utilized in Corollary 1.4, Corollary 1.5, Corollary 1.6, but it is less convenient.) Although $X_{\text{res},0}^{\text{sp}}$ is diffeomorphic to the mwc used in [Vas21c][Hin22] to study the low energy resolvent in the short-range case, it differs from their mwc as a compactification of $(0, 1]_\sigma \times X$. Indeed, besides

³We will omit the comma between “sc,b.”

⁴In this paper, we use ‘ \circ ’ and the term “interior” in the mwc-theoretic sense; for instance, for any boundary hypersurface f , f° is the collection of points in f that are not in any other boundary hypersurface.

having a different smooth structure at $\{0\} \times X^\circ$, our blow-up resolves the ratio E/x rather than σ/x , the latter being the ratio parametrizing the front face of the blow-up in [Vas21c][Hin22]. ■

Remark 7. When

- $g = g_0$ is an exactly conic metric and
- $f, V \in C_c^\infty(X^\circ)$,

the radial dependence of u_\pm near infinity can be solved for (up to a multiplicative constant) exactly via separation of variables. When separating variables, we need only solve the “radial ODE,” which ends up being a Whittaker ODE for $\sigma > 0$ and a Bessel ODE for $\sigma = 0$. This motivating example is discussed in Appendix §A, where the consequences of Theorem 1.2 for solutions to Whittaker’s ODE are discussed. The necessity of the factor of $(\sigma^2 + Zx)^{-1/4}$ in eq. (19) can be read off of the asymptotics of the Bessel functions (or can alternatively be deduced from the structure of the radial ODE, with slightly more work). See eq. (723). Even in this exactly conic case, which was studied in [Yaf82], Theorem 1.2 (or at least the part specifying the existence of a full asymptotic expansion) seems to be novel. ■

Remark 8. At first, it may seem somewhat paradoxical that at zero energy an attractive (rather than repulsive) Coulomb-like Schrödinger operator has scattering behavior. The paradox is resolved if we realize that, if a classical particle traveling in an attractive Coulomb force field has zero energy, then it must have more kinetic energy than if the force field were absent (the point being that “zero energy” refers to energy as measured *with* the potential present, not without).

It may also seem somewhat surprising that the presence of the attractive Coulomb potential allows us to avoid the b-analysis of the Laplacian in [Vas21c], but such considerations are expected to be necessary in order to understand the $E \rightarrow 0^-$ limit. Roughly speaking, the presence of an attractive Coulomb potential has moved the locus of b-analysis infinitesimally negativewards – more specifically negativewards on the front face of $[\mathbb{R}_E \times X; \{0\} \times \partial X]$, to its intersection with the lift of $\{E = -Zx\}$ – where it is irrelevant to the $E \rightarrow 0^+$ limit studied here. ■

The proof of Theorem 1.2, which is a special case of Proposition 6.3, is spread throughout §6, with a key estimate proven in §5 (and other lemmas appearing in §2, §3, §4). The conormality component of the theorem and smoothness at zf° are deduced in Proposition 6.15. From there, classicality is deduced via an inductive argument using an explicit parametrix for the “b,leC-normal operator” of the conjugated spectral family — see Proposition 6.16. Under only partial classicality assumptions, the inductive step of the argument can only be carried out finitely many times, and the upshot is Proposition 6.3. The analysis in §6 is essentially the analysis of a family of ODEs and applies equally well in the $n = 1$ case. The analysis in §5 is a uniform version of the sc-analysis of the Laplacian carried out in [Mel94][Vas21a], where we rely on the presence of the attractive Coulomb potential to prevent a full degeneration (from the perspective of the b-calculus) of the spectral family to the Laplacian at zero energy.

We note four deficiencies of our treatment, which can form the basis for further work:

- (1) while we discuss the output and mapping properties of the low energy resolvent, we do not discuss its Schwartz kernel (as Guillarmou, Hassell, and Sikora [GHS13] do in the case of a short-range potential),
- (2) we only treat the case when the Coulomb-like potential has the form Zx for constant $Z > 0$, while more general $Z \in C^\infty(\partial X; \mathbb{R}^+)$ (and $a_{00} \in C^\infty(\partial X; \mathbb{R})$) may be of interest [Mel94, §14] (if $Z \in C^\infty(\partial X; \mathbb{R}^+)$, a change of coordinates from x to $x_0 = Zx$ does not necessarily preserve the form of the metric we require, so we cannot just reduce the case of nonconstant Z to the constant case via a change of bdf),
- (3) it should be possible to study the resolvent for σ in a cone $\{\sigma \in \mathbb{C} : \arg(\sigma) \in [0, \theta)\}$ for $\theta \in (0, \pi/2)$, in accordance with the work of Skibsted and Fournais [FS04][Ski13] in the Euclidean case – see [Ski13, Theorem 1.2] in particular – thus giving a version of the limiting

absorption principle that is uniform down to zero energy (as opposed to the *ad hoc* use of $Z \pm i0$ in Theorem 1.2 to describe the zero energy limit), and

- (4) it should be possible to extend Theorem 1.2 to the case where the Coulomb potential and manifold possess conic singularities, so as to handle an exact Coulomb potential on $[\mathbb{R}^d; \{0\}]$ as a special case. (Indeed, blowing up the origin of Euclidean space and using r as a bdf, the Coulomb potential Z/r and spectral term σ^2 are both lower order than the Euclidean Laplacian with respect to the b-calculus at $r = 0$ in terms of both decay and regularity.) The exact 1D model problem discussed in §A is an example of this.

Along with some analysis of the $E < 0$ case, the present work can serve as input to the study of the Klein–Gordon equation on (not necessarily axially symmetric) asymptotically Schwarzschild spacetimes (away from the event horizon) in the spirit of Hintz’s recent treatment of Price’s law [Hin22] on asymptotically subextremal Kerr spacetimes. This problem was the original motivation for the present work. Some heuristic investigations in this direction have been undertaken by physicists [Det80][HP98][KT01][KT02][BK04][KZM07][Bar+12][Bar+14], usually in the case of a spherically or axially symmetric black hole spacetime (e.g. Schwarzschild, Reissner–Nordström, Kerr, etc.). Indeed, [KT01] treat their problem (pointwise temporal decay rates) using an uncontrolled approximation of the radial part of their PDE as a Whittaker ODE, this being equivalent to the radial part of the time-independent Schrödinger equation in an attractive Coulomb potential. This problem can be solved exactly in terms of special functions (see §A), but we do not need to do so in order to understand the asymptotics.

Moreover, while the asymptotics for fixed $E > 0$ can be read off of the exact solution (using the large argument expansions of Whittaker functions), the exact solution does not even help much in understanding the $E \rightarrow 0$ limit. Instead of relying on the Whittaker approximation, it is possible to exactly solve the Klein–Gordon equation on the Schwarzschild exterior (at the level of individual spherical harmonics) using confluent Heun functions [Bar+14, §IV]. When this paper was first written, it was an open problem to rigorously establish temporal decay at a polynomial rate (as this involves knowing the asymptotics of the confluent Heun functions in some regime where the argument and a parameter are both taken to infinity). Since then, [PSRM23][SRM23] proved the required result. For more complicated spacetimes, no exact solution is possible. A more robust analysis is therefore necessary.

For other mathematical work on Price’s law and similar problems, see Donninger–Schlag–Soffer [DSS11][DSS12], Metcalfe–Tataru–Tohaneanu [MTT12][Tat13], Morgan–Wunsch [Mor19][MW21], and Looi [Loo21]. The Klein–Gordon equation on Kerr-like spacetimes is subtle — for instance, on subextremal Kerr, Shlapentokh–Rothman [SR14] has constructed finite energy solutions that grow exponentially in time by exploiting superradiant instability.

We remark finally that there does not appear to be much overlap between the present work and similarly titled physics papers, e.g. [BOW67][BT67][BTO67][MB67][Mac67], because these concern the scattering of electrons off *atomic* hydrogen, not ionic hydrogen. The former consists of one hydrogen nucleus and one electron (different from the incoming electron) which is tightly bound to the nucleus. Thus, atomic hydrogen is electrically neutral, and the incoming electron feels at most a dipole interaction, which is short-range.

1.2. Preliminaries and Outline. Recall that an exactly conic manifold X (according to one of several essentially equivalent definitions) consists of the following data: an arbitrary (smooth, connected) compact manifold-with-boundary X_0 , a Riemannian metric $g_{\partial X} \in C^\infty(X; \text{Sym}^2 T^* \partial X_0)$ on the closed manifold ∂X_0 , and an embedding

$$\iota : \hat{X} = [0, \bar{x}]_x \times \partial X_0 \rightarrow X_0 \quad (46)$$

(the “boundary-collar”) which is a diffeomorphism between the cylinder $[0, \bar{x}]_x \times \partial X_0$ for some $\bar{x} > 0$ and a neighborhood of ∂X_0 . Note that ι satisfies $\text{id}|_{\partial X_0} = \iota(0, -)$. We notationally conflate X with

the underlying X_0 . We also conflate the image of ι with $[0, \bar{x}]_x \times \partial X_0$. (This is an alternative, if a somewhat crude one, to Melrose’s use in [Mel94] of the inward pointing normal bundle of ∂X_0 .) Given an exactly conic manifold X , a Riemannian metric g_0 on X° is called an exactly conic metric (this notion being defined relative to ι) if

$$g_0 \in C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X), \quad g_0 = x^{-4} dx^2 + x^{-2} g_{\partial X} \text{ near } \partial X. \quad (47)$$

Below, we write $x : X \rightarrow [0, \infty)$ to denote a globally defined boundary-defining function, compatible with the boundary-collar ι in the sense that $x(\iota(x_0, y)) = x_0$ for all $x_0 \in [0, \bar{x}]$. (Euclidean space fits into this framework, where by “Euclidean space” we mean the radial compactification of \mathbb{R}^d , endowed with the induced metric. Here $x = 1/\langle r \rangle$, r being the Euclidean radial coordinate. See [Mel94, §1].) We will also consider (symbolically) asymptotically conic metrics, which – according to our use of the term “asymptotically” – are Riemannian metrics g on X° of the form

$$g - g_0 \in x C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) + S^{-1-\delta}(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \quad (48)$$

for some $\delta > 0$. Hence, g is conic up to a classical subleading term – decaying faster than g_0 by one order of x – and a merely symbolic error decaying slightly faster than that. More generally, for $\alpha_1 > 1$, we will say that the asymptotically conic metric g is “classical to α_1 th order” if eq. (48) can be strengthened to

$$g - g_0 \in x C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) + S^{-\alpha_1}(X; {}^{\text{sc}}\text{Sym}^2 T^* X). \quad (49)$$

We will call g “(fully) classical” if the symbolic term in eq. (48) can be dropped entirely, i.e. if eq. (49) holds for all $\alpha_1 > 1$. (Any sc-metric in the sense of [Mel94] is fully classical in this sense.) The term *asymptotically conic manifold* will be used to refer to a conic manifold equipped with an asymptotically conic metric, so a triple (X, ι, g) . This is a more general notion than that of an “sc-manifold” [Mel94][Mel95][JSB99][HW05], which are also sometimes called asymptotically conic — besides allowing a symbolic term in eq. (48), we also allow a classical $O(x^{-3})dx^2$ term (see a_{00} in Theorem 1.2). (Rewriting the metric in terms of the tortoise coordinate x_* , the $O(x^{-3})dx^2$ term can be removed, but x_* is not a smooth function of x and so x_* is not a bdf on the mwb X .) We will assume that this term is constant on ∂X . Thus, the metrics we consider are Riemannian metrics on X° of the form

$$g = g_0 + a_{00} \frac{dx^2}{x^3} + \frac{\Gamma_{1, \partial X} \odot dx}{x^2} + \frac{h_{1, \partial X}}{x} + x^2 C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) + S^{-\alpha_1}(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \quad (50)$$

for some $a_{00} \in \mathbb{R}$, 1-form $\Gamma_{1, \partial X} \in \Omega^1(\partial X)$, and symmetric 2-cotensor $h_{1, \partial X} \in C^\infty(\partial X; \text{Sym}^2 T^* \partial X)$. A fully classical g is then of the form

$$g = g_0 + a_{00} \frac{dx^2}{x^3} + \frac{\Gamma_{1, \partial X} \odot dx}{x^2} + \frac{h_{1, \partial X}}{x} + x^2 C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \quad (51)$$

near ∂X .

Before we specialize to the case of spectral families of Schrödinger operators for the proof of Theorem 1.2, we study in §3, §5 smooth families $\{P(\sigma)\}_{\sigma \geq 0} \subset \text{Diff}^2(X^\circ)$ of elliptic 2nd-order differential operators on X° of the form

$$P(\sigma) = P_0(\sigma) + P_1(\sigma) + P_2(\sigma), \quad (52)$$

where, for each $i = 1, 2, 3$, $P_i = \{P_i(\sigma)\}_{\sigma \geq 0} \subset \text{Diff}^2(X^\circ)$ is a family of differential operators (depending smoothly on $\sigma^2 \geq 0$, all the way down to $\sigma = 0$), which near ∂X can be written as

$$\bullet \quad P_0(\sigma) = -(1 + x a_{00}(\sigma))(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} + x^3 [a(\sigma) + n - 1] \partial_x - Zx - \sigma^2 \quad (53)$$

for

$$\begin{aligned}
& - a_{00} \in C^\infty([0, \infty)_{\sigma^2}; \mathbb{R}), Z > 0 \text{ satisfying the } \textit{attractivity condition} \\
& Z - \sigma^2 a_{00}(\sigma) > 0
\end{aligned} \tag{54}$$

$$\begin{aligned}
& \text{for each } \sigma \geq 0, \\
& - a \in C^\infty([0, \infty)_{\sigma^2}),
\end{aligned}$$

•

$$P_1(\sigma) = \sum_{j=1}^J (x^4 P_{\perp,j}(\sigma) b_j(x; \sigma) \partial_x + x^3 b'_j(x; \sigma) P_{\partial X,j}(\sigma) + x^2 b''_j(x; \sigma) Q_{\partial X,j}(\sigma)) \tag{55}$$

for some $J \in \mathbb{N}$, where

$$- b_j(-; \sigma), b'_j(-; \sigma), b''_j(-; \sigma) \in S^0(X), \text{ more specifically}$$

$$b_j, b'_j, b''_j \in C^\infty([0, \infty)_{\sigma^2}; S^0(X)), \tag{56}$$

$$- P_{\perp,j}(\sigma), Q_{\partial X,j}(\sigma) \in \text{Diff}^1(\partial X) \text{ and } P_{\partial X,j}(\sigma) \in \text{Diff}^2(\partial X) \text{ for each } \sigma \geq 0, \text{ so that}$$

$$P_{\perp,j}, Q_{\partial X,j} \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^1(\partial X)), \tag{57}$$

$$P_{\partial X,j} \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^2(\partial X)), \tag{58}$$

for each $j = 1, \dots, J$,

$$\bullet P_2 \in C^\infty([0, \infty)_{\sigma^2}; S \text{Diff}_{\text{scb}}^{2,-1-\delta,-3/2-\delta}(X)) \text{ for some } \delta > 0.$$

Here $\text{Diff}_{\text{scb}}^{m,s,l}(X) = \text{Diff}_b^{m,l}(X) \cap \text{Diff}_{\text{sc}}^{m,s}(X)$ and $S \text{Diff}_{\text{scb}}^{m,s,l}(X) = S \text{Diff}_b^{m,l}(X) \cap S \text{Diff}_{\text{sc}}^{m,s}(X)$. We may take $\delta < 1/2$, as this will be useful in reducing casework later on. Note that we are requiring smoothness with respect to $E = \sigma^2$, rather than with respect to σ , in compact subsets of X° . This is not strictly necessary for the analysis in §5, for which even smoothness with respect to σ is mostly unnecessary, but since we use smoothness in E in §6 we will work with it from the outset and develop Ψ DO calculi accordingly in §2. We will say that P_1 is “classical to β_1 th order” for $\beta_1 > 0$ if we can replace eq. (56) by

$$b_j, b'_j, b''_j \in C^\infty([0, \infty)_{\sigma^2}; C^\infty(X)) + C^\infty([0, \infty)_{\sigma^2}; x^{\beta_1} S^0(X)). \tag{59}$$

Similarly, we say that P_2 is “classical to (β_2, β_3) th order” for $\beta_2 > \beta_3 > 0$ if

$$\begin{aligned}
P_2 & \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}_{\text{scb}}^{2,-2,-2}(X) + S \text{Diff}_{\text{scb}}^{2,-1-\beta_2,-3/2-\beta_2}(X) + x^{3/2+\beta_3} S^0(X)) \\
& = C^\infty([0, \infty)_{\sigma^2}; \text{Diff}_{\text{sc}}^{2,-2}(X) + S \text{Diff}_{\text{scb}}^{2,-1-\beta_2,-3/2-\beta_2}(X) + x^{3/2+\beta_3} S^0(X))
\end{aligned} \tag{60}$$

(note that $\text{Diff}_{\text{scb}}^{2,-2,-2}(X) = \text{Diff}_{\text{sc}}^{2,-2}(X)$). We could work with even finer measures of classicality, but for the proof of Theorem 1.2 (and eventually Proposition 6.3) it suffices to keep track of $\beta_1, \beta_2, \beta_3$.

We restrict attention to the case of constant Z , though the case of variable Z may also be of interest. (The $x^3(n-1)\partial_x$ term in eq. (55) is merely conventional and can be parametrized away upon redefining a , but it is convenient because $x^2\partial_x - x(n-1)/2$ is formally anti-self-adjoint with respect to the $L^2 = L^2([0, \bar{x}) \times \partial X, g_0) = L^2_{\text{sc}}([0, \bar{x}) \times \partial X)$ inner product.) The “main” and “subleading” terms of P are collected in P_0 . Although subleading by only one order in the sense of decay, the terms in P_1 are all negligible in the analysis of §5, as are the terms in P_2 (which is why they are allowed to be symbolic). (The key point here is that the principal symbol of P_1 vanishes to one additional order at the radial sets of P ’s Hamiltonian flow because of the tangential derivatives. On the other hand, P_2 is just sufficiently low order everywhere.) Note that we are only requiring $P(\sigma), P_0(\sigma)$ to be elliptic in the traditional sense; the operators above can be considered as sc-operators, but they will be nonelliptic in the sc-calculus. The ‘ S ’ in front of $S \text{Diff}$ indicates that the differential operators therein have coefficients which are required only to be conormal functions on X , hence are “symbolic.”

We do not concern ourselves with uniformity as $\sigma \rightarrow \infty$, this already being handled in [Vas21a, §5] for a suitably large class of operators.

By “attractive Coulomb-like Schrödinger operator (on X),” we mean an operator of the form

$$P(0) = \Delta_g - Zx + V(x) \quad (61)$$

for some asymptotically conic metric g on X , $Z > 0$, and real-valued $V \in S^{-3/2-\delta}(X)$ for $\delta > 0$. The spectral family of $P(0)$ is the family $\{P(\sigma) = P(0) - \sigma^2\}_{\sigma \geq 0}$. The spectral families of attractive Coulomb-like Schrödinger operators are therefore included in the class of families we study here if we restrict attention to sufficiently small σ such that the attractivity condition eq. (54) is satisfied. See Proposition 6.2. From the explicit formula for the Laplacian in local coordinates, we see that a_{00} is the restriction to ∂X of $-x^3(\delta g)_{00}$, where $\delta g = g - g_0$. Although it is only obvious upon changing coordinates from x to $x/(1 + xa_{00})$ – see §A – a nonzero value of a_{00} has essentially the same effect as a nonzero Coulomb potential on the asymptotics of solutions to Helmholtz’s equation for positive σ , and it turns out that eq. (54) – taking into account both the actual Coulomb potential and the effective Coulomb potential from a_{00} – is the correct notion of “attractivity” in this generality. Note that if P is the spectral family of a Coulomb-like Schrödinger operator with $Z > 0$ on an asymptotically conic manifold, then we can modify P by cutting off a_{00} outside of some sufficiently small neighborhood of $\sigma = 0$ such that the resulting family of operators satisfies the attractivity condition. We only care about behavior in a neighborhood of $\sigma = 0$, so this does not restrict applicability.

For $\sigma \gg 0$, attractive Coulomb-like operators do not need to be distinguished from non-attractive Coulomb-like operators as far as the limiting absorption principle is concerned, except with regards to the precise logarithmic corrections to spherical waves. The difference matters only in the $\sigma \rightarrow 0$ limit. That this limit is delicate can be seen already in the case when $X = [0, 1]_x$, equipped with an exactly conic metric (x now only denoting a bdf for one end). Then, when the potential is an exact Coulomb potential near the $x = 0$ boundary, $P(\sigma)u = 0$ is an ODE,

$$-(x^2 \partial_x)^2 u - \sigma^2 u - Zxu = 0 \quad (62)$$

for x sufficiently close to zero. This is a form of Whittaker’s ODE for $\sigma > 0$, degenerating to a Bessel ODE

$$-(x^2 \partial_x)^2 u - Zxu = 0 \quad (63)$$

at $\sigma = 0$. Hence, smooth families $\{u(-; \sigma)\}_{\sigma \geq 0}$ of solutions $u = u(-; \sigma)$ to $P(\sigma)u(-; \sigma) = f$ for $f \in C_c^\infty(X^\circ)$ are near $x = 0$ linear combinations of Whittaker functions (evaluated along the imaginary axis) for $\sigma > 0$ degenerating in some way to linear combinations of Bessel functions as $\sigma \rightarrow 0^+$. See Proposition A.3 for a precise statement. An analysis of the low-energy limit of attractive Coulomb-like Schrödinger operators must therefore – at the very least – describe in appropriate asymptotic regimes the degeneration of Whittaker functions to Bessel functions (a classical topic, but nonetheless delicate). The former oscillate like

$$\exp\left(\pm i \frac{\sigma}{x}\right) x^{\mp iZ/2\sigma} = \exp\left(\pm i \left(\frac{\sigma}{x} - \frac{Z}{2\sigma} \log x\right)\right) \quad (64)$$

as $x \rightarrow 0^+$, while the latter oscillate like $\exp(\pm 2iZ^{1/2}/x^{1/2})$. While the leading order term of the phase in eq. (64) is suppressed as $\sigma \rightarrow 0^+$, the logarithmic correction $\mp i(Z/\sigma) \log x$ blows up as $\sigma \rightarrow 0^+$. This, along with $O_\sigma(1)$, $O_\sigma(x)$, etc. contributions which were omitted from eq. (64) (and include terms like $1/\sigma$, $1/\sigma^2$, \dots), end up contributing as $\sigma \rightarrow 0^+$ to the phase $\pm 2iZ^{1/2}/x^{1/2}$ relevant at zero energy. Clearly, eq. (64) paints an incomplete picture of the oscillations of solutions to eq. (62) in the $\sigma \rightarrow 0^+$ limit. This is the first indication that the analysis of the low energy limit will need to be done on a resolution of $[0, \infty)_\sigma \times X$. We refer the reader to §A for a further discussion of the ODE case.

The “conjugated perspective” involves studying the conjugated spectral family $\tilde{P} = \{\tilde{P}(\sigma)\}_{\sigma \geq 0} \subset \text{Diff}^2(X^\circ)$,

$$\tilde{P}(\sigma) = \exp(-i\Phi)P(\sigma)\exp(+i\Phi) : \mathcal{S}'(X) \ni u \mapsto \exp(-i\Phi)(P(\sigma)(\exp(+i\Phi)u)), \quad (65)$$

where the “phase” $\Phi = \Phi(-; \sigma)$ is an appropriate function on $[0, \infty)_\sigma \times X$ such that $\exp(+i\Phi)$ captures the asymptotics of outgoing solutions u to the PDE $Pu = f$ to some desired level of precision. (Thus, we deal with the ‘+’ case of Theorem 1.2, the ‘−’ case being analogous.) At the level of the phase space of the sc-calculus, this conjugation corresponds to a symplectomorphism moving one of the two sets of radial points (the “selected radial set”) to the zero section of the sc-fibers, which upon second microlocalization gets blown up to the fibers of ${}^b\bar{T}^*_{\partial X}X$, the phase space of the b-calculus over ∂X (see [Vas21a]). The following choice of Φ is sufficiently detailed:

$$\Phi(x; \sigma) = \frac{1}{x} \sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} + \frac{1}{\sigma} (Z - \sigma^2 a_{00}) \operatorname{arcsinh} \left(\frac{\sigma}{x^{1/2}} \frac{1}{(Z - \sigma^2 a_{00})^{1/2}} \right) - \frac{i}{2} a \log x, \quad (66)$$

(extended from $\sigma > 0$ to $\sigma = 0$ by continuity), hence the appearance of Φ in Theorem 1.2 (for which we have $a = 0$). The square roots in eq. (66) are well-defined by the attractivity condition eq. (54). See the beginning of §3 for a motivation of eq. (66). Given compact $K \subset (0, \infty)$, we can write, for $\sigma \in K$,

$$\Phi = \sigma x^{-1} - \frac{Z - \sigma^2 a_{00}}{2\sigma} \log x - \frac{i}{2} a \log x + O_K(1) \quad (67)$$

as $x \rightarrow 0^+$. This is in accordance with eq. (64), [Vas21a, Theorem 1.1]; Φ differs from Vasy’s phase modulo a logarithmic plus smooth correction for each individual $\sigma > 0$. The O_K term in eq. (67), when formally expanded around $\sigma = 0^+$, contains terms proportional to σ^{-2N} for all $N \in \mathbb{N}^+$. The estimate eq. (67) is therefore not uniform as $\sigma \rightarrow 0$. Similarly, while we can write $(\sigma^2 + Zx)^{1/2} = Z^{1/2}x^{1/2}(1 + \sigma^2/2Zx + O(\sigma^4/Z^2x^2))$ for x bounded away from zero, this obviously cannot be used to understand the asymptotics of outgoing solutions to $Pu = f$ for any $\sigma > 0$. Hence, the precise form of $(\sigma^2 + Zx)^{1/2}$ in eq. (66) (and the precise form of the other terms appearing in eq. (66)) is important (at least modulo functions which are well-behaved on $X_{\text{res}}^{\text{sp}}$). This is the second indication that we will need to resolve the $x, \sigma \rightarrow 0^+$ regime, and apparently in doing so we had ought to resolve the ratio x/σ^2 .

This is accomplished by the mwc $X_{\text{res}}^{\text{sp}}$. While, as already remarked upon in Remark 6, this resolution is reminiscent of that used in [Vas21c][Hin22] in order to resolve the low energy behavior of interest there, it performs a somewhat different role here; theirs was used to study the degeneration of $\Delta_g - \sigma^2$ as $\sigma \rightarrow 0^+$ to an elliptic element

$$\Delta_g \in \operatorname{Diff}_b^{2,-2}(X) \quad (68)$$

of the b-calculus, but

$$\Delta_g - Zx \in \operatorname{Diff}_b^{2,-1}(X) \quad (69)$$

is not elliptic in the b-calculus, hence the oscillating solutions to the ODE eq. (63) noted above. Instead, we use the resolution to interpolate between Vasy’s analysis in [Vas21a] of $\Delta_g - \sigma^2 - Zx$ performed in $\Psi_{\text{scb}}(X)$ for $\sigma > 0$ and the analysis of the “zero-energy operator” $\Delta_g - Zx$ performed in $\Psi_{\text{scb}}(X_{1/2})$ along *similar lines* (see §4). According to the previous paragraph, what makes the Φ in Theorem 1.2 the “right” choice is the asymptotics of Φ near the boundary $\partial X_{\text{res}}^{\text{sp}}$ of this resolution (or more accurately near $\partial X_{\text{res}}^{\text{sp}} \setminus \text{zf}^0 = \text{bf} \cup \text{tf}$). Indeed, for $\sigma = 0$, we can write

$$\Phi = 2\sqrt{\frac{Z}{x}} - \frac{i}{2} a \log x. \quad (70)$$

Hence, eq. (66) interpolates between the oscillations eq. (67) seen in solutions of the ODE eq. (62) at positive energy and the oscillations eq. (70) seen in solutions of the ODE eq. (63) at zero energy.

We refer to §3 for a further discussion of the conjugated operator family and §4 for a discussion of the situation at zero energy.

As our main technical tool, we situate the family $\tilde{P} = \{\tilde{P}(\sigma)\}_{\sigma \geq 0}$ in a pseudodifferential calculus,

$$\Psi_{\text{leC}}(X) = \bigcup_{m,s,\varsigma,l,\ell \in \mathbb{R}} \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X), \quad (71)$$

which we will call the “leC-calculus” (“low energy Coulomb”-calculus, for lack of a better name), the elements of which can be interpreted as particular families of b- Ψ DOs on X . The calculus comes with a refined symbol calculus tailored to the problem at hand. Compared to the calculus of b- Ψ DOs with parameters (i.e. $\Psi_b(X)$ -valued symbols on some parameter space), the symbol calculus here is refined in the sense of being second-microlocalized à la Vasy (so as to keep track both of b-decay and sc-decay orders) and “resolved at the corner” (so as to keep track of the asymptotic regime when both $\sigma^2 \rightarrow 0$ and $x \rightarrow 0$ at compatible rates). The corresponding symbols are conormal functions on the “leC-phase space,” an iterated blow-up of $[0, \infty)_\sigma \times {}^b\overline{T}^*X$. This mwc has six boundary faces – df, sf, ff, bf, tf, zf – and is described in the next section. In eq. (71), m is the “differential order” (order at fiber infinity, df), s is the sc-decay order at positive energy (order at sf), ς is the sc-decay order at zero energy (order at ff), l is the b-decay order at positive energy (order at bf), and ℓ is the b-decay order at zero energy (order at tf). We remark that the scattering calculus with respect to x^2 (rather than $x^{1/2}$) has been used by Wunsch [Wun99][HW05], who called it the “quadratic scattering calculus.” The leC-calculus is discussed in §2.

We now say a word about our usage of the symbol ‘ \preceq ’ (the usage of ‘ \succeq ’ being analogous). When stating a proposition involving an estimate, we will be explicit about the dependence of the constants involved on parameters. In order to avoid a proliferation of symbols denoting different but unimportant constants, when proving a proposition of the form

- for all $p_1 \in \mathcal{P}_2, \dots, p_M \in \mathcal{P}_M$, there exists a constant $C(p_1, \dots, p_M) > 0$ such that $r_1(p_1, \dots, p_N) \leq C(p_1, \dots, p_M)r_2(p_1, \dots, p_N)$ for all $p_{M+1} \in \mathcal{P}_{M+1}, \dots, p_N \in \mathcal{P}_N$,

(where $M, N \in \mathbb{N}$, $N \geq M$, $\mathcal{P}_2, \dots, \mathcal{P}_N$ are some sets, $r_i : \mathcal{P}_2 \times \dots \times \mathcal{P}_N \rightarrow \mathbb{R}$ for $i = 1, 2$) we will write an intermediate estimate of the form

$$r_3(p_1, \dots, p_N) \leq C'(p_1, \dots, p_M)r_4(p_1, \dots, p_N) \quad (72)$$

as $r_3 \preceq r_4$, with the key point being that, according to eq. (72), C' depends only on the parameters that C depends on (so that the estimate eq. (72) is “uniform” in p_{M+1}, \dots, p_N). All constants below depend on the geometric data in the setup of Theorem 1.2, so we will not be explicit about that dependence.

2. THE leC-CALCULUS

We now turn to our discussion of the leC-calculus. This calculus is, in many ways, similar to Vasy’s second microlocalized calculi $\Psi_{\text{scb}}(X) = \cup_{m,s,l \in \mathbb{R}} \Psi_{\text{scb}}^{m,s,l}(X)$ [Vas21a, §2] and $\Psi_{\text{scb, res}}(X) = \cup_{m,s,l,\ell \in \mathbb{R}} \Psi_{\text{scb, res}}^{m,s,l,\ell}(X)$ [Vas21c, §3], with the main novel feature of the leC-calculus (besides the relatively unimportant alteration of the smooth structure at $\sigma = 0$) being another resolution of the phase space (which we actually carry out before resolving the scattering face for $\sigma > 0$). This can be seen at a glance, comparing Figure 3 to Figure 4, [Vas21c, Figure 4]. The zero face zf of the leC-phase space is identifiable with the phase space ${}^{\text{scb}}\overline{T}^*X_{1/2}$ of the calculus $\Psi_{\text{scb}}(X_{1/2})$, while for $\sigma_0 > 0$ the $\{\sigma = \sigma_0\}$ cross-section of the leC-phase space is identifiable with the phase space ${}^{\text{scb}}\overline{T}^*X$ of the scb-calculus. Thus, the leC-calculus interpolates between these two calculi as $\sigma \rightarrow 0^+$, as $\Psi_{\text{scb, res}}(X)$ interpolates between $\Psi_{\text{scb}}(X)$ and $\Psi_b(X)$ in the same limit.

The leC-phase space ${}^{\text{leC}}\overline{T}^*X$ is introduced in §2.1, along with corresponding algebras of symbols. Calculi are discussed in §2.2, and the corresponding leC-Sobolev spaces (which are really families of scb-Sobolev spaces) are discussed in §2.3. As many of the results in this section are either consequences of standard results for the b-calculus or derivable via similar arguments, some details are omitted. Still, we’ve made an effort to give a relatively complete list of the results needed later, and in the process we give the leC analogues of some standard arguments. In this section we mainly write $\lambda = \sigma^2$ in place of the parameter E used in the introduction.

When we write ‘ $\Psi_b(X)$,’ and more generally when we refer to “b- Ψ DOs,” we mean the *conormal* (a.k.a. “symbolic”) b-algebra [Vas18, Definition 5.15], rather than the closely related, slightly

smaller calculus defined in [Mel93, Definition 4.22]. The latter calculus has symbols that have some additional classicality at the boundary. Our convention follows [Vas21a]. In [Vas21b], the notation ‘ $\Psi_{bc}(X)$ ’ is used instead.

2.1. Phase Spaces and Symbols. Recall that we can identify ${}^bT^*X$ over the boundary collar with $[0, \bar{x})_x \times \mathbb{R}_{\xi_b} \times (T^*\partial X)_{\eta_b}$ via the map

$$(0, \bar{x})_x \times \mathbb{R}_{\xi_b} \times (T^*\partial X)_{\eta_b} \ni (x, \xi_b, \eta_b) \mapsto \xi_b \frac{dx}{x} + \text{plr}^*(\eta_b), \quad (73)$$

where $\text{plr} = \pi_R \circ \iota^{-1} : \iota(\hat{X}) \rightarrow \partial X$. This defines a diffeomorphism between $[0, \bar{x})_x \times \mathbb{R}_{\xi_b} \times (T^*\partial X)_{\eta_b}$ and ${}^bT^*_{\iota(\hat{X})}X$.

Let

$$\begin{aligned} {}^{b,\text{sp}}\overline{T}^*X &= [[0, \infty)_\lambda \times {}^b\overline{T}^*X; \{\lambda = x = 0\}] \\ &= [[0, \infty)_\lambda \times {}^b\overline{T}^*X; \{0\} \times {}^b\overline{T}^*_{\partial X}X] \end{aligned} \quad (74)$$

denote the phase space of the resolved calculus of 1-parameter families of b-ΨDOs ([Vas21c, Figure 4]), and letting $\beta : {}^{b,\text{sp}}\overline{T}^*X \rightarrow [0, \infty)_\lambda \times {}^b\overline{T}^*X$ denote the blowdown map, let

$$\begin{aligned} \text{zf}_{00} &= \text{cl } \beta^{-1}(\{\lambda = 0, x > 0\}), \\ \text{tf}_{00} &= \beta^{-1}(\{\lambda = 0 = x\}), \\ \text{bf}_{00} &= \text{cl } \beta^{-1}(\{\lambda > 0, x = 0\}), \end{aligned} \quad (75)$$

and

$$\text{df}_{00} = \text{cl}((\partial^{\text{sp},b}\overline{T}^*X) \setminus (\text{zf}_{00} \cup \text{tf}_{00} \cup \text{bf}_{00})) \quad (76)$$

denote its (closed) boundary faces. For each $f \in \{\text{zf}_{00}, \text{tf}_{00}, \text{bf}_{00}, \text{df}_{00}\}$, let $\varrho_{0,f} \in C^\infty({}^{b,\text{sp}}\overline{T}^*X; [0, \infty))$ denote a bdf of the respective face, which we can take to be equal near $\{x = 0\}$ to

$$\varrho_{0,\text{zf}_{00}} = \frac{\lambda}{\lambda + x}, \quad \varrho_{0,\text{tf}_{00}} = \lambda + Zx, \quad \varrho_{0,\text{bf}_{00}} = \frac{x}{\lambda + Zx}, \quad \varrho_{0,\text{df}_{00}} = (1 + \xi_b^2 + \eta_b^2)^{-1/2} \quad (77)$$

(defined initially in the interior of ${}^{b,\text{sp}}\overline{T}^*X$, these then extending to smooth functions on ${}^{b,\text{sp}}\overline{T}^*X$). (Below, we conflate smooth functions on mwc with their restrictions to the interior when such a conflation does not cause trouble.) In eq. (77) and below, we write ξ_b for the b-cofiber coordinate dual to x , and $\eta_b^2 = g_{\partial X}^{-1}(\eta_b, \eta_b)$ for $\eta_b \in T^*\partial X$.

There exists a unique mwc ${}^{b,\text{leC}}\overline{T}^*X = [[0, \infty)_\lambda \times {}^b\overline{T}^*X; \{0\} \times {}^b\overline{T}^*_{\partial X}X; \frac{1}{2}] = [{}^{b,\text{sp}}\overline{T}^*X; \text{tf}_{00}; \frac{1}{2}]$, depicted in Figure 2, with the following properties:

- (1) as a set, ${}^{b,\text{leC}}\overline{T}^*X$ is equal to ${}^{b,\text{sp}}\overline{T}^*X$ (a convenient convention),
- (2) ${}^{b,\text{leC}}\overline{T}^*X$ has the same smooth structure as ${}^{b,\text{sp}}\overline{T}^*X$ away from tf_{00} , so that if $\varphi \in C^\infty({}^{b,\text{sp}}\overline{T}^*X)$ is supported away from tf_{00} , then $\varphi \in C^\infty({}^{b,\text{leC}}\overline{T}^*X)$; moreover,
- (3) ${}^{b,\text{leC}}\overline{T}^*X$ has four faces, equal as sets to $\text{zf}_{00}, \text{tf}_{00}, \text{bf}_{00}, \text{df}_{00}$, for which $\varrho_{0,\text{zf}_{00}}, \varrho_{0,\text{tf}_{00}}^{1/2}, \varrho_{0,\text{bf}_{00}}, \varrho_{0,\text{df}_{00}}$ serve as bdfs (respectively).

We will refer to ${}^{b,\text{leC}}\overline{T}^*X$ as the ${}^{b,\text{leC}}$ -phase space. See Figure 2.

We will refer to the bdfs of ${}^{b,\text{leC}}\overline{T}^*X$ as $\varrho_{\text{zf}_{00}} = \varrho_{0,\text{zf}_{00}}, \varrho_{\text{tf}_{00}} = \varrho_{0,\text{tf}_{00}}^{1/2}, \varrho_{\text{bf}_{00}} = \varrho_{0,\text{bf}_{00}}$, and $\varrho_{\text{df}_{00}} = \varrho_{0,\text{df}_{00}}$. Thus, in terms of $\sigma = \lambda^{1/2}$,

$$\varrho_{\text{tf}_{00}} = \sqrt{\sigma^2 + Zx}, \quad \varrho_{\text{bf}_{00}} = \frac{x}{\sigma^2 + Zx}, \quad \varrho_{\text{zf}_{00}} = \frac{\sigma^2}{\sigma^2 + Zx}. \quad (78)$$

Note that the notions of zeroth order conormality on ${}^{\text{b,leC}}\overline{T}^*X$ and ${}^{\text{b,sp}}\overline{T}^*X$ agree, as does the notion of smoothness at zf_{00}° . For each $m, l, \ell \in \mathbb{R}$, we let

$$S_{\text{b,leC}}^{m,l,\ell}(X) = \mathcal{A}_{\text{loc}}^{-m,-l,-\ell/2,(0,0)}(\text{sp}, {}^{\text{b}}\overline{T}^*X) = \mathcal{A}_{\text{loc}}^{-m,-l,-\ell,(0,0)}({}^{\text{b,leC}}\overline{T}^*X), \quad (79)$$

$$S_{\text{b,leC}}(X) = \cup_{m,l,\ell \in \mathbb{R}} S_{\text{b,leC}}^{m,l,\ell}(X), \quad (80)$$

where we are enumerating the faces of the b,leC -phase space in the order $\text{df}_{00}, \text{bf}_{00}, \text{tf}_{00}, \text{zf}_{00}$. Thus, m is the order at df_{00} , l is the order at bf_{00} , ℓ is the order at tf_{00} , and the order at zf_{00} is just zero (and we have a full Taylor series there, with the terms in the Taylor series elements of $\mathcal{A}^{-m,-\ell/2}(\text{zf}_{00})$). (The ‘loc’ subscript in “ \mathcal{A}_{loc} ” refers to the fact that we do not require L^∞ bounds in the $\sigma \rightarrow \infty$ direction. That is, we only have bounds in compact subsets (which can include boundary points) of the mwc ${}^{\text{b,leC}}\overline{T}^*X$, which is only noncompact because of the $\sigma \rightarrow \infty$ direction.) We also define

$$S_{\text{cl,b,leC}}^{m,l,\ell}(X) = \varrho_{\text{df}_{00}}^{-m} \varrho_{\text{bf}_{00}}^{-l} \varrho_{\text{tf}_{00}}^{-\ell} C^\infty({}^{\text{b,sp}}\overline{T}^*X) \subset S_{\text{b,leC}}^{m,l,\ell}(X). \quad (81)$$

Given $a \in S_{\text{b,leC}}^{m,l,\ell}(X)$, we may restrict a to zf_{00} , giving an element $a(-; 0) \in S_{\text{b}}^{m,\ell/2}(X) = S_{\text{b}}^{m,\ell}(X_{1/2})$.

Note that the elements of $S_{\text{b,leC}}(X)$ are all symbols on $[0, \infty)_\lambda \times {}^{\text{b}}\overline{T}^*X$, i.e. elements of

$$\bigcup_{m,l,\nu \in \mathbb{R}} x^{-l} \lambda^{-\nu} \varrho_{\text{df}_{00}}^{-m} \mathcal{A}^0([0, \infty)_\lambda \times {}^{\text{b}}\overline{T}^*X). \quad (82)$$

Specifically, for all $m, l, \ell \in \mathbb{R}$,

$$S_{\text{b,leC}}^{m,l,\ell}(X) \subset x^{-l_0} \lambda^{-\nu} \varrho_{\text{df}_{00}}^{-m} \mathcal{A}^0([0, \infty)_\lambda \times {}^{\text{b}}\overline{T}^*X) \quad (83)$$

whenever $l_0 \geq l$ and $2\nu + 2l_0 \geq \ell$. Consequently, if $l_0 \geq l, \ell/2$,

$$S_{\text{b,leC}}^{m,l,\ell}(X) \subseteq \mathcal{A}^0([0, \infty)_\lambda; S_{\text{b}}^{m,l_0}(X)). \quad (84)$$

Lemma 2.1. *For any $\ell \in \mathbb{R}$, $(\sigma^2 + \text{Z}x)^{-\ell/2} \in C^0([0, \infty)_\sigma; S_{\text{b}}^{0, \max\{0, \ell/2 + \varepsilon\}}(X))$ for any $\varepsilon > 0$.* ■

Proof. Continuity at $\sigma > 0$ is clear. Let $l = \max\{0, \ell/2 + \varepsilon\}$. It suffices to restrict attention to $x < \bar{x}$, so we have to prove that, for each $k \in \mathbb{N}$, $x^l (x\partial_x)^k (\sigma^2 + \text{Z}x)^{-\ell/2} \rightarrow x^l (x\partial_x)^k (\sigma^2 + \text{Z}x)^{-\ell/2}|_{\sigma=0}$ in $L^\infty[0, \bar{x}]$. We compute

$$x^l (x\partial_x)^k (\sigma^2 + \text{Z}x)^{-\ell/2} = x^l \sum_{j=0}^k c_{j,k} x^j (\sigma^2 + \text{Z}x)^{-\ell/2-j} \quad (85)$$

for some $c_{j,k}(\text{Z}) \in \mathbb{R}$. Observe that $x^{\max\{0, \ell/2\}} x^j (\sigma^2 + \text{Z}x)^{-\ell/2-j} \in L^\infty([0, \bar{x}]_x \times [0, 1]_\sigma)$ for every $j \in \mathbb{N}$. Consequently, if $\ell > 0$, the extra factor of x^ε in eq. (85) in conjunction with the uniform convergence as $\sigma \rightarrow 0^+$ of $x^{j+\ell/2} (\sigma^2 + \text{Z}x)^{-\ell/2-j} \rightarrow \text{Z}^{-\ell/2-j}$ in compact subsets of $(0, \bar{x}]_x$ implies that

$$x^l \sum_{j=0}^k c_{j,k} x^j (\sigma^2 + \text{Z}x)^{-\ell/2-j} \rightarrow x^\varepsilon \sum_{j=0}^k c_{j,k} \text{Z}^{-\ell/2-j} \quad (86)$$

uniformly in all of $[0, \bar{x}]$ as $\sigma \rightarrow 0^+$.

If $\ell \leq 0$, then we can instead use that $x^j (\sigma^2 + \text{Z}x)^{-\ell/2-j} \rightarrow (\text{Z}x)^{-\ell/2} \text{Z}^{-j}$ uniformly in $[0, \bar{x}]$. □

Proposition 2.2. *If $a \in S_{\text{b,leC}}^{m,l,\ell}(X)$ and $l_0 \in \mathbb{R}$ satisfies $l_0 \geq l$ and $l_0 > \ell/2$, then $\{a(-; \sigma)\}_{\sigma \geq 0} \in C^0([0, \infty)_\sigma; S_{\text{b}}^{m,l_0}(X))$.* ■

Proof. We first reduce to the case $m, l, \ell = 0$:

- For any $m, l, \ell \in \mathbb{R}$,

$$x^{-l} (\sigma^2 + \text{Z}x)^{l-\ell/2} \varrho_{\text{df}_{00}}^{-m} \in C^0([0, \infty)_\sigma; S_{\text{b}}^{m,l_1}(X)) \quad (87)$$

if $l_1 \geq l, l_1 > \ell/2$ by the previous lemma.

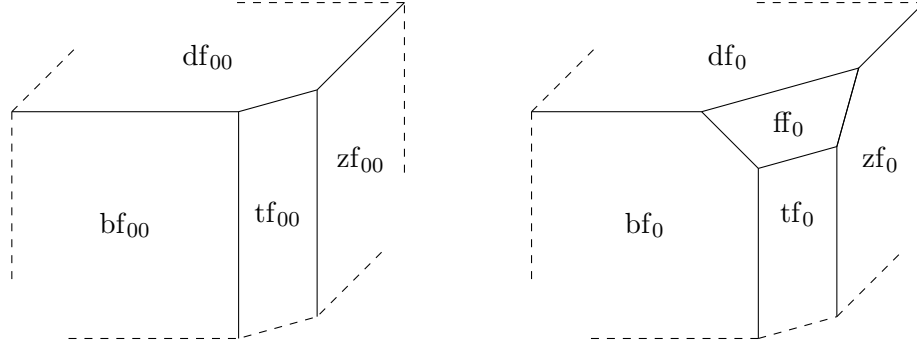


FIGURE 2. The phase spaces ${}^{\text{b,leC}}\overline{T}^*X$ (left, cf. [Vas21c, Fig. 1]) and $[{}^{\text{b,leC}}\overline{T}^*X; \text{df}_{00} \cap \text{tf}_{00}]$ (right), with the degrees of freedom associated with $T^*\partial X$ omitted. (In other words, if we were to consider the case $\dim X = 1$, then the figures above would depict the phases spaces.) For simplicity, we only depict the $\xi_{\text{b}} > 0$ half of phase space.

Any $a \in S_{\text{b,leC}}^{m,l,\ell}(X)$ can be written as $a = x^{-l}(\sigma^2 + \mathbb{Z}x)^{l-\ell/2} \varrho_{\text{df}_{00}}^{-m} a_0$ for $a_0 \in S_{\text{b,leC}}^{0,0,0}(X)$, so if we know that $a_0 \in C^0([0, \infty)_{\sigma}; S_{\text{b}}^{0,\varepsilon}(X))$ for $\varepsilon \geq 0$, then we can conclude that

$$a \in C^0([0, \infty)_{\sigma}; S_{\text{b}}^{m,l_1}(X)) C^0([0, \infty)_{\sigma}; S_{\text{b}}^{0,\varepsilon}(X)) = C^0([0, \infty)_{\sigma}; S_{\text{b}}^{m,l_1+\varepsilon}(X)). \quad (88)$$

Taking $l_1 < l_0$ and $\varepsilon \in (0, l_0 - l_1)$, we get $a \in C^0([0, \infty)_{\sigma}; S_{\text{b}}^{m,l_0}(X))$.

To prove the proposition in the case $m, l, \ell = 0$:

- Continuity at $\sigma > 0$ is clear, so we only need to check that, for $a \in S_{\text{b,leC}}^{0,0,0}(X)$, $a(-; \sigma) \rightarrow a(-; 0)$ in $S_{\text{b}}^{0,\varepsilon}(X)$ for every $\varepsilon > 0$. Note that $La \in L^\infty([0, 1]_{\sigma} \times [0, \bar{x}]_x \times \partial X)$ for any $L \in \text{Diff}_{\text{b}}(X)$. Since, as $\sigma \rightarrow 0^+$, $La \rightarrow La|_{\sigma=0}$ uniformly in compact subsets of $(0, \bar{x}]_x \times \partial X$, we can conclude that

$$x^\varepsilon La \rightarrow x^\varepsilon La|_{\sigma=0} \quad (89)$$

uniformly in all of $[0, \bar{x}] \times \partial X$. Thus, $a(-; \sigma) \rightarrow a(-; 0)$ in $S_{\text{b}}^{0,\varepsilon}(X)$.

□

We now introduce the full leC- phase space ${}^{\text{leC}}\overline{T}^*X$. This is the mwc gotten from ${}^{\text{b,leC}}\overline{T}^*X$ by first blowing up the edge $\text{df}_{00} \cap \text{tf}_{00}$, resulting in a mwc with five faces – $\text{df}_0, \text{ff}_0, \text{bf}_0, \text{tf}_0, \text{zf}_0$ (Figure 2, right), where ff_0 is the front face of the blow up – with bdfs

$$\varrho_{\text{df}_0} = \frac{\varrho_{\text{df}_{00}}}{\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}}}, \quad \varrho_{\text{ff}_0} = \varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}}, \quad \varrho_{\text{tf}_0} = \frac{\varrho_{\text{tf}_{00}}}{\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}}}, \quad (90)$$

$\varrho_{\text{zf}_0} = \varrho_{\text{zf}_{00}}$, $\varrho_{\text{bf}_0} = \varrho_{\text{bf}_{00}}$, and then performing a polar blowup of the edge $\text{df}_0 \cap \text{bf}_0$, resulting in a mwc with six faces – $\text{df}, \text{sf}, \text{ff}, \text{bf}, \text{tf}, \text{zf}$, where sf is front face of the second blowup – with bdfs

$$\begin{aligned} \varrho_{\text{df}} &= \frac{\varrho_{\text{df}_0}}{\varrho_{\text{df}_0} + \varrho_{\text{bf}_0}} = \frac{\varrho_{\text{df}_{00}}}{\varrho_{\text{df}_{00}} + \varrho_{\text{bf}_{00}}(\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}})}, \\ \varrho_{\text{sf}} &= \varrho_{\text{df}_0} + \varrho_{\text{bf}_0} = \frac{\varrho_{\text{df}_{00}}}{\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}}} + \varrho_{\text{bf}_{00}}, \\ \varrho_{\text{bf}} &= \frac{\varrho_{\text{bf}_0}}{\varrho_{\text{df}_0} + \varrho_{\text{bf}_0}} = \frac{\varrho_{\text{bf}_{00}}(\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}})}{\varrho_{\text{df}_{00}} + \varrho_{\text{bf}_{00}}(\varrho_{\text{df}_{00}} + \varrho_{\text{tf}_{00}})}, \end{aligned} \quad (91)$$

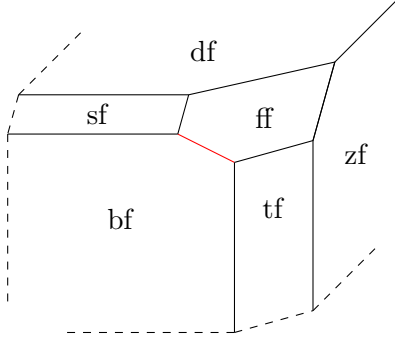


FIGURE 3. The leC-phase space ${}^{\text{leC}}\overline{T}^*X$, with the degrees of freedom associated with $T^*\partial X$ omitted. The edge $\text{ff} \cap \text{bf}$ has been highlighted in red.

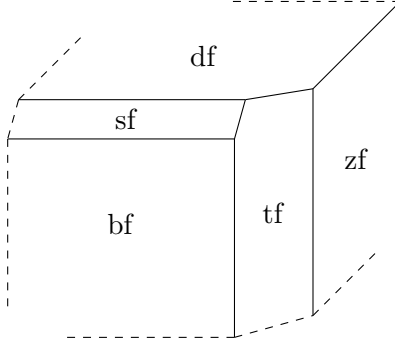


FIGURE 4. The phase space of Vasy's resolved calculus. Cf. [Vas21c, Figure 4] (which also depicts the $\xi < 0$ half of this phase space).

$\varrho_{\text{tf}} = \varrho_{\text{tf}0}$, $\varrho_{\text{ff}} = \varrho_{\text{ff}0}$, $\varrho_{\text{zf}} = \varrho_{\text{zf}0}$. See Figure 3. Now let, for each $m, s, \varsigma, l, \ell \in \mathbb{R}$,

$$\begin{aligned} S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) &= \mathcal{A}^{-m,-s,-\varsigma,-l,-\ell,(0,0)}({}^{\text{leC}}\overline{T}^*X) = \varrho_{\text{df}}^{-m} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell} S_{\text{b,leC}}^{0,0,0}(X), \\ S_{\text{leC}}(X) &= \cup_{m,s,\varsigma,l,\ell \in \mathbb{R}} S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X), \\ S_{\text{cl,leC}}^{m,s,\varsigma,l,\ell}(X) &= \varrho_{\text{df}}^{-m} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell} C^\infty({}^{\text{leC}}\overline{T}^*X) \subset S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X), \\ S_{\text{cl,leC}}(X) &= \cup_{m,s,\varsigma,l,\ell \in \mathbb{R}} S_{\text{cl,leC}}^{m,s,\varsigma,l,\ell}(X). \end{aligned} \tag{92}$$

At the level of sets (and at the level of \mathbb{C} -algebras), $S_{\text{leC}}(X)$ is equal to $S_{\text{b,leC}}(X) = \cup_{m,l,\ell \in \mathbb{R}} S_{\text{b,leC}}^{m,l,\ell}(X)$, but the filtration above presents $S_{\text{leC}}(X)$ as a multigraded \mathbb{C} -algebra. The isomorphism

$$\times \varrho_{\text{df}}^{-m} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell} : S_{\text{b,leC}}^{0,0,0}(X) \rightarrow S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) \tag{93}$$

of vector spaces allows us to consider each $S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ as a Fréchet space, and likewise for $S_{\text{cl,leC}}^{m,s,\varsigma,l,\ell}(X)$. The \mathbb{C} -algebras $S_{\text{leC}}(X)$ and $S_{\text{cl,leC}}(X)$ are then multigraded Fréchet algebras, as pointwise multiplication of symbols is jointly continuous with respect to the relevant topologies.

Observe that if $a \in S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ then $a(-; 0) \in S_{\text{scb}}^{m,\varsigma,\ell}(X_{1/2})$. For $\sigma > 0$, $a(-; \sigma) \in S_{\text{scb}}^{m,s,l}(X)$.

2.2. Calculi. After recalling some preliminary notions in §2.2.1, we discuss the b,leC-calculus in §2.2.2 and the full leC-calculus in §2.2.3. Since the b,leC-calculus is essentially $\Psi_{\text{b,sp,res}}(X)$, except that we enforce classicality at zf_{00} (which introduces no complications), and since all references to the b,leC-calculus in the rest of the paper (i.e. §3, §5, §6) could be replaced by references to the leC-calculus with only notational complications, we will only sketch the arguments in §2.2.2. (However, it is important to understand the residual operators $\Psi_{\text{b,leC}}^{-\infty,l,\ell}(X)$, as these are the “non-symbolic”

parts of leC-operators. But – once again – this is essentially $\Psi_{b,sp,res}^{-\infty,l,\ell}(X)$ except for additional classicality at the “zero face” of Vasy’s double space.)

2.2.1. $\Psi_{b,sp,res}(X)$. We now recall the notion of the b-calculus with conormal dependence on parameters: for any (compact) mwb X and a (connected) mwc M (the “parameter space”), we have a multigraded \mathbb{C} -algebra

$$\Psi_{b;M}(X) = \bigcup_{m,l \in \mathbb{R}} \Psi_{b;M}^{m,l}(X), \quad (94)$$

$$\Psi_{b;M}^{m,l}(X) = \mathcal{A}_{loc}^0(M; \Psi_b^{m,l}(X)), \quad (95)$$

the members of which are the families $\{A_\lambda\}_{\lambda \in M^\circ}$ of b- Ψ DOs on X depending conormally on a parameter $\lambda \in M^\circ$ (as above, the ‘loc’ refers to the fact that we only require uniform bounds in compact subsets of M). So, letting $C^\infty \cap L_{loc}^\infty(M; \Psi_b^{m,l}(X))$ denote the Fréchet space of smooth maps $M^\circ \rightarrow \Psi_b^{m,l}(X)$ that are (but whose derivatives are not necessarily) uniformly bounded with respect to each Fréchet seminorm of $\Psi_b^{m,l}(X)$ in every compact subset of M ,

$$\Psi_{b;M}^{m,l}(X) = \{A_\bullet \in C^\infty \cap L_{loc}^\infty(M; \Psi_b^{m,l}(X)) : [LA_\bullet]_{a.e.} \in L_{loc}^\infty(M; \Psi_b^{m,l}(X)) \forall L \in \text{Diff}_b(M)\}. \quad (96)$$

Note that each $\Psi_{b;M}^{m,l}(X)$ is a Fréchet space, and $\Psi_{b;M}(X)$ can be regarded as a multigraded Fréchet algebra. Relevant to the study of spectral families of operators is the case $M = [0, \infty)_\lambda$. In this case, we write “,sp” in place of “; M ” in the notation.

In [Vas21c, §3], Vasy (using slightly different notation) defines a particular “refinement” of $\Psi_{b,sp}(X) = \Psi_{b;M}(X)$,

$$\Psi_{b,sp,res}(X) = \Psi_{b,sp}(X), \quad \Psi_{b,sp,res}(X) = \bigcup_{m,l,\ell \in \mathbb{R}} \Psi_{b,sp,res}^{m,l,\ell}(X), \quad (97)$$

a multigraded \mathbb{C} -algebra which is equal, at the level of \mathbb{C} -algebras, to $\Psi_{b,sp}(X)$, but with a 3-parameter multigrading (and associated symbol calculus) such that

- $\Psi_{b,sp,res}^{m,l,\ell}(X) \subset \Psi_{b,sp}^{m,\max\{l,\ell\}}(X)$
- and $(\lambda + x)^{-l}$, considered as a multiplication operator, is in $\Psi_{b,sp,res}^{0,0,l}(X)$.

The three indices m, l, ℓ in $\Psi_{b,sp,res}^{m,l,\ell}(X)$ keep track of three notions of order, roughly the “differential order” m , the b-decay order away from zero energy l – that is at bf_{00} – and the b-decay order ℓ at tf_{00} (with respect to $\varrho_{0,\text{tf}_{00}}$). In [Vas21c], Vasy keeps track of an additional order, the order at zf_{00} , which for our purposes can be set to zero throughout (as in eq. (79)).

Define $\sigma_{b,sp,res}^{m,l,\ell} : \Psi_{b,sp,res}^{m,l,\ell}(X) \rightarrow \mathcal{A}_{loc}^{-m,-l,-\ell,0}(\text{b}, \text{sp} \overline{T}^* X) / \mathcal{A}_{loc}^{-m+1,-l,-\ell,0}(\text{b}, \text{sp} \overline{T}^* X)$ by

$$\sigma_{b,sp,res}^{m,l,\ell}(a) = (\lambda + x)^{l-\ell} \sigma_{b,sp}^{m,l}((\lambda + x)^{\ell-l} a) \quad (98)$$

for $a \in S_{b,sp,res}^{m,l,\ell}(X)$, where $\sigma_{b,sp}(a)$ denotes the b-principal symbol map applied λ -wise to a considered as an element of the family b-algebra $\Psi_{b,sp}(X)$. Then, for all $m, l, \ell \in \mathbb{R}$,

$$0 \rightarrow \Psi_{b,sp,res}^{m-1,l,\ell}(X) \hookrightarrow \Psi_{b,sp,res}^{m,l,\ell}(X) \xrightarrow{\sigma_{b,sp,res}} \mathcal{A}_{loc}^{-m,-l,-\ell,0}(\text{sp}, \text{b} \overline{T}^* X) / \mathcal{A}_{loc}^{-m+1,-l,-\ell,0}(\text{sp}, \text{b} \overline{T}^* X) \rightarrow 0 \quad (99)$$

is a short exact sequence and, for all $m', l', \ell' \in \mathbb{R}$,

$$\sigma_{b,sp,res}^{m,l,\ell}(A) \sigma_{b,sp,res}^{m',l',\ell'}(B) = \sigma_{b,sp,res}^{m+m',l+l',\ell+\ell'}(AB) \quad (100)$$

$$\{\sigma_{b,sp,res}^{m,l,\ell}(A), \sigma_{b,sp,res}^{m',l',\ell'}(B)\} = -i \sigma_{b,sp,res}^{m+m'-1,l+l',\ell+\ell'}([A, B]). \quad (101)$$

for all $A \in \Psi_{b,sp,res}^{m,l,\ell}$, $B \in \Psi_{b,sp,res}^{m',l',\ell'}$. We will compute Poisson brackets using the convention that momentum derivatives of the first entry have positive sign. (The sign in eq. (101) depends on the choice of sign used in the Fourier transform used in defining the calculus.) The bdfs $(\sigma^2 + \text{Z}x)^{1/2}$

and $x/(\sigma^2 + Zx)$ of $X_{\text{res}}^{\text{sp}}$, considered as multiplication operators, are representatives of their own principal symbols:

$$\sigma_{\text{b,sp,res}}^{0,l,0}(x^{-l}(\sigma^2 + Zx)^l) = x^{-l}(\sigma^2 + Zx)^l \mod S_{\text{b,sp,res}}^{-1,l,0}(X) \quad (102)$$

$$\sigma_{\text{b,sp,res}}^{0,0,\ell}((\sigma^2 + Zx)^{-\ell}) = (\sigma^2 + Zx)^{-\ell} \mod S_{\text{b,sp,res}}^{-1,0,\ell}(X). \quad (103)$$

More generally, if $a \in S_{\text{b,sp,res}}^{m,l,\ell}(X)$, then $a \in \sigma_{\text{b,sp,res}}^{m,l,\ell}(a)$.

It is very convenient to make use of a “(left) quantization” map (right quantization working equally well):

$$\text{Op} : S_{\text{b}}(X) \rightarrow \Psi_{\text{b}}(X) \quad (104)$$

$$: S_{\text{b}}^{m,l}(X) \rightarrow \Psi_{\text{b}}^{m,l}(X), \quad (105)$$

discussed e.g. in [Vas18] among other places, given by the left quantization of symbols in local coordinates. This will be noncanonical, depending on a choice of atlas on X , among other things. While not surjective (missing out on the remainder term R' in [Vas18, Definition 5.15]), it will be modulo $\Psi_{\text{b}}^{-\infty,l}(X)$.

Applied λ -wise to an element of $\mathcal{A}^0([0, \infty)_{\lambda}; S_{\text{b}}^{m,l}(X))$, the result is an element of $\Psi_{\text{b,sp}}^{m,l}(X) = \mathcal{A}^0([0, \infty)_{\lambda}; \Psi_{\text{b}}^{m,l}(X))$. Some elementary properties of Op which we can arrange are:

- for any $f \in \cup_{l \in \mathbb{R}} \mathcal{A}^l(X)$,

$$\text{Op}(f) = f \quad (106)$$

(this property distinguishing left quantization from right), where the f on the right-hand side denotes the multiplication operator $u \mapsto fu$,

- Op is \mathbb{C} -linear,
- eq. (105) is continuous for any $m, l \in \mathbb{R}$,
- $\sigma_{\text{b}}^{m,l}(\text{Op}(a)) = a \mod S_{\text{b}}^{m-1,l}(X)$ for all $a \in S_{\text{b}}^{m,l}(X)$.

(Equation (106) holds for the Kohn–Nirenberg calculus Ψ_{∞} under left quantization. Since the Schwartz kernel of the multiplication operator $u \mapsto fu$ is supported on the diagonal, it is unaltered by the cutoff $\psi(t - t')$ in [Vas18, Definition 5.15]. As a consequence, eq. (106) holds also for Ψ_{b} .)

2.2.2. $\Psi_{\text{b,leC}}(X)$. Let $\Psi_{\text{b,leC}}^{-\infty,l,\ell}(X)$ denote the elements of $\Psi_{\text{b,sp,res}}^{-\infty,l,\ell/2}(X)$ whose Schwartz kernels are smooth at the face $\text{cl}\{\sigma = 0, x' > 0, x > 0\}$ of the double space $X^{2\text{b,sp,res}}$ [Vas21c, Figure 2] (with the terms in the Taylor series being elements of $\Psi_{\text{b}}^{-\infty,\ell}(X_{1/2})$):

$$\text{SK } \Psi_{\text{b,leC}}^{-\infty,l,\ell}(X) = \mathcal{A}_{\text{loc}}^{-l,-\ell/2,-\infty,-\infty,(0,0)}(X^{2\text{b,sp,res}}), \quad (107)$$

where we are listing the boundary faces of $X^{2\text{b,sp,res}}$ in the order $\text{bf} = \text{cl}\{x, x' = 0, \sigma > 0\}$, $\text{tf} = \{x, x' = 0, \sigma = 0\}$, $\text{lb} = \text{cl}\{x = 0, x' > 0, \sigma > 0\}$, $\text{rb} = \text{cl}\{x' = 0, x > 0, \sigma > 0\}$, and $\text{zf} = \text{cl}\{\sigma = 0, x > 0, x' > 0\}$. Thus, $\Psi_{\text{b,leC}}^{-\infty,l,\ell}(X)$ inherits from eq. (107) a Fréchet space structure, and it can be shown that

$$\Psi_{\text{b,leC}}^{-\infty,\infty,\infty}(X) = \cup_{l,\ell \in \mathbb{R}} \Psi_{\text{b,leC}}^{-\infty,l,\ell}(X) \quad (108)$$

is then a multigraded Fréchet algebra. (So operator composition defines a jointly continuous map $\Psi_{\text{b,leC}}^{-\infty,l,\ell}(X) \times \Psi_{\text{b,leC}}^{-\infty,l',\ell'}(X) \rightarrow \Psi_{\text{b,leC}}^{-\infty,l+l',\ell+\ell'}(X)$ for all $l, \ell, l', \ell' \in \mathbb{R}$.)

An argument similar to that used to prove Proposition 2.2 yields:

Proposition 2.3. *Given an element $K \in \text{SK } \Psi_{\text{b,leC}}^{-\infty,l,\ell}(X)$, if $l_0 \geq l$ and $l_0 > \ell/2$, then it is the case that $\{K(-; \sigma)\}_{\sigma \geq 0} \in C^0([0, \infty)_{\sigma}; \text{SK } \Psi_{\text{b}}^{m,l_0}(X))$ for any $m \in \mathbb{R}$. \blacksquare*

Consequently, elements of $\Psi_{b,\text{leC}}^{-\infty,\infty,\infty}(X)$ can be considered as continuous families of b- Ψ DOs indexed either by \mathbb{R}_σ^+ or by $[0, \infty)_\sigma$. We now define, for each $m, l, \ell \in \mathbb{R}$,

$$\Psi_{b,\text{leC}}^{m,l,\ell}(X) = \text{Op}(S_{b,\text{leC}}^{m,l,\ell}(X)) + \Psi_{b,\text{leC}}^{-\infty,l,\ell}(X). \quad (109)$$

Thus, by eq. (84), if $l_0 \geq l, \ell/2$, then $\Psi_{b,\text{leC}}^{m,l,\ell}(X) \subseteq \Psi_{b,\text{sp}}^{m,l_0}(X)$. In addition:

Proposition 2.4. *If $a = \{a(-; \sigma)\}_{\sigma \geq 0} \in S_{b,\text{leC}}^{m,l,\ell}(X)$, then, if $l_0 \geq l$ and $l_0 > \ell/2$,*

$$\{\text{Op}(a(-; \sigma))\}_{\sigma \geq 0} \in C^0([0, \infty)_\sigma; \Psi_b^{m,l_0}). \quad (110)$$

Consequently, if $A = \{A(\sigma)\}_{\sigma > 0} \in \Psi_{b,\text{leC}}^{m,l,\ell}(X)$, then there exists some $A(0) \in \Psi_b^{m,l_0}(X)$ such that $\{A(\sigma)\}_{\sigma \geq 0} \in C^0([0, \infty)_\sigma; \Psi_b^{m,l_0})$. ■

Proof. Using the continuity of $\text{Op} : S_b^{m,l_0}(X) \rightarrow \Psi_b^{m,l_0}(X)$, the first statement follows from Proposition 2.2. The second statement follows from the first in conjunction with Proposition 2.3. □

Since Op is linear, $\Psi_{b,\text{leC}}^{m,l,\ell}(X)$ is a vector space, and it inherits a Fréchet space structure from $S_{b,\text{leC}}^{m,l,\ell}(X)$ and $\Psi_{b,\text{leC}}^{-\infty,l,\ell}(X)$, more specifically the quotient topology associated to the definitional surjection

$$S_{b,\text{leC}}^{m,l,\ell}(X) \times \Psi_{b,\text{leC}}^{-\infty,l,\ell}(X) \rightarrow \Psi_{b,\text{leC}}^{m,l,\ell}(X). \quad (111)$$

From the definition of $\Psi_{b,\text{sp, res}}^{m,l,\ell/2}(X)$ given in [Vas21c, §3], $\Psi_{b,\text{leC}}^{m,l,\ell}(X) \subset \Psi_{b,\text{sp, res}}^{m,l,\ell/2}(X)$.

Just as the set of classical Kohn–Nirenberg Ψ DOs is a subalgebra of the calculus of all Kohn–Nirenberg Ψ DOs,

$$\Psi_{b,\text{leC}}(X) = \bigcup_{m,l,\ell \in \mathbb{R}} \Psi_{b,\text{leC}}^{m,l,\ell}(X) \quad (112)$$

is a subalgebra of $\Psi_{b,\text{sp, res}}(X)$, with composition of Ψ DOs defining jointly continuous products

$$\Psi_{b,\text{leC}}^{m,l,\ell}(X) \times \Psi_{b,\text{leC}}^{m',l',\ell'}(X) \rightarrow \Psi_{b,\text{leC}}^{m+m',l+l',\ell+\ell'}(X) \quad (113)$$

for all $m, l, \ell, m', l', \ell' \in \mathbb{R}$. The key observation here, in addition to the continuity of $\Psi_{b,\text{leC}}^{m,l,\ell}(X) \times \Psi_{b,\text{leC}}^{-\infty,l',\ell'}(X) \rightarrow \Psi_{b,\text{leC}}^{-\infty,l+l',\ell+\ell'}(X)$ for all $m, l, \ell, l', \ell' \in \mathbb{R}$, is that the reduction formula for full symbols in local coordinates respects classicality at zf_{00} .

From $\sigma_{b,\text{sp, res}}$, we get a set $\{\sigma_{b,\text{leC}}^{m,l,\ell}\}_{m,l,\ell \in \mathbb{R}}$ of maps $\sigma_{b,\text{leC}}^{m,l,\ell} : \Psi_{b,\text{leC}}^{m,l,\ell} \rightarrow S_{b,\text{leC}}^{m,l,\ell}(X)/S_{b,\text{leC}}^{m-1,l,\ell}(X)$ such that

$$0 \rightarrow \Psi_{b,\text{leC}}^{m-1,l,\ell}(X) \hookrightarrow \Psi_{b,\text{leC}}^{m,l,\ell}(X) \xrightarrow{\sigma_{b,\text{leC}}} S_{b,\text{leC}}^{m,l,\ell}(X)/S_{b,\text{leC}}^{m-1,l,\ell}(X) \rightarrow 0 \quad (114)$$

is a short exact sequence and such that

$$\sigma_{b,\text{leC}}^{m,l,\ell}(A) \sigma_{b,\text{sp, res}}^{m',l',\ell'}(B) = \sigma_{b,\text{leC}}^{m+m',l+l',\ell+\ell'}(AB) \quad (115)$$

$$\{\sigma_{b,\text{leC}}^{m,l,\ell}(A), \sigma_{b,\text{leC}}^{m',l',\ell'}(B)\} = -i \sigma_{b,\text{leC}}^{m+m'-1,l+l',\ell+\ell'}([A, B]). \quad (116)$$

for all $A \in \Psi_{b,\text{leC}}^{m,l,\ell}, B \in \Psi_{b,\text{leC}}^{m',l',\ell'}$, with each of eq. (114), eq. (115), eq. (116) following from each of eq. (99), eq. (100), eq. (101) respectively.

Let $\mathcal{A}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}})$ denote the set of distributions on $X_{\text{res}}^{\text{sp}}$ which are conormal to all boundaries and smooth at zf (in particular smooth everywhere except possibly at tf, bf).

Proposition 2.5. *For any $f \in \mathcal{A}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}})$ and $l, \ell \in \mathbb{R}$, the multiplication operator given by multiplication by $x^{-l}(\sigma^2 + Zx)^{l-\ell/2}f(x; \sigma)$ defines an element of $\Psi_{b,\text{leC}}^{0,l,\ell}(X)$. ■*

Proof. The given multiplication operator $M = \{M(\sigma)\}_{\sigma > 0}$ is given by $M(\sigma) = \text{Op}(x^{-l}(\sigma^2 + Zx)^{l-\ell/2}f(x; \sigma))$ (using eq. (106) for each individual $\sigma > 0$), so the proposition follows from $f \in S_{b,\text{leC}}^{0,0,0}(X)$. □

2.2.3. $\Psi_{\text{leC}}(X)$. In order to define the full leC-calculus, we will use the following properties of Op :

- $\text{Op}(a) \in \Psi_{\text{b,leC}}^{m,l,\ell}(X)$ (if and) only if $a \in S_{\text{b,leC}}^{m,l,\ell}(X)$.
-

$$\sigma_{\text{b,leC}}^{m,l,\ell}(\text{Op}(a)) = a \bmod S_{\text{b,leC}}^{m-1,l,\ell}(X) \quad (117)$$

whenever $\text{Op}(a) \in \Psi_{\text{b,leC}}^{m,l,\ell}(X)$,

- there exists a function $\sharp : S_{\text{b,leC}}(X)^2 \rightarrow S_{\text{b,leC}}(X)$ (given by the “reduction formula” for Ψ_∞ , related to Ψ_{b} via [Vas18, §6]) such that, for any $m, m', s, s', \varsigma, \varsigma', l, l', \ell, \ell' \in \mathbb{R}$, $a \in S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ and $b \in S_{\text{leC}}^{m',s',\varsigma',l',\ell'}(X)$,

$$a\sharp b \in S_{\text{leC}}^{m+m',s+s',\varsigma+\varsigma',l+l',\ell+\ell'}(X), \quad (118)$$

$$\text{esssupp}_{\text{leC}}(a\sharp b) \subseteq \text{esssupp}_{\text{leC}}(a) \cap \text{esssupp}_{\text{leC}}(b), \quad (119)$$

$$\text{Op}(a)\text{Op}(b) = \text{Op}(a\sharp b) + E, \quad (120)$$

for some $E \in \Psi_{\text{b,leC}}^{-\infty,l+l',\ell+\ell'}(X)$ which depends continuously on a, b . Moreover,

$$a\sharp b = ab \bmod S_{\text{leC}}^{m+m'-1,s+s'-1,\varsigma+\varsigma'-1,l+l',\ell+\ell'}(X), \quad (121)$$

$$a\sharp b - b\sharp a = i\{a, b\} \bmod S_{\text{leC}}^{m+m'-2,s+s'-2,\varsigma+\varsigma'-2,l+l',\ell+\ell'}(X) \quad (122)$$

for all such a, b ,

- there exists another continuous (\mathbb{C} -antilinear) function $\flat : S_{\text{b,leC}}(X) \rightarrow S_{\text{b,leC}}(X)$ (which can also be written in local coordinates in terms of the reduction formula) such that for all $a \in S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$,

$$\text{Op}(\flat a) = \text{Op}(a)^* + E \quad (123)$$

for some $E \in \Psi_{\text{b,leC}}^{-\infty,l,\ell}(X)$ which depends continuously on a , where the asterisk denotes an $L^2(X, g_0)$ -based adjoint (for arbitrary exactly conic g_0), and

$$\flat a = a^* \bmod S_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X) \quad (124)$$

$$\text{esssupp}_{\text{leC}}(\flat a) = \text{esssupp}_{\text{leC}}(a). \quad (125)$$

Here, for $s \in S_{\text{b,leC}}(X)$, $\text{esssupp}_{\text{leC}}(s)$ consists of those points in $\text{df} \cup \text{sf} \cup \text{ff}$ failing to possess a neighborhood in which s vanishes to infinite order at the boundary of the leC- phase space.

As in [Vas21a][Vas21c], these properties follow from the relation between $\Psi_{\text{b}}(X)$ and $\Psi_\infty(\mathbb{R}^n)$, as explained in [Vas18, §6], and the basic properties of $\Psi_\infty(\mathbb{R}^n)$ (in particular the reduction formula), for which the standard reference is [H07]. It is crucial for us that \sharp satisfies the equations eq. (119), eq. (121), eq. (122) above and not just the weaker b,leC- analogues. This fundamental fact can be read off of the reduction formula for full symbols in local coordinates, in terms of which \sharp can be written.

We can now define, for each $m, s, \varsigma, l, \ell \in \mathbb{R}$,

$$\Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) = \text{Op}(S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)) + \Psi_{\text{b,leC}}^{-\infty,l,\ell} \subseteq \Psi_{\text{b,leC}}^{\infty,l,\ell}. \quad (126)$$

Evidently, eq. (126) endows $\Psi_{\text{b,leC}}^{m,s,\varsigma,l,\ell}(X)$ with a topology, so that it becomes a Fréchet space. Consider the graded vector space

$$\Psi_{\text{leC}}(X) = \bigcup_{m,s,\varsigma,l,\ell \in \mathbb{R}} \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X). \quad (127)$$

At the level of vector spaces, this is just $\Psi_{\text{b,leC}}(X)$.

Moreover, since $S_{\text{leC}}^{m,m+l,m+\ell,l,\ell}(X) = S_{\text{b,leC}}^{m,l,\ell}(X)$ at the level of sets, $\Psi_{\text{leC}}^{m,m+l,m+\ell,l,\ell}(X) = \Psi_{\text{b,leC}}^{m,l,\ell}(X)$ for all $m, l, \ell \in \mathbb{R}$.

Proposition 2.6. $\Psi_{\text{leC}}(X)$ is a multigraded \mathbb{C} -algebra: for any $m, m', s, s', \varsigma, \varsigma', l, l', \ell, \ell' \in \mathbb{R}$, $A \in \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ and $B \in \Psi_{\text{leC}}^{m', s', \varsigma', l', \ell'}(X)$,

$$AB \in \Psi_{\text{leC}}^{m+m', s+s', \varsigma+\varsigma', l+l', \ell+\ell'}(X). \quad (128)$$

■

Proof. We can write $A = \text{Op}(a) + E$ and $B = \text{Op}(b) + F$ for $a \in S_{\text{leC}}^{m, s, \varsigma, l, \ell}$, $b \in S_{\text{leC}}^{m', s', \varsigma', l', \ell'}$, $E \in \Psi_{\text{b,leC}}^{-\infty, l, \ell}$, $F \in \Psi_{\text{b,leC}}^{-\infty, l', \ell'}$. Thus,

$$AB = \text{Op}(a) \text{Op}(b) + E \text{Op}(b) + \text{Op}(a)F + EF \quad (129)$$

$$= \text{Op}(a \# b) + E \text{Op}(b) + \text{Op}(a)F + EF + G \quad (130)$$

for some $G \in \Psi_{\text{b,leC}}^{-\infty, l+l', \ell+\ell'}$. Since $a \in S_{\text{b,leC}}^{M, l, \ell}$, $b \in S_{\text{b,leC}}^{M', l', \ell'}$ for $M = \max\{m, s - l, \varsigma - \ell\}$ and $M' = \max\{m', s' - l', \varsigma' - \ell'\}$,

$$\text{Op}(a) \in \Psi_{\text{b,leC}}^{M, l, \ell}(X), \quad \text{Op}(b) \in \Psi_{\text{b,leC}}^{M', l', \ell'}(X), \quad (131)$$

which implies that $E \text{Op}(b), \text{Op}(a)F \in \Psi_{\text{b,leC}}^{-\infty, l+l', \ell+\ell'}$, and likewise $EF \in \Psi_{\text{b,leC}}^{-\infty, l+l', \ell+\ell'}$.

Since $a \# b \in S_{\text{leC}}^{m+m', s+s', \varsigma+\varsigma', l+l', \ell+\ell'}$, we deduce that eq. (128) holds. \square

In fact (as can be proven using finite-order truncations of the reduction formula), operator composition defines a jointly continuous map

$$\Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X) \times \Psi_{\text{leC}}^{m', s', \varsigma', l', \ell'}(X) \rightarrow \Psi_{\text{leC}}^{m+m', s+s', \varsigma+\varsigma', l+l', \ell+\ell'}(X), \quad (132)$$

so we can say that the leC-calculus is a multigraded Fréchet algebra.

Lemma 2.7. Suppose that $A \in \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ can be written either as $A = \text{Op}(a_1) + E_1$ or $A = \text{Op}(a_2) + E_2$ for some symbols $a_1, a_2 \in S_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ and $E_1, E_2 \in \Psi_{\text{b,leC}}^{-\infty, l, \ell}$.

Then $a_1 - a_2 \in S_{\text{b,leC}}^{-\infty, l, \ell}(X)$. \blacksquare

Proof. By the linearity of Op , $\text{Op}(a_1 - a_2) \in \Psi_{\text{b,leC}}^{-\infty, l, \ell}(X)$. Then, using the $\sigma_{\text{b,leC}}$ -short exact sequence and the property eq. (117) of Op : for all $N \in \mathbb{N}$,

$$0 = \sigma_{\text{b,leC}}^{-N, l, \ell}(\text{Op}(a_1 - a_2)) = a_1 - a_2 \bmod S_{\text{b,leC}}^{-N-1, l, \ell}(X), \quad (133)$$

which means that $a_1 - a_2 \in S_{\text{b,leC}}^{-N-1, l, \ell}(X)$. \square

Thus, for any $A \in \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$, we get notions

$$\text{Ell}_{\text{leC}}^{m, s, \varsigma, l, \ell}(A) = \text{ell}_{\text{leC}}^{m, s, \varsigma, l, \ell}(a) \subset \text{df} \cup \text{sf} \cup \text{ff}, \quad (134)$$

$$\text{Char}_{\text{leC}}^{m, s, \varsigma, l, \ell}(A) = \text{char}_{\text{leC}}^{m, s, \varsigma, l, \ell}(a) = \text{df} \cup \text{sf} \cup \text{ff} \setminus \text{ell}_{\text{leC}}^{m, s, \varsigma, l, \ell}(a), \quad (135)$$

$$\text{WF}_{\text{leC}}^{l, \ell}(A) = \text{esssupp}_{\text{leC}}(a), \quad (136)$$

for any $a \in S_{\text{leC}}^{m, s, \varsigma, l, \ell}$ with $A = \text{Op}(a) + E$ for $E \in \Psi_{\text{b,leC}}^{-\infty, l, \ell}(X)$. All three are subsets of $\text{df} \cup \text{sf} \cup \text{ff}$. Another useful notion, which we will only apply to $A \in \cap_{l, \ell \in \mathbb{R}} \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X) = \Psi_{\text{leC}}^{m, s, \varsigma, -\infty, -\infty}(X)$, is

$$\text{WF}'_{\text{leC}}(A) = \cap_{l, \ell \in \mathbb{R}} \text{WF}_{\text{leC}}^{l, \ell}(A). \quad (137)$$

Observe that if $a \in S_{\text{leC}}^{m, s, \varsigma, -\infty, -\infty}(X)$, then $\text{Op}(A) \in \Psi_{\text{leC}}^{m, s, \varsigma, -\infty, -\infty}(X)$ (tautologically) and

$$\text{WF}'_{\text{leC}}(A) = \text{esssupp}_{\text{leC}}(a). \quad (138)$$

The $a \in S_{\text{leC}}^{m, s, \varsigma, -\infty, -\infty}(X)$ that we consider are all supported away from $\text{bf} \cup \text{tf}$, in which case $\text{WF}'_{\text{leC}}(A)$ is disjoint from $\text{bf} \cup \text{tf}$.

In addition, we get the leC- principal symbol maps

$$\{\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}\}_{m,s,\varsigma,l,\ell \in \mathbb{R}}, \quad \sigma_{\text{leC}}^{m,s,\varsigma,l,\ell} : \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) \rightarrow S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)/S_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X),$$

$$\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}(A) = a \bmod S_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell} \text{ for any } a \in S_{\text{leC}}^{m,s,\varsigma,l,\ell} \text{ with } A = \text{Op}(a) + E \text{ for } E \in \Psi_{\text{b,leC}}^{-\infty,l,\ell}.$$

Unlike Op, which is not canonical and depends on a particular choice of local coordinate charts, these notions are all canonical.

Proposition 2.8. *For every $m, s, \varsigma, l, \ell \in \mathbb{R}$, we have a short exact sequence*

$$0 \rightarrow \Psi_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X) \hookrightarrow \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) \rightarrow S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)/S_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X) \rightarrow 0, \quad (139)$$

where the second-to-last map is $\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}$. ■

Proof. The surjectivity of $\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}$ follows from the properties of Op listed above.

If, on the other hand, $A \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ satisfies $\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}(A) = 0$, then $A = \text{Op}(a) + E$ for $a \in S_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}$ and $E \in \Psi_{\text{b,leC}}^{-\infty,l,\ell}$. Then, by the definition of $\Psi_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X)$, $A \in \Psi_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}(X)$. □

Proposition 2.9. *For every $m, s, \varsigma, l, \ell, m', s', \varsigma', l', \ell' \in \mathbb{R}$ and pair of $A \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ and $B \in \Psi_{\text{leC}}^{m',s',\varsigma',l',\ell'}(X)$,*

$$\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}(A)\sigma_{\text{leC}}^{m',s',\varsigma',l',\ell'}(B) = \sigma_{\text{leC}}^{m+m',s+s',\varsigma+\varsigma',l+l',\ell+\ell'}(AB). \quad (140)$$

Moreover, $\sigma_{\text{leC}}^{m+m'-1,s+s'-1,\varsigma+\varsigma'-1,l+l',\ell+\ell'}([A, B])$ is equal to

$$i\{a, b\} \bmod S_{\text{leC}}^{m+m'-2,s+s'-2,\varsigma+\varsigma'-2,l+l',\ell+\ell'}(X), \quad (141)$$

where a, b are any representatives of $\sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}(A)$ and $\sigma_{\text{leC}}^{m',s',\varsigma',l',\ell'}(B)$. ■

Proof. Write $A = \text{Op}(a) + E$ and $B = \text{Op}(b) + F$ for $a \in S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$, $b \in S_{\text{leC}}^{m',s',\varsigma',l',\ell'}(X)$, $E \in \Psi_{\text{leC}}^{-\infty,l,\ell}(X)$, $F \in \Psi_{\text{leC}}^{-\infty,l',\ell'}(X)$. Then

$$\begin{aligned} \sigma_{\text{leC}}^{m+m',s+s',\varsigma+\varsigma',l+l',\ell+\ell'}(AB) &= \sigma_{\text{leC}}^{m+m',s+s',\varsigma+\varsigma',l+l',\ell+\ell'}(\text{Op}(a)\text{Op}(b)) \\ &= \sigma_{\text{leC}}^{m+m',s+s',\varsigma+\varsigma',l+l',\ell+\ell'}(\text{Op}(a\sharp b)) \\ &= a\sharp b \bmod S_{\text{leC}}^{m+m'-1,s+s'-1,\varsigma+\varsigma'-1,l+l',\ell+\ell'}(X) \\ &= ab \bmod S_{\text{leC}}^{m+m'-1,s+s'-1,\varsigma+\varsigma'-1,l+l',\ell+\ell'}(X) \\ &= \sigma_{\text{leC}}^{m,s,\varsigma,l,\ell}(A)\sigma_{\text{leC}}^{m',s',\varsigma',l',\ell'}(B). \end{aligned} \quad (142)$$

The proof of eq. (141) is similar. □

Proposition 2.10. *Suppose that $L \in S\text{Diff}_{\text{scb}}^{m,s,l}(X)$, $m \in \mathbb{N}$, $s, l \in \mathbb{R}$. Then, the constant family $\{L(\sigma)\}_{\sigma>0}$, $L(\sigma) = L$, (which we conflate with L) defines an element of $\Psi_{\text{leC}}^{m,s,s+l,2l}(X)$.*

Given any $f \in \mathcal{A}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}})$ and $l_0, \ell_0 \in \mathbb{R}$, $(x/(\sigma^2 + \text{Z}x))^{l_0}(\sigma^2 + \text{Z}x)^{\ell_0/2}fL$ defines an element of $\Psi_{\text{leC}}^{m,s-l_0,s+l-\ell_0,l-l_0,2l-\ell_0}(X)$. ■

Proof. First consider the case when $s = m + l$, so that $L \in S\text{Diff}_{\text{b}}^{m,l}(X)$. Then,

$$L \in \Psi_{\text{b,leC}}^{m,l,2l}(X) = \Psi_{\text{leC}}^{m,m+l,m+2l,l,2l}(X) = \Psi_{\text{leC}}^{m,s,s+l,l,2l}(X). \quad (143)$$

To handle the general case, we note that we may write any $L \in S\text{Diff}_{\text{scb}}^{m,s,l}(X)$ – where now $m, s, l \in \mathbb{R}$ are arbitrary – as

$$L = \sum_{j=0}^m L_j \quad (144)$$

for $L_j \in S \text{Diff}_b^{j, \min\{s-j, l\}}(X) \subset S \text{Diff}_{\text{scb}}^{j, \min\{s, l+j\}, \min\{s-j, l\}}(X) \subset S \text{Diff}_{\text{scb}}^{m, s, l}(X)$. (Indeed, it suffices to construct such a decomposition on $[0, \bar{x}]_x \times \partial X$, where L_j can be written as a linear combination of elements of $S(X) \partial_x^{j_1} \text{Diff}^{j_2}(\partial X)$ for $j_1 + j_2 = j$, $j_1, j_2 \in \mathbb{N}$.) Since

$$\begin{aligned} L_j &\in \Psi_{b, \text{leC}}^{j, \min\{s-j, l\}, 2 \min\{s-j, l\}}(X) \\ &= \Psi_{\text{leC}}^{j, \min\{s, l+j\}, \min\{2s-j, j+2l\}, \min\{s-j, l\}, 2 \min\{s-j, l\}}(X), \end{aligned} \quad (145)$$

we deduce that

$$L \in \Psi_{\text{leC}}^{m, \min\{s, l+m\}, s+l, \min\{s, l\}, \min\{2s, 2l\}}(X) \subset \Psi_{\text{leC}}^{m, s, s+l, l, 2l}(X). \quad (146)$$

From Proposition 2.5, for any $f \in \mathcal{A}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}})(X_{\text{res}}^{\text{sp}})$ and $l_0, \ell_0 \in \mathbb{R}$, the multiplication operator $x^{l_0}(\sigma^2 + Zx)^{-l_0+\ell_0/2} f(x; \sigma)$ defines an element of $\Psi_{b, \text{leC}}^{0, -l_0, -\ell_0}(X) = \Psi_{\text{leC}}^{0, -l_0, -\ell_0, -l_0, -\ell_0}(X)$. The second statement of the proposition therefore follows from the first via an application of Proposition 2.6. \square

For each $m, s, \varsigma, l, \ell \in \mathbb{R}$, we let $S \text{Diff}_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ denote the set of elements of $\Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ which are families of differential operators (all of which arise from the construction in Proposition 2.10). Likewise, we let

$$\text{Diff}_{\text{leC}}^{m, s, \varsigma, l, \ell}(X) \subset S \text{Diff}_{\text{leC}}^{m, s, \varsigma, l, \ell}(X) \quad (147)$$

denote the subset of families of differential operators which can be written as a linear combination of elements of $\text{Diff}_b(X)$ times classical symbols on $X_{\text{res}}^{\text{sp}}$.

The elliptic parametrix construction, applied to the leC-calculus, yields the following: for any $m, s, \varsigma, l, \ell \in \mathbb{R}$ and (totally) elliptic $A \in \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ – that is A with

$$\text{Ell}_{\text{leC}}^{m, s, \varsigma, \ell}(A) = \text{df} \cup \text{sf} \cup \text{ff} \quad (148)$$

– there exists, for each $N \in \mathbb{N}$, some $B \in \Psi_{\text{leC}}^{-m, -s, -\varsigma, -l, -\ell}(X)$ such that $AB - 1, BA - 1 \in \Psi_{b, \text{leC}}^{-N, 0, 0}(X)$. (To elaborate, the leC-symbol calculus analogue of the construction of left and right parametrices via Neumann series yields B_L, B_R such that $AB_R - 1, B_L A - 1 \in \Psi_{b, \text{leC}}^{-N, 0, 0}(X)$. Then, setting $E_L = B_L A - 1$ and $E_R = AB_R - 1$,

$$B_L = B_L(AB_R - E_R) = (E_L + 1)B_R - B_L E_R = B_R + (E_L B_R - B_L E_R), \quad (149)$$

so taking either $B = B_L$ or $B = B_R$, both of $AB - 1, BA - 1 \in \Psi_{b, \text{leC}}^{-N, 0, 0}(X)$ hold. Hence we do not need to distinguish left vs. right parametrices.)

Lemma 2.11. *Given any $m, s, \varsigma, l, \ell \in \mathbb{R}$, $A \in \Psi_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)$ and $\alpha \in \text{df} \cup \text{sf} \cup \text{ff}$, $\alpha \notin \text{WF}_{\text{leC}}^{l, \ell}(A)$ only if there exists some $B \in \Psi_{\text{leC}}^{0, 0, 0, 0, 0}(X)$ that is elliptic at α and satisfies $AB, BA \in \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, l, \ell}(X)$.*

The same statement applies if we replace α with the intersection of a finite union of closed balls within $\text{df} \cup \text{sf} \cup \text{ff}$. \blacksquare

Proof. We write $A = \text{Op}(a)$ for $a \in S_{\text{leC}}^{\infty, \infty, \infty, l, \ell}(X)$, so that $\text{WF}_{\text{leC}}^{l, \ell}(A) = \text{esssupp}_{\text{leC}}(a)$.

Given $B \in \Psi_{\text{leC}}^{0, 0, 0, 0, 0}(X)$, $B = \text{Op}(b)$, $b \in S_{\text{leC}}^{0, 0, 0, 0, 0}(X)$, define $E \in \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, l, \ell}(X)$ by

$$AB = \text{Op}(a \sharp b) + E. \quad (150)$$

If $\alpha \notin \text{esssupp}_{\text{leC}}(a)$, then (since essential supports are closed) we can choose $b \in S_{\text{cl, leC}}^{0, 0, 0, 0, 0}(X)$ that is identically equal to one in a neighborhood of α but supported away from $\text{esssupp}_{\text{leC}}(a)$, so that

$$\text{esssupp}_{\text{leC}}(a) \cap \text{esssupp}_{\text{leC}}(b) = \emptyset.$$

Then, by eq. (119), $a \sharp b$ has empty essential support, which implies that $AB \in \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, l, \ell}(X)$.

That handles AB , and BA is analogous. If we replace α in the previous argument with the intersection of a finite union of closed balls within $\text{df} \cup \text{sf} \cup \text{ff}$, the argument goes through the same. \square

We briefly discuss uniform families of leC- Ψ DOs (i.e. one-parameter families of b- Ψ DOs which are uniformly bounded in a sense appropriate for the leC-calculus). For each $m, s, \varsigma, l, \ell \in \mathbb{R} \cup \{-\infty\}$,

$$\Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) = \bigcap_{\substack{m' \geq m, \dots, \ell' \geq \ell \\ m', \dots, \ell' \in \mathbb{R}}} \Psi_{\text{leC}}^{m',s',\varsigma',l',\ell'}(X) \quad (151)$$

is a Fréchet space, so for any nonempty interval $I \subset \mathbb{R}$, we have a Fréchet space $L^\infty(I; \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$ whose elements are a.e. equivalence classes of measurable functions $I \rightarrow \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ which are uniformly bounded with respect to each of our countably many Fréchet seminorms on the codomain. For I closed, we can safely conflate an element of $\mathcal{A}^0(I; \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$, which we can consider as a smooth family $\{A_t\}_{t \in I^\circ}$, with the corresponding element of $L^\infty(I; \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$.

For $A = \{A_t\}_{t \in I} \in L^\infty(I; \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$, we define subsets $\text{WF}'_{L^\infty, \text{leC}}(A)$, $\text{WF}^{n,\ell}_{L^\infty, \text{leC}}(A) \subset \text{df} \cup \text{sf} \cup \text{ff}$ by stipulating that a point $\alpha \in \text{df} \cup \text{sf} \cup \text{ff}$ does not lie in

- $\text{WF}'_{L^\infty, \text{leC}}(A)$ if and only if there exists some $B \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ that is elliptic at α and satisfies $\{BA_t\}_{t \in I} \in L^\infty(I; \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, -\infty, -\infty}(X))$,
- $\text{WF}^{n,\ell}_{L^\infty, \text{leC}}(A)$ if and only if there exists some $B \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ that is elliptic at α and satisfies $\{BA_t\}_{t \in I} \in L^\infty(I; \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, l, \ell}(X))$.

A uniform version of the elliptic parametrix construction goes through. We likewise have a notion of $\text{esssupp}_{L^\infty, \text{leC}}(a)$ for $a \in L^\infty(I; S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$. One definition is that $\alpha \in \text{df} \cup \text{sf} \cup \text{ff}$ is not in $\text{esssupp}_{L^\infty, \text{leC}}(a)$ if and only if there exists some $b \in S_{\text{leC}}^{0,0,0,0,0}(X)$ that is elliptic at α and satisfies $ab \in S_{\text{leC}}^{-\infty, -\infty, -\infty, -\infty, -\infty}(X)$. Quantizing: for any interval $I \subset \mathbb{R}$ and any $m, s, \varsigma, l, \ell \in \mathbb{R}$, given $a \in L^\infty(I; S_{\text{leC}}^{m,s,\varsigma,l,\ell}(X))$, letting $A = \{\text{Op}(A_t)\}_{t \in I}$,

$$\text{WF}^{n,\ell}_{L^\infty, \text{leC}}(A) = \text{esssupp}_{L^\infty, \text{leC}}(a). \quad (152)$$

2.3. Sobolev Spaces. For each $m, s, l \in \mathbb{R}$, let $H_{\text{scb}}^{m,s,l}(X)$ denote the Sobolev space of differential order m , sc-decay order s , and b-decay order l associated to the scb-calculus in [Vas21a]. Associated to the leC-calculus is a five-parameter family

$$\{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)\}_{m,s,\varsigma,l,\ell \in \mathbb{R}}$$

of families

$$H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) = \{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)\}_{\sigma \geq 0} \quad (153)$$

of “leC-based Sobolev spaces,” where

- for each $\sigma > 0$, $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)$ is a Hilbertizable Banach space equal to $H_{\text{scb}}^{m,s,l}(X) \subset \mathcal{S}'(X)$ at the level of TVSs (i.e. equivalent at the level of Banach spaces),
- $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(0) = H_{\text{scb}}^{m,\varsigma+n/2,\ell+n/2}(X_{1/2})$ (at the level of TVSs).

The family $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ had ought to be thought of as interpolating between the Sobolev spaces $H_{\text{scb}}^{m,s,l}(X)$ and $H_{\text{scb}}^{m,\varsigma+n/2,\ell+n/2}(X_{1/2})$ as $\sigma \rightarrow 0^+$.

Besides using the leC-Sobolev spaces to relate $H_{\text{scb}}^{\bullet,\bullet,\bullet}(X)$ and $H_{\text{scb}}^{\bullet,\bullet,\bullet}(X_{1/2})$ we can, more crudely, observe that

$$H_{\text{scb}}^{m,m+l,l}(X) = H_{\text{b}}^{m,l}(X) = H_{\text{b}}^{m,2l+n/2}(X_{1/2}) = H_{\text{scb}}^{m,m+2l+n/2,2l+n/2}(X_{1/2}). \quad (154)$$

The spaces $H_{\text{scb}}^{\bullet,\bullet,\bullet}(X)$ are more refined than $H_{\text{b}}^{\bullet,\bullet}(X)$, in that the decay rate of terms like $\exp(i/x)$ need not be treated the same as the decay rate of the constant function as measured by the former. Similarly, the spaces $H_{\text{scb}}^{\bullet,\bullet,\bullet}(X_{1/2})$ are more refined than $H_{\text{b}}^{\bullet,\bullet}(X_{1/2})$ in that the decay rate of terms like $\exp(i/x^{1/2})$ need not be the same as the decay rate of the constant function as measured by the former. As shown in [Vas21a], the scb-Sobolev spaces are well-suited for formulating a conjugated version of the Sommerfeld radiation condition, and for similar reasons the leC-Sobolev spaces are well-suited to the study of attractive Coulomb-like Schrödinger operators down to zero energy.

In order to define the σ -dependent norm on $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)$, we first note:

Lemma 2.12. *For any two elliptic $A, B \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}$ and $N, M \in \mathbb{Z}$ with $N, M \geq -\min\{m, s-l, \varsigma-\ell\}$, and for any $\Sigma > 0$, there exist $c, C > 0$ such that*

$$c(\|B(\sigma)u\|_{L^2(X)} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)}) \leq \|A(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \leq C(\|B(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)}) \quad (155)$$

holds for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. ■

Proof. It suffices to prove $\|A(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \leq C(\|B(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)})$, the other inequality following by symmetry.

By the elliptic parametrix construction, for arbitrary $N_0 \in \mathbb{N}$, we can find $\Lambda \in \Psi_{\text{leC}}^{-m,-s,-\varsigma,-l,-\ell}$ and $R \in \Psi_{\text{b,leC}}^{-N_0,0,0}$ with $\Lambda B = 1 + R$. Now, setting $m_0 = \min\{m, s-l, \varsigma-\ell\}$,

$$\begin{aligned} \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} &= \|(\Lambda(\sigma)B(\sigma) - R(\sigma))u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \\ &\leq \|\Lambda(\sigma)B(\sigma)u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} + \|R(\sigma)u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \\ &\leq \|\Lambda(\sigma)B(\sigma)u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)} \\ &\leq \|B(\sigma)u\|_{H_{\text{b,leC}}^{-N-m_0,0,0}(\sigma)} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)} \leq \|B(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)} \end{aligned} \quad (156)$$

for sufficiently large N_0 , where we used that $\Lambda \in \Psi_{\text{b,leC}}^{-m_0,-l,-\ell}$.

Similarly, for sufficiently large N_0 ,

$$\begin{aligned} \|A(\sigma)u\|_{L^2} &= \|A(\sigma)(\Lambda(\sigma)B(\sigma) - R(\sigma))u\|_{L^2} \leq \|A(\sigma)\Lambda(\sigma)B(\sigma)u\|_{L^2} + \|A(\sigma)R(\sigma)u\|_{L^2} \\ &\leq \|A(\sigma)\Lambda(\sigma)B(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)} \\ &\leq \|B(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l,\ell}(\sigma)}. \end{aligned} \quad (157)$$

Combining eq. (156) with eq. (157) yields the desired inequality. □

We can now define a Hilbertizable norm on $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)$, for each $\sigma \geq 0$, by writing

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)} = \|\Lambda(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \quad (158)$$

for $N \in \mathbb{N}$ with $N \geq -\min\{m, s-l, \varsigma-\ell\}$ and arbitrary elliptic $\Lambda \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$. (Such Λ can be constructed by the symbol calculus.) The previous lemma suffices to guarantee that the estimates we prove do not depend on the particular choice of Λ and N used in defining the norm eq. (158) except with regards to the particular constants involved (which we do not keep track of anyways). However, it will be convenient to fix

$$\Lambda_{m,s,\varsigma,l,\ell} = (1/2)(\text{Op}(\varrho_{\text{df}}^{-m} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell}) + \text{Op}(\varrho_{\text{df}}^{-m} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell})^*), \quad (159)$$

which is certainly an elliptic element of $\Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$. Since $\Lambda_{m,s,\varsigma,l,\ell}(\sigma)$ is an elliptic element of $\Psi_{\text{scb}}^{m,s,l}(X)$ for $\sigma > 0$ and $\Psi_{\text{scb}}^{m,\varsigma,\ell}(X_{1/2})$ for $\sigma = 0$, we see that $H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)$ is indeed equivalent to

- $H_{\text{scb}}^{m,s,l}(X)$ for $\sigma > 0$ and
- $H_{\text{scb}}^{m,\varsigma+n/2,\ell+n/2}(X_{1/2})$ for $\sigma = 0$. (The shift of $n/2$ orders occurs because a sc-density on X is a weighted sc-density on $X_{1/2}$, and vice versa; see Lemma 2.18.)

On the other hand:

Lemma 2.13. *If $m, s, \varsigma, l, \ell \in \mathbb{R}$ satisfy $\varsigma = m + \ell$ and $s = m + l$, for each $\Sigma > 0$ there exist constants $c(m, s, \varsigma, l, \ell, \Sigma), C(m, s, \varsigma, l, \ell, \Sigma) > 0$ such that*

$$c\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)} \leq \|u\|_{H_{\text{b,leC}}^{m,l,\ell}(X)(\sigma)} \leq C\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)} \quad (160)$$

for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. In particular, if $\ell = 2l$, $c' \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)} \leq \|u\|_{H_{\text{b}}^{m,l}(X)} \leq C' \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)}$ for some other constants $c'(m, s, \varsigma, l, \ell, \Sigma), C'(m, s, \varsigma, l, \ell, \Sigma) > 0$. ■

Proof. We have that $\Lambda_{m,s,\varsigma,l,\ell} \in \Psi_{\text{b,leC}}^{m,l,\ell}(X)$, so

$$\|u\|_{H_{\text{leC}}^{m,m+l,m+l,\ell}(X)(\sigma)} = \|\Lambda_{m,m+l,m+l,\ell}(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}(\sigma)} \preceq \|u\|_{H_{\text{b,leC}}^{m,l,\ell}(X)(\sigma)} \quad (161)$$

for sufficiently large N (by the boundedness of the elements of the resolved family b-calculus).

The reverse inequality follows from the ellipticity of $\Lambda_{m,m+l,m+l,\ell}$ as an element of $\Psi_{\text{b,leC}}^{m,l,\ell}(X)$. □

We now check the boundedness of leC-ΨDOs acting on leC-Sobolev spaces (from the L^2 -boundedness of zeroth order b-ΨDOs):

Proposition 2.14. *For any $m, s, \varsigma, l, \ell, m_0, s_0, \varsigma_0, l_0, \ell_0 \in \mathbb{R}$, $A \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$, and $\Sigma > 0$, there exists some constant $C = C(m, s, \varsigma, l, \ell, m_0, s_0, \varsigma_0, l_0, \ell_0, A, \Sigma) > 0$ such that*

$$\|A(\sigma)u\|_{H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}(X)(\sigma)} \leq C \|u\|_{H_{\text{leC}}^{m+m_0,s+s_0,\varsigma+\varsigma_0,l+l_0,\ell+\ell_0}(X)(\sigma)} \quad (162)$$

for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. ■

Proof. Pick arbitrary elliptic $\Lambda_0 \in \Psi_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}(X)$. Then

$$\|A(\sigma)u\|_{H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}(X)(\sigma)} \preceq \|\Lambda_0(\sigma)A(\sigma)u\|_{L^2} + \|A(\sigma)u\|_{H_{\text{b,leC}}^{-N,l_0,\ell_0}(X)(\sigma)} \quad (163)$$

$$\|u\|_{H_{\text{leC}}^{m+m_0,s+s_0,\varsigma+\varsigma_0,l+l_0,\ell+\ell_0}(X)(\sigma)} \succeq \|\Lambda_1(\sigma)u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l+l_0,\ell+\ell_0}(X)(\sigma)} \quad (164)$$

for N, M not too negative and arbitrary elliptic $\Lambda_1 \in \Psi_{\text{leC}}^{m+m_0,s+s_0,\varsigma+\varsigma_0,l+l_0,\ell+\ell_0}(X)$.

For N sufficiently large compared to M (dependent on m, \dots, ℓ_0), $\|A(\sigma)u\|_{H_{\text{b,leC}}^{-N,l_0,\ell_0}(X)(\sigma)} \preceq \|u\|_{H_{\text{b,leC}}^{-M,l+l_0,\ell+\ell_0}(X)(\sigma)}$.

For each $M' \in \mathbb{R}$, an elliptic parametrix for Λ_1 can be used to construct $A_0 \in \Psi_{\text{b,leC}}^{0,0,0}(X) \subset L^\infty([0, \infty); \Psi_{\text{b}}^0(X))$ and $R \in \Psi_{\text{b,leC}}^{-M',l+l_0,\ell+\ell_0}(X)$ such that $\Lambda_0 A = A_0 \Lambda_1 + R$. Then, if M' is sufficiently large,

$$\|\Lambda_0 A(\sigma)u\|_{L^2} \leq \|A_0 \Lambda_1 u\|_{L^2} + \|Ru\|_{L^2} \preceq \|\Lambda_1 u\|_{L^2} + \|u\|_{H_{\text{b,leC}}^{-M,l+l_0,\ell+\ell_0}(X)(\sigma)}, \quad (165)$$

where we have used the uniform L^2 -boundedness of A_0 . Combining the estimates above, we deduce eq. (162). □

In the rest of the paper, we will abbreviate

$$H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma) = H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) = H_{\text{leC}}^{m,s,\varsigma,l,\ell}, \quad (166)$$

leaving the dependence on σ implicit (and likewise for other σ -dependent notions). In particular, what we call “estimates” will really be 1-parameter families of estimates with multiplicative constants that – for any fixed $\Sigma > 0$ – are uniform for $\sigma \in [0, \Sigma]$.

The following two lemmas will be used mostly without comment in §5:

Lemma 2.15. *For any $m, s, \varsigma, l, \ell \in \mathbb{R}$, there exists a constant $C = C(m, s, \varsigma, l, \ell)$ such that, for all $u, v \in \mathcal{S}'(X)$ and $\bar{\delta} > 0$,*

$$|\langle u, v \rangle_{L^2}| \leq C \cdot (\bar{\delta}^{-1} \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{leC}}^{-m,-s,-\varsigma,-l,-\ell}}^2) \quad (167)$$

for all $\sigma \geq 0$ for which the right-hand side is finite (in the strong sense that if the right-hand side is finite, then the left-hand side makes sense using the duality pairing for scb-Sobolev spaces and obeys the stated inequality). ■

Proof. By the parametrix construction, we can find, for each $N_0 \in \mathbb{N}$, elliptic $V = V_{N_0} \in \Psi_{\text{leC}}^{-m,-s,-\varsigma,-l,-\ell}(X)$ such that $V^* \Lambda_{m,s,\varsigma,l,\ell} = 1 + Y$ for $Y = Y_{N_0} \in \Psi_{\text{leC}}^{-N_0,-N_0,-N_0,0,0}(X)$. Then,

$$\begin{aligned} 2|\langle u, v \rangle_{L^2}| &\leq 2|\langle \Lambda_{m,s,\varsigma,l,\ell} u, Vv \rangle_{L^2}| + 2|\langle Yu, v \rangle_{L^2}| \\ &\leq \bar{\delta}^{-1} \|\Lambda_{m,s,\varsigma,l,\ell} u\|_{L^2}^2 + \bar{\delta} \|Vv\|_{L^2}^2 + 2|\langle Yu, v \rangle_{L^2}| \\ &\preceq \bar{\delta}^{-1} \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{leC}}^{-m,-s,-\varsigma,-l,-\ell}}^2 + 2|\langle Yu, v \rangle_{L^2}|. \end{aligned} \quad (168)$$

On the other hand, for any $N_1 \in \mathbb{N}$,

$$\begin{aligned} 2|\langle Yu, v \rangle_{L^2}| &= 2|\langle x^{-l}(\sigma^2 + Zx)^{-\ell/2+l} Yu, x^l(\sigma^2 + Zx)^{\ell/2-l} v \rangle_{L^2}| \\ &\leq \bar{\delta}^{-1} \|x^{-l}(\sigma^2 + Zx)^{-\ell/2+l} Yu\|_{H_{\text{b}}^{N_1,0}}^2 + \bar{\delta} \|x^l(\sigma^2 + Zx)^{\ell/2-l} v\|_{H_{\text{b}}^{-N_1,0}}^2 \\ &\preceq \bar{\delta}^{-1} \|x^{-l}(\sigma^2 + Zx)^{-\ell/2+l} Yu\|_{H_{\text{b,leC}}^{N_1,0,0}}^2 + \bar{\delta} \|x^l(\sigma^2 + Zx)^{\ell/2-l} v\|_{H_{\text{b,leC}}^{-N_1,0,0}}^2 \\ &\preceq \bar{\delta}^{-1} \|Yu\|_{H_{\text{b,leC}}^{N_1,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{b,leC}}^{N_1,-l,-\ell}}^2 \\ &\preceq \bar{\delta}^{-1} \|Yu\|_{H_{\text{leC}}^{N_1,N_1+l,N_1+\ell,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{leC}}^{-N_1,-N_1-l,-N_1-\ell,-l,-\ell}}^2 \\ &\preceq \bar{\delta}^{-1} \|u\|_{H_{\text{leC}}^{N_1-N_0,N_1+l-N_0,N_1+\ell-N_0,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{leC}}^{-N_1,-N_1-l,-N_1-\ell,-l,-\ell}}^2. \end{aligned} \quad (169)$$

Taking N_1 sufficiently large, and then taking N_0 sufficiently large relative to that, we get $|\langle Yu, v \rangle_{L^2}| \preceq \bar{\delta}^{-1} \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}}^2 + \bar{\delta} \|v\|_{H_{\text{leC}}^{-m,-s,-\varsigma,-l,-\ell}}^2$. Combining this with eq. (168), we get eq. (167). \square

Lemma 2.16. *Let $m, s, \varsigma, l, \ell, m_0, s_0, \varsigma_0, l_0, \ell_0 \in \mathbb{R}$. Suppose that $A \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$ and that we have some $J \in \mathbb{N}$ and $G_1, \dots, G_J \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ such that*

$$\text{WF}_{\text{leC}}^{\prime\prime,\ell}(A) \subseteq \bigcup_{j=1}^J \text{Ell}_{\text{leC}}^{0,0,0,0,0}(G_j). \quad (170)$$

Then, for each $\Sigma > 0$ and $N \in \mathbb{N}$, there exists some $C = C(A, G_1, \dots, G_J, \Sigma, N) > 0$ such that

$$\|Au\|_{H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}} \leq C \left[\|u\|_{H_{\text{leC}}^{-N,-N,-N,l+l_0,\ell+\ell_0}} + \sum_{j=1}^J \|G_j u\|_{H_{\text{leC}}^{m+m_0,\dots,\ell+\ell_0}} \right] \quad (171)$$

for all $\sigma \in [0, \Sigma]$ and $u \in \mathcal{S}'(X)$. \blacksquare

Proof. It suffices to consider the case of $u \in \mathcal{S}(X)$, the general estimate eq. (171) following from this case via continuity.

Via quantizing some explicit symbols, there exist some $\bar{G}_1, \dots, \bar{G}_J \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ such that $\text{WF}^{0,0}(\bar{G}_j) \subset \text{Ell}_{\text{leC}}^{0,0,0,0,0}(G_j)$ and

$$\text{WF}_{\text{leC}}^{\prime\prime,\ell}(A) \subseteq \bigcup_{j=1}^J \text{Ell}_{\text{leC}}^{0,0,0,0,0}(\bar{G}_j). \quad (172)$$

We now apply the leC-analogue of the Gårding's inequality- type argument. We can choose $E \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ such that

$$\text{Ell}_{\text{leC}}^{0,0,0,0,0}(E) \cup \bigcup_{j=1}^J \text{Ell}_{\text{leC}}^{0,0,0,0,0}(\bar{G}_j) = \text{df} \cup \text{sf} \cup \text{ff} \quad (173)$$

$$\text{WF}_{\text{leC}}^{\prime\prime}(E) \cap \text{WF}_{\text{leC}}^{\prime\prime,\ell}(A) = \emptyset. \quad (174)$$

Thus, every representative of $\sigma_{\text{leC}}^{0,0,0,0}(E^*E + \bar{G}_1^*\bar{G}_1 + \cdots + \bar{G}_J^*\bar{G}_J)$ is nonvanishing on $\text{df} \cup \text{sf} \cup \text{ff}$. We may assume without loss of generality that $E, \bar{G}_1, \dots, \bar{G}_J$ are constant for $\sigma \geq 2\Sigma$, in which case (since $(\text{df} \cup \text{sf} \cup \text{ff}) \cap \{\sigma \leq 2\Sigma\}$ is compact) there exists some $c > 0$ such that $\sigma_{\text{leC}}^{0,0,0,0}(E^*E + \bar{G}_1^*\bar{G}_1 + \cdots + \bar{G}_J^*\bar{G}_J) \geq 2c$ in some neighborhood of $\text{df} \cup \text{sf} \cup \text{ff}$ (in the sense that every representative of the principal symbol has this property, for different neighborhoods). Via an iterative symbolic construction: for each $N_0 \in \mathbb{N}$ there exists some elliptic $B = B_{N_0,c} \in \Psi_{\text{leC}}^{0,0,0,0}(X)$ such that

$$E^*E + \bar{G}_1^*\bar{G}_1 + \cdots + \bar{G}_J^*\bar{G}_J - c - B^*B \in \Psi_{\text{leC}}^{-N_0,-N_0,-N_0,0,0}(X). \quad (175)$$

Then, for $\mathcal{X} = H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}(X)$,

$$c\|Au\|_{\mathcal{X}}^2 \leq -\|BAu\|_{\mathcal{X}}^2 + \|EAu\|_{\mathcal{X}}^2 + \sum_{j=1}^J \|\bar{G}_j Au\|_{\mathcal{X}}^2 \leq \|EAu\|_{\mathcal{X}}^2 + \sum_{j=1}^J \|\bar{G}_j Au\|_{\mathcal{X}}^2 \quad (176)$$

$$\leq \|u\|_{H_{\text{leC}}^{-N,-N,-N,l+l_0,\ell+\ell_0}}^2 + \sum_{j=1}^J \|G_j u\|_{H_{\text{leC}}^{m+m_0,\dots,\ell+\ell_0}}^2 \quad (177)$$

for N_0 sufficiently large, from which eq. (171) follows. \square

Lemma 2.17. *If $m, s, \varsigma, l, \ell, m_0, s_0, \varsigma_0, l_0, \ell_0, m_1, s_1, \varsigma_1, l_1, \ell_1 \in \mathbb{R}$ satisfy $m_1 > m > m_0, \dots, \ell_1 > \ell > \ell_0$, then, for each $\epsilon > 0$ and $\Sigma > 0$, there exists a $C(\epsilon) = C(\epsilon, m, \dots, \ell_1, \Sigma) > 0$ such that*

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq \epsilon \|u\|_{H_{\text{leC}}^{m_1,s_1,\varsigma_1,l_1,\ell_1}} + C(\epsilon) \|u\|_{H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}} \quad (178)$$

for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. \blacksquare

Proof. Fix $\Sigma > 0$. Suppose, to the contrary, that there exists some $\epsilon > 0$, $\{\sigma_k\}_{k \in \mathbb{N}} \subset [0, \Sigma]$, and $\{u_k\}_{k \in \mathbb{N}} \subset \mathcal{S}'(X)$ such that

$$1 = \|u_k\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_k)} \geq \epsilon \|u_k\|_{H_{\text{leC}}^{m_1,s_1,\varsigma_1,l_1,\ell_1}(X)(\sigma_k)} + k \|u_k\|_{H_{\text{leC}}^{m_0,s_0,\varsigma_0,l_0,\ell_0}(X)(\sigma_k)}, \quad (179)$$

i.e., for sufficiently large $N, M \in \mathbb{N}$,

$$\begin{aligned} 1 &= \|\Lambda_{m,s,\varsigma,l,\ell}(\sigma_k)u_k\|_{L^2} + \|x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\|_{H_b^{-M}} \\ &\geq \epsilon \|\Lambda_{m_1,s_1,\varsigma_1,l_1,\ell_1}(\sigma_k)u_k\|_{L^2} + \epsilon \|x^{-l_1}(\sigma_k^2 + Zx)^{l_1-\ell_1/2}u_k\|_{H_b^{-N}} \\ &\quad + k \|\Lambda_{m_0,s_0,\varsigma_0,l_0,\ell_0}(\sigma_k)u_k\|_{L^2} + k \|x^{-l_0}(\sigma_k^2 + Zx)^{l_0-\ell_0/2}u_k\|_{H_b^{-M}}. \end{aligned} \quad (180)$$

Passing to a subsequence if necessary, we may assume without loss of generality that $\sigma_k \rightarrow \sigma_\infty$ for some $\sigma_\infty \in [0, \Sigma]$.

By the Banach–Alaoglu theorem, by passing to a further subsequence if necessary, we can arrange that $\Lambda_{m_i,s_i,\varsigma_i,l_i,\ell_i}(\sigma_k)u_k \rightarrow v_i$ weakly in L^2 for some $v_0, v_1 \in L^2(X)$ and that

$$x^{-l_i}(\sigma_k^2 + Zx)^{l_i-\ell_i/2}u_k \rightarrow w_i \quad (181)$$

weakly in $H_b^{-M}(X), H_b^{-N}(X)$ for some $w_0, w_1 \in H_b^{-M}(X), H_b^{-N}(X)$, respectively. Since

$$\Lambda_{m-m_1,s-s_1,\varsigma-\varsigma_1,l-l_1,\ell-\ell_1}\Lambda_{m_1,s_1,\varsigma_1,l_1,\ell_1} \in \Psi_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) \quad (182)$$

is elliptic, Lemma 2.12 yields

$$\begin{aligned} 1 &= \|u_k\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_k)} \preceq \\ &\|\Lambda_{m-m_1,s-s_1,\varsigma-\varsigma_1,l-l_1,\ell-\ell_1}(\sigma_k)\Lambda_{m_1,s_1,\varsigma_1,l_1,\ell_1}(\sigma_k)u_k\|_{L^2} + \|x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\|_{H_b^{-M}}. \end{aligned} \quad (183)$$

Taking $M > N$, we will produce a contradiction by showing that both terms on the right-hand side of eq. (183) converge to zero as $k \rightarrow \infty$.

(1) For $M > N$,

$$\begin{aligned} x^{-l_0}(\sigma_k^2 + Zx)^{l_0-\ell_0/2}u_k - x^{l_1-l_0}(\sigma_k^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}w_1 \\ = x^{l_1-l_0}(\sigma_k^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}(x^{-l_1}(\sigma_k^2 + Zx)^{l_1-\ell_1/2}u_k - w_1) \rightarrow 0 \end{aligned} \quad (184)$$

strongly in $H_b^{-M}(X)$. It follows that $x^{l_1-l_0}(\sigma_k^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}w_1 \rightarrow w_0$ weakly in $H_b^{-M}(X)$, and thus in $\mathcal{S}'(X)$. But, it is also the case that

$$x^{l_1-l_0}(\sigma_k^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}w_1 \rightarrow x^{l_1-l_0}(\sigma_\infty^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}w_1 \quad (185)$$

in $\mathcal{S}'(X)$. So, in fact,

$$w_0 = x^{l_1-l_0}(\sigma_\infty^2 + Zx)^{l_0-l_1+\ell_1/2-\ell_0/2}w_1, \quad (186)$$

and

$$x^{-l_0}(\sigma_k^2 + Zx)^{l_0-\ell_0/2}u_k \rightarrow w_0 \quad (187)$$

strongly in $H_b^{-M}(X)$. We can therefore deduce from the family of inequalities eq. (180) that $\|w_0\|_{H_b^{-M}} = 0$, i.e. $w_0 = 0$, from which and eq. (186) it follows that $w_1 = 0$.

But then, by the joint continuity of the multiplication operator

$$x^{l_0}(\sigma^2 + Zx)^{\ell_0/2-l_0} : [0, \Sigma]_\sigma \times \mathcal{S}'(X) \rightarrow \mathcal{S}'(X), \quad (188)$$

$u_k \rightarrow x^{l_0}(\sigma_\infty^2 + Zx)^{\ell_0/2-l_0}w_0 = 0$ in $\mathcal{S}'(X)$. We conclude that $v_0, v_1 = 0$.

Via the same initial argument, we deduce that $x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k$ converges strongly in H_b^{-M} to $x^{l_1-l}(\sigma_\infty^2 + Zx)^{l-l_1+\ell_1/2-\ell/2}w_1 = 0$. Thus, $\|x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\|_{H_b^{-M}} \rightarrow 0$ as $k \rightarrow \infty$.

(2) We now consider

$$\Lambda = \Lambda_{m-m_1, s-s_1, \varsigma-\varsigma_1, l-l_1, \ell-\ell_1} \in C^0([0, \infty)_\sigma; \Psi_b^{-\varepsilon, -\varepsilon}(X)), \quad (189)$$

where $\varepsilon > 0$ is sufficiently small. Thus, we can write Λ as the composition of a fixed compact operator on $L^2(X)$ and a continuous family of bounded operators on $L^2(X)$. It follows (since $\Lambda_{m_1, s_1, \varsigma_1, l_1, \ell_1}(\sigma_k)u_k \rightarrow v_1$ weakly in L^2) that

$$\Lambda(\sigma_k)\Lambda_{m_1, s_1, \varsigma_1, l_1, \ell_1}(\sigma_k)u_k - \Lambda(\sigma_k)v_1 \rightarrow 0 \quad (190)$$

strongly in $L^2(X)$. But $v_1 = 0$, as proven above, so

$$\|\Lambda_{m-m_1, s-s_1, \varsigma-\varsigma_1, l-l_1, \ell-\ell_1}(\sigma_k)\Lambda_{m_1, s_1, \varsigma_1, l_1, \ell_1}(\sigma_k)u_k\|_{L^2} \rightarrow 0 \quad (191)$$

as $k \rightarrow \infty$.

□

Lemma 2.18. *The $L_{\text{sc}}^2(X), L_{\text{sc}}^2(X_{1/2})$ -based b-Sobolev spaces on X and $X_{1/2}$ are related by*

$$H_b^{m, l}(X) = H_b^{m, 2l+n/2}(X_{1/2}). \quad (192)$$

■

Proof. For $f, g \in L_{\text{sc}}^2([0, \bar{x})_x \times \partial X)$,

$$\begin{aligned} \langle f, g \rangle_{L_{\text{sc}}^2([0, \bar{x})_x \times \partial X)} &= \int_{\partial X} \left(\int_0^{\bar{x}} f^*(x)g(x) \frac{dx}{x^{n+1}} \right) d\text{Vol}_{g_{\partial X}}(y) \\ &= 2 \int_{\partial X} \left(\int_0^{\bar{x}^{1/2}} f^*(\rho^2)g(\rho^2) \frac{d\rho}{\rho^{2n+1}} \right) d\text{Vol}_{g_{\partial X}}(y) \\ &= 2 \int_{\partial X} \left(\int_0^{\bar{x}^{1/2}} \frac{f^*(\rho^2)}{\rho^{n/2}} \frac{g(\rho^2)}{\rho^{n/2}} \frac{d\rho}{\rho^{n+1}} \right) d\text{Vol}_{g_{\partial X}}(y). \end{aligned} \quad (193)$$

Thus, $L_{\text{sc}}^2([0, \bar{x})_x \times \partial X) \ni f(x) \mapsto \sqrt{2}f(\rho^2) \in \rho^{n/2}L_{\text{sc}}^2([0, \bar{x}^{1/2})_\rho \times \partial X)$ defines an isomorphism of Hilbert spaces.

This implies that

$$L_{\text{sc}}^2(X) = x^{n/4}L_{\text{sc}}^2(X_{1/2}). \quad (194)$$

As follows from the definition [Vas18, Definition 5.15] (see also [Mel93, Definition 4.22] for the case of classicality at the front face of the b-double space), $\Psi_b^{m,l}(X) = \Psi_b^{m,2l}(X_{1/2})$ for all $m, l \in \mathbb{R}$. In conjunction with eq. (194), this implies that

$$H_b^{m,l}(X) = \Psi_b^{-m,-l}(X)L^2(X) = \Psi_b^{-m,-2l-n/2}(X_{1/2})L^2(X_{1/2}) = H_b^{m,2l+n/2}(X_{1/2}). \quad (195)$$

□

For use in §6, we briefly recall the connection between spaces of conormal distributions on $X_{\text{res}}^{\text{sp}}$, which are defined using L^∞ -based spaces, and the b-Sobolev spaces, which are defined using L^2 . We have

$$\mathcal{A}_{\text{loc}}^{\alpha,\beta,0}(X_{\text{res}}^{\text{sp}}) = x^\alpha(\sigma^2 + Zx)^{-\alpha+\beta/2}\mathcal{A}_{\text{loc}}^{0,0}([0, \infty) \times X), \quad (196)$$

and

$$\begin{aligned} \mathcal{A}_{\text{loc}}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}}) &= \{u \in \mathcal{A}^0(X_{\text{res}}^{\text{sp}}) : [\hat{E} \mapsto u|_{E/x=\hat{E}}] \in C^\infty([0, \infty)_{\hat{E}}; \mathcal{A}^0(X_{1/2}))\}, \\ \mathcal{A}_{\text{loc}}^{\alpha,\beta,(0,0)}(X_{\text{res}}^{\text{sp}}) &= x^\alpha(\sigma^2 + Zx)^{-\alpha+\beta/2}\mathcal{A}_{\text{loc}}^{0,0,(0,0)}(X_{\text{res}}^{\text{sp}}), \end{aligned} \quad (197)$$

where $\hat{E} = E/x$ and we are identifying level sets of \hat{E} with $\text{zf} \cong X_{1/2}$. Note that greater indices α, β means greater decay, the convention opposite of that used for symbols and Ψ DOs. We also use

$$\mathcal{A}_{\text{loc}}^{\alpha-, \beta-, 0}(X_{\text{res}}^{\text{sp}}) = \cap_{\alpha' < \alpha, \beta' < \beta} \mathcal{A}_{\text{loc}}^{\alpha', \beta', 0}(X_{\text{res}}^{\text{sp}}) \quad (198)$$

$$\mathcal{A}_{\text{loc}}^{\alpha-, \beta-, (0,0)}(X_{\text{res}}^{\text{sp}}) = \cap_{\alpha' < \alpha, \beta' < \beta} \mathcal{A}_{\text{loc}}^{\alpha', \beta', (0,0)}(X_{\text{res}}^{\text{sp}}). \quad (199)$$

for $\alpha, \beta \in \mathbb{R}$. These are all Fréchet spaces of (locally) conormal distributions on $X_{\text{res}}^{\text{sp}}$ that are smooth at zf° (where, to reiterate, “local” just means that we do not require uniformity as $\sigma \rightarrow \infty$, only as $\sigma \rightarrow 0^+$). Via Sobolev embedding,

$$\cup_{l' > l} H_b^{\infty, l'}(X) = H_b^{\infty, l+}(X) \subseteq \mathcal{A}^{l+n/2}(X) \subseteq H_b^{\infty, l-}(X) = \cap_{l' < l} H_b^{\infty, l'}(X), \quad (200)$$

and, for each $\epsilon > 0$, each seminorm of $\mathcal{A}^{l+n/2}(X)$ can be controlled using only finitely many of the norms in the family $\{\|\cdot\|_{H_b^{m, l+\epsilon}}\}_{m \in \mathbb{R}}$.

Since, for each $\alpha \geq 0$, $\mathcal{A}_{\text{loc}}^0([0, \infty)_\sigma; \mathcal{A}^{\alpha-}(X)) \subseteq \mathcal{A}_{\text{loc}}^{\alpha-, 2\alpha-, 0}(X_{\text{res}}^{\text{sp}})$, the preceding observation implies that

Proposition 2.19. *For any function $m_0 : \mathbb{R} \rightarrow \mathbb{R}$,*

$$\bigcap_{l < -1/2} \bigcap_{m > m_0(l)} \mathcal{A}_{\text{loc}}^0([0, \infty)_\sigma; H_b^{m, l+\alpha}(X)) \subseteq \mathcal{A}_{\text{loc}}^{(\alpha+(n-1)/2)-, (2\alpha+(n-1))-, 0}(X_{\text{res}}^{\text{sp}}) \quad (201)$$

holds for each $\alpha \in \mathbb{R}$. ■□

Remark. Implicit in the statements above is the identification of elements of $\mathcal{A}_{\text{loc}}^0([0, \infty)_\sigma; \mathcal{A}^{\alpha-}(X))$ with extendable distributions on $X_{\text{res}}^{\text{sp}}$, which occurs via

$$\begin{aligned} \mathcal{A}_{\text{loc}}^0([0, \infty)_\sigma; \mathcal{A}^{\alpha-}(X)) &\ni \{u(-; \sigma)\}_{\sigma > 0} \mapsto \\ &\left[\dot{C}^\infty(X_{\text{res}}^{\text{sp}}) \ni \chi \mapsto \int_0^\infty \langle u(-; \sigma), \chi(-; \sigma) \rangle d\sigma \right] \in \mathcal{D}'(X_{\text{res}}^{\text{sp}}). \end{aligned} \quad (202)$$

We defer to [Mel92][Mel93] for more about conormal distributions. See also [Hin22] for the particular case of $X_{\text{res},0}^{\text{sp}}$ (where the notation X_{res}^+ is used instead).

Proposition 2.20. *Fix $\chi \in C_c^\infty(X_{\text{res}}^{\text{sp}})$ supported away from bf and nonvanishing near zf. Then, if $(\chi x \partial_E)^k u \in \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}})$ for all $k \in \mathbb{N}$, then $u \in \mathcal{A}_{\text{loc}}^{0-,0-,0,(0,0)}(X_{\text{res}}^{\text{sp}})$. ■*

Proof. We want to show that $[\hat{E} \mapsto u|_{E/x=\hat{E}}] \in C^\infty([0, \infty)_{\hat{E}}; \mathcal{A}^0(X_{1/2}))$. Observe that $x \partial_E = \partial_{\hat{E}}$ away from bf, where the partial derivative on the right-hand side is taken with x held constant.

Consequently, the $k = 0, \dots, K+1$ cases of $(\chi x \partial_E)^k u \in \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}})$ together show that $[\hat{E} \mapsto u|_{\hat{E}}] \in C^K([0, \infty)_{\hat{E}}; \mathcal{A}^{0-}(X_{1/2}))$. Here we are using that, given $v \in \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}})$,

$$[\hat{E} \mapsto v|_{E/x=\hat{E}}] \in \mathcal{A}^0([0, 1)_{\hat{E}}; \mathcal{A}^{0-}(X_{1/2})) = \mathcal{A}^0([0, 1)_{\hat{E}}; \mathcal{A}^{0-}(X)). \quad (203)$$

Since K can be taken arbitrarily large, we conclude the claim. \square

Remark. As the argument shows, each Fréchet seminorm of $u \in \mathcal{A}_{\text{loc}}^{0-,0-,0,(0,0)}(X_{\text{res}}^{\text{sp}})$ is controlled by finitely many Fréchet seminorms of $(\chi x \partial_E)^k u \in \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}})$ for finitely many k . In other words, the map

$$\begin{aligned} \left(\prod_{k=0}^{\infty} \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}}) \right) \cap \{ \{ (\chi x \partial_E)^k u \}_{k=0}^{\infty} : u \in \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}}) \} &\ni \{ (\chi x \partial_E)^k u \}_{k=0}^{\infty} \\ &\mapsto u \in \mathcal{A}_{\text{loc}}^{0-,0-,0,(0,0)}(X_{\text{res}}^{\text{sp}}) \end{aligned} \quad (204)$$

is continuous when we endow the domain with the topology of $\prod_{k=0}^{\infty} \mathcal{A}_{\text{loc}}^{0-,0-,0}(X_{\text{res}}^{\text{sp}})$.

3. THE CONJUGATED PERSPECTIVE

We now construct the “conjugated” operator $\tilde{P} = \{\tilde{P}(\sigma)\}_{\sigma \geq 0}$. Given some

$$\{f(-; \sigma)\}_{\sigma \geq 0}, \{g(-; \sigma)\}_{\sigma \geq 0} \subset C^\infty(X^\circ)$$

and a family of differential operators $\{D(\sigma)\}_{\sigma \geq 0} \subset \text{Diff}(X^\circ)$, we use the somewhat abusive notation $fDg = \{f(-; \sigma)D(\sigma)g(-; \sigma)\}_{\sigma \geq 0} \subset \text{Diff}(X^\circ)$ to denote the family of differential operators

$$M_f D M_g = \{M_{f(-; \sigma)} D(\sigma) M_{g(-; \sigma)}\}_{\sigma \geq 0}, \quad (205)$$

where for $h \in C^\infty(X^\circ)$, $M_h : C_c^\infty(X^\circ) \rightarrow C_c^\infty(X^\circ)$ denotes the multiplication operator $C_c^\infty(X^\circ) \ni \varphi \mapsto h\varphi$. For us, f, g will be of the form $f(-; \sigma) = \exp(-i\Phi(-; \sigma))$ and $g(-; \sigma) = \exp(+i\Phi(-; \sigma))$ for some $\Phi = \{\Phi(-; \sigma)\}_{\sigma \geq 0} \subset C^\infty(X^\circ)$. The conjugated operator \tilde{P} is defined by

$$\tilde{P} = M_{\exp(-i\Phi)} P M_{\exp(+i\Phi)} = e^{-i\Phi} P e^{+i\Phi} \quad (206)$$

for our eventual choice of Φ .

As discussed in the introduction, Φ had ought to be determined to an order or two (including logarithmic terms) on $X_{\text{res}}^{\text{sp}}$ by the actual asymptotics of solutions of $P(\sigma)u = f$ for $f \in \mathcal{S}(X)$. Rather than determine what Φ should be in this manner (that is by solving the PDE, or a model thereof, to a sufficient degree of accuracy), it is actually easier to work backwards, meaning to find the asymptotics of solutions to the given PDE by first finding a choice of Φ for which \tilde{P} has a workable form, where “workable” roughly means qualitatively similar to the conjugated operator in [Vas21a]. (We will then have to actually show that this choice describes the asymptotics of solutions to the PDE.) Such a choice has already been stated in the introduction, §1, eq. (66), and the resultant conjugated operator is computed below. We now motivate that choice. Consider the model operator $P_{\text{Model}} \in \text{Diff}_{\text{sc}}^2([0, \infty)_x)$ given by

$$P_{\text{Model}}(\sigma) = -(1 + xa_{00})(x^2 \partial_x)^2 + (a + n - 1)x^3 \partial_x - \sigma^2 - Zx, \quad (207)$$

where $a_{00}, a, Z \in \mathbb{R}$, $Z > 0$, and where we require $a_{00} < 0$ for simplicity. This captures the leading and subleading terms of $P(\sigma)$ in $\Psi_{\text{scb}}(X)$, modulo the terms involving nonradial derivatives (which

can be expected to be unimportant – i.e. under symbolic control – based on considerations similar to those in [Vas21a]). For $\Phi(-; \sigma) \in C^\infty(\mathbb{R}_x^+ \times \mathbb{R}_\sigma^+)$, let

$$\tilde{P}_{\text{Model}} = M_{\exp(-i\Phi)} P_{\text{Model}} M_{\exp(+i\Phi)} = \exp(-i\Phi) P_{\text{Model}} \exp(+i\Phi). \quad (208)$$

This will be qualitatively similar (at least with regards to the sc-calculus) to Vasy's conjugated operator family (with $\Im \alpha_\pm = 0$) if \tilde{P}_{Model} is equal to

$$\begin{aligned} P_{\text{Goal}} = & -(1 + xa_{00}) \left(x^2 \partial_x - \frac{x(n-1)}{2} \right)^2 + 2i(1 + xa_{00}) \sqrt{\sigma^2 + Zx - \sigma^2 a_{00} x} (x^2 \partial_x) \\ & + 2i(1 + xa_{00}) \left(-\frac{(n-1)}{2} x \sqrt{\sigma^2 + Zx} + \frac{Zx^2}{4} \frac{1}{\sqrt{\sigma^2 + Zx}} \right) \end{aligned} \quad (209)$$

modulo terms which are two sc-decay orders subleading at sf and ff (in particular in $x^2 \text{Diff}_{\text{sc}}^{*,0}(X)$ for $\sigma > 0$ and $x^2 \text{Diff}_{\text{sc}}^{*,0}(X_{1/2})$ for $\sigma = 0$, where $X = [0, \infty)_x$). The specific form of the lower order terms in eq. (209) bears comment: the operators $x^2 \partial_x - x(n-1)/2$ and

$$(\sigma^2 + Zx)^{1/2} x^2 \partial_x - x(\sigma^2 + Zx)^{1/2} \frac{(n-1)}{2} + \frac{Zx^2}{4} \frac{1}{(\sigma^2 + Zx)^{1/2}} \quad (210)$$

are both formally anti- self-adjoint with respect to the $L^2([0, 1], x^{-(n+1)} dx) = L_{\text{sc}}^2[0, 1]$ inner product. Indeed, for any $f, g \in \dot{\mathcal{S}}([0, \infty))$,

$$\int_0^\infty f^*(x) \partial_x g(x) x^{-(n+1)} dx = - \int_0^\infty (\partial_x f)^* g(x) x^{-(n+1)} dx + (n+1) \int_0^\infty \left(\frac{f(x)}{x} \right)^* g(x) x^{-(n+1)} dx, \quad (211)$$

so, as bilinear forms $\dot{\mathcal{S}}([0, \infty))^2 \rightarrow \mathbb{C}$,

$$\partial_x^* = -\partial_x + (n+1)/x, \quad (212)$$

$$(x^2 \partial_x)^* = x^2 \partial_x^* + [\partial_x^*, x^2] = -x^2 \partial_x + x(n-1) \quad (213)$$

and

$$\begin{aligned} (\sqrt{\sigma^2 + Zx} x^2 \partial_x)^* &= \sqrt{\sigma^2 + Zx} (x^2 \partial_x)^* + [(x^2 \partial_x)^*, \sqrt{\sigma^2 + Zx}] \\ &= -\sqrt{\sigma^2 + Zx} x^2 \partial_x + x(n-1) \sqrt{\sigma^2 + Zx} - \frac{Zx^2}{2} \frac{1}{\sqrt{\sigma^2 + Zx}}, \end{aligned} \quad (214)$$

which implies the claimed anti- self-adjointness. The $(1 + xa_{00})$ terms in eq. (209), along with the $-\sigma^2 a_{00} x$ under the square root, spoil anti- self-adjointness or self-adjointness, but only negligibly. Thus, the terms in eq. (209) have definite adjointness modulo negligible errors.

As a preliminary step towards P_{Goal} , we can conjugate away the $(a+n-1)x^3 \partial_x$ term in eq. (207), getting

$$x^{-(a+n-1)/2} P_{\text{Model}} x^{(a+n-1)/2} = -(1 + xa_{00}) (x^2 \partial_x)^2 - \sigma^2 - Zx \text{ mod } x^2 \text{Diff}_{\text{sc}}^{1,0}(X). \quad (215)$$

The $x^2 \text{Diff}_{\text{sc}}^{1,0}$ remainder may contain some terms in eq. (209), for example the square of the $x(n-1)/2$ terms, but we can conjugate these back in at the end of the computation. Conjugating eq. (215) by $\exp(+i\varphi)$ for a to-be-decided $\varphi(x; \sigma) \in C^\infty(\mathbb{R}_x^+ \times \mathbb{R}_\sigma^+)$, we get

$$e^{-i\varphi} x^{-(a+n-1)/2} P_{\text{Model}} x^{(a+n-1)/2} e^{+i\varphi} = -(1 + xa_{00}) (x^2 \partial_x + ix^2 \varphi')^2 - \sigma^2 - Zx \text{ mod } x^2 \text{Diff}_{\text{sc}}^{1,0}(X), \quad (216)$$

assuming the contribution from conjugating the $x^2 \text{Diff}_{\text{sc}}^{1,0}$ in eq. (215) is negligible. It is not unreasonable to expect (perhaps based on the $Z = 0$ case) that, for our eventual choice of φ , the leading order new contribution to eq. (216) is $x^4 (\varphi')^2$.

Since the Coulomb term in eq. (216) is subleading order relative to σ^2 in the sense of decay, we should really be keeping track of the new contributions to one subleading order. To this order, the

new contribution to the effective potential is $x^4(1 + xa_{00})(\varphi')^2$ (assuming the terms with second derivatives of φ are negligible). Thus, we seek to arrange

$$x^4(1 + xa_{00})(\varphi')^2 = \sigma^2 + Zx \bmod x^2 C^\infty(X). \quad (217)$$

Multiplying through by $(1 + xa_{00})^{-1} = 1 - xa_{00} \bmod x^2 C^\infty(X)$, this suggests setting $x^4(\varphi')^2 = \sigma^2 + Zx - \sigma^2 a_{00}x$, the solution of which (up to an arbitrary additive constant and conventional choice of sign) is

$$\varphi(x; \sigma) = \frac{1}{x} \sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} + \frac{1}{\sigma} (Z - \sigma^2 a_{00}) \operatorname{arcsinh} \left(\frac{\sigma}{x^{1/2}} \frac{1}{(Z - \sigma^2 a_{00})^{1/2}} \right). \quad (218)$$

Recall that $\operatorname{arcsinh}(z) = \log(z + (1 + z^2)^{1/2})$ for all $z \geq 0$. Expanding $\operatorname{arcsinh}(z)$ in Taylor series around $z = 0$, we see that the apparent singularity in eq. (218) at $\sigma = 0$ is removable (to all orders), and hence $\varphi(x; \sigma)$ defines a smooth function on $\mathbb{R}_x^+ \times \mathbb{R}_\sigma$, and it is even in σ .

We observe that, given eq. (217), the $x^4\varphi''$ term in eq. (216) is indeed negligible except at ff, where there is one non-negligible contribution, and it is precisely the final term in eq. (209) (modulo negligible terms).

Given this definition of φ , $e^{-i\varphi}x^{-(a+n-1)/2}P_{\text{Model}}x^{(a+n-1)/2}e^{+i\varphi}$ is given, modulo $x^2 \operatorname{Diff}_{\text{sc}}^{1,0}(X)$, by

$$\begin{aligned} & -(1 + xa_{00})((x^2\partial_x)^2 + 2ix^4\varphi'\partial_x - x^4\varphi'\varphi' + i(x^2\partial_x)^2\varphi) - \sigma^2 - Zx = -(1 + xa_{00})((x^2\partial_x)^2 \\ & - 2ix^2\sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x}\partial_x + i(x^2\partial_x)^2\varphi) \bmod C^\infty(\mathbb{R}_{\sigma^2}; x^2 C^\infty(X)), \end{aligned} \quad (219)$$

plus the term in $\operatorname{Diff}_{\text{leC}}^{0,-2,-5,-2,-5}$ that results from applying the first order operator in the $x^2 \operatorname{Diff}_{\text{sc}}^{1,0}(X)$ remainder in eq. (216) to φ , which we will work out later.

We can now add back in the $x(n-1)/2$ terms:

$$e^{-i\varphi}x^{-a/2}P_{\text{Model}}e^{+i\varphi}x^{a/2} = P_{\text{Goal}} \bmod \operatorname{Diff}_{\text{leC}}^{1,-2,-4,-2,-4}. \quad (220)$$

So, at least in this model case, conjugation by $e^{+i\Phi} = e^{i\varphi+(1/2)a \log x}$ has the required properties.

Returning to the full problem, we consider the family $\Phi = \{\Phi(-; \sigma)\}_{\sigma \geq 0}$ of $\Phi(-; \sigma) \in C^\infty(X^\circ \times \mathbb{R}_\sigma^+)$ given by

$$\Phi(x; \sigma) = \frac{1}{x} \sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} + \frac{1}{\sigma} (Z - \sigma^2 a_{00}) \operatorname{arcsinh} \left(\frac{\sigma}{x^{1/2}} \frac{1}{(Z - \sigma^2 a_{00})^{1/2}} \right) - \frac{i}{2} a \log x. \quad (221)$$

Observe that (after removing the removable singularity at $\sigma = 0$) eq. (70) holds. We then define, for each $\sigma \geq 0$,

$$\tilde{P}(\sigma) = M_{e^{-i\Phi}(-; \sigma)} P(\sigma) M_{e^{+i\Phi}(-; \sigma)}, \quad (222)$$

i.e. $\tilde{P} = \exp(-i\Phi)P \exp(+i\Phi)$. Thus, $\tilde{P}(\sigma) \in S \operatorname{Diff}_{\text{sc}}(X)$ for each $\sigma \geq 0$, and the coefficients all depend smoothly on σ all the way down to $\sigma = 0$ in compact subsets of X° . Of course, this does not mean that $[0, \infty)_\sigma \ni \sigma \mapsto \tilde{P}(\sigma) \in S \operatorname{Diff}_{\text{sc}}^{2,0}(X)$ is smooth all the way down to $\sigma = 0$. This map is continuous for $\sigma > 0$ but discontinuous at $\sigma = 0$, as can be verified by computing sc-principal symbols.

Let $\tilde{P}_0 = e^{-i\Phi}P_0e^{+i\Phi}$, $\tilde{P}_1 = e^{-i\Phi}P_1e^{+i\Phi}$, and $\tilde{P}_2 = e^{-i\Phi}P_2e^{+i\Phi}$, where P_0, P_1, P_2 are as in §1. We have:

Proposition 3.1. *For each $\sigma \geq 0$, $\tilde{P}_0(\sigma)$ is given with respect to the boundary-collar ι by*

$$\tilde{P}_0(\sigma) = -(1 + xa_{00})(x^2\partial_x)^2 + x^2\Delta_{\partial X} + (n-1)x^3\partial_x + L(\sigma) + V_{\text{eff}}(x; \sigma), \quad (223)$$

where

$$L(\sigma) = 2ix(1 + xa_{00})\sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} \left(x\partial_x - \frac{n-1}{2} + \frac{Z}{4} \frac{x}{\sigma^2 + Zx} \right) - aa_{00}x^4\partial_x \quad (224)$$

and $V_{\text{eff}} \in x^2 C^\infty(X_{\text{res}}^{\text{sp}})$. ■

Proof. We work on $\hat{X} = [0, \bar{x}]_x \times \partial X$.

We may write $\tilde{P}_0 = P_0 + e^{-i\Phi}[P_0, e^{+i\Phi}]$. Observe that

$$[P_0, e^{+i\Phi}] = -(1 + xa_{00})[(x^2\partial_x)^2, e^{i\Phi}] + x^3(a + n - 1)[\partial_x, e^{i\Phi}] + x^2[\Delta_{\partial X}, e^{i\Phi}] \quad (225)$$

$$e^{-i\Phi}[(x^2\partial_x)^2, e^{i\Phi}] = 2ix^4\Phi'\partial_x - x^4\Phi'\Phi' + i(x^2\partial_x)^2\Phi \quad (226)$$

$$e^{-i\Phi}[\partial_x, e^{i\Phi}] = i\Phi' \quad (227)$$

$$e^{-i\Phi}[\Delta_{\partial X}, e^{i\Phi}] = 0, \quad (228)$$

where the primes denote differentiation in x . (If a, a_{00} were nonconstant functions on ∂X , then eq. (228) would not hold. This is ultimately the reason for assuming that a, a_{00} are constant.) Thus, if we set

$$L(\sigma) = ax^3\partial_x - 2ix^4(1 + xa_{00})\Phi'\partial_x + V_{\text{SA}}, \quad (229)$$

$$V_{\text{eff}}(x; \sigma) = -(1 + xa_{00})(-x^4\Phi'\Phi' + i(x^2\partial_x)^2\Phi) + ix^3(a + n - 1)\Phi' - V_{\text{SA}} - \sigma^2 - Zx \quad (230)$$

for

$$V_{\text{SA}} = -2ix(1 + xa_{00})\sqrt{\sigma^2 + Zx - \sigma^2a_{00}x}\left(\frac{n-1}{2} - \frac{Z}{4}\frac{x}{\sigma^2 + Zx}\right), \quad (231)$$

then eq. (223) holds, and it only remains to verify eq. (224) and the fact that $V_{\text{eff}} \in x^2C^\infty(X_{\text{res}}^{\text{sp}})$.

We compute that

$$\begin{aligned} +x^4\Phi'\Phi' &= x^4\left(\frac{1}{x^2}\sqrt{\sigma^2 + Zx - \sigma^2a_{00}x} + \frac{ia}{2x}\right)^2 \\ &= \sigma^2 + Zx - \sigma^2a_{00}x + iax\sqrt{\sigma^2 + Zx - \sigma^2a_{00}x} - \frac{x^2a^2}{4}, \end{aligned} \quad (232)$$

$$-i(x^2\partial_x)^2\Phi = -\frac{ax^2}{2} + \frac{i}{2}x^2\frac{(Z - \sigma^2a_{00})}{(\sigma^2 + Zx - \sigma^2a_{00}x)^{1/2}}, \quad (233)$$

$$+ix^3\Phi' = \frac{ax^2}{2} - ix\sqrt{\sigma^2 + Zx - \sigma^2a_{00}x}. \quad (234)$$

From eq. (234) and eq. (229), we get eq. (224).

When adding up the various contributions to V_{eff} , as written in eq. (230), a few key cancellations happen by design:

- (I) the first σ^2 in eq. (230) (coming from $x^4\Phi'\Phi'$, eq. (232)) cancels with the last $-\sigma^2$ in eq. (230),
- (II) the first Zx term, also coming from $x^4\Phi'\Phi'$, cancels with the $-Zx$ in eq. (230), so that the original Coulomb-like term has been “conjugated away,”
- (III) multiplying eq. (232) by $(1 + xa_{00})$ the terms in $(1 + xa_{00})(\sigma^2 - \sigma^2a_{00}x)$ linear in a_{00} cancel per difference-of-squares,
- (IV) the $iax(\sigma^2 + Zx - \sigma^2a_{00}x)^{1/2}$ term in eq. (232) cancels with the term in $iax^3\Phi'$ coming from the last term in eq. (234).
- (V) The $ix(\sigma^2 + Zx - \sigma^2a_{00}x)^{1/2}(n-1)$ term in V_{eff} coming from V_{SA} cancels with the identical term in $i(n+1)x^3\Phi'$ coming from the last term in eq. (234).

All in all, V_{eff} is given by

$$\begin{aligned} a_{00}(Z - \sigma^2a_{00})x^2 + ia_{00}(a + n - 1)x^2\sqrt{\sigma^2 + Zx - \sigma^2a_{00}x} - \left[(1 + xa_{00})\left(\frac{a^2}{4} + \frac{a}{2}\right) - \frac{a}{2}(a + n - 1)\right]x^2 \\ + \frac{ix}{2}(1 + xa_{00})\left[\frac{Zx - \sigma^2a_{00}x}{(\sigma^2 + Zx - \sigma^2a_{00}x)^{1/2}} - (\sigma^2 + Zx - \sigma^2a_{00}x)^{1/2}\frac{Zx}{\sigma^2 + Zx}\right]. \end{aligned} \quad (235)$$

It is clear that the first line of eq. (235) defines an element of $x^2 C^\infty(X_{\text{res}}^{\text{sp}})$. On the other hand, $-(ix/2)(1 + xa_{00})\sigma^2 a_{00}x/(\sigma^2 + Zx - \sigma^2 a_{00}x)^{1/2}$ is in $x^2 C^\infty(X_{\text{res}}^{\text{sp}})$ as well. Thus, it remains to verify that

$$\frac{Zx}{(\sigma^2 + Zx)^{1/2}} - (\sigma^2 + Zx - \sigma^2 a_{00}x)^{1/2} \frac{Zx}{\sigma^2 + Zx} \in x C^\infty(X_{\text{res}}^{\text{sp}}), \quad (236)$$

i.e. that

$$1 - \left(1 - \frac{\sigma^2 a_{00}x}{\sigma^2 + Zx}\right)^{1/2} \in (\sigma^2 + Zx)^{1/2} C^\infty(X_{\text{res}}^{\text{sp}}). \quad (237)$$

This is of course not true for each term on the left-hand side individually, but we can expand

$$\left(1 - \frac{\sigma^2 a_{00}x}{\sigma^2 + Zx}\right)^{1/2} = 1 \bmod \frac{\sigma^2 a_{00}x}{\sigma^2 + Zx} C^\infty(X_{\text{res}}^{\text{sp}}) = 1 \bmod x C^\infty(X_{\text{res}}^{\text{sp}}) \quad (238)$$

(since $f(\zeta) = \zeta^{-1}((1 - \zeta)^{1/2} - 1) \in C^\infty(-\infty, 1)_\zeta$, $f(\sigma^2 a_{00}x/(\sigma^2 + Zx)) \in C^\infty(X_{\text{res}}^{\text{sp}})$, which implies eq. (238)), so in fact eq. (237) is true with some room to spare. We can then conclude that V_{eff} is in $x^2 C^\infty(X_{\text{res}}^{\text{sp}})$. \square

Proposition 3.2. *The family $L = \{L(\sigma)\}_{\sigma \geq 0} \in \text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X)$ satisfies*

$$L = 2i(1 + xa_{00}) \left(1 - \frac{\sigma^2 a_{00}x}{2} \frac{1}{\sigma^2 + Zx}\right) x \sqrt{\sigma^2 + Zx} \left(x \partial_x - \frac{n-1}{2} + \frac{Z}{4} \frac{x}{\sigma^2 + Zx}\right) \bmod \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X) \quad (239)$$

near ∂X . \blacksquare

Proof. It suffices to restrict attention to $\hat{X} = [0, \bar{x}) \times \partial X$.

We refer to eq. (224). Expanding $(1 - \sigma^2 a_{00}x/(\sigma^2 + Zx))^{1/2} = 1 - (1/2)\sigma^2 a_{00}x/(\sigma^2 + Zx) + O(\sigma^4 x^2/(\sigma^2 + Zx)^2)$ in Taylor series, we deduce that

$$2ix(1 + xa_{00}) \left(\left(1 - \frac{\sigma^2 a_{00}x}{2(\sigma^2 + Zx)}\right) \sqrt{\sigma^2 + Zx} - \sqrt{\sigma^2 + Zx - \sigma^2 a_{00}x} \right) \left(x \partial_x - \frac{n-1}{2} + \frac{Z}{4} \frac{x}{\sigma^2 + Zx}\right) \in \text{Diff}_{\text{leC}}^{1,-2,-6,-3,-7}(X). \quad (240)$$

Thus,

$$L = 2i(1 + xa_{00}) \left(1 - \frac{\sigma^2 a_{00}x}{2} \frac{1}{\sigma^2 + Zx}\right) x \sqrt{\sigma^2 + Zx} \left(x \partial_x - \frac{n-1}{2} + \frac{Z}{4} \frac{x}{\sigma^2 + Zx}\right) - aa_{00}x^4 \partial_x \bmod \text{Diff}_{\text{leC}}^{1,-2,-6,-3,-7}(X). \quad (241)$$

On the other hand, $aa_{00}x^4 \partial_x \in \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X)$. We conclude eq. (239) from the above. \square

Proposition 3.3. $\tilde{P}_0 \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X)$, and

$$\tilde{P}_0 = -(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} + 2ix \sqrt{\sigma^2 + Zx} \left(x \partial_x - \frac{n-1}{2} + \frac{Z}{4} \frac{x}{\sigma^2 + Zx}\right) \bmod \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X). \quad (242)$$

Proof. We have $L(\sigma) \in \text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X)$ and

$$\begin{aligned} -(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} &\in \text{Diff}_{\text{leC}}^{2,0,-2,-2,-4}(X) \\ xa_{00}(x^2 \partial_x)^2 &\in \text{Diff}_{\text{leC}}^{2,-1,-4,-3,-6}(X) \\ (n-1)x^3 \partial_x &\in \text{Diff}_{\text{leC}}^{1,-1,-3,-2,-4}(X) \\ V_{\text{eff}} &\in \text{Diff}_{\text{leC}}^{0,-2,-4,-2,-4}(X). \end{aligned} \quad (243)$$

Thus, by Proposition 3.1, $\tilde{P}_0 \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X)$, and

$$\tilde{P}_0 = -(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} + L \bmod \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X). \quad (244)$$

Simplifying L modulo $\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X)$ using Proposition 3.2, we get eq. (242). \square

Recall that $b_j, P_{\perp,j}$ were defined in eq. (55).

Proposition 3.4. *For some $\Upsilon_1, \dots, \Upsilon_J \in S \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X)$ which near ∂X are given by $\Upsilon_j = ix^4 \Phi' b_j P_{\perp,j}$, we have*

$$\tilde{P}_1(\sigma) = P_1(\sigma) + \sum_{j=1}^J \Upsilon_j + R \quad (245)$$

for some $R \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^2(X))$ which is supported outside of some neighborhood $U \subset X$ of ∂X . Thus, $\tilde{P}_1 = \{\tilde{P}_1(\sigma)\}_{\sigma \geq 0} \in S \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X)$. If P_1 is classical to order $\beta_1 > 0$, then

$$\tilde{P}_1 \in \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + S \text{Diff}_{\text{leC}}^{2,-1-\beta_1,-3-2\beta_1,-2-\beta_1,-4-2\beta_1}(X). \quad (246)$$

■

Proof. First observe that $\chi \tilde{P}_1 \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^2(X^\circ))$ for any $\chi \in C_c^\infty(X^\circ)$. Now let $P_{j,\text{ext}} \in C^\infty([0, \infty)_{\sigma^2}; S \text{Diff}_{\text{scb}}^{2,-1,-3}(X))$ be equal to $x^4 P_{\perp,j} b_j \partial_x$ near ∂X . Now define

$$\Upsilon_j = e^{-i\Phi}[P_{j,\text{ext}}, e^{+i\Phi}] = e^{-i\Phi}[(1-\chi)P_{j,\text{ext}}, e^{+i\Phi}] + e^{-i\Phi}[\chi P_{j,\text{ext}}, e^{+i\Phi}], \quad (247)$$

where χ is identically equal to one in a sufficiently large open set such that $1-\chi$ is supported in a neighborhood for which eq. (55) applies. Evidently, we have $e^{-i\Phi}[\chi P_{j,\text{ext}}, e^{+i\Phi}] \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^2(X \setminus U))$ for some neighborhood $U \subset X$ of ∂X . On the other hand, $e^{-i\Phi}[(1-\chi)P_{j,\text{ext}}, e^{+i\Phi}] = i(1-\chi)x^4 \Phi' b_j P_{\perp,j}$, so $\Upsilon_j = e^{-i\Phi}[P_{j,\text{ext}}, e^{+i\Phi}]$ near ∂X .

We now write

$$\begin{aligned} \tilde{P}_1 &= P_1 + e^{-i\Phi}[(1-\chi)P_1, e^{+i\Phi}] + e^{-i\Phi}[\chi P_1, e^{+i\Phi}] \\ &= P_1 + \sum_{j=1}^J \Upsilon_j + e^{-i\Phi}[\chi P_1, e^{+i\Phi}]. \end{aligned} \quad (248)$$

Set $R = e^{-i\Phi}[\chi P_1, e^{+i\Phi}]$. Then eq. (245) holds, and $R \in C^\infty([0, \infty)_{\sigma^2}; \text{Diff}^2(X))$ is supported outside of some neighborhood $U \subset X$ of ∂X .

We observe, from eq. (55) (and Proposition 2.10), that

$$P_1 \in S \text{Diff}_{\text{leC}}^{2,-1,-4,-3,-6}(X) + S \text{Diff}_{\text{leC}}^{1,-1,-3,-2,-4}(X) \subseteq S \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X). \quad (249)$$

Since $\Upsilon_1, \dots, \Upsilon_J \in S \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X)$ (and the same holds for R , trivially), we conclude that

$$\tilde{P}_1 \in S \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X), \quad (250)$$

as claimed.

If b_j, b'_j, b''_j satisfy eq. (59), then the conclusion is similar, except now we can write $P_1 \in \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + x^{\beta_1} S \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X)$ and

$$\Upsilon_1, \dots, \Upsilon_J \in \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) + x^{\beta_1} S \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X), \quad (251)$$

which leads to the conclusion eq. (248) as a strengthening of eq. (250). \square

Proposition 3.5. *Given $\delta > 0$ such that $P_2 \in C^\infty([0, \infty)_{\sigma^2}; S \text{Diff}_{\text{scb}}^{2,-1-\delta,-3/2-\delta}(X))$,*

$$\tilde{P}_2 = \{\tilde{P}_2(\sigma)\}_{\sigma \geq 0} \in S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X). \quad (252)$$

If P_2 is classical to order (β_2, β_3) , then we can write

$$\tilde{P}_2 = \{\tilde{P}_2(\sigma)\}_{\sigma \geq 0} \in \text{Diff}_{\text{leC}}^{2,-2,-4,-2,-4}(X) + S \text{Diff}_{\text{leC}}^{2,-1-\beta_2,-3-2\beta_2,-1-\beta_2,-3-2\beta_2}(X) + x^{3/2+\beta_3} S^0(X). \quad (253)$$

Proof. We restrict attention to $\hat{X} = [0, \bar{x})_x \times \partial X$, which suffices by an argument similar to that in the proof of Proposition 3.4.

Let $y = (y_1, \dots, y_{n-1})$ denote local coordinates on ∂X . In terms of these, we can write

$$P_2(\sigma) = x^\delta \left[x^5 c \partial_x^2 + x^4 \sum_{j=1}^{n-1} c_j \partial_x \partial_{y_j} + x^3 \sum_{j,k=1}^{n-1} c_{j,k} \partial_{y_j} \partial_{y_k} + dx^3 \partial_x + \sum_{j=1}^{n-1} d_j x^2 \partial_{y_j} + x^{3/2} e \right] \quad (254)$$

where $\{c, d, e\} \cup \{c_j, d_j, c_{j,k}\}_{j,k=1}^{n-1} \subset C^\infty([0, \infty)_{\sigma^2}; S^0(X))$ and $c_{j,k} = c_{k,j}$. We then have

$$\begin{aligned} P_2(\sigma) - \tilde{P}_2(\sigma) &= x^\delta \left[2ix^5 c \Phi' \partial_x + ix^4 \sum_{j=1}^{n-1} c_j (\Phi' \partial_{y_j} + \partial_{y_j} \Phi \partial_x) + 2ix^3 \sum_{j,k=1}^{n-1} c_{j,k} \partial_{y_j} \Phi \partial_{y_k} - x^5 c \Phi' \Phi' \right] \\ &\quad + x^\delta \left[ix^5 c \Phi'' + x^4 \sum_{j=1}^{n-1} c_j (i \partial_{y_j} \Phi' - \Phi' \partial_{y_j} \Phi) + x^3 \sum_{j,k=1}^{n-1} c_{j,k} (i \partial_{y_j} \partial_{y_k} \Phi - \partial_{y_j} \Phi \partial_{y_k} \Phi) \right] \\ &\quad + x^\delta \left[id x^3 \Phi' + \sum_{j=1}^{n-1} id_j x^2 \partial_{y_j} \Phi \right], \quad (255) \end{aligned}$$

i.e., since Φ does not depend on tangential coordinates,

$$P_2(\sigma) - \tilde{P}_2(\sigma) = x^\delta \left[2ix^5 c \Phi' \partial_x + ix^4 \sum_{j=1}^{n-1} c_j \Phi' \partial_{y_j} - x^5 c \Phi' \Phi' + ix^5 c \Phi'' + id x^3 \Phi' \right]. \quad (256)$$

It follows from eq. (234) that

$$x^{3+\delta} d \Phi' \in \Psi_{\text{leC}}^{0,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X), \quad (257)$$

$x^{5+\delta} c \Phi' \Phi' \in \Psi_{\text{leC}}^{0,-1-\delta,-4-2\delta,-1-\delta,-4-2\delta}(X)$ by eq. (232), and $x^{5+\delta} c \Phi'' \in \Psi_{\text{leC}}^{0,-2-\delta,-5-2\delta,-2-\delta,-5-2\delta}(X)$ by eq. (234). On the other hand,

$$x^{5+\delta} \Phi' \partial_x, x^{4+\delta} \Phi' \partial_{y_j} \in \Psi_{\text{leC}}^{1,-1-\delta,-4-2\delta,-2-\delta,-5-2\delta}(X) \quad (258)$$

by eq. (234). Combining the observations above, in particular eq. (257) and eq. (258), we conclude that

$$P_2 - \tilde{P}_2 \in S \text{Diff}_{\text{leC}}^{1,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X). \quad (259)$$

By eq. (254) (and Proposition 2.10, eq. (145)),

$$\begin{aligned} P_2 &\in S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-2-\delta,-4-2\delta}(X) + S \text{Diff}_{\text{leC}}^{0,-3/2-\delta,-3-2\delta,-3/2-\delta,-3-2\delta}(X) \\ &\subset S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-3/2-\delta,-3-2\delta}(X). \quad (260) \end{aligned}$$

We conclude eq. (252) from eq. (259) and eq. (260).

If eq. (60) holds, then instead of eq. (259) we conclude

$$P_2 - \tilde{P}_2 \in \text{Diff}_{\text{leC}}^{1,-2,-5,-2,-5}(X) + S \text{Diff}_{\text{leC}}^{1,-1-\beta_2,-3-2\beta_2,-1-\beta_2,-3-2\beta_2}(X), \quad (261)$$

and instead of eq. (260) we have

$$P_2 \in \text{Diff}_{\text{leC}}^{2,-2,-4,-2,-4}(X) + S \text{Diff}_{\text{leC}}^{2,-1-\beta_2,-3-2\beta_2,-3/2-\beta_2,-3-2\beta_2}(X) + x^{3/2+\beta_3} S^0(X). \quad (262)$$

Equation (253) follows from eq. (261) and eq. (262). \square

Proposition 3.6. $\tilde{P} \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X) + S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X)$, with $\tilde{P} = \tilde{P}_0 \bmod S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X)$. If P_1, P_2 are classical to orders β_1 and (β_2, β_3) respectively, then

$$\tilde{P} \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X) + x^{\beta_1} S\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + x^{\beta_2} S\text{Diff}_{\text{leC}}^{2,-1,-3,-1,-3}(X) + x^{3/2+\beta_3} S^0(X). \quad (263)$$

Moreover,

$$\begin{aligned} \tilde{P} = & -(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} + 2ix\sqrt{\sigma^2 + \mathbb{Z}x} \left(x\partial_x - \frac{n-1}{2} + \frac{\mathbb{Z}}{4} \frac{x}{\sigma^2 + \mathbb{Z}x} \right) \\ & \bmod \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X). \end{aligned} \quad (264)$$

Thus,

$$\begin{aligned} \tilde{P} = & -(x^2 \partial_x)^2 + x^2 \Delta_{\partial X} + 2ix(\sigma^2 + \mathbb{Z}x)^{1/2} x\partial_x \\ & \bmod \text{Diff}_{\text{leC}}^{2,-1,-3,-1,-3}(X) + S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X). \end{aligned} \quad (265)$$

■

Proof. We have $\tilde{P}(\sigma) = \tilde{P}_0 + \tilde{P}_1 + \tilde{P}_2$. As seen in Proposition 3.4 and Proposition 3.5, $\tilde{P}_1 \in S\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X)$ and $\tilde{P}_2 \in S\text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X)$, so

$$\tilde{P}_1 + \tilde{P}_2 \in S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X) \quad (266)$$

(where we are using $\delta < 1/2$). Likewise, if P_1, P_2 are classical to orders β_1 and (β_2, β_3) then Proposition 3.4 and Proposition 3.5 yield

$$\begin{aligned} \tilde{P}_1 + \tilde{P}_2 \in & \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + x^{\beta_1} S\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + x^{\beta_2} S\text{Diff}_{\text{leC}}^{2,-1,-3,-1,-3}(X) \\ & + x^{3/2+\beta_3} S^0(X). \end{aligned} \quad (267)$$

By Proposition 3.3, $\tilde{P}_0 \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X)$, so \tilde{P} is in the claimed spaces.

Furthermore, by eq. (266), $\tilde{P} = \tilde{P}_0 \bmod S\text{Diff}_{\text{leC}}^{2,-1,-3-2\delta,-1-\delta,-3-2\delta}(X)$. Combining with Proposition 3.3, we get eq. (264). \square

To conclude this discussion, we let

$$N(\tilde{P}) = \{N(\tilde{P})(\sigma)\}_{\sigma \geq 0}, \quad N(\tilde{P})(\sigma) = 2ix\sqrt{\sigma^2 + \mathbb{Z}x} \left(x\partial_x - \frac{n-1}{2} + \frac{\mathbb{Z}}{4} \frac{x}{\sigma^2 + \mathbb{Z}x} \right) \quad (268)$$

denote the *leC-normal operator*, defined initially near ∂X . To avoid technicalities, we extend $N(\tilde{P})(\sigma)$ to a differential operator on X° such that, in any compact subset of X° , $N(\tilde{P})(\sigma)$ depends smoothly on $E = \sigma^2$, all the way down to $\sigma = 0$. Thus:

Proposition 3.7. $N(\tilde{P}) \in \text{Diff}_{\text{b,leC}}^{1,-1,-3}(X)$. $\blacksquare \square$

The following proposition justifies the term “leC-normal operator:”

Proposition 3.8. $N(\tilde{P}) - \tilde{P} \in S\text{Diff}_{\text{leC}}^{2,0,-2,-1-\delta,-3-2\delta}(X) \subseteq S\text{Diff}_{\text{b,leC}}^{2,-1-\delta,-3-2\delta}(X)$. If P_1, P_2 are classical to orders β_1 and (β_2, β_3) respectively, then

$$\begin{aligned} N(\tilde{P}) - \tilde{P} \in & \text{Diff}_{\text{leC}}^{2,0,-2,-2,-4}(X) + x^{\beta_1} S\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + x^{\beta_2} S\text{Diff}_{\text{leC}}^{2,-1,-3,-1,-3}(X) \\ & + x^{3/2+\beta_3} S^0(X). \end{aligned} \quad (269)$$

■

Proof. We have $N(\tilde{P}) - \tilde{P}(\sigma) = (N(\tilde{P}) - \tilde{P}_0) - \tilde{P}_1 - \tilde{P}_2$. We first check that $N(\tilde{P}) - \tilde{P} \in S\text{Diff}_{\text{leC}}^{2,0,-2,-1-\delta,-3-2\delta}(X)$.

- By eq. (266), $\tilde{P}_1 + \tilde{P}_2 \in S \text{Diff}_{\text{leC}}^{2,0,-2,-1-\delta,-3-2\delta}(X)$, and by eq. (243) the same holds for $\tilde{P}_0 - L$, so it suffices to check that

$$N(\tilde{P}) - L \in \text{Diff}_{\text{leC}}^{2,0,-2,-1-\delta,-3-2\delta}(X). \quad (270)$$

Indeed, by Proposition 3.2,

$$\begin{aligned} N(\tilde{P}) - L &\in \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X) + x \text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X) + \sigma^2 x / (\sigma^2 + Zx) \text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X) \\ &\subseteq \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X) + \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) + \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) \\ &= \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) \\ &\subset \text{Diff}_{\text{leC}}^{2,0,-2,-1-\delta,-3-2\delta}(X). \end{aligned} \quad (271)$$

If P_1, P_2 are classical to orders $\beta_1, (\beta_2, \beta_3)$, then we instead get eq. (269). \square

We now consider the $L_{\text{sc}}^2(X) = L^2(X, g_0)$ -based adjoint \tilde{P}^* , defined such that

$$\int_X f^* \tilde{P} g \, d\text{Vol}_{g_0} = \int_X (\tilde{P}^* f)^* g \, d\text{Vol}_{g_0} \quad (272)$$

for all $f, g \in \mathcal{S}(X)$. This is a family of differential operators and, by Proposition 2.10, an element of $S \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X) \subset \Psi_{\text{leC}}^{2,0,-2,-1,-3}(X)$. We form the differential operators

$$\Re \tilde{P} = \frac{1}{2}(\tilde{P} + \tilde{P}^*), \quad \Im \tilde{P} = \frac{1}{2i}(\tilde{P} - \tilde{P}^*), \quad (273)$$

the self-adjoint and anti-self-adjoint parts of \tilde{P} .

Proposition 3.9. *For any exactly conic metric g_0 , there exists a differential operator $R = R_{g_0} \in S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X)$ such that*

$$\tilde{P}^* = \tilde{P} + (P_1^* - P_1) + \sum_{j=1}^J (\Upsilon_j^* - \Upsilon_j) + R, \quad (274)$$

and, near ∂X ,

$$\Upsilon_j^* \in x^2 \varrho_{\text{tf}} S^0(X_{\text{res}}^{\text{sp}}) \text{Diff}^1(\partial X) \subset S \text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) \quad (275)$$

$$P_1^* = \sum_{j=1}^J [-x^2 P_{\perp,j}^* (b_j^* x^2 \partial_x + x b_{*,j}) + x^3 b_j'^* P_{\partial X,j}^* + x^2 b_j''^* Q_{\partial X,j}^*] \in S \text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) \quad (276)$$

for some $b_{*,1}, \dots, b_{*,J} \in S^0(X)$. \blacksquare

Proof. It clearly suffices to restrict attention to $\hat{X} = [0, \bar{x}]_x \times \partial X$, that is to compute the formal adjoint of \tilde{P} with respect to

$$L_{\text{sc}}^2(\hat{X}) = L^2([0, \bar{x}]_x \times \partial X_y, x^{-(n+1)} dx d\text{Vol}_{g_{\partial X}}(y)) \quad (277)$$

up to the required order.

Begin with \tilde{P}_0 , which we rewrite as

$$\tilde{P}_0 = -(1 + x a_{00}) \left(x^2 \partial_x - \frac{x(n-1)}{2} \right)^2 + x^2 \Delta_{\partial X} - x^4 (n-1) a_{00} \partial_x + L(\sigma) + W \quad (278)$$

for $W \in x^2 C^\infty(X_{\text{res}}^{\text{sp}})$.

- We see that $x^2 \Delta_{\partial X} = x^2 \Delta_{g_{\partial X}}$ is formally self-adjoint on $L_{\text{sc}}^2([0, \bar{x}] \times \partial X)$, and
- the adjoint of W is its complex conjugate $W^* \in x^2 C^\infty(X_{\text{res}}^{\text{sp}})$.
- We also have $a_{00} x^4 \partial_x \in S \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X)$, thus $(a_{00} x^4 \partial_x)^* \in S \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X)$.

- On the other hand, $x^2\partial_x - x(n-1)/2$ is formally anti-self-adjoint on $L_{\text{sc}}^2([0, \bar{x}) \times \partial X)$, so

$$\begin{aligned} \left[(1 + xa_{00}) \left(x^2\partial_x - \frac{x(n-1)}{2} \right)^2 \right]^* &= \left(x^2\partial_x - \frac{x(n-1)}{2} \right)^2 (1 + xa_{00}) \\ &= (1 + xa_{00}) \left(x^2\partial_x - \frac{x(n-1)}{2} \right)^2 + a_{00} \left[\left(x^2\partial_x - \frac{x(n-1)}{2} \right)^2, x \right] \\ &= (1 + xa_{00}) \left(x^2\partial_x - \frac{x(n-1)}{2} \right)^2 + 2a_{00}x^4\partial_x + (3-n)a_{00}x^3. \end{aligned}$$

- By the same computation opening this subsection, $N(\tilde{P})$ is formally self-adjoint. So, by Proposition 3.2,

$$\begin{aligned} L^* &= L + \left[N(\tilde{P}), (1 + xa_{00}) \left(1 - \frac{\sigma^2 a_{00} x}{2(\sigma^2 + \mathbb{Z}x)} \right) \right] \bmod \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X) \\ &= L + 2ix\sqrt{\sigma^2 + \mathbb{Z}x} \left[x\partial_x, (1 + xa_{00}) \left(1 - \frac{\sigma^2 a_{00} x}{2(\sigma^2 + \mathbb{Z}x)} \right) \right] \bmod \text{Diff}_{\text{leC}}^{1,-2,-5,-3,-6}(X) \\ &= L \bmod \text{Diff}_{\text{leC}}^{1,-2,-5,-2,-5}(X) = L \bmod \text{Diff}_{\text{leC}}^{1,-2,-4,-2,-4}(X). \end{aligned} \quad (279)$$

So,

$$\tilde{P}_0^* = \tilde{P}_0 \bmod S \text{Diff}_{\text{leC}}^{1,-2,-4,-2,-4}(X). \quad (280)$$

On the other hand, we trivially have from Proposition 3.5 that

$$\tilde{P}_2^* \in S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X). \quad (281)$$

Now define $R = \tilde{P}^* - \tilde{P} - (P_1^* - P_1) - \sum_{j=1}^J (\Upsilon_j^* - \Upsilon_j)$, so that eq. (274) holds by construction. Rearranging this definition,

$$\begin{aligned} R &= (\tilde{P}_0^* - \tilde{P}_0) + (\tilde{P}_1^* - \tilde{P}_1) + (\tilde{P}_2^* - \tilde{P}_2) - (P_1^* - P_1) - \sum_{j=1}^J (\Upsilon_j^* - \Upsilon_j) \\ &= (\tilde{P}_0^* - \tilde{P}_0) + (\tilde{P}_2^* - \tilde{P}_2) + \left([\tilde{P}_1 - P_1 - \sum_{j=1}^J \Upsilon_j]^* - [\tilde{P}_1 - P_1 - \sum_{j=1}^J \Upsilon_j] \right). \end{aligned} \quad (282)$$

Thus, using eq. (280), eq. (281), Proposition 3.5, and Proposition 3.4,

$$\begin{aligned} R &\in S \text{Diff}_{\text{leC}}^{1,-2,-4,-2,-4}(X) + S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X) \\ &= S \text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X), \end{aligned} \quad (283)$$

as claimed.

Equation (275) follows from the observation that $\Upsilon_j^* = -ix^4\Phi'b_j^*P_{\perp,j}^*$ near ∂X , where $P_{\perp,j}^*$ is computed using the $L^2(\partial X, g_{\partial X})$ -inner product. On the other hand,

$$\begin{aligned} P_1^* &= \sum_{j=1}^J \left[P_{\perp,j}^* (x^2 b_j^* (x^2 \partial_x)^* + [(x^2 \partial_x)^*, x^2 b_j^*]) + x^3 b_j^* P_{\partial X,j}^* + x^2 b_j^{''*} Q_{\partial X,j}^* \right] \\ &= \sum_{j=1}^J \left[P_{\perp,j}^* (x^2 b_j^* (-x^2 \partial_x + x(n-1)) - x[x\partial_x, x^2 b_j^*]) + x^3 b_j^* P_{\partial X,j}^* + x^2 b_j^{''*} Q_{\partial X,j}^* \right] \\ &= \sum_{j=1}^J \left[P_{\perp,j}^* (x^2 b_j^* (-x^2 \partial_x + x(n-1)) - 2x^3 b_j^* - x^3 (x\partial_x b_j^*)) + x^3 b_j^* P_{\partial X,j}^* + x^2 b_j^{''*} Q_{\partial X,j}^* \right] \end{aligned} \quad (284)$$

near ∂X , where b_j, b_j', b_j'' are as in eq. (55). Equation (276) follows from this, for some choice of $b_{*,j}$. \square

We see from the above that

$$\begin{aligned} \Im \tilde{P} &\in S\text{Diff}_{\text{leC}}^{1,-1,-4,-2,-5}(X) + S\text{Diff}_{\text{leC}}^{2,-1,-3,-2,-4}(X) + S\text{Diff}_{\text{leC}}^{2,-1-\delta,-3-2\delta,-1-\delta,-3-2\delta}(X) \\ &\subset S\text{Diff}_{\text{leC}}^{2,-1,-3,-1-\delta,-3-2\delta}(X). \end{aligned} \quad (285)$$

Thus, $\Im \tilde{P}$ is one order lower than \tilde{P} at sf and ff and slightly lower order at bf and tf, and we have been entirely explicit about the leading terms of $\Im \tilde{P}$ (namely $\Im P_1$ and $\Im \Upsilon_j$) at sf and ff, the remainder $(2i)^{-1}R$ being slightly more than one order lower than \tilde{P} at both faces.

4. THE SITUATION AT ZERO ENERGY

In this section we consider $P(0)$ and $\tilde{P}(0)$ in some detail. Specifically, we apply [Vas21a, Theorem 1.1] to study the strong limit

$$R(E = 0; Z \pm i0) = \text{slim}_{\epsilon \rightarrow 0+} R(E = 0; Z \pm i\epsilon) \quad (286)$$

used in the statement of Theorem 1.2 to characterize the resolvent output at zero energy. The mapping properties of this operator will be used in §6 in order to prove the smoothness of the output of the conjugated resolvent at positive energy all the way down to zero energy (as used e.g. in Corollary 1.4).

Recall that $x_{1/2} = 2^{-1/2}x^{1/2}$. Then, from the form eq. (53) of P_0 ,

$$\begin{aligned} P_0(0) &= -(1 + 2x_{1/2}^2 a_{00}(0))(x_{1/2}^3 \partial_{x_{1/2}})^2 + 4x_{1/2}^4 \Delta_{\partial X} + 2x_{1/2}^5 [a(0) + n - 1] \partial_{x_{1/2}} - 2Zx_{1/2}^2 \\ &= -x_{1/2}^2 (1 + 2x_{1/2}^2 a_{00}(0))(x_{1/2}^2 \partial_{x_{1/2}})^2 + 4x_{1/2}^4 \Delta_{\partial X} + 2x_{1/2}^5 \left[a(0) + n - \frac{3}{2} - x_{1/2}^2 a_{00}(0) \right] \partial_{x_{1/2}} - 2Zx_{1/2}^2. \end{aligned} \quad (287)$$

(The extra $-x_{1/2}^3 \partial_{x_{1/2}}$ in eq. (287) is the source of the $(\sigma^2 + Zx)^{-1/4}$ term in eq. (19).) Therefore

$$\begin{aligned} x_{1/2}^{-1-n/2} P_0(0) x_{1/2}^{(n-2)/2} &= -(1 + 2x_{1/2}^2 a_{00}(0))(x_{1/2}^2 \partial_{x_{1/2}})^2 + 4x_{1/2}^2 \Delta_{\partial X} - 2Z \\ &+ x_{1/2}^3 \left[2a(0) + n - 1 - 2(n-1)x_{1/2}^2 a_{00}(0) \right] \partial_{x_{1/2}} + (n-2) \left[a(0) + \frac{3n}{4} - \frac{3}{2} - \frac{n+2}{2} x_{1/2}^2 a_{00}(0) \right] x_{1/2}^2. \end{aligned} \quad (288)$$

Thus, $x_{1/2}^{-1-n/2} \tilde{P}_0(0) x_{1/2}^{-1+n/2} \in \text{Diff}_b(X_{1/2})$ has the same form as the conjugated spectral family at positive energy, except with respect to $x_{1/2}$ instead of x (and with an extra short-range potential in eq. (288)).

We generalize this observation:

Proposition 4.1. *If g is an asymptotically conic metric on X , and if we set $g_{1/2} = x_{1/2}^2 g$, then there exists some $V_{\text{eff}} \in x^2 S^0(X)$ such that*

$$x_{1/2}^{1-n/2} \Delta_g x_{1/2}^{-1+n/2} = x_{1/2}^2 \Delta_{g_{1/2}} + V_{\text{eff}}. \quad (289)$$

holds. ■

Proof. We first want to show that $x_{1/2}^{1-n/2} \Delta_g x_{1/2}^{-1+n/2} - x_{1/2}^2 \Delta_{g_{1/2}}$ is zeroth order (and therefore a function on X°). Indeed, $\Delta_g = x_{1/2}^2 \Delta_{g_{1/2}} + (n-2)x_{1/2} \nabla_{g_{1/2}} x_{1/2}$, so

$$x_{1/2}^{1-n/2} \Delta_g x_{1/2}^{-1+n/2} = x_{1/2}^2 \Delta_{g_{1/2}} + x_{1/2}^{3-n/2} \Delta_{g_{1/2}} x_{1/2}^{-1+n/2} + (n-2)x_{1/2} g_{1/2} (dx_{1/2}, dx_{1/2}^{(n-2)/2}). \quad (290)$$

It is therefore the case that $x_{1/2}^{1-n/2} \Delta_g x_{1/2}^{-1+n/2} - x_{1/2}^2 \Delta_{g_{1/2}} = V_{\text{eff}}$ for

$$V_{\text{eff}} = x_{1/2}^{3-n/2} \Delta_{g_{1/2}} x_{1/2}^{-1+n/2} + (n-2)x_{1/2} g_{1/2} (dx_{1/2}, dx_{1/2}^{(n-2)/2}). \quad (291)$$

Since $g_{1/2}$ is an asymptotically conic metric on $X_{1/2}$ (see below), the fact that $V_{\text{eff}} \in x^2 S^0(X)$ can be read off eq. (291), but we check in local coordinates.

It suffices to restrict attention to a neighborhood of ∂X . Let y_1, \dots, y_{n-1} denote a local system of coordinates on ∂X . Then, using $g = x_{1/2}^{-2} g_{1/2}$,

$$\begin{aligned} \Delta_g &= -\frac{x_{1/2}^n}{|g_{1/2}|^{1/2}} \partial_i (x_{1/2}^{2-n} |g_{1/2}|^{1/2} g_{1/2}^{ij} \partial_j) \\ &= -\frac{x_{1/2}^2}{|g_{1/2}|^{1/2}} \partial_i (|g_{1/2}|^{1/2} g_{1/2}^{ij} \partial_j) + (n-2) x_{1/2} g_{1/2}^{0i} \partial_i, \end{aligned} \quad (292)$$

where $\partial_0 = \partial_{x_{1/2}}$ and $\partial_i = \partial_{y_i}$ for $i = 1, \dots, n-1$.

Conjugating the right-hand side of eq. (292) by $x_{1/2}^{(n-2)/2}$, we see that

$$x_{1/2}^{1-n/2} \Delta_g x_{1/2}^{-1+n/2} = x_{1/2}^2 \Delta_{g_{1/2}} - \frac{x_{1/2}}{|g_{1/2}|^{1/2}} \frac{n-2}{2} \partial_i (|g_{1/2}|^{1/2} g_{1/2}^{i0}) + 4^{-1} n(n-2) g_{1/2}^{00}. \quad (293)$$

Since $g_{1/2}^{00} \in x_{1/2}^4 S^0(X_{1/2}) = x^2 S^0(X)$, $4^{-1} n(n-2) g_{1/2}^{00} \in x_{1/2}^4 S^0(X_{1/2})$.

Likewise, we see that $|g_{1/2}|^{-1/2} \partial_0 |g_{1/2}|^{1/2} \in x_{1/2}^{-1} S^0(X_{1/2})$, $|g_{1/2}|^{-1/2} \partial_i |g_{1/2}|^{1/2} \in S^0(X_{1/2})$ for $i \neq 0$, and

$$\partial_0 g_{1/2}^{00}, \partial_i g_{1/2}^{i0} \in x_{1/2}^3 S^0(X_{1/2}) \quad (294)$$

for $i \neq 0$. So,

$$V_{\text{eff}} = -\frac{x_{1/2}}{|g_{1/2}|^{1/2}} \frac{n-2}{2} \partial_i (|g_{1/2}|^{1/2} g_{1/2}^{i0}) + 4^{-1} n(n-2) g_{1/2}^{00} \in x_{1/2}^4 S^0(X_{1/2}) = x^2 S^0(X). \quad (295)$$

□

Observe that if g is an asymptotically conic metric on X , then $g_{1/2}$ is an asymptotically conic metric on $X_{1/2}$. Indeed, $g_{1/2}$ is certainly a Riemannian metric on $X^\circ = X_{1/2}^\circ$, and it is a sum of $x_{1/2}^2 g_0$, which is exactly conic on $X_{1/2}$, and terms in

$$x^2 C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \subset x_{1/2}^2 C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^* X_{1/2}), \quad (296)$$

$$x^{2+\delta} S^0(X; {}^{\text{sc}}\text{Sym}^2 T^* X) \subset x_{1/2}^{2+2\delta} S^0(X; {}^{\text{sc}}\text{Sym}^2 T^* X_{1/2}). \quad (297)$$

Consequently, by Proposition 4.1, if P is the spectral family of an attractive Coulomb-like Schrödinger operator, then

$$x_{1/2}^{-1-n/2} P(0) x_{1/2}^{-1+n/2} = x_{1/2}^{-2} (x_{1/2}^{-1+n/2})^{-1} P(0) x_{1/2}^{-1+n/2} \quad (298)$$

is a member $P_{\text{zero}}(2Z) = \Delta_{g_{1/2}} - 2Z + W$ of the spectral family $\{P_{\text{zero}}(\zeta) = P_{\text{zero}}(0) - \zeta\}_{\zeta \geq 0}$ of a Schrödinger operator

$$P_{\text{zero}}(0) = \Delta_{g_{1/2}} + W \quad (299)$$

on $X_{1/2}$, where the potential $W \in x S^0(X) = x_{1/2}^2 S^0(X_{1/2})$ is short-range.

Thus, $\tilde{P}_0(2Z)$ satisfies the hypotheses of [Vas21a, §3], with $2^{1/2} Z^{1/2}$ in place of σ and $X_{1/2}$ in place of X . Moreover, as seen from eq. (70) with $a = 0$, the phase $\Phi(-; 0)$ is just that used by Vasy's in his conjugation. Thus, $\tilde{P}(0)$ has the form of Vasy's conjugated operator (with $2^{1/2} Z^{1/2}$ in place of σ and $X_{1/2}$ in place of X). In order to denote the Z dependence of $P(0)$ and $\tilde{P}(0)$, we write

$$P(0) = P(0; Z) = x_{1/2}^{1-n/2} P_{\text{zero}}(2Z) x_{1/2}^{1-n/2} \quad (300)$$

and $\tilde{P}(0) = \tilde{P}(0; Z)$. For $\epsilon > 0$, let $P(0; Z + i\epsilon)$ denote $P(0)$ with Z replaced by $Z + i\epsilon$. Since $\mathcal{S}(X_{1/2}) = \mathcal{S}(X)$ and $\mathcal{S}'(X_{1/2}) = \mathcal{S}'(X)$, the limiting absorption principle (as in [Mel94]) applies in the following form: for $\epsilon > 0$, the resolvent

$$R_0(2Z + 2i\epsilon) = x_{1/2}^{-1-n/2} R(0; Z + i\epsilon) x_{1/2}^{(n-2)/2} : \mathcal{S}(X) \rightarrow \mathcal{S}'(X) \quad (301)$$

of $P_{\text{zero}}(0)$ (evaluated at “energy” $\zeta = 2Z + 2i\epsilon$) – defined e.g. via the functional calculus – admits a strong limit $x_{1/2}^{-1-n/2} R(0; Z + i0) x_{1/2}^{-1+n/2} : \mathcal{S}(X) \rightarrow \mathcal{S}'(X)$. Using [Vas21a, Theorem 1.1], we can construct this resolvent as a map between suitable Sobolev spaces using the conjugated perspective:

Proposition 4.2. *If $P(\sigma)$ is the spectral family of an attractive Coulomb-like Schrödinger operator on X , then, for any $m, \varsigma, \ell \in \mathbb{R}$ satisfying $\ell < -3/2 < \varsigma$,*

$$\begin{aligned} \tilde{P}(0; Z) : \{u \in H_{\text{scb}}^{m, \varsigma+n/2, \ell+n/2}(X_{1/2}) : \tilde{P}(0)u \in H_{\text{scb}}^{m-2, \varsigma+3+n/2, \ell+3+n/2}(X_{1/2})\} \\ \rightarrow H_{\text{scb}}^{m-2, \varsigma+3+n/2, \ell+3+n/2}(X_{1/2}) \end{aligned} \quad (302)$$

is invertible, and the inverse

$$\tilde{R}_+(0; Z) : H_{\text{scb}}^{m-2, \varsigma+3+n/2, \ell+3+n/2}(X_{1/2}) \rightarrow H_{\text{scb}}^{m, \varsigma+n/2, \ell+n/2}(X_{1/2}) \quad (303)$$

is related to $R(0; Z + i0)$ by the formula $e^{+i\Phi(-;0)} \tilde{R}_+(0; Z) e^{-i\Phi(-;0)} f = R(0; Z + i0) f$, which holds for all $f \in \mathcal{S}(X)$. ■

Proof. Suppose that we are given $m, \varsigma_{1/2}, \ell_{1/2} \in \mathbb{R}$ with $\ell_{1/2} < -1/2 < \varsigma_{1/2}$. Then, combining the observations above and [Vas21a, Theorem 1.1],

$$\begin{aligned} e^{-i\Phi(-;0)} P_{\text{zero}}(2Z) e^{+i\Phi(-;0)} &= x_{1/2}^{-1-n/2} \tilde{P}(0; Z) x_{1/2}^{-1+n/2} : \\ \{u \in H_{\text{scb}}^{m, \varsigma_{1/2}, \ell_{1/2}}(X_{1/2}) : x_{1/2}^{-1-n/2} \tilde{P}(0; Z) x_{1/2}^{-1+n/2} u \in H_{\text{scb}}^{m-2, \varsigma_{1/2}+1, \ell_{1/2}+1}(X_{1/2})\} \\ &\rightarrow H_{\text{scb}}^{m-2, \varsigma_{1/2}+1, \ell_{1/2}+1}(X_{1/2}) \end{aligned} \quad (304)$$

is invertible, defining a continuous linear map

$$\tilde{R}_0(Z) : H_{\text{scb}}^{m-2, \varsigma_{1/2}+1, \ell_{1/2}+1}(X_{1/2}) \rightarrow H_{\text{scb}}^{m, \varsigma_{1/2}, \ell_{1/2}}(X_{1/2}), \quad (305)$$

and this is related to the limiting resolvent $R_0(2Z + i0)$ of the spectral family $P_{\text{zero}} = \{P_{\text{zero}}(0) - \zeta\}_{\zeta \geq 0}$ of the Schrödinger operator $P_{\text{zero}}(0)$ on $X_{1/2}$ by

$$e^{+i\Phi(-;0)} \tilde{R}_0(0; Z) e^{-i\Phi(-;0)} f = R_0(2Z + i0) f = x_{1/2}^{-1-n/2} R(0; Z + i0) x_{1/2}^{-1+n/2} f, \quad (306)$$

which holds for all $f \in \mathcal{S}(X)$.

Call the (Z -dependent) domain and codomain of eq. (304)

$$\begin{aligned} \tilde{\mathcal{X}}_{m, \varsigma_{1/2}, \ell_{1/2}} &= \{u \in H_{\text{scb}}^{m, \varsigma_{1/2}, \ell_{1/2}}(X_{1/2}) : x_{1/2}^{-1-n/2} \tilde{P}(0; Z) x_{1/2}^{-1+n/2} u \in H_{\text{scb}}^{m-2, \varsigma_{1/2}+1, \ell_{1/2}+1}(X_{1/2})\} \\ &= \{u \in H_{\text{scb}}^{m, \varsigma_{1/2}, \ell_{1/2}}(X_{1/2}) : \tilde{P}(0; Z) x_{1/2}^{-1+n/2} u \in H_{\text{scb}}^{m-2, \varsigma_{1/2}+2+n/2, \ell_{1/2}+2+n/2}(X_{1/2})\} \\ \tilde{\mathcal{Y}}_{m, \varsigma_{1/2}, \ell_{1/2}} &= H_{\text{scb}}^{m-2, \varsigma_{1/2}+1, \ell_{1/2}+1}(X_{1/2}). \end{aligned}$$

Setting

$$\begin{aligned} \mathcal{X}_{m, \varsigma, \ell} &= x_{1/2}^{-1+n/2} \tilde{\mathcal{X}}_{m, \varsigma_{1/2}, \ell_{1/2}} \\ &= \{u \in H_{\text{scb}}^{m, \varsigma+n/2, \ell+n/2}(X_{1/2}) : \tilde{P}(0; Z) u \in H_{\text{scb}}^{m-2, \varsigma+3+n/2, \ell+3+n/2}(X_{1/2})\} \end{aligned} \quad (307)$$

$$\mathcal{Y}_{m, \varsigma, \ell} = H_{\text{scb}}^{m-2, \varsigma+3+n/2, \ell+3+n/2}(X_{1/2}) \quad (308)$$

for $\varsigma = \varsigma_{1/2} - 1$ and $\ell = \ell_{1/2} - 1$, we have a commutative diagram

$$\begin{array}{ccc} \tilde{\mathcal{X}}_{m, \varsigma_{1/2}, \ell_{1/2}} & \xrightarrow{x_{1/2}^{-1-n/2} \tilde{P}(0) x_{1/2}^{-1+n/2}} & \tilde{\mathcal{Y}}_{m, \varsigma_{1/2}, \ell_{1/2}} \\ \downarrow \times x_{1/2}^{-1+n/2} & & \downarrow \times x_{1/2}^{1+n/2} \\ \mathcal{X}_{m, \varsigma, \ell} & \xrightarrow{\tilde{P}(0)} & \mathcal{Y}_{m, \varsigma, \ell} \end{array} \quad (309)$$

in the category of Banach spaces. The vertical arrows are manifestly isomorphisms, and as observed the top horizontal arrow is as well.

Hence, $\tilde{P}(0) = \tilde{P}(0; Z)$ is invertible, and the inverse $\tilde{R}_+(0) = \tilde{R}_+(0; Z)$ has the properties specified in the proposition. \square

And for the b-Sobolev spaces:

Proposition 4.3. *If $P(\sigma)$ is the spectral family of an attractive Coulomb-like Schrödinger operator on X , for any $m, \ell = 2l \in \mathbb{R}$ satisfying $\ell < -3/2 < m + \ell$,*

$$\begin{aligned} \tilde{P}(0; Z) : \{u \in H_b^{m, \ell+n/2}(X_{1/2}) : \tilde{P}(0; Z)u \in H_b^{m, \ell+3+n/2}(X_{1/2})\} &\rightarrow H_b^{m, \ell+3+n/2}(X_{1/2}) \\ : \{u \in H_b^{m, l}(X) : \tilde{P}(0; Z)u \in H_b^{m, l+3/2}(X)\} &\rightarrow H_b^{m, l+3/2}(X) \end{aligned} \quad (310)$$

is invertible. \blacksquare

Proof. The equality

$$\{u \in H_b^{m, \ell+n/2}(X_{1/2}) : \tilde{P}(0; Z)u \in H_b^{m, \ell+3+n/2}(X_{1/2})\} = \{u \in H_b^{m, l}(X) : \tilde{P}(0; Z)u \in H_b^{m, l+3/2}(X)\} \quad (311)$$

follows from Lemma 2.18.

First of all, setting $\varsigma = m + \ell$,

$$\{u \in H_b^{m, \ell+n/2}(X_{1/2}) : \tilde{P}(0; Z)u \in H_b^{m, \ell+3+n/2}(X_{1/2})\} \subset \mathcal{X}_{m, \varsigma, \ell}, \quad (312)$$

and so eq. (310) is injective (by Proposition 4.2). Conversely,

$$H_b^{m, l+3/2}(X) \subset \tilde{\mathcal{Y}}_{m, \varsigma, \ell}, \quad (313)$$

so given any $f \in H_b^{m, l+3/2}(X)$ there exists a $u \in \tilde{\mathcal{X}}_{m, \varsigma, \ell}$ such that $\tilde{P}(0; Z)u = f$. Thus, $u \in H_b^{m, \ell+n/2}(X_{1/2})$, and we already know that $f = \tilde{P}(0; Z)u \in H_b^{m, \ell+3+n/2}(X_{1/2})$, so u is actually in the codomain of eq. (310). Thus, eq. (310) is surjective. \square

5. SYMBOLIC ESTIMATES

We now proceed to establish quantitative control of $u \in \mathcal{S}'(X)$ in terms of $\tilde{P}u$ microlocally in the symbolic region of the leC-phase space ${}^{\text{leC}}\overline{T}^*X$, meaning at $\text{df} \cup \text{sf} \cup \text{ff}$. Since we make no attempt to be uniform in the $\sigma \rightarrow \infty$ limit, we simply restrict attention to $\sigma \in [0, \Sigma]$ for some arbitrary $\Sigma > 0$, and the estimates will all depend on Σ in some unexamined way. The main result of this section, duplicated below as Proposition 5.14, says:

- for every $\Sigma > 0$, $N \in \mathbb{N}$, and $m, s, \varsigma, l, \ell, s_0, \varsigma_0 \in \mathbb{R}$ satisfying $l < -1/2 < s_0 < s$ and $\ell < -3/2 < \varsigma_0 < \varsigma \leq \ell + s - l$, there exists a constant $C = C(\tilde{P}, \Sigma, N, m, s, \varsigma, l, \ell) > 0$ such that

$$\|u\|_{H_{\text{leC}}^{m, s, \varsigma, l, \ell}} \leq C(\|\tilde{P}u\|_{H_{\text{leC}}^{m-2, s+1, \varsigma+3, l+1, \ell+3}} + \|u\|_{H_{\text{b, leC}}^{-N, l, \ell}}) \quad (314)$$

holds for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$ such that $\|u\|_{H_{\text{leC}}^{-N, s_0, \varsigma_0, -N, -N}(X)(\sigma)} < \infty$.

(As mentioned in the introduction, [Vas21a, §5] suffices for the analysis of the $\sigma \rightarrow \infty$ regime.) The contents of this section should be compared to the contents of [Vas21a, §4], as our argument below is very similar to some of the symbolic computations there. Using Lemma 2.16, in order to prove eq. (314) it suffices to establish quantitative control of u within each member of some finite collection of open subsets of the leC-phase space covering $\text{df} \cup \text{sf} \cup \text{ff}$ — see Figure 5.

By eq. (265),

$$\tilde{P} = -(x^2 \partial_x)^2 + x^2 \triangle_{\partial X} + 2ix(\sigma^2 + Zx)^{1/2} x \partial_x \bmod S \text{Diff}_{\text{leC}}^{2, -1, -3, -1, -3}(X) \quad (315)$$

(near ∂X). Thus, the leC-principal symbol $\sigma_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \in S_{\text{leC}}^{2,0,-2,-1,-3}(X)/S_{\text{leC}}^{1,-1,-3,-1,-3}(X)$ of \tilde{P} has a representative of the form

$$\tilde{p} = \tilde{p}_0 + \tilde{p}_{1,2} \quad (316)$$

for $\tilde{p}_{1,2} \in S_{\text{leC}}^{2,-1,-3,-1,-3}(X)$, where \tilde{p}_0 is a representative of $\sigma_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}_0)$. Thus,

$$\tilde{p}_0 = x^2 \xi_b^2 + x^2 \eta_b^2 - 2x(\sigma^2 + Zx)^{1/2} \xi_b \quad (317)$$

near ∂X . (Recall from §2 that ξ_b is the b-cofiber coordinate dual to x and $\eta = \eta_b \in T^*\partial X$.) The ellipticity of $\tilde{P} \in S\text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X)$ (see Proposition 5.2) at and near df makes establishing control there trivial, so we work on ${}^{\text{leC}}T^*X = {}^{\text{leC}}\overline{T}^*X \setminus \text{df}$ and establish control at $\{x=0\} \subset {}^{\text{leC}}T^*X$. Thus, $\tilde{p}_{1,2}$ is irrelevant for the symbolic considerations below (except those in §5.1), all of which are restricted away from df.

In order to investigate the dynamics away from $\text{bf} \cup \text{tf}$, we introduce new coordinates on $\mathbb{R}_\sigma^+ \times T^*X^\circ$ (over some collar neighborhood of ∂X , not including the boundary itself) by

$$\xi_{\text{sc,leC}} = \xi_b \varrho_{\text{bf}00} \varrho_{\text{tf}00} = \xi_b \varrho_{\text{sf}} \varrho_{\text{ff}} \varrho_{\text{bf}} \varrho_{\text{tf}} \quad (318)$$

$$\eta_{\text{sc,leC}} = \eta \varrho_{\text{bf}00} \varrho_{\text{tf}00} = \eta \varrho_{\text{sf}} \varrho_{\text{ff}} \varrho_{\text{bf}} \varrho_{\text{tf}} \in T^*\partial X, \quad (319)$$

for $\xi_b \in \mathbb{R}$ and $\eta = \eta_b \in T^*\partial X$. More explicitly, $\xi_{\text{sc,leC}} = \xi_b x(\sigma^2 + Zx)^{-1/2}$, $\eta_{\text{sc,leC}} = \eta x(\sigma^2 + Zx)^{-1/2}$. These extend to fiber coordinates on the bundle $[0, \bar{x}]_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}} \rightarrow [0, \bar{x}]_{\text{res}}^{\text{sp}}$, and we write

$$\begin{aligned} {}^{\text{sc,leC}}T_{\partial X}^*X &= (\partial[0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus \text{zf}^\circ) \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}} \\ &= \partial([0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus \text{zf}^\circ) \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}}. \end{aligned} \quad (320)$$

(We do not endow ${}^{\text{sc,leC}}T_{\partial X}^*X$ with any more structure than that of a set.) Let o denote the zero section of $T^*\partial X$. While ${}^{\text{sc,leC}}T_{\partial X}^*X$ is not a mwc, $[0, \bar{x}]_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times T^*\partial X$ is.

The punctured space ${}^{\text{sc,leC}}T_{\partial X}^*X \setminus o_\partial$, $o_\partial = \{x=0, \xi_{\text{sc,leC}}=0 \text{ and } \eta_{\text{sc,leC}} \in o\}$, can be identified with $\text{sf} \cup \text{ff} \setminus (\text{df} \cup \text{bf} \cup \text{tf})$:

Proposition 5.1. *The identity map $i : \{0 \leq x < \bar{x}\} \cap (\mathbb{R}_\sigma^+ \times T^*X^\circ) \rightarrow \{0 \leq x < \bar{x}\} \cap (\mathbb{R}_\sigma^+ \times T^*X^\circ)$ composed with $(D\iota)^* : \{0 \leq x < \bar{x}\} \cap (\mathbb{R}_\sigma^+ \times T^*X^\circ) \rightarrow \mathbb{R}_\sigma^+ \times (0, \bar{x})_x \times \mathbb{R}_\xi \times T^*\partial X$ (where ι is the boundary collar) extends (uniquely) to a smooth map*

$$\bar{i} : \{0 \leq x < \bar{x}\} \cap {}^{\text{leC}}T^*X \rightarrow [0, \bar{x}]_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}} \quad (321)$$

restricting to a diffeomorphism $\{0 \leq x < \bar{x}\} \cap {}^{\text{leC}}\overline{T}^*X \setminus (\text{df} \cup \text{bf} \cup \text{tf}) \rightarrow ([0, \bar{x}]_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}}) \setminus o_\partial$. ■

Proof. Here we are using the boundary collar ι (really $(D\iota)^*$) to identify $\{0 < x < \bar{x}\} \cap (\mathbb{R}_\sigma^+ \times T^*X^\circ)$ with $\mathbb{R}_\sigma^+ \times (0, \bar{x})_x \times \mathbb{R}_\xi \times T^*\partial X$.

Using ι , we see that $\xi_{\text{sc,leC}}$ is a smooth function on the domain of eq. (321), as is $\eta_{\text{sc,leC}} \in T^*\partial X$. Likewise, the (smooth) projection ${}^{\text{leC}}T^*X \rightarrow X_{\text{res}}^{\text{sp}}$ fits into a composition

$$\{0 \leq x < \bar{x}\} \cap {}^{\text{leC}}T^*X \rightarrow \{0 \leq x < \bar{x}\} \cap X_{\text{res}}^{\text{sp}} \xrightarrow{\iota^{-1}} \hat{X}_{\text{res}}^{\text{sp}} = [0, \bar{x}]_{\text{res}}^{\text{sp}} \times \partial X \rightarrow [0, \bar{x}]_{\text{res}}^{\text{sp}}, \quad (322)$$

where $\hat{X} = [0, \bar{x}] \times \partial X$, which shows that the $[0, \bar{x}]_{\text{res}}^{\text{sp}}$ component of eq. (321) is smooth. Thus, \bar{i} is smooth.

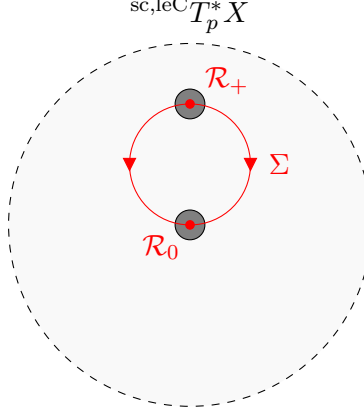


FIGURE 5. Schematic of the proof of eq. (314) when $\dim \partial X = 1$. The situation over a point $p \in (\partial[0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus \text{zf}^\circ) \times \partial X$ in the sc,leC -phase space is illustrated, with $\Sigma = \text{Char}_{\text{sc,leC}}^{2,0,-2}(\tilde{P})$ in red. There are four subsets of interest: the dark gray neighborhood of \mathcal{R}_+ , controlled using a high regularity radial point estimate, the dark gray neighborhood of $\mathcal{R}_0 = o_\partial$, controlled using a low regularity radial point estimate (and an elliptic estimate), a neighborhood of the rest of the characteristic set which stays away from $\mathcal{R}_+, \mathcal{R}_0$ over which we can propagate regularity, and the rest of the cofiber (light gray background), where elliptic estimates apply. Cf. [Mel94, Figure 2]. The direction of the Hamiltonian flow of \tilde{P} is indicated with arrows. (Note that only the “vertical” components of the flow are drawn — the Hamiltonian flow also changes p .)

The diffeomorphism clause follows from the inversion formulas

$$\varrho_{\text{df}} = \left[1 + \frac{x}{\sigma^2 + Zx} + \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{1/2} \right]^{-1}, \quad (323)$$

$$\varrho_{\text{sf}} = \frac{x}{\sigma^2 + Zx} \left[1 + \left(\frac{x}{\sigma^2 + Zx} + \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{1/2} \right)^{-1} \right], \quad (324)$$

$$\varrho_{\text{ff}} = (\sigma^2 + Zx)^{1/2} \left(\frac{x}{\sigma^2 + Zx} \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{-1/2} + 1 \right), \quad (325)$$

$$\varrho_{\text{tf}} = \left[1 + \frac{x}{\sigma^2 + Zx} \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{-1/2} \right]^{-1}, \quad (326)$$

and

$$\begin{aligned} \varrho_{\text{bf}} = & \left[1 + \left(1 + \frac{x}{\sigma^2 + Zx} \right) \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{-1/2} \right]^{-1} \\ & \times \left[1 + \frac{x}{\sigma^2 + Zx} \left(\frac{x^2}{\sigma^2 + Zx} + \xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 \right)^{-1/2} \right], \end{aligned} \quad (327)$$

(holding on $\{0 < x < \bar{x}\} \cap \mathbb{R}_\sigma^+ \times T^*X^\circ$), which certainly suffice to define a smooth two-sided inverse $([0, \bar{x}]_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{\eta_{\text{sc,leC}}}) \setminus o_\partial \rightarrow \{0 \leq x < \bar{x}\} \cap {}^{\text{leC}}T^*X \setminus (\text{df} \cup \text{bf} \cup \text{tf})$ of \tilde{i} . \square

In terms of $\xi_{\text{sc,leC}}$ and $\eta_{\text{sc,leC}}$, \tilde{p}_0 can be written as

$$\tilde{p}_0 = (\sigma^2 + Zx)(\xi_{\text{sc,leC}}^2 + g_{\partial X}^{-1}(\eta_{\text{sc,leC}}, \eta_{\text{sc,leC}}) - 2\xi_{\text{sc,leC}}) \quad (328)$$

$$= (\sigma^2 + Zx)(\xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2 - 2\xi_{\text{sc,leC}}) \quad (329)$$

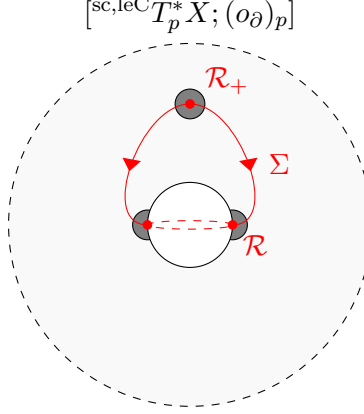


FIGURE 6. The result of blowing up the zero section of ${}^{\text{sc, leC}}T_{\partial X}^*X$ over a fixed point $p \in (\partial[0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus (\text{zf}^\circ \cup (\text{bf} \cap \text{tf}))) \times \partial X$, with $\Sigma = \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$ (which, according to \bar{i} , agrees with $\text{Char}_{\text{sc, leC}}^{2,0,-2}(\tilde{P})$ away from $\text{bf} \cup \text{tf}$ but is deformed by the blow-up) and the vertical Hamiltonian flow on it indicated. Cf. [Vas21a, Figure 3]. The inner black circle depicts either $\text{bf} \cap \text{sf}$ or $\text{tf} \cap \text{ff}$, depending on whether $\sigma > 0$ or $\sigma = 0$. Over each p , it is a copy of \mathbb{S}^{n-1} . Its equator is the portion of \mathcal{R} over p . The radial point estimate in §5.3 is used to control the dark gray neighborhood, while elliptic estimates suffice for the rest of $\text{bf} \cap \text{sf}$ or $\text{tf} \cap \text{ff}$ (which is away from \mathcal{R}).

in T^*X° . Weighting, $\tilde{p}\varrho_{\text{df}}^2\varrho_{\text{ff}}^{-2}\varrho_{\text{bf}}^{-1}\varrho_{\text{tf}}^{-3} = \tilde{p}^{2,0,-2,-1,-3}$ induces a well-defined function on $\partial^{\text{leC}}T^*X$ and thus on ${}^{\text{sc, leC}}T_{\partial X}^*X \setminus o_\partial$. Note that this restriction does not depend on $\tilde{p}_{1,2}$. The portion of the leC-characteristic set of \tilde{P} disjoint from $\text{bf} \cup \text{tf}$ can be written as

$$\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \setminus (\text{bf} \cup \text{tf}) = (\tilde{p}^{2,0,-2,-1,-3}|_{{}^{\text{sc, leC}}T_{\partial X}^*X \setminus o_\partial})^{-1}(\{0\}). \quad (330)$$

This is equal to $\{\xi_{\text{sc, leC}}^2 + \eta_{\text{sc, leC}}^2 - 2\xi_{\text{sc, leC}} = 0\} \setminus o_\partial \subset {}^{\text{sc, leC}}T_{\partial X}^*X$. For convenience, we add in o_∂ to give a new set

$$\text{Char}_{\text{sc, leC}}^{2,0,-2}(\tilde{P}) = \{\xi_{\text{sc, leC}}^2 + g_{\partial X}^{-1}(\eta_{\text{sc, leC}}, \eta_{\text{sc, leC}}) - 2\xi_{\text{sc, leC}} = 0\} \subset {}^{\text{sc, leC}}T_{\partial X}^*X, \quad (331)$$

which consists of an off-center sphere over each point of $(\partial[0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus \text{zf}^\circ) \times \partial X$ (in terms of the trivialization eq. (320)), of radius one and centered at $\xi_{\text{sc, leC}} = 1$ and $\eta_{\text{sc, leC}} \in o$. See Figure 5, which is a translated version of [Mel94, Figure 2]. The portion of

$$\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) = (\tilde{p}^{2,0,-2,-1,-3})^{-1}(\{0\}) \cap (\text{df} \cup \text{sf} \cup \text{ff}) \quad (332)$$

over a point $p \in (\partial[0, \bar{x}]_{\text{res}}^{\text{sp}} \setminus (\text{zf}^\circ \cup (\text{bf} \cap \text{tf}))) \times \partial X$ is depicted in Figure 6. Note that $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$ only contains a codimension one subset of $\text{sf} \cap \text{bf}$, $\text{ff} \cap \text{tf}$, that corresponding to $\xi_{\text{sc, leC}} = 0$, while the lift of $\text{Char}_{\text{sc, leC}}^{2,0,-2}(\tilde{P})$ contains the whole sets.

Elliptic estimates control u in terms of $\tilde{P}u$ away from this set — see §5.1.

As shown in §5.2, the Hamiltonian flow $H_{\tilde{p}} \in \mathcal{V}(\mathbb{R}_\sigma^+ \times T^*X^\circ)$ associated to \tilde{p} is given in terms of the sc, leC- coordinates above by

$$\begin{aligned} H_{\tilde{p}_0} &= x(\sigma^2 + \text{Z}x)^{1/2} \left[2(\xi_{\text{sc, leC}} - 1)x\partial_x + 2g_{\partial X}^{-1}(\eta_{\text{sc, leC}}, -) \right. \\ &\quad \left. + \frac{2\sigma^2 + \text{Z}x}{\sigma^2 + \text{Z}x} \left((\xi_{\text{sc, leC}} - 1)\eta_{\text{sc, leC}}\partial_{\eta_{\text{sc, leC}}} + \xi_{\text{sc, leC}}(\xi_{\text{sc, leC}} - 2)\partial_{\xi_{\text{sc, leC}}} \right) - \frac{2\tilde{p}_0}{\sigma^2 + \text{Z}x}\partial_{\xi_{\text{sc, leC}}} \right] \end{aligned} \quad (333)$$

to leading order at ∂X , where $\eta_{\text{sc},\text{leC}}\partial\eta_{\text{sc},\text{leC}} \in \mathcal{V}(T^*\partial X) \subset \mathcal{V}(\mathbb{R}_\sigma^+ \times [0, \bar{x})_x \times \mathbb{R}_{\xi_{\text{sc},\text{leC}}} \times T^*\partial X)$ denotes the vector field on $T^*\partial X$ given in local coordinates y_1, \dots, y_{n-1} for ∂X by

$$\eta_{\text{sc},\text{leC}}\partial\eta_{\text{sc},\text{leC}} = \sum_{j=1}^{n-1} \eta_{\text{sc},\text{leC},j} \partial\eta_{\text{sc},\text{leC},j}, \quad (334)$$

where $\eta_{\text{sc},\text{leC},j}$ is the cofiber component of $\eta_{\text{sc},\text{leC}}$ dual to y_j . It will be convenient to work with the weighted Hamiltonian vector fields $H_{\tilde{p}_0}^{2,0,-2} = \varrho_{\text{df}} x^{-1}(\sigma^2 + Zx)^{-1/2} H_{\tilde{p}_0}$, $H_{\tilde{p}}^{2,0,-2} = \varrho_{\text{df}} x^{-1}(\sigma^2 + Zx)^{-1/2} H_{\tilde{p}}$.

The vector field $H_{\tilde{p}}$, as given by eq. (333), defines on the sc,leC-characteristic set $\text{Char}_{\text{sc},\text{leC}}^{2,0,-2}(\tilde{P})$ a source-to-sink flow, with

- $\{\xi_{\text{sc},\text{leC}} = \eta_{\text{sc},\text{leC}} = 0\} = \mathcal{R}_0 \subset {}^{\text{sc},\text{leC}}T_{\partial X}^*X$, the “selected radial set,” and
- $\{\xi_{\text{sc},\text{leC}} = 2, \eta_{\text{sc},\text{leC}} = 0\} = \mathcal{R}_+ \subset \text{sf} \cup \text{ff}$, the “unselected radial set,”

the sink and source (respectively, under our sign conventions) of the flow. Hence, leC-analogues of standard propagation and radial point estimates apply away from o_∂ (the dependence on the one parameter σ being unimportant), and the proofs are straightforward modifications of the analogous estimates in [Mel94]. These estimates are proven in §5.2. Note that the $(\xi_{\text{sc},\text{leC}} - 1)x\partial_x$ term in eq. (333) does not actually vanish on the portion of $\mathcal{R}_0, \mathcal{R}_+$ over the interior of tf — rather, it induces a flow from the bf side to the zf side. Thus, it is also possible to prove a propagation result in which regularity at $\mathcal{R}_+ \cap \text{sf}$ is propagated into $\mathcal{R}_+ \cap \text{zf}$. We will not need (and hence will not prove) such an estimate, as a much weaker argument in §6 suffices to show that for nice f the output $\tilde{R}_+(E)f$ of the conjugated resolvent $\tilde{R}_+(E)$ converges weakly to something satisfying the conjugated version of the Sommerfeld radiation condition as $E \rightarrow 0^+$. (But it is still worth pointing out that $\mathcal{R}_+ \cap \text{zf}$ is not a source of the flow in every direction — it is a sink in the $\varsigma = \sigma/x^{1/2}$ direction.)

The radial point estimate at the selected radial set is somewhat more nonstandard. Rather than \mathcal{R}_0 , we work with $\mathcal{R} = \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \cap (\text{bf} \cup \text{tf})$. See [Vas21a, Figure 3] for the $\sigma > 0$ case. The face ff meets bf at an edge — see Figure 3, Figure 7 — and so $\text{sf} \cup \text{ff} \setminus \text{df}$ is not just the boundary of

$$[[0, \bar{x})_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc},\text{leC}}} \times T^*\partial X; [0, \bar{x})_{\text{res}}^{\text{sp}} \times \{0\} \times o], \quad (335)$$

which is why Figure 6 depicts the situation only over $p \in [0, \bar{x})_{\text{sp}}^{\text{res}} \times \partial X$ not on the zero face or corner. Away from that edge, our situation looks (even at ff) very much like that in [Vas21a]. An argument similar to that used to prove the radial point estimate there suffices to prove an estimate here that is uniform down to $\sigma = 0$, albeit without the desired number of independent orders. In order to prove an estimate with the desired number of independent orders, it will be necessary to take into account the aberrant edge. See §5.3.

The elliptic, propagation, and radial point estimates in §5.1, §5.2, §5.3 are combined in a short epilogue §5.4, which contains Theorem 5.12, as well as the corollary Proposition 5.14 stated above.

5.1. Ellipticity.

Proposition 5.2. $\tilde{P} \in \text{Diff}_{\text{leC}}^{2,0,-2,-1,-3}(X)$ is elliptic in some neighborhood of df . ■

Proof. By assumption, \tilde{P} is elliptic at every point of df° . By eq. (264), \tilde{P} will be elliptic at a point of $\text{sf} \cup \text{ff}$ (including the points of $\partial \text{df} \subset \text{sf} \cup \text{ff}$) if and only if \tilde{P}_0 is. $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}_0)$ is contained away from df , so the same then holds for $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$. □

Via the leC-calculus analogue of the usual argument via parametrix:

Proposition 5.3. Let $A, A_0, B \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$, with $\text{WF}_{\text{leC}}^{0,0}(A), \text{WF}'_{\text{leC}}(A_0) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0,0}(B) \cap \text{Ell}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$.

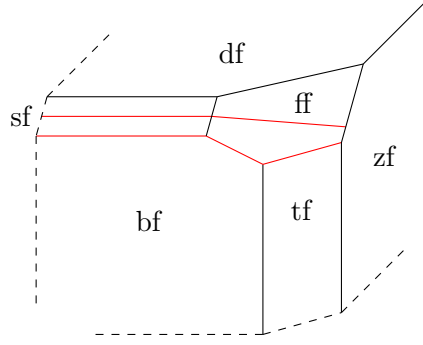


FIGURE 7. The radial sets $\mathcal{R}_+, \mathcal{R}$ (in red) as subsets of the leC-phase space (ignoring the degrees of freedom associated with $T^*\partial X$). The top line represents \mathcal{R}_+ , while the bottom line depicts \mathcal{R} . (Note that on $\text{bf} \cap \text{ff}$, $x/(\sigma^2 + \mathbb{Z}x) = \sigma^2 + \mathbb{Z}x = 0$.)

Then, for each $\Sigma > 0$, $m, s, \varsigma, l, \ell \in \mathbb{R}$, and $N \in \mathbb{N}$, there exists a

$$C = C(\tilde{P}, A, A_0, B, m, s, \varsigma, l, \ell, N, \Sigma) > 0$$

such that, for any $u \in \mathcal{S}'(X)$,

$$\|Au\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C[\|B\tilde{P}u\|_{H_{\text{leC}}^{m-2,s,\varsigma+2,l+1,\ell+3}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,l,\ell}}] \quad (336)$$

$$\|A_0u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C[\|B\tilde{P}u\|_{H_{\text{leC}}^{m-2,s,\varsigma+2,l+1,\ell+3}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,-N,-N}}] \quad (337)$$

for all $\sigma \in (0, \Sigma]$ (in the strong sense that if the right-hand side is finite then the left-hand side is as well). \blacksquare

Proof. Let $b \in S_{\text{leC}}^{0,0,0,0,0}(X) = S_{\text{b,leC}}^{0,0,0}(X)$ denote a representative of $\sigma_{\text{leC}}^{0,0,0,0,0}(B)$. Let $\varphi \in C_c^\infty(\mathbb{R})$ be identically equal to one in some neighborhood of $[0, \Sigma]$ and supported in $(-\infty, 2\Sigma)$.

The set

$$K = \text{WF}_{\text{leC}}^{0,0}(A) \cap \{\sigma \leq 2\Sigma\} \subseteq \text{df} \cup \text{sf} \cup \text{ff} \subseteq \partial^{\text{leC}} \overline{T^*X} \quad (338)$$

is a compact subset of ${}^{\text{leC}}\overline{T^*X}$, so we can find some $\chi \in S_{\text{leC}}^{0,0,0,0,0}(X)$ such that $\chi = 1$ identically in some neighborhood of K and such that $\chi = 0$ identically in some neighborhood of $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(B\tilde{P}) = \text{Char}_{\text{leC}}^{0,0,0,0,0}(B) \cup \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$. We can moreover choose χ such that $\chi = 0$ identically in some neighborhood of $b^{-1}(\{0\}) \cup \tilde{p}^{-1}(\{0\})$. Consider $f = \chi/b\tilde{p} \in S_{\text{leC}}^{-2,0,2,1,3}(X)$.

Quantizing, we get some $F = \text{Op}(f) \in \Psi_{\text{leC}}^{-2,0,2,1,3}(X)$ such that – via the leC-principal symbol short exact sequence –

$$\varphi A(1 - FB\tilde{P}) \in \Psi_{\text{leC}}^{-1,-1,-1,0,0}(X). \quad (339)$$

Indeed, $\varphi A(1 - FB\tilde{P}) \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ by Proposition 2.6. By Proposition 2.9, for any representative $a \in S_{\text{leC}}^{0,0,0,0,0}(X)$ of the principal symbol $\sigma_{\text{leC}}^{0,0,0,0,0}(A)$,

$$\begin{aligned} \sigma_{\text{leC}}^{0,0,0,0,0}(\varphi A(1 - FB\tilde{P})) &= \varphi a(1 - \chi) \bmod S_{\text{leC}}^{-1,-1,-1,0,0}(X) \\ &= 0 \bmod S_{\text{leC}}^{-1,-1,-1,0,0}(X), \end{aligned} \quad (340)$$

which implies eq. (339). So, for $\sigma \in [0, \Sigma]$,

$$\begin{aligned} \|Au\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} &\leq \|\varphi AFB\tilde{P}u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} + \|\varphi A(1 - FB\tilde{P})u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \\ &\leq \|B\tilde{P}u\|_{H_{\text{leC}}^{m-2,s,\varsigma+2,l+1,\ell+3}} + \|\varphi A_1u\|_{H_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,l,\ell}} \\ &= \|B\tilde{P}u\|_{H_{\text{leC}}^{m-2,s,\varsigma+2,l+1,\ell+3}} + \|A_1u\|_{H_{\text{leC}}^{m-1,s-1,\varsigma-1,l,\ell}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,l,\ell}} \end{aligned} \quad (341)$$

for A_1 satisfying the same hypotheses as A but also elliptic on the essential support of A . Inducting, we conclude the estimate eq. (336).

The second estimate, eq. (337), is proven in a completely analogous manner. \square

5.2. Propagation. The Hamiltonian vector field $H_{\tilde{p}} \in \mathcal{V}(\mathbb{R}_\sigma^+ \times T^*X^\circ)$ associated with the symbol $\tilde{p} \in C^\infty((0, \infty)_\sigma \times T^*X^\circ)$ is given near ∂X by

$$H_{\tilde{p}} = (\partial_{\xi_b} \tilde{p}) x \partial_x - (x \partial_x \tilde{p}) \partial_{\xi_b} + \sum_{j=1}^{n-1} ((\partial_{\eta_j} \tilde{p}) \partial_{y_j} - (\partial_{y_j} \tilde{p}) \partial_{\eta_j}) \in C^\infty((0, \infty)_\sigma; \mathcal{V}(T^*X^\circ)) \quad (342)$$

with respect to any set of local coordinates $y = (y_1, \dots, y_{n-1})$ on ∂X . (We will alternatively identify $H_{\tilde{p}}$ as a smooth family of elements of $\mathcal{V}(T^*X^\circ)$ and as a vector field on $\mathbb{R}_\sigma^+ \times T^*X^\circ$ without ∂_σ component.) Together, $x, y, \xi_{\text{sc,leC}}, \eta_{\text{sc,leC}}$ constitute a coordinate chart for T^*X° , so we can rewrite $H_{\tilde{p}}(\sigma) \in C^\infty(T^*X^\circ; TT^*X^\circ)$ in terms of them, and the result (patching together the various coordinate charts for ∂X) can be interpreted as a weighted b-vector field on $[0, \bar{x})_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc,leC}}} \times (T^*\partial X)_{y, \eta_{\text{sc,leC}}}$ (after restricting attention to a small collar neighborhood of ∂X). To perform this rewrite, we need the following substitutions:

$$\begin{aligned} x \partial_x &\rightarrow x \partial_x + \frac{\partial \xi_{\text{sc,leC}}}{\partial x} x \partial_{\xi_{\text{sc,leC}}} + \sum_{j=1}^{n-1} \frac{\partial \eta_{\text{sc,leC},j}}{\partial x} x \partial_{\eta_{\text{sc,leC},j}} \\ &= x \partial_x + \frac{2\sigma^2 + Zx}{2(\sigma^2 + Zx)} \xi_{\text{sc,leC}} \partial_{\xi_{\text{sc,leC}}} + \frac{2\sigma^2 + Zx}{2(\sigma^2 + Zx)} \sum_{j=1}^{n-1} \eta_{\text{sc,leC},j} \partial_{\eta_{\text{sc,leC},j}} \end{aligned} \quad (343)$$

and

$$\partial_{\xi_b} \rightarrow \frac{\partial \xi_{\text{sc,leC}}}{\partial \xi_b} \partial_{\xi_{\text{sc,leC}}} = \frac{x}{\sqrt{\sigma^2 + Zx}} \partial_{\xi_{\text{sc,leC}}}, \quad \partial_{\eta_j} \rightarrow \frac{\partial \eta_{\text{sc,leC},j}}{\partial \eta_j} \partial_{\eta_{\text{sc,leC},j}} = \frac{x}{\sqrt{\sigma^2 + Zx}} \partial_{\eta_{\text{sc,leC},j}} \quad (344)$$

for $j = 1, \dots, n-1$, where the partial derivatives are taken with respect to the coordinate system x, y, ξ_b, η . In other words, letting $(x \partial_x)_{\text{old}}, (\partial_{y_j})_{\text{old}}$ denote the local vector fields defined using the coordinate system x, y, ξ_b, η , we have

$$(\partial_{y_j})_{\text{old}} = \partial_{y_j}, \quad (345)$$

$$(x \partial_x)_{\text{old}} = x \partial_x + \frac{2\sigma^2 + Zx}{2(\sigma^2 + Zx)} \xi_{\text{sc,leC}} \partial_{\xi_{\text{sc,leC}}} + \frac{2\sigma^2 + Zx}{2(\sigma^2 + Zx)} \sum_{j=1}^{n-1} \eta_{\text{sc,leC},j} \partial_{\eta_{\text{sc,leC},j}}, \quad (346)$$

where the partial derivatives on the right-hand side are defined using the coordinate system $x, y, \xi_{\text{sc,leC}}, \eta_{\text{sc,leC}}$. In terms of this new notation, eq. (342) says

$$H_{\tilde{p}} = (\partial_{\xi_b} \tilde{p}) (x \partial_x)_{\text{old}} - ((x \partial_x)_{\text{old}} \tilde{p}) \partial_{\xi_b} + \sum_{j=1}^{n-1} ((\partial_{\eta_j} \tilde{p}) (\partial_{y_j})_{\text{old}} - ((\partial_{y_j})_{\text{old}} \tilde{p}) \partial_{\eta_j}). \quad (347)$$

The same holds for \tilde{p}_0 in place of \tilde{p} . The ∂_{y_j} component of $H_{\tilde{p}_0} = H_{\tilde{p}} - H_{\tilde{p}_{1,2}}$ is given by

$$\partial_{\eta_j} \tilde{p}_0 = \frac{x}{\sqrt{\sigma^2 + Zx}} \frac{\partial \tilde{p}_0}{\partial \eta_{\text{sc,leC},j}} = 2x \sqrt{\sigma^2 + Zx} \sum_{k=1}^{n-1} g^{kj} \eta_{\text{sc,leC},k}. \quad (348)$$

On the other hand, the $\partial_{\eta_{\text{sc,leC},j}}$ component of $H_{\tilde{p}_0}$ is given by

$$\begin{aligned} \frac{2\sigma^2 + Zx}{2\sigma^2 + 2Zx} \eta_{\text{sc,leC},j} \frac{\partial \tilde{p}_0}{\partial \xi_b} - \frac{x}{(\sigma^2 + Zx)^{1/2}} \frac{\partial \tilde{p}_0}{\partial y_j} &= \frac{2\sigma^2 + Zx}{2\sigma^2 + 2Zx} \frac{x}{(\sigma^2 + Zx)^{1/2}} \eta_{\text{sc,leC},j} \frac{\partial \tilde{p}_0}{\partial \xi_{\text{sc,leC}}} \\ &= \frac{2\sigma^2 + Zx}{(\sigma^2 + Zx)^{1/2}} x \eta_{\text{sc,leC},j} (\xi_{\text{sc,leC}} - 1), \end{aligned} \quad (349)$$

while the $\partial_{\xi_{\text{sc},\text{leC}}}$ component is given by

$$\begin{aligned} & \frac{2\sigma^2 + Zx}{2\sigma^2 + 2Zx} \xi_{\text{sc},\text{leC}} \frac{\partial \tilde{p}_0}{\partial \xi_{\text{b}}} - \frac{x^2}{(\sigma^2 + Zx)^{1/2}} \left(\frac{\partial \tilde{p}_0}{\partial x} \right)_{\text{old}} \\ &= \frac{2\sigma^2 + Zx}{(\sigma^2 + Zx)^{1/2}} x \xi_{\text{sc},\text{leC}} (\xi_{\text{sc},\text{leC}} - 1) - \frac{x}{(\sigma^2 + Zx)^{1/2}} \left[2\tilde{p}_0 + \xi_{\text{sc},\text{leC}} (2\sigma^2 + Zx) \right] \\ &= \frac{2\sigma^2 + Zx}{(\sigma^2 + Zx)^{1/2}} x \xi_{\text{sc},\text{leC}} (\xi_{\text{sc},\text{leC}} - 2) - \frac{2x\tilde{p}_0}{(\sigma^2 + Zx)^{1/2}}, \quad (350) \end{aligned}$$

and the $x\partial_x$ component is $x(\sigma^2 + Zx)^{-1/2} \partial_{\xi_{\text{sc},\text{leC}}} \tilde{p}_0 = 2x(\sigma^2 + Zx)^{1/2} (\xi_{\text{sc},\text{leC}} - 1)$. To summarize:

Proposition 5.4. *In terms of the coordinates $x, y, \xi_{\text{sc},\text{leC}}, \eta_{\text{sc},\text{leC}}$,*

$$\begin{aligned} H_{\tilde{p}_0} &= x(\sigma^2 + Zx)^{1/2} \left[2(\xi_{\text{sc},\text{leC}} - 1)x\partial_x + 2g_{\partial X}^{-1}(\eta_{\text{sc},\text{leC}}, -) \right. \\ &\quad \left. + \frac{2\sigma^2 + Zx}{\sigma^2 + Zx} \left(\xi_{\text{sc},\text{leC}} (\xi_{\text{sc},\text{leC}} - 2) \partial_{\xi_{\text{sc},\text{leC}}} + \sum_{j=1}^{n-1} \eta_{\text{sc},\text{leC},j} (\xi_{\text{sc},\text{leC}} - 1) \partial_{\eta_{\text{sc},\text{leC},j}} \right) - \frac{2\tilde{p}_0}{\sigma^2 + Zx} \partial_{\xi_{\text{sc},\text{leC}}} \right] \quad (351) \end{aligned}$$

(in the relevant neighborhood of ${}^{\text{leC}}T^*X$). ■□

Letting $H_{\tilde{p}}^{-,0,-2} = x^{-1}(\sigma^2 + Zx)^{-1/2} H_{\tilde{p}} \in \mathcal{V}(\mathbb{R}_\sigma^+ \times T^*X^\circ)$, $H_{\tilde{p}}^{-,0,-2}$ defines a b-vector field on

$${}^{\text{sc},\text{leC}}T^*X = [0, \bar{x})_{\text{res}}^{\text{sp}} \times \mathbb{R}_{\xi_{\text{sc},\text{leC}}} \times (T^*\partial X)_{y,\eta_{\text{sc},\text{leC}}}. \quad (352)$$

Likewise for \tilde{p}_0 . Note that $H_{\tilde{p}}^{-,0,-2}$ and $H_{\tilde{p}_0}^{-,0,-2}$ agree at every point of $\text{sf} \cup \text{ff}$. Restricting to $\{x = 0\}$, $H_{\tilde{p}}^{-,0,-2}$ can be considered as a family

$$H : \partial([0, \bar{x})_{\text{res}}^{\text{sp}} \setminus \text{zf}^\circ) \rightarrow \mathcal{V}(\mathbb{R}_{\xi_{\text{sc},\text{leC}}} \times (T^*\partial X)_{y,\eta_{\text{sc},\text{leC}}}) \quad (353)$$

of vector fields on the fiber $\mathbb{R} \times T^*\partial X$. In order to understand H , consider $H_{\tilde{p}}^{-,0,-2}$ over a neighborhood of a subset of the interior of the transition face tf of $[0, \bar{x})_{\text{res}}^{\text{sp}}$, which we can parametrize in terms of $x \in [0, \bar{x})$ and $\lambda \in \mathbb{R}^+$ by writing $\sigma^2 = Z\lambda x$. In this neighborhood, $x^{1/2}$ is a bdf for the transition face. Then, H can be thought of as a family of vector fields on $\mathbb{R}_{\xi_{\text{sc},\text{leC}}} \times T^*\partial X$ dependent on the parameter λ , with explicit formula

$$H = \frac{2\lambda + 1}{\lambda + 1} \left[(\xi_{\text{sc},\text{leC}} - 1) \eta_{\text{sc},\text{leC}} \partial_{\eta_{\text{sc},\text{leC}}} + (\xi_{\text{sc},\text{leC}} - 2) \xi_{\text{sc},\text{leC}} \partial_{\xi_{\text{sc},\text{leC}}} \right] + 2g_{\partial X}^{-1}(\eta_{\text{sc},\text{leC}}, -) \quad (354)$$

when $\tilde{p}_0 = 0$. The ∂_y -component of H is independent of λ and comes from the projection of the geodesic flow on $T^*\partial X$ down to ∂X . The cofiber components of H (bracketed) depend on λ only in the form of an overall factor that – crucially – is nonzero for all $\lambda \in [0, \infty)$. (In fact, eq. (354) makes sense as a family of vector fields on $\mathbb{R} \times T^*\partial X$ for all $\lambda > -1$, and it is non-vanishing for $\lambda > -1/2$, but we do not consider this “extended transition face” here. The vanishing at $\lambda = -1/2$ is one sign that the consideration of negative λ would require solving a b-problem analogous to the b-problem encountered in the low-energy analysis of Coulomb-free Schrödinger operators. Cf. Remark 8.) From eq. (354), we read off the following crucial observation: over the transition face of $[0, \bar{x})_{\text{res}}^{\text{sp}}$, parametrized as above, H , when restricted to the characteristic set of \tilde{p} and away from $\text{bf} \cup \text{tf}$, vanishes if and only if $\eta_{\text{sc},\text{leC}} = 0$ and $\xi_{\text{sc},\text{leC}} = 2$, i.e. on \mathcal{R}_+ . Between this submanifold and the zero section \mathcal{R}_0 , H has (over each point in $\partial[0, \bar{x})_{\text{res}}^{\text{sp}}$) a source-to-sink flow within $\text{Char}_{\text{sc},\text{leC}}^{2,0,-2}(\tilde{P})$. In order to see that $\{\xi_{\text{sc},\text{leC}} = \eta_{\text{sc},\text{leC}} = 0\}$ is a sink of the flow, observe that

$$H(\xi_{\text{sc},\text{leC}}^2 + \eta_{\text{sc},\text{leC}}^2) = \frac{4\lambda + 2}{\lambda + 1} \left[\eta_{\text{sc}}^2 (\xi_{\text{sc},\text{leC}} - 1) + \xi_{\text{sc}}^2 (\xi_{\text{sc},\text{leC}} - 2) \right] - \frac{4\tilde{p}_0 \xi_{\text{sc},\text{leC}}}{\sigma^2 + Zx} \quad (355)$$

(note that $\tilde{p}_0/(\sigma^2 + Zx) = \xi_{\text{sc},\text{leC}}^2 + \eta_{\text{sc},\text{leC}}^2 - 2\xi_{\text{sc},\text{leC}}$ is well-defined).

The same computation yields:

Proposition 5.5. $H_{\tilde{p}_0}^{-,0,-2}(\xi_{\text{sc},\text{leC}}^2 + \eta_{\text{sc},\text{leC}}^2) = \beta_{0,1}(\xi_{\text{sc},\text{leC}}^2 + \eta_{\text{sc},\text{leC}}^2) + F_{0,2} + F_{0,3}$ for

$$\beta_{0,1} = \frac{4\sigma^2 + 2Zx}{\sigma^2 + Zx}(\xi_{\text{sc},\text{leC}} - 1), \quad F_{0,2} = -\xi_{\text{sc},\text{leC}}^2 \frac{4\sigma^2 + 2Zx}{\sigma^2 + Zx}, \quad F_{0,3} = -\frac{4\tilde{p}_0 \xi_{\text{sc},\text{leC}}}{\sigma^2 + Zx}. \quad (356)$$

These extend to symbols on the *leC*-phase space. The first two are nonpositive in the vicinity of $\text{bf} \cup \text{tf}$, while $F_{0,3}$ vanishes cubically there. $\blacksquare \square$

Note a similar statement holds if we replace \tilde{p}_0 by \tilde{p} , if instead of $F_{0,3} = -4\tilde{p}_0 \xi_{\text{sc},\text{leC}}/(\sigma^2 + Zx)$ we use

$$H_{\tilde{p}_{1,2}}^{-,0,-2}(\xi_{\text{sc},\text{leC}}^2 + \eta_{\text{sc},\text{leC}}^2) - 4\tilde{p}_0 \xi_{\text{sc},\text{leC}}/(\sigma^2 + Zx), \quad (357)$$

which is the sum of a term suppressed by a factor of $x(\sigma^2 + Zx)^{-1/2}$ and a cubically vanishing term.

Consider, for each pair of $\Theta_1, \Theta_2 \in (0, \pi)$ with $\Theta_1 < \Theta_2$, the set $\mathcal{P}[\Theta_1, \Theta_2] \subset \text{sf} \cup \text{ff}$ defined by $\mathcal{P}[\Theta_1, \Theta_2] = \text{Char}_{\text{leC}}^{2,0,-2}(\tilde{P}) \cap \{\arccos(\xi_{\text{sc},\text{leC}} - 1) \in [\Theta_1, \Theta_2]\}$. The following proposition is a symbolic version of the statement that the Hamiltonian flow is source-to-sink, \mathcal{R}_+ to \mathcal{R}_0 .

Proposition 5.6. Let $\Theta \in S_{\text{cl},\text{leC}}^{0,0,0,0}(X)$ satisfy $\Theta = \arccos(\xi_{\text{sc},\text{leC}} - 1)$ in some neighborhood of $\mathcal{P}[\Theta_1, \Theta_2]$.

For any pair if $\Theta_1, \Theta_2 \in (0, \pi)$ with $\Theta_1 < \Theta_2$, the symbol $\alpha \in S_{\text{leC}}^{0,0,0,0}(X)$ defined by $H_{\tilde{p}}^{2,0,-2}\Theta = \alpha$ satisfies $\alpha > 0$ on $\mathcal{P}[\Theta_1, \Theta_2]$. \blacksquare

Proof. Given such $\Theta \in S_{\text{cl},\text{leC}}^{0,0,0,0}(X)$, Θ is equal to $\arccos(\xi_{\text{sc},\text{leC}} - 1)$ in some neighborhood $U \subset {}^{\text{leC}}\overline{T}^*X$ of $\mathcal{P}[\Theta_1, \Theta_2]$. Thus,

$$H_{\tilde{p}_0}^{-,0,-2}\Theta = \frac{2\sigma^2 + Zx}{\sigma^2 + Zx} \xi_{\text{sc},\text{leC}}^{1/2} (2 - \xi_{\text{sc},\text{leC}})^{1/2} + \frac{2\tilde{p}_0}{\sigma^2 + Zx} \frac{1}{\xi_{\text{sc},\text{leC}}^{1/2} (2 - \xi_{\text{sc},\text{leC}}^2)^{1/2}} \quad (358)$$

in some neighborhood of $\mathcal{P}[\Theta_1, \Theta_2]$, where we are taking positive square roots. Since \tilde{p}_0 vanishes on $\mathcal{P}[\Theta_1, \Theta_2]$, the expression on the right-hand side is positive on $\mathcal{P}[\Theta_1, \Theta_2]$.

Since ϱ_{df} is positive on $\mathcal{P}[\Theta_1, \Theta_2]$, the same statement applies to $H_{\tilde{p}_0}^{2,0,-2}\Theta$. And since $H_{\tilde{p}_{1,2}}^{2,0,-2}\Theta$ vanishes at $\mathcal{P}[\Theta_1, \Theta_2]$, the same statement applies to $\alpha = H_{\tilde{p}}^{2,0,-2}\Theta$. \square

Clearly, there exist $\Theta = \Theta_{\Theta_1, \Theta_2} \in S_{\text{cl},\text{leC}}^{0,0,0,0}(X)$ satisfying the hypotheses of the previous proposition.

Proposition 5.7. Suppose that $G_1, G_2, G_3 \in \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty}(X)$ satisfy the following: there exist

$$\Theta_1, \Theta_2, \Theta_3, \Theta_4, \Theta_5 \in (0, \pi) \quad (359)$$

satisfying $\Theta_1 < \Theta_2 < \Theta_3 < \Theta_4 < \Theta_5$ and

- $\text{WF}'_{\text{leC}}(G_1) \cap \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subseteq \mathcal{P}[\Theta_3, \Theta_4]$,
- $\mathcal{P}[\Theta_1, \Theta_2] \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_3)$,
- $\mathcal{P}[\Theta_1, \Theta_5] \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$,

and $\text{WF}'_{\text{leC}}(G_1) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$. Then, for every $\Sigma > 0$, $N \in \mathbb{N}$, $m, s, \varsigma, l, \ell \in \mathbb{R}$, there exists a constant

$$C = C(\tilde{P}, G_1, G_2, G_3, \Theta_1, \Theta_2, \Theta_3, \Theta_4, \Theta_5, \Sigma, N, m, s, \varsigma, l, \ell) > 0 \quad (360)$$

such that

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C \left[\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{-N,s+1,\varsigma+3,-N,-N}} + \|G_3 u\|_{H_{\text{leC}}^{-N,s,\varsigma,-N,-N}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,-N,-N}} \right] \quad (361)$$

holds for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$ (in the strong sense that if the right-hand side is finite, then the left-hand side is as well, and the stated inequality holds). \blacksquare

Proof. Throughout the argument below, we can take N to be sufficiently large such that any of the finitely many functions of m, s, ς, l, ℓ that arise can be bounded below $-N$.

We may assume without loss of generality that G_2, G_3 are essentially supported away from $\text{bf} \cup \text{tf}$ and that G_2 is essentially supported away from \mathcal{R}_+ :

$$\mathcal{R}_+ \cap \text{WF}'_{\text{leC}}(G_2) = \emptyset. \quad (362)$$

Let $\varphi \in C_c^\infty(\mathbb{R})$ satisfy $\text{supp } \varphi \subset [\Theta_1, \Theta_5]$ and

$$\varphi'(\theta) = \varphi_0(\theta)^2 + \varphi_1(\theta) \quad (363)$$

for $\varphi_0 \in C_c^\infty(\mathbb{R})$ nonvanishing on $[\Theta_3, \Theta_4]$ and $\varphi_1 \in C_c^\infty(\mathbb{R})$ supported within (Θ_1, Θ_2) . The construction of such φ is standard. Moreover, for any closed interval $I \subset ((\Theta_1 + 3\Theta_2)/4, \Theta_5)$ and any $\epsilon > 0$, we can construct φ such that $\epsilon\varphi_0^2 \geq \varphi$ in I . This construction is also standard — we consider $\varphi_{00} \in C^\infty(\mathbb{R})$ given by

$$\varphi_{00} = \begin{cases} e^{-F/(\Theta_5 - \Theta)} & (\Theta \leq \Theta_5) \\ 0 & (\Theta > \Theta_5) \end{cases} \quad (364)$$

for a parameter $F = F(\sigma) > 0$ and $\varphi_{01} \in C^\infty(\mathbb{R})$ that is identically equal to one in some neighborhood of $[\Theta_2, \infty)$ and identically zero in some neighborhood of $(-\infty, \Theta_1]$; setting $\varphi = -\varphi_{00}\varphi_{01}^2$ we have

$$\varphi' = -2\varphi_{00}\varphi_{01}\varphi_{01}' - \varphi_{00}'\varphi_{01}^2 = -2\varphi_{00}\varphi_{01}\varphi_{01}' + F(\Theta_5 - \Theta)^{-2}\varphi_{00}\varphi_{01}^2. \quad (365)$$

Setting $\varphi_0 = F^{1/2}(\Theta_5 - \Theta)^{-1}\varphi_{00}^{1/2}\varphi_{01}$ and $\varphi_1 = -2\varphi_{00}\varphi_{01}\varphi_{01}'$, we see that φ, φ' have the desired form.

Fix $\Theta \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ that is equal to $\arccos(\xi_{\text{sc,leC}} - 1)$ in some neighborhood of $\mathcal{P}[\Theta_1, \Theta_5]$. Pick a neighborhood U_0 of $\mathcal{P}[\Theta_1, \Theta_5]$ on which Θ is identically $\arccos(\xi_{\text{sc,leC}} - 1)$ and such that α is bounded below on U_0 (in compact sets worth of σ). Let

$$\varphi(\Theta) \in S_{\text{cl,leC}}^{-\infty,0,0,-\infty,-\infty}(X) \quad (366)$$

denote a symbol equal to $\varphi \circ \Theta$ on some neighborhood $U \Subset U_0$ of $\mathcal{P}[\Theta_1, \Theta_5]$, and let $\psi \in S_{\text{cl,leC}}^{-\infty,0,0,-\infty,-\infty}(X)$ have $\text{supp } \psi \Subset U$ and be identically equal to 1 on some neighborhood of $\mathcal{P}[\Theta_1, \Theta_5]$. We can take both of these to be supported away from $\text{bf} \cup \text{tf}$. We can choose ψ such that

$$\text{supp}(\varphi(\Theta)H_{\tilde{p}}^{2,0,-2}\psi) \cap (\tilde{p}^{2,0,-2,-1,-3})^{-1}(\{0\}) = \emptyset \quad (367)$$

and $\text{supp } \psi \cap (\text{df} \cup \mathcal{R}_+) = \emptyset$. Consider the symbol

$$a_0 = \varphi(\Theta)\psi^2 \in S_{\text{cl,leC}}^{-\infty,0,0,-\infty,-\infty}(X). \quad (368)$$

We then compute

$$H_{\tilde{p}}^{2,0,-2}a_0 = (\varphi_0(\Theta)^2 + \varphi_1(\Theta))\psi^2\alpha + 2\psi\varphi(\Theta)H_{\tilde{p}}^{2,0,-2}\psi, \quad (369)$$

where α is as in Proposition 5.6.

Now set, for to-be-decided $K > 0$, set $\phi_\epsilon(x) = (1 + \epsilon x^{-1})^{-K}$, for each $\epsilon \geq 0$. Now set $a_0^{(\epsilon)} = \phi_\epsilon^2 a_0$.

In §3, we checked that $\Im \tilde{P} \in S\text{Diff}_{\text{leC}}^{2,-1,-3,-1-\delta,-3-2\delta}(X) \subset \Psi_{\text{leC}}^{2,-1,-3,-1-\delta,-3-2\delta}(X)$. Let p_1 denote a representative of $\sigma_{\text{leC}}^{2,-1,-3,-1-\delta,-3-2\delta}(-2\Im \tilde{P})$. Then,

$$H_{\tilde{p}}^{2,0,-2}a_0^{(\epsilon)} = (\varphi_0(\Theta)^2 + \varphi_1(\Theta))\phi_\epsilon^2\psi^2\alpha + 2\psi\varphi(\Theta)H_{\tilde{p}}^{2,0,-2}\psi + 2K\psi^2\phi_\epsilon^2(\epsilon x^{-1})(1 + \epsilon x^{-1})^{-1}\varphi\beta_1, \quad (370)$$

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2}a_0^{(\epsilon)} + \varrho_{\text{df}}\varrho_{\text{sf}}^{-1}\varrho_{\text{ff}}^{-3}\varrho_{\text{bf}}^{-1}\varrho_{\text{tf}}^{-3}p_1a_0^{(\epsilon)} &= (\varphi_0(\Theta)^2 + \varphi_1(\Theta))\phi_\epsilon^2\psi^2\alpha + 2\psi\varphi(\Theta)H_{\tilde{p}}^{2,0,-2}\psi \\ &\quad + 2K\psi^2\phi_\epsilon^2(\epsilon x^{-1})(1 + \epsilon x^{-1})^{-1}\varphi\beta_1 + \psi^2p_1\varrho_{\text{df}}\varrho_{\text{sf}}^{-1}\varrho_{\text{ff}}^{-3}\varrho_{\text{bf}}^{-1}\varrho_{\text{tf}}^{-3}\phi_\epsilon^2\varphi, \end{aligned} \quad (371)$$

where $\beta_1 \in S_{\text{leC}}^{0,0,0,0}(X)$ is defined by $H_{\tilde{p}}^{2,0,-2}x = \beta_1x$.

We now let $a = \varrho_{\text{sf}}^{-2s-1} \varrho_{\text{ff}}^{-2\varsigma-3} \varrho_{\text{bf}}^{-2l-1} \varrho_{\text{tf}}^{-2\ell-3} a_0$, $a^{(\varepsilon)} = \varrho_{\text{sf}}^{-2s-1} \varrho_{\text{ff}}^{-2\varsigma-3} \varrho_{\text{bf}}^{-2l-1} \varrho_{\text{tf}}^{-2\ell-3} a_0^{(\varepsilon)}$. Then, $a^{(\bullet)} \in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, 2s+1, 2\varsigma+3, -\infty, -\infty}(X))$ and

$$H_{\tilde{p}}^{2,0,-2} a^{(\varepsilon)} = \varrho_{\text{sf}}^{-2s-1} \varrho_{\text{ff}}^{-2\varsigma-3} \varrho_{\text{bf}}^{-2l-1} \varrho_{\text{tf}}^{-2\ell-3} (H_{\tilde{p}}^{2,0,-2} a_0^{(\varepsilon)} + a_0^{(\varepsilon)} p_2) \quad (372)$$

for some $p_2 \in S_{\text{leC}}^{0,0,0,0}(X)$. Thus,

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2} a^{(\varepsilon)} + \varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} p_1 a^{(\varepsilon)} &= \varrho_{\text{sf}}^{-2s-1} \varrho_{\text{ff}}^{-2\varsigma-3} \varrho_{\text{bf}}^{-2l-1} \varrho_{\text{tf}}^{-2\ell-3} \left[(\varphi_0(\Theta)^2 + \varphi_1(\Theta)) \phi_\varepsilon^2 \psi^2 \alpha \right. \\ &\quad + 2\psi \varphi(\Theta) H_{\tilde{p}}^{2,0,-2} \psi + 2K \psi^2 \phi_\varepsilon^2 (\varepsilon x^{-1}) (1 + \varepsilon x^{-1})^{-1} \varphi \beta_1 + \phi_\varepsilon^2 \varphi(\Theta) \psi^2 p_2 \\ &\quad \left. + \psi^2 p_1 \varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} \phi_\varepsilon^2 \varphi \right]. \quad (373) \end{aligned}$$

Dividing by $\varrho_{\text{df}} x^{-1} (\sigma^2 + \mathbf{Z}x)^{-1/2}$,

$$\begin{aligned} H_{\tilde{p}} a^{(\varepsilon)} + p_1 a^{(\varepsilon)} &= \varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^{-2s} \varrho_{\text{ff}}^{-2\varsigma} \varrho_{\text{bf}}^{-2l} \varrho_{\text{tf}}^{-2\ell} \left[(\varphi_0(\Theta)^2 + \varphi_1(\Theta)) \phi_\varepsilon^2 \psi^2 \alpha \right. \\ &\quad + 2\psi \varphi(\Theta) H_{\tilde{p}}^{2,0,-2} \psi + 2K \psi^2 \phi_\varepsilon^2 (\varepsilon x^{-1}) (1 + \varepsilon x^{-1})^{-1} \varphi \beta_1 + \phi_\varepsilon^2 \varphi(\Theta) \psi^2 p_2 \\ &\quad \left. + \psi^2 p_1 \varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} \phi_\varepsilon^2 \varphi \right]. \quad (374) \end{aligned}$$

For each $K, \delta > 0$, we may choose $\psi = \psi_{K,\delta}$ (perhaps dilating it if necessary) such that its essential support is a subset of $\text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$ and such that, taking $F = F_{K,\delta,\psi} \in C^\infty([0, \infty)_\sigma; \mathbb{R}^+)$ sufficiently large,

$$\begin{aligned} b_\varepsilon &= \varrho_{\text{df}}^{-1/2} \varrho_{\text{sf}}^{-s} \varrho_{\text{ff}}^{-\varsigma} \varrho_{\text{bf}}^{-l} \varrho_{\text{tf}}^{-\ell} \phi_\varepsilon \varphi_{00}^{1/2} \varphi_{01} \psi \left[F(\Theta_5 - \Theta)^{-2} \alpha - 2K \frac{\varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \beta_1 - p_2 \right. \\ &\quad \left. - p_1 \varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} - 2\delta \phi_\varepsilon^2 \psi^2 \varphi \right]^{1/2} \quad (375) \end{aligned}$$

is a well-defined uniform family of leC-symbols, specifically $b_\bullet \in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, s, \varsigma, -\infty, -\infty}(X))$. In addition, we set

$$e_\varepsilon = \varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^{-2s} \varrho_{\text{ff}}^{-2\varsigma} \varrho_{\text{bf}}^{-2l} \varrho_{\text{tf}}^{-2\ell} \varphi_1(\Theta) \phi_\varepsilon^2 \psi^2 \alpha, \quad (376)$$

$$f_\varepsilon = 2\varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^{-2s} \varrho_{\text{ff}}^{-2\varsigma} \varrho_{\text{bf}}^{-2l} \varrho_{\text{tf}}^{-2\ell} \varphi(\Theta) \tilde{p}^{-1} \psi H_{\tilde{p}}^{2,0,-2} \psi \quad (377)$$

(the division by \tilde{p} in eq. (377) being unproblematic by eq. (367)). Thus,

$$e_\bullet \in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, 2s, 2\varsigma, -\infty, -\infty}(X)), \quad (378)$$

$$f_\bullet \in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, 2s, 2\varsigma+2, -\infty, -\infty}(X)). \quad (379)$$

In terms of these,

$$H_{\tilde{p}} a^{(\varepsilon)} + p_1 a^{(\varepsilon)} = 2\delta \varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^{2s+2} \varrho_{\text{ff}}^{2\varsigma+6} \varrho_{\text{bf}}^{2l+2} \varrho_{\text{tf}}^{2\ell+6} a^{(\varepsilon)2} + b_\varepsilon^2 + e_\varepsilon + f_\varepsilon \tilde{p}. \quad (380)$$

Setting $A_\varepsilon = (1/2)(\text{Op}(a^{(\varepsilon)}) + \text{Op}(a^{(\varepsilon)})^*)$, $B_\varepsilon = \text{Op}(b_\varepsilon)$, $E_\varepsilon = \text{Op}(e_\varepsilon)$, $F_\varepsilon = \text{Op}(f_\varepsilon)$,

$$-i[\Re \tilde{P}, A_\varepsilon] - \{\Im \tilde{P}, A_\varepsilon\} = 2\delta A_\varepsilon \Lambda_{1/2, -s-1, -\varsigma-3, -l-1, -\ell-3}^2 A_\varepsilon + B_\varepsilon^* B_\varepsilon + E_\varepsilon + F_\varepsilon^* \tilde{P} + R_\varepsilon \quad (381)$$

for some $R_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, 2s-1, 2\varsigma-1, 2l, 2\ell}(X))$. We have

$$\begin{aligned} A_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, 2s+1, 2\varsigma+3, -\infty, -\infty}(X)), \\ F_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, 2s, 2\varsigma+2, -\infty, -\infty}(X)), \\ B_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, s, \varsigma, -\infty, -\infty}(X)), \\ E_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, 2s, 2\varsigma, -\infty, -\infty}(X)), \end{aligned} \quad (382)$$

and

$$\begin{aligned} \text{WF}'_{L^\infty, \text{leC}}(A_\bullet), \text{WF}'_{L^\infty, \text{leC}}(B_\bullet), \text{WF}'_{L^\infty, \text{leC}}(E_\bullet), \\ \text{WF}'_{L^\infty, \text{leC}}(F_\bullet), \text{WF}'_{L^\infty, \text{leC}}(R_\bullet) \subset \text{supp } \psi\varphi(\Theta), \end{aligned} \quad (383)$$

where the last of these inclusions (the one for R_\bullet) follows from the one for A_\bullet and $\text{WF}'_{L^\infty}([\tilde{P}, A_\bullet]) \subset \text{WF}'_{L^\infty}(A_\bullet)$.

Now, for each $m_0, s_0, \varsigma_0, l_0, \ell_0 \in \mathbb{R}$, there exist $K > 0$ (dependent on $m_0, s_0, \varsigma_0, l_0, \ell_0$ and m, l, ℓ but nothing else) such that, given $\{u(-; \sigma)\}_{\sigma \geq 0} \in L^\infty([0, 2\Sigma]; H_{\text{leC}}^{m_0, s_0, \varsigma_0, l_0, \ell_0}(X))$, it is the case that, for any $\varepsilon > 0$ (and for each $\sigma > 0$, implicit in the notation below),

$$2\Im \langle \tilde{P}u, A_\varepsilon u \rangle_{L^2} = -\langle \{\Im \tilde{P}, A_\varepsilon\}u, u \rangle_{L^2} + i\langle [\Re \tilde{P}(\sigma), A_\varepsilon]u, u \rangle_{L^2}, \quad (384)$$

where the pairings above are well-defined (and where we are using the convention that $\langle -, - \rangle_{L^2}$ is antilinear in the first slot). Indeed, A_ε , $\{\Im \tilde{P}, A_\varepsilon\}$, and $[\Re \tilde{P}(\sigma), A_\varepsilon]$ are all smoothing operators – i.e. lying in $\Psi_{\text{scb}}^{-\infty, \infty, \infty}(X)$ if $\sigma > 0$ and $\Psi_{\text{scb}}^{-\infty, \infty, \infty}(X_{1/2})$ if $\sigma = 0$ (in a uniform sense made precise by the leC-calculus, but since we simply need to justify some integration by parts σ -wise the uniformity is not important here) – and by taking K large they can be made to induce an arbitrarily large amount of decay for each $\varepsilon > 0$. Given N , we fix $m_0, s_0, \varsigma_0, l_0, \ell_0 \in \mathbb{R}$ such that $m_0, s_0, \varsigma_0, l_0, \ell_0 < -N$. Then, we can take K depending on m, l, ℓ, N and nothing else such that eq. (384) holds for all $\{u(-; \sigma)\}_{\sigma \geq 0} \in L^\infty([0, \Sigma]; H_{\text{leC}}^{-N, -N, -N, -N, -N}(X))$.

Applying eq. (384) to $\{u(-; \sigma)\}_{\sigma > 0} \in H_{\text{leC}}^{-N, -N, -N, -N, -N}(X)$ and pairing against u (after taking K large enough), we have

$$\begin{aligned} 2\Im \langle \tilde{P}u, A_\varepsilon u \rangle_{L^2} &= \|B_\varepsilon u\|_{L^2}^2 + \langle u, E_\varepsilon u \rangle_{L^2} + \langle \tilde{P}u, F_\varepsilon u \rangle_{L^2} + \langle R_\varepsilon u, u \rangle_{L^2} \\ &\quad + 2\delta \|\Lambda_{1/2, -s-1, -\varsigma-3, -l-1, -\ell-3} A_\varepsilon u\|_{L^2}^2 \end{aligned} \quad (385)$$

for $\varepsilon > 0$. So,

$$\begin{aligned} \|B_\varepsilon u\|_{L^2}^2 + 2\delta \|\Lambda_{1/2, -s-1, -\varsigma-3, -l-1, -\ell-3} A_\varepsilon u\|_{L^2}^2 &\leq \\ 2|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| + |\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| + |\langle R_\varepsilon u, u \rangle_{L^2}| + |\langle u, E_\varepsilon u \rangle_{L^2}|. \end{aligned} \quad (386)$$

Fix self-adjoint $G \in \Psi_{\text{leC}}^{-\infty, 0, 0, -\infty, -\infty}(X)$ (constructed via Op) such that $\text{WF}'_{\text{leC}}{}^{0,0}(1-G)$ is disjoint from a neighborhood of the L^∞ -esssupp of a, b, f, e and such that $\text{WF}'_{\text{leC}}(G)$ is disjoint from \mathcal{R}_+ and satisfies $\text{WF}'_{\text{leC}}(G) \subset \text{Ell}_{\text{leC}}^{0,0,0,0,0}(G_2)$. (Both a, f are supported away from \mathcal{R}_+ and $\text{Char}_{\text{leC}}^{0,0,0,0,0}(G_2)$, so such a G exists.) Then (for K large enough):

- Writing $\tilde{P} = (1-G)\tilde{P} + G\tilde{P}$ and noting that

$$(1-G)A_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, -\infty, -\infty}(X)), \quad (387)$$

we have, for each $N \in \mathbb{N}$,

$$\begin{aligned} 2|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| &\leq 2|\langle G\tilde{P}u, A_\varepsilon u \rangle_{L^2}| + 2|\langle \tilde{P}u, (1-G)A_\varepsilon u \rangle_{L^2}| \\ &\leq 2^{-1}\bar{\delta}^{-1}\|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + 2\bar{\delta}\|\Lambda_{1/2, -s-1, -\varsigma-3, -l-1, -\ell-3} A_\varepsilon u\|_{L^2}^2 + \|u\|_{\mathcal{E}_N}^2, \end{aligned} \quad (388)$$

for any $\bar{\delta} > 0$, where the constant in eq. (388) is independent of $\bar{\delta}$. We have abbreviated $\mathcal{E}_N = H_{\text{leC}}^{-N, -N, -N, -N, -N}(X)$,

$$\mathcal{Y}_N = H_{\text{leC}}^{-N, s+1, \varsigma+3, l+1, \ell+3}(X). \quad (389)$$

We also set $\mathcal{Y}_{*,N} = H_{\text{leC}}^{-N, -(s+1), -(\varsigma+3), -(l+1), -(\ell+3)}(X)$ (so dual in all orders except that at df).

- Similarly, we can chose self-adjoint $\bar{G}_3 \in \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty}(X)$ with

$$\text{WF}_{\text{leC}}^{'0,0}(1 - \bar{G}_3) \cap \text{WF}_{L^\infty, \text{leC}}'(E_\bullet) = \emptyset, \quad (390)$$

$\text{WF}_{\text{leC}}'(\bar{G}_3) \cap \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subseteq \mathcal{P}[\Theta_1, \Theta_2]$, and $\text{WF}_{\text{leC}}'(\bar{G}_3) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$. Then

$$\begin{aligned} |\langle u, E_\varepsilon u \rangle_{L^2}| &\leq |\langle \bar{G}_3 u, E_\varepsilon u \rangle_{L^2}| + |\langle u, (1 - \bar{G}_3) E_\varepsilon u \rangle_{L^2}| \\ &\leq \|\bar{G}_3 u\|_{\mathcal{X}_N}^2 + \|E_\varepsilon u\|_{\mathcal{X}_N^*}^2 + \|u\|_{\mathcal{E}_N}^2, \\ &\leq \|\bar{G}_3 u\|_{\mathcal{X}_N}^2 + \|E_\varepsilon u\|_{\mathcal{X}_{*,N}}^2 + \|u\|_{\mathcal{E}_N}^2, \end{aligned} \quad (391)$$

where $\mathcal{X}_N = H_{\text{leC}}^{-N,s,\varsigma,l,\ell}(X)$, $\mathcal{X}_{*,N} = H_{\text{leC}}^{-N,-s,-\varsigma,-l,-\ell}(X)$.

The bound $\|E_\varepsilon u\|_{\mathcal{X}_N^*}^2 \leq \|E_\varepsilon u\|_{\mathcal{X}_{*,N}}^2 + \|u\|_{\mathcal{E}_N}^2$ follows (using eq. (382)) from the construction via Op of $H_1 \in \Psi_{\text{leC}}^{-1,0,0,0,0}(X)$ and $H_2 \in \Psi_{\text{leC}}^{0,0,0,0,0}(X)$ such that

$$\text{WF}_{L^\infty, \text{leC}}'(E_\bullet) \cap \text{WF}_{\text{leC}}^{'0,0}(H_2) = \emptyset \quad (392)$$

and $1 = H_1 + H_2$. Then, we can compute

$$\begin{aligned} \|E_\varepsilon u\|_{\mathcal{X}_N^*}^2 &\leq \|H_1 E_\varepsilon u\|_{\mathcal{X}_N^*}^2 + \|H_2 E_\varepsilon u\|_{\mathcal{X}_N^*}^2 \\ &\leq \|E_\varepsilon u\|_{H_{\text{leC}}^{-N-1,s,\varsigma,l,\ell}}^2 + \|H_2 E_\varepsilon u\|_{\mathcal{X}_N^*}^2 \\ &\leq \|E_\varepsilon u\|_{H_{\text{leC}}^{-N-1,s,\varsigma,l,\ell}}^2 + \|u\|_{\mathcal{E}_N}^2. \end{aligned} \quad (393)$$

Proceeding inductively, we deduce $\|E_\varepsilon u\|_{\mathcal{X}_N^*}^2 \leq \|E_\varepsilon u\|_{\mathcal{X}_{*,N}}^2 + \|u\|_{\mathcal{E}_N}^2$. This argument will be used below without further comment.

•

$$\begin{aligned} |\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| &\leq |\langle G\tilde{P}u, F_\varepsilon u \rangle_{L^2}| + |\langle \tilde{P}u, (1 - G)F_\varepsilon u \rangle_{L^2}| \\ &\leq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|F_\varepsilon u\|_{\mathcal{Y}_{*,N}}^2 + \|u\|_{\mathcal{E}_N}^2. \end{aligned} \quad (394)$$

We have $\text{WF}_{L^\infty, \text{leC}}'(F_\bullet) \cap \text{Char}^{2,0,-2,-1,-3}(G\tilde{P}) = \emptyset$ (since, by eq. (367), f_\bullet is supported away from the characteristic set), so we can deduce via the elliptic parametrix construction that

$$\begin{aligned} \|F_\varepsilon u\|_{\mathcal{Y}_{*,N}} &\leq \|F_\varepsilon u\|_{\mathcal{Y}_N^*} \leq \|G\tilde{P}u\|_{H_{\text{leC}}^{-N,s-1,\varsigma-1,-N,-N}} + \|u\|_{\mathcal{E}_N} \\ &\leq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|u\|_{\mathcal{E}_N}^2 \end{aligned} \quad (395)$$

for N sufficiently large, so

$$|\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| \leq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|u\|_{\mathcal{E}_N}^2. \quad (396)$$

- Writing $R_\varepsilon = (1 - G^2)R_\varepsilon + G^2 R_\varepsilon$, since $1 - G^2 = (1 - G)(1 + G)$ implies $\text{WF}_{L^\infty, \text{leC}}^{'0,0}(1 - G^2) \subset \text{WF}_{L^\infty, \text{leC}}^{'0,0}(1 - G)$, for each $N \in \mathbb{N}$ we have

$$\begin{aligned} |\langle R_\varepsilon u, u \rangle_{L^2}| &\leq (\|GR_\varepsilon u\|_{\mathcal{Z}_{*,N}} \|Gu\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N}^2) \leq (\|GR_\varepsilon u\|_{\mathcal{Z}_{*,N}}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|u\|_{\mathcal{E}_N}^2) \\ &\leq (\|\tilde{R}_\varepsilon Gu\|_{\mathcal{Z}_{*,N}}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|u\|_{\mathcal{E}_N}^2) \\ &\leq (\|Gu\|_{\mathcal{Z}_N}^2 + \|u\|_{\mathcal{E}_N}^2) \end{aligned} \quad (397)$$

for some $\tilde{R}_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, 2s-1, 2\varsigma-1, 2l, 2\ell}(X))$, where

$$\begin{aligned} \mathcal{Z}_N &= H_{\text{leC}}^{-N, (2s-1)/2, (2\varsigma-1)/2, l, \ell}(X), \\ \mathcal{Z}_{*,N} &= H_{\text{leC}}^{-N, -(2s-1)/2, -(2\varsigma-1)/2, -l, -\ell}(X). \end{aligned} \quad (398)$$

Combining eq. (386), eq. (388) with $\bar{\delta}$ sufficiently small, eq. (396), eq. (397), we have proven that

$$\|B_\varepsilon u\|_{L^2}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|\bar{G}_3u\|_{\mathcal{X}_N}^2 + \|E_\varepsilon u\|_{\mathcal{X}_{*,N}}^2 + \|u\|_{\mathcal{E}_N}^2, \quad (399)$$

i.e.

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|\bar{G}_3u\|_{\mathcal{X}_N}^2 + \|\hat{E}_\varepsilon u\|_{\mathcal{X}_N}^2 + \|u\|_{\mathcal{E}_N}^2, \quad (400)$$

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N} + \|Gu\|_{\mathcal{Z}_N} + \|\bar{G}_3u\|_{\mathcal{X}_N} + \|\hat{E}_\varepsilon u\|_{\mathcal{X}_N} + \|u\|_{\mathcal{E}_N}, \quad (401)$$

where $\hat{B}_\varepsilon = \Lambda_{0,-s,-\varsigma,0,0}B_\varepsilon$ and $\hat{E}_\varepsilon = \Lambda_{-2s,-2\varsigma,0,0}E_\varepsilon$. Hence, $\hat{B}_\bullet, \hat{E}_\bullet \in \Psi_{\text{leC}}^{0,0,0,0}(X)$. By the choice of ψ , $\text{WF}'_{L^\infty, \text{leC}}(B_\bullet), \text{WF}'_{L^\infty, \text{leC}}(E_\bullet) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$. Since $\hat{E}_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty})$,

$$\|\hat{E}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|\hat{E}_\varepsilon u\|_{\bar{\mathcal{X}}_N} \quad (402)$$

for $\bar{\mathcal{X}}_N = H_{\text{leC}}^{-N,s,\varsigma,-N,-N}(X)$. Since $\text{WF}'_{L^\infty, \text{leC}}(E_\bullet) \cap \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subset \mathcal{P}[\Theta_1, \Theta_2]$,

$$\text{WF}'_{L^\infty, \text{leC}}(E_\bullet) \subset \text{Ell}_{\text{leC}}^{2,0,-2,-1,-3}(G\tilde{P}) \cup \text{Ell}_{\text{leC}}^{0,0,0,0}(G_3). \quad (403)$$

Consequently, $\|\hat{E}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\bar{\mathcal{Y}}_N} + \|G_3u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}$ for sufficiently large N , where $\bar{\mathcal{Y}}_N = H_{\text{leC}}^{-N,s+1,\varsigma+3,-N,-N}(X)$. Likewise,

$$\|\bar{G}_3u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\bar{\mathcal{Y}}_N} + \|G_3u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (404)$$

Letting $\bar{\mathcal{Z}}_N = H_{\text{leC}}^{-N,s-1/2,\varsigma-1/2,-N,-N}(X)$, since $G \in \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty}(X)$, $\|Gu\|_{\mathcal{Z}_N} \preceq \|Gu\|_{\bar{\mathcal{Z}}_N}$.

Because $\text{WF}'_{\text{leC}}(G) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$, $\|G\tilde{P}u\|_{\bar{\mathcal{Y}}_N} \preceq \|G_2\tilde{P}u\|_{\bar{\mathcal{Y}}_N} + \|u\|_{\mathcal{E}_N}$ and $\|Gu\|_{\bar{\mathcal{Z}}_N} \preceq \|G_2u\|_{\bar{\mathcal{Z}}_N} + \|u\|_{\mathcal{E}_N}$ for sufficiently large N . We have therefore shown that

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G_2\tilde{P}u\|_{\bar{\mathcal{Y}}_N} + \|G_2u\|_{\bar{\mathcal{Z}}_N} + \|G_3u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (405)$$

Thus, for each $\sigma \geq 0$, $\hat{B}_\varepsilon u(-; \sigma)$ is uniformly bounded in $\mathcal{X}_N(\sigma)$ as $\varepsilon \rightarrow 0^+$.

For each $\sigma \geq 0$, given any sequence $\{\varepsilon_k\}_{k \in \mathbb{N}} \subset (0, 1]$ with $\varepsilon_k \rightarrow 0$ as $k \rightarrow \infty$, there exists – by the Banach–Alaoglu theorem – a subsequence ε_{k_κ} thereof such that

$$\{\hat{B}_{\varepsilon_{k_\kappa}}(\sigma)u(-; \sigma)\}_{\kappa \in \mathbb{N}} \subset \mathcal{X}_N(\sigma) \quad (406)$$

is weakly convergent in the scb-Sobolev space $\mathcal{X}_N(\sigma)$ and in fact in any closed ball in $\mathcal{X}_N(\sigma)$ in which $\hat{B}_{\varepsilon_k}u(-; \sigma)$ is eventually contained. Call the weak limit $v = v(N, \sigma, \{\varepsilon_{k_\kappa}\}_{\kappa \in \mathbb{N}}) \in \mathcal{X}_N(\sigma)$. The preceding clause means that

$$\|v\|_{\mathcal{X}_N(\sigma)} \leq \limsup_{\kappa \rightarrow \infty} \|\hat{B}_{\varepsilon_{k_\kappa}}(\sigma)u(-; \sigma)\|_{\mathcal{X}_N(\sigma)}. \quad (407)$$

The family $\{B_\varepsilon(\sigma)\}_{\varepsilon \in [0,1]}$ was constructed so that it is continuous down to $\varepsilon = 0$ with respect to the topology of some space of high order Ψ DOs. (This follows from the analogous observation for $b_\bullet(\sigma)$ and the continuity of the quantization map.) Consequently,

$$\hat{B}_{\varepsilon_k}(\sigma)u(-; \sigma) \rightarrow \hat{B}_0(\sigma)u(-; \sigma) \quad (408)$$

in the topology of $\mathcal{S}'(X)$ as $k \rightarrow \infty$. But $\hat{B}_{\varepsilon_{k_\kappa}}(\sigma)u(-; \sigma) \rightarrow v$ in $\mathcal{S}'(X)$, so $v = \hat{B}_0(\sigma)u(-; \sigma)$. The sequence $\{\varepsilon_k\}_{k \in \mathbb{N}}$ was arbitrary, so we can actually conclude from eq. (407) that

$$\|\hat{B}_0(\sigma)u(-; \sigma)\|_{\mathcal{X}_N(\sigma)} \leq \liminf_{\varepsilon \rightarrow 0^+} \|\hat{B}_\varepsilon u(-; \sigma)\|_{\mathcal{X}_N(\sigma)}. \quad (409)$$

This applies for each $\sigma \geq 0$, so

$$\|\hat{B}_0u\|_{\mathcal{X}_N} \preceq \|G_2\tilde{P}u\|_{\bar{\mathcal{Y}}_N} + \|G_2u\|_{\bar{\mathcal{Z}}_N} + \|G_3u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (410)$$

Since φ is nonvanishing on $[\Theta_3, \Theta_4]$, we have $\text{Ell}_{\text{leC}}^{0,0,0,0}(\hat{B}_0) \supseteq \mathcal{P}[\Theta_3, \Theta_4] \supseteq \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \cap \text{WF}'_{\text{leC}}(G_1)$, so (via elliptic regularity) $\|G_1 u\|_{\mathcal{X}} \preceq \|\hat{B}_0 u\|_{\mathcal{X}_N} + \|G_2 \tilde{P} u\|_{\tilde{\mathcal{Y}}_N} + \|u\|_{\mathcal{E}_N}$, where $\mathcal{X} = H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$. This yields

$$\|G_1 u\|_{\mathcal{X}} \preceq \|G_2 \tilde{P} u\|_{\tilde{\mathcal{Y}}_N} + \|G_2 u\|_{\tilde{\mathcal{Z}}_N} + \|G_3 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (411)$$

Observe that the leC-Sobolev space \mathcal{Z}_N is lower order than \mathcal{X} at sf and ff. Since the leC-Sobolev spaces $\mathcal{X}_N, \mathcal{Y}_N$ get bigger as s, ς decrease, an inductive argument (which we can carry out because eq. (362)) upgrades eq. (411) to eq. (361). \square

Since $H_{\tilde{p}}^{2,0,-2}$, viewed as a vector field on ${}^{\text{sc,leC}}T^*X$, vanishes at the two radial sets, in order to carry out a positive commutator argument we must take into account the previously negligible radial component

$$(H_{\tilde{p}_0}^{-,0,-2} x) \partial_x = 2(\xi_{\text{sc,leC}} - 1)x \partial_x \quad (412)$$

of the rescaled flow. For $\xi_{\text{sc,leC}} = 0, 2$, this is $\pm 2x \partial_x$, which (projecting down to X) is the ur-example of a nondegenerate radial b-vector field on X . Moreover, \mathcal{R}_+ is a source in the radial direction (as well as the other directions, as seen earlier), so \mathcal{R}_+ is a source for the Hamiltonian flow in all directions. (Similarly, \mathcal{R}_0 is a sink in all directions.) It is therefore straightforward to prove a radial point estimate at \mathcal{R}_+ .

To begin:

Proposition 5.8. *There exist $\beta_1, \beta_2 \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ such that*

$$H_{\tilde{p}}^{2,0,-2} x = \beta_1 x \quad (413)$$

$$H_{\tilde{p}}^{2,0,-2}(\sigma^2 + Zx)^{1/2} = \beta_2(\sigma^2 + Zx)^{1/2}, \quad (414)$$

with $\beta_1, \beta_2 > 0$ on \mathcal{R}_+ . \blacksquare

Proof. It suffices to work near ∂X . There, applying Proposition 5.4,

$$H_{\tilde{p}_0}^{-,0,-2} x = 2(\xi_{\text{sc,leC}} - 1)x, \quad (415)$$

$$H_{\tilde{p}_0}^{-,0,-2}(\sigma^2 + Zx)^{1/2} = (\xi_{\text{sc,leC}} - 1)Zx(\sigma^2 + Zx)^{-1/2}. \quad (416)$$

So, defining β_1, β_2 by eq. (413) and eq. (414), $\beta_1, \beta_2 \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ and are given at $\{x = 0\}$ by $\beta_1 = 2\varrho_{\text{df}}(\xi_{\text{sc,leC}} - 1)$ and $\beta_2 = \varrho_{\text{df}}(\xi_{\text{sc,leC}} - 1)Zx(\sigma^2 + Zx)^{-1}$. Thus, $\beta_1, \beta_2 > 0$ on \mathcal{R}_+ . \square

Proposition 5.9. *Suppose that $G_1, G_2, G_3 \in \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty}(X)$ satisfy*

- (1) $\text{WF}'_{\text{leC}}(G_1) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$,
- (2) $\mathcal{R}_+ \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_3), \text{Ell}_{\text{leC}}^{0,0,0,0}(G_1)$,
- (3) *there exists some $\Theta \in (0, \pi)$ such that*

$$\text{WF}'_{\text{leC}}(G_1) \cap \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subseteq \mathcal{R}_+ \cup \bigcup_{0 < \Theta' < \Theta} \mathcal{P}[\Theta', \Theta] \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2). \quad (417)$$

Then, for every $\Sigma > 0$ and $N \in \mathbb{N}$, $m, s, \varsigma, l, \ell, s_0, \varsigma_0 \in \mathbb{R}$ such that $s > s_0 > -1/2$ and $\varsigma > \varsigma_0 > -3/2$, there exists some constant

$$C = C(\tilde{P}, G_1, G_2, \Sigma, m, s, \varsigma, l, \ell, s_0, \varsigma_0, N) > 0 \quad (418)$$

such that, for all $u \in \mathcal{S}'(X)$,

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C \left[\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{-N,s+1,\varsigma+3,-N,-N}} + \|G_3 u\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}} + \|u\|_{H_{\text{leC}}^{-N,-N,-N,-N,-N}} \right] \quad (419)$$

holds (in the usual strong sense, i.e. the left-hand side is finite if the right-hand side is) for all $\sigma \in [0, \Sigma]$. \blacksquare

Proof. The symbolic constructions will only be specified for $\sigma \in [0, \Sigma]$, which is evidently unproblematic. Also, it suffices to take N to be sufficiently large such that any of the finitely many functions of m, s, ς, l, ℓ that arise below can be bounded below $-N$.

We can assume without loss of generality that $\text{WF}'_{\text{leC}}(G_3) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$ and

$$\text{WF}'_{\text{leC}}(G_3) \cap \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subseteq \mathcal{R}_+ \cup \cup_{0 < \Theta' < \Theta} \mathcal{P}[\Theta', \Theta] \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2). \quad (420)$$

We read off Proposition 5.8 that $\sup_{\mathcal{R}_+} \beta_1, \sup_{\mathcal{R}_+} \beta_2 > 0$.

Now choose nonnegative $\rho \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ equal to $(\xi_{\text{sc,leC}} - 2)^2 + \eta_{\text{sc,leC}}^2$ in some neighborhood of $\mathcal{R}_+ = \{\xi_{\text{sc}} = 2, \eta_{\text{sc}} = 0\} \subset \text{sf} \cup \text{ff}$.

There exist some symbols, which we call $\tilde{\beta}_0, \tilde{F}_2, \tilde{F}_3, \tilde{F}_4 \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ such that

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2} \rho &= \tilde{\beta}_0 \rho + \tilde{F}_2 + \tilde{F}_3 + x(\sigma^2 + Zx)^{-1/2} \tilde{F}_4 \\ \inf \tilde{\beta}_0|_{\mathcal{R}_+} &> 0, \end{aligned} \quad (421)$$

$\tilde{\beta}_0, \tilde{F}_2 \geq 0$ everywhere, and \tilde{F}_3 vanishes cubically at \mathcal{R}_+ (uniformly in σ). This computation is completely analogous to the one in the proof of Proposition 5.5. We now consider the weight (for to-be-decided $l', \ell' \in \mathbb{R}$)

$$a_0 = x^{-l'}(\sigma^2 + Zx)^{-\ell'/2+l'} \in S_{\text{cl,leC}}^{0,l',\ell',l'}(X). \quad (422)$$

Then, the symbol $\beta \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ defined by

$$\beta = -l'\beta_1 + (2l' - \ell')\beta_2 = a_0^{-1} H_{\tilde{p}}^{2,0,-2} a_0 \quad (423)$$

has a definite sign near \mathcal{R}_+ if $l', -l' + \ell'/2 \neq 0$ and have the same sign. Using the explicit formula for β_1, β_2 in the proof of Proposition 5.8,

$$\beta = -\varrho_{\text{df}} \left(2l' + (\ell' - 2l) \frac{Zx}{\sigma^2 + Zx} \right) (\xi_{\text{sc,leC}} - 1) \quad (424)$$

at $\{x = 0\}$. Negativity on $\mathcal{R}_+ \cap \text{sf}$ requires that $l' > 0$. negativity on $\mathcal{R}_+ \cap \text{ff}$ requires that $\ell' > 0$. And this suffices; for $l', \ell' > 0$, $\beta < 0$ in some neighborhood of \mathcal{R}_+ .

There exists $\chi \in C_c^\infty(\mathbb{R})$ such that $-\text{sgn}(t)\chi'(t)\chi(t) = \chi_0^2(t)$ for some $\chi_0 \in C_c^\infty(\mathbb{R})$ and such that $\chi = 1$ identically in some neighborhood of the origin. (The construction is standard and uses a translate of $\exp(-1/t)$.) Replacing χ with $\chi \circ \text{dil}_\lambda = \chi(\lambda \bullet)$ for sufficiently large $\lambda = \lambda(l', \ell', \chi)$ if necessary, choose χ such that

$$\beta|_{\text{supp } \chi(\rho)} < 0 \quad (425)$$

and such that $\text{supp } \chi(\tilde{p}^{2,0,-2})\chi(\rho)$ is disjoint from $\text{df} \cup \text{bf} \cup \text{tf}$, where $\tilde{p}^{2,0,-2} = (\sigma^2 + Zx)^{-1}\tilde{p}$.

Choose $\psi \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ such that ψ is identically equal to one in some neighborhood of $\{x = 0\}$ and such that the formula $(H_{\tilde{p}}^{2,0,-2}\rho)^{1/2}\psi\chi_0(\rho)$ defines a symbol:

$$\chi_0(\rho)\psi\sqrt{H_{\tilde{p}}^{2,0,-2}}\rho \in S_{\text{cl,leC}}^{0,0,0,0}(X). \quad (426)$$

(The existence of such a ψ follows from eq. (421).) Now set

$$a = a_0\psi^2\chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2 \in S_{\text{cl,leC}}^{-\infty,l',\ell',-\infty,-\infty}(X). \quad (427)$$

We compute

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2} a &= \chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2 H_{\tilde{p}}^{2,0,-2} a_0 - 2a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2})^2 H_{\tilde{p}}^{2,0,-2} \rho \\ &\quad + 2a_0\chi(\tilde{p}^{2,0,-2})\chi'(\tilde{p}^{2,0,-2})\chi(\tilde{p})^2\tilde{p}^{2,0,-2}\tilde{q} + 2a_0\chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2\psi H_{\tilde{p}}^{2,0,-2}\psi. \end{aligned} \quad (428)$$

Here $\tilde{q} \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ is defined by $\tilde{q} = (\sigma^2 + Zx)H_{\tilde{p}}^{2,0,-2}(\sigma^2 + Zx)^{-1}$, so that

$$H_{\tilde{p}}^{2,0,-2}\tilde{p}^{2,0,-2} = \tilde{q}\tilde{p}^{2,0,-2}. \quad (429)$$

In terms of β , eq. (428) says

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2}a &= \chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2\beta a_0 - 2a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2})^2H_{\tilde{p}}^{2,0,-2}\rho \\ &\quad + 2a_0\chi(\tilde{p}^{2,0,-2})\chi'(\tilde{p}^{2,0,-2})\chi(\rho)^2\tilde{p}^{2,0,-2}\tilde{q} + 2a_0\chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2\psi H_{\tilde{p}}^{2,0,-2}\psi. \end{aligned} \quad (430)$$

The first two terms have definite sign (the same sign for $l', \ell' > 0$), while the third and fourth terms are unproblematic (being controllable by elliptic or propagation estimates).

Set $\phi_\varepsilon = (1 + \varepsilon x^{-1})^{-K_1}(1 + \varepsilon(\sigma^2 + Zx)^{-1/2})^{-K_2}$ and $a^{(\varepsilon)} = \phi_\varepsilon^2 a \in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, l', \ell', -\infty, -\infty}(X))$ for to-be-decided $K_1, K_2 \in \mathbb{R}$. The replacement for eq. (430) is

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2}a^{(\varepsilon)} &= \phi_\varepsilon^2 \left[-\chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2a_0\psi^2 \left(\beta_1 \left(l' - \frac{K_1\varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) + \beta_2 \left(\ell' - 2l' - \frac{K_2\varepsilon(\sigma^2 + Zx)^{-1/2}}{1 + \varepsilon(\sigma^2 + Zx)^{-1/2}} \right) \right) \right. \\ &\quad - 2a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2})^2\psi^2H_{\tilde{p}}^{2,0,-2}\rho + 2a_0\chi'(\tilde{p}^{2,0,-2})\chi(\tilde{p}^{2,0,-2})\chi(\rho)^2\psi^2\tilde{p}^{2,0,-2}\tilde{q} \\ &\quad \left. + 2a_0\chi(\tilde{p}^{2,0,-2})^2\chi(\rho)^2\psi H_{\tilde{p}}^{2,0,-2}\psi \right]. \end{aligned} \quad (431)$$

Rewriting the first parenthetical,

$$\begin{aligned} &\beta_1 \left(l' - \frac{K_1\varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) + \beta_2 \left(\ell' - 2l' - \frac{K_2\varepsilon(\sigma^2 + Zx)^{-1/2}}{1 + \varepsilon(\sigma^2 + Zx)^{-1/2}} \right) \\ &= \varrho_{\text{df}} \left(2 \left(l' - \frac{K_1\varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) + \frac{Zx}{\sigma^2 + Zx} \left(\ell' - 2l' - \frac{K_2\varepsilon(\sigma^2 + Zx)^{-1/2}}{1 + \varepsilon(\sigma^2 + Zx)^{-1/2}} \right) \right) (\xi_{\text{sc,leC}} - 1). \end{aligned} \quad (432)$$

So, we will require that

$$K_1 < l', \quad K_2 < \ell', \quad (433)$$

and then the quantity in eq. (432) is positive in some neighborhood of \mathcal{R}_+ . Thus, only a limited amount of “regularization” can be performed. This is a standard technicality, and we can deal with it via citing the standard arguments used to handle it elsewhere — see [Vas18]. We do not even need to worry about uniformity: we can justify the formal integrations-by-parts below σ -wise, by citing essentially verbatim the arguments in [Vas18] for the $\sigma > 0$ and applying the argument with $X_{1/2}$ in place of X to handle the $\sigma = 0$ case. (In fact, since the $\sigma = 0$ case of the proposition follows from the estimates in [Vas21a] applied on $X_{1/2}$, to prove the proposition here it suffices to prove estimates that are uniform as $\sigma \rightarrow 0^+$, and thus to restrict attention to the $\sigma > 0$ case, for which we can take $K_2 = 0$ and apply [Vas18] essentially verbatim.)

Given that the inequalities eq. (433) are satisfied, we can (perhaps dilating χ or shrinking the support of ψ if necessary) find

$$\underline{\delta} = \underline{\delta}(K_1, K_2, l', \ell', \chi) > 0 \quad (434)$$

sufficiently small such that there exist uniform families of leC-symbols

$$\begin{aligned} b_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, (l'-1)/2, (\ell'-3)/2, -\infty, -\infty}(X)), \\ e_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, (l'-1)/2, (\ell'-3)/2, -\infty, -\infty}(X)), \\ f_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, l'-1, \ell'-1, -\infty, -\infty}(X)), \\ r_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, -\infty, -\infty, -\infty, -\infty}(X)). \end{aligned} \quad (435)$$

given by

$$\begin{aligned}
b_\varepsilon &= \varrho_{\text{df}}^{-1/2} \varrho_{\text{sf}}^{1/2} \varrho_{\text{ff}}^{3/2} \varrho_{\text{bf}}^{1/2} \varrho_{\text{tf}}^{3/2} a_0^{1/2} \chi(\tilde{p}^{2,0,-2}) \chi(\rho) \phi_\varepsilon \psi \left[-\varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} p_1 + \beta_1 \left(l' - \frac{K_1 \varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) \right. \\
&\quad \left. + \beta_2 \left(\ell' - 2l' - \frac{K_2 \varepsilon (\sigma^2 + \mathbf{Z}x)^{-1/2}}{1 + \varepsilon (\sigma^2 + \mathbf{Z}x)^{-1/2}} \right) - 2\delta \phi_\varepsilon^2 \chi(\tilde{p}^{2,0,-2})^2 \chi(\rho)^2 \right]^{1/2} \\
e_\varepsilon &= \varrho_{\text{df}}^{-1/2} \varrho_{\text{sf}}^{1/2} \varrho_{\text{ff}}^{3/2} \varrho_{\text{bf}}^{1/2} \varrho_{\text{tf}}^{3/2} a_0^{1/2} \phi_\varepsilon \psi \chi_0(\rho) \chi(\tilde{p}^{2,0,-2}) \sqrt{2H_{\tilde{p}}^{2,0,-2} \rho} \\
f_\varepsilon &= 2\varrho_{\text{df}} \varrho_{\text{sf}} \varrho_{\text{ff}} \varrho_{\text{bf}} \varrho_{\text{tf}} \phi_\varepsilon^2 a_0 \chi(\tilde{p}^{2,0,-2}) \chi'(\tilde{p}^{2,0,-2}) \chi(\rho)^2 \psi^2 \tilde{q} \\
r_\varepsilon &= 2\varrho_{\text{df}} \varrho_{\text{sf}} \varrho_{\text{ff}}^3 \varrho_{\text{bf}} \varrho_{\text{tf}}^3 \phi_\varepsilon^2 a_0 \chi(\tilde{p}^{2,0,-2})^2 \chi(\rho)^2 \psi H_{\tilde{p}}^{2,0,-2} \psi.
\end{aligned} \tag{436}$$

Here $p_1 \in S_{\text{leC}}^{2,-1,-3,-1-\delta,-3-2\delta}(X)$ is as in the proof of the propagation estimate, and we are using Proposition 3.9, which shows that $\varrho_{\text{df}}^2 \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} p_1$ vanishes to some fractional order at \mathcal{R}_+ and therefore does not spoil the sign of the quantity under the first square root in eq. (436) for an appropriate choice of χ, ψ .

In terms of these new symbols, we can write

$$\begin{aligned}
H_{\tilde{p}} a^{(\varepsilon)} + p_1 a^{(\varepsilon)} &= -2\delta \varrho_{\text{df}}^{-1} \varrho_{\text{sf}} \varrho_{\text{ff}}^3 \varrho_{\text{bf}} \varrho_{\text{tf}}^3 a^{(\varepsilon)} \phi_\varepsilon^2 \chi(\tilde{p}^{2,0,-2})^2 \chi(\rho)^2 - b_\varepsilon^2 - e_\varepsilon^2 + f_\varepsilon \tilde{p} + r_\varepsilon \\
&= -2\delta \varrho_{\text{df}}^{-1} \varrho_{\text{sf}} \varrho_{\text{ff}}^3 \varrho_{\text{bf}} \varrho_{\text{tf}}^3 a_0^{-1} a^{(\varepsilon)2} - b_\varepsilon^2 - e_\varepsilon^2 + f_\varepsilon \tilde{p} + r_\varepsilon
\end{aligned} \tag{437}$$

We apply the quantization map Op . Setting $A_\varepsilon = (1/2)(\text{Op}(a^{(\varepsilon)}) + \text{Op}(a^{(\varepsilon)})^*)$, $B_\varepsilon = \text{Op}(b_\varepsilon)$, $E_\varepsilon = \text{Op}(e_\varepsilon)$, $F_\varepsilon = \text{Op}(f_\varepsilon)$,

$$\begin{aligned}
A_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, l', \ell', -\infty, -\infty}(X)), \\
B_\bullet, E_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, (l'-1)/2, (\ell'-3)/2, -\infty, -\infty}), \\
F_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, l'-1, \ell'-1, -\infty, -\infty}(X)),
\end{aligned} \tag{438}$$

and

$$-i[\Re \tilde{P}, A_\varepsilon] - \{\Im \tilde{P}, A_\varepsilon\} = -2\delta A_\varepsilon \Lambda_{1/2, -(l'+1)/2, -(\ell'+3)/2, -(l'+1)/2, -(\ell'+3)/2} A_\varepsilon - B_\varepsilon^* B_\varepsilon - E_\varepsilon^* E_\varepsilon + F_\varepsilon^* \tilde{P} + R_\varepsilon \tag{439}$$

for some $R_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, l'-2, \ell'-4, -\infty, -\infty}(X))$. We have

$$\text{WF}'_{L^\infty, \text{leC}}(A_\bullet), \text{WF}'_{L^\infty, \text{leC}}(B_\bullet), \text{WF}'_{L^\infty, \text{leC}}(F_\bullet), \text{WF}'_{L^\infty}(E_\bullet), \text{WF}'_{L^\infty, \text{leC}}(R_\bullet) \subset \text{supp } \chi(\tilde{p}^{2,0,-2}) \chi(\rho) \psi. \tag{440}$$

We now set the parameters l', ℓ' in the definition eq. (422) of a_0 to $l' = 2s+1 > 0$ and $\ell' = 2\varsigma+3 > 0$. Fix $K_1 \in (0, l')$, $K_2 \in (0, \ell')$ such that $-1/2 < s - K_1 < s_0$ and $-3/2 < \varsigma - K_2 < \varsigma_0$. Suppose now that $u \in \mathcal{S}'(X)$, $\sigma \in [0, \Sigma]$ are such that

$$\|G_3 u\|_{H_{\text{leC}}^{-N, s_0, \varsigma_0, -N, -N}} < \infty. \tag{441}$$

The argument in [Vas18, §4.7, above Proposition 5.27] justifies the computation

$$2\Im \langle \tilde{P} u, A_\varepsilon u \rangle_{L^2} = -\langle \{\Im \tilde{P}, A_\varepsilon\} u, u \rangle_{L^2} + i \langle [\Re \tilde{P}(\sigma), A_\varepsilon] u, u \rangle_{L^2} \tag{442}$$

$$\begin{aligned}
&= -\|B_\varepsilon u\|_{L^2}^2 - \|E_\varepsilon u\|_{L^2}^2 + \langle \tilde{P} u, F_\varepsilon u \rangle_{L^2} + \langle R_\varepsilon u, u \rangle_{L^2} \\
&\quad - 2\delta \|\Lambda_{1/2, -s-1, -\varsigma-3, -s-1, -\varsigma-3} A_\varepsilon u\|_{L^2}^2
\end{aligned} \tag{443}$$

where the individual terms above are all well-defined distributional pairings (in the sense of Hörmander) or (finite) norms. Thus,

$$\|B_\varepsilon u\|_{L^2}^2 + \|E_\varepsilon u\|_{L^2}^2 + 2\delta\|\Lambda_{1/2,-s-1,-\varsigma-3,-s-1,-\varsigma-3}A_\varepsilon u\|_{L^2}^2 \leq 2|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| + |\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| + |\langle R_\varepsilon u, u \rangle_{L^2}|. \quad (444)$$

We estimate each of the terms on the right-hand side as in the proof of the propagation estimate: for self-adjoint $G \in \Psi_{\text{leC}}^{-\infty,0,0,-\infty,-\infty}(X)$ with $1 - G$ essentially supported away from the L^∞ -essential support of a, f, e, b and with $\text{WF}'_{\text{leC}}(G) \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$,

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \leq \|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} + \|\hat{E}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N} + \|Gu\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N}, \quad (445)$$

where $\mathcal{E}_N = H_{\text{leC}}^{-N,-N,-N,-N}(X)$, $\mathcal{Z}_N = H_{\text{leC}}^{-N,(2s-1)/2,(2\varsigma-1)/2,-N,-N}(X)$, $\mathcal{X}_N = H_{\text{leC}}^{-N,s,\varsigma,-N,-N}(X)$, $\mathcal{Y}_N = H_{\text{leC}}^{-N,s+1,\varsigma+3,-N,-N}(X)$, and \hat{B}_ε and \hat{E}_ε are given by $\hat{B}_\varepsilon = \Lambda_{0,-s,-\varsigma,0,0}B_\varepsilon$ and

$$\hat{E}_\varepsilon = \Lambda_{0,-s,-\varsigma,0,0}E_\varepsilon. \quad (446)$$

Via elliptic regularity, we can estimate

$$\|G\tilde{P}u\|_{\mathcal{Y}_N} \preceq \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|u\|_{\mathcal{E}_N}. \quad (447)$$

By shrinking the support of χ, ψ if necessary, we can arrange that the L^∞ -esssupp of a, b, f, e is a subset $\text{Ell}_{\text{leC}}^{0,0,0,0}(G_3)$, and then we can choose G such that $\text{WF}'_{\text{leC}}(G) \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_3)$, so that the estimate eq. (445) implies

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|G_3u\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N}. \quad (448)$$

Using the Banach–Alaoglu theorem, applied as during the proof of the propagation estimate, we can take $\varepsilon \rightarrow 0^+$ to conclude

$$\|\hat{B}_0 u\|_{\mathcal{X}_N} \preceq \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|G_3u\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N}. \quad (449)$$

Let $\mathcal{X} = H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)$. Since $\text{Ell}_{\text{leC}}^{0,0,0,0}(\hat{B}_0) \supset \mathcal{R}_+$, eq. (449) implies

$$\|G_1 u\|_{\mathcal{X}} \preceq \|\hat{B}_0 u\|_{\mathcal{X}_N} + \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|u\|_{\mathcal{E}_N}, \quad (450)$$

where we used the propagation estimate to control $G_1 u$ on $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P})$ away from \mathcal{R}_+ . So,

$$\begin{aligned} \|G_1 u\|_{\mathcal{X}} &\preceq \|\hat{B}_0 u\|_{\mathcal{X}_N} + \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|u\|_{\mathcal{E}_N} \preceq \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|G_3u\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N} \\ &\preceq \|G_2\tilde{P}u\|_{\mathcal{Y}_N} + \|G_3u\|_{H_{\text{leC}}^{-N,s-1/2,\varsigma-1/2,-N,-N}} + \|u\|_{\mathcal{E}_N}. \end{aligned} \quad (451)$$

An inductive argument (using eq. (420)) estimating $\|G_3u\|_{H_{\text{leC}}^{-N,s-1/2,\varsigma-1/2,-N,-N}}$ finishes the proof. \square

5.3. The Radial “Point” \mathcal{R} . In order to analyze matters uniformly near the corners of leC-phase space (in particular the highlighted edge $\text{bf} \cup \text{ff}$ in Figure 7), we work with the bdfs $\varrho_{\text{sf}}, \varrho_{\text{bf}}, \varrho_{\text{ff}}, \varrho_{\text{tf}}$ defined in eq. (91) rather than the leC-adapted momentum coordinates $\xi_{\text{sc,leC}}, \eta_{\text{sc,leC}}$ used in the previous section.

In terms of the bdfs,

$$\tilde{p}_0 = \varrho_{\text{df}}^{-2} \varrho_{\text{ff}}^2 \varrho_{\text{bf}}^2 \varrho_{\text{tf}}^4 (1 - \varrho_{\text{df}}^2 \varrho_{\text{sf}}^2 \varrho_{\text{ff}}^2) - 2 \left(\frac{\xi_{\text{b}}}{(1 + \xi_{\text{b}}^2 + g_{\partial X}^{-1}(\eta_{\text{b}}, \eta_{\text{b}}))^{1/2}} \right) \varrho_{\text{df}}^{-1} \varrho_{\text{ff}}^2 \varrho_{\text{bf}} \varrho_{\text{tf}}^3 \quad (452)$$

$$= \varrho_{\text{df}}^{-2} \varrho_{\text{ff}}^2 \varrho_{\text{bf}}^2 \varrho_{\text{tf}}^4 (1 - \varrho_{\text{df}}^2 \varrho_{\text{sf}}^2 \varrho_{\text{ff}}^2) - 2 \varrho_{\text{df}}^{-1} \varrho_{\text{ff}}^2 \varrho_{\text{bf}} \varrho_{\text{tf}}^3 \quad (453)$$

near $\{x = 0\} \subset {}^{\text{leC}}\overline{T}^*X$, where $\beth \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ is given by $\xi_b(1 + \xi_b^2 + g_{\partial X}^{-1}(\eta_b, \eta_b))^{-1/2}$ near $\{x = 0\}$. Factoring out common powers of bdfs from eq. (453), we are left with

$$\tilde{p}_0^{2,0,-2,-1,-3} = \varrho_{\text{bf}}\varrho_{\text{tf}}(1 - \varrho_{\text{df}}^2\varrho_{\text{sf}}^2\varrho_{\text{ff}}^2) - 2\beth\varrho_{\text{df}} \quad (454)$$

$$= \varrho_{\text{bf}}\varrho_{\text{tf}} - 2\beth\varrho_{\text{df}} \bmod S_{\text{leC}}^{-1,-1,-1,-1}(X). \quad (455)$$

On $\text{sf} \cup \text{ff}$, $\tilde{p}_0^{2,0,-2,-1,-3}$ and thus $\tilde{p}^{2,0,-2,-1,-3}$ vanishes if and only if $2\beth = \varrho_{\text{bf}}\varrho_{\text{tf}}/\varrho_{\text{df}}$. Therefore

$$\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) = \{2\beth = \varrho_{\text{bf}}\varrho_{\text{tf}}/\varrho_{\text{df}}\} \subset {}^{\text{leC}}T^*X. \quad (456)$$

In particular, the portion of the characteristic set that is on the boundary of $\text{bf} \cup \text{tf}$ is precisely $\{\beth = 0\} \cap (\text{bf} \cup \text{tf}) \cap (\text{sf} \cup \text{ff})$.

Proposition 5.10. *We have $H_{\tilde{p}_0}^{2,0,-2}\varrho_{\text{df}00} = F_{0,1}\varrho_{\text{df}00}$ for $F_{0,1} \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ given by*

$$F_{0,1} = 2\varrho_{\text{bf}}\varrho_{\text{tf}}\beth(1 - \varrho_{\text{df}}^2\varrho_{\text{sf}}^2\varrho_{\text{ff}}^2) - 2\varrho_{\text{df}}\beth^2 - Z\varrho_{\text{df}}\varrho_{\text{sf}}\varrho_{\text{bf}}\beth^2 \quad (457)$$

near $\{x = 0\} \subset {}^{\text{leC}}\overline{T}^*X$. ■

Proof. Applying eq. (342) to $\varrho_{\text{df}00}$,

$$H_{\tilde{p}_0}\varrho_{\text{df}00} = -(x\partial_x\tilde{p}_0)\partial_{\xi_b}\varrho_{\text{df}00} \quad (458)$$

near $\{x = 0\}$. Writing $\tilde{p}_0 = x^2\varrho_{\text{df}00}^{-2}(1 - \varrho_{\text{df}00}^2) - 2x\xi_b(\sigma^2 + Zx)^{1/2}$,

$$\begin{aligned} \partial_x\tilde{p}_0 &= 2x\varrho_{\text{df}00}^{-2}(1 - \varrho_{\text{df}00}^2) - 2\xi_b(\sigma^2 + Zx)^{1/2} - Zx\xi_b(\sigma^2 + Zx)^{-1/2} \\ \partial_{\xi_b}\varrho_{\text{df}00} &= -\xi_b\varrho_{\text{df}00}^3. \end{aligned} \quad (459)$$

Thus,

$$H_{\tilde{p}_0}\varrho_{\text{df}00} = 2x^2\xi_b\varrho_{\text{df}00}(1 - \varrho_{\text{df}00}^2) - 2x\xi_b^2\varrho_{\text{df}00}^3(\sigma^2 + Zx)^{1/2} - Z\xi_b^2\varrho_{\text{df}00}^3x^2(\sigma^2 + Zx)^{-1/2}. \quad (460)$$

In terms of bdfs ${}^{\text{leC}}\overline{T}^*X$, $x^2\xi_b\varrho_{\text{df}00} = \varrho_{\text{sf}}^2\varrho_{\text{ff}}^4\varrho_{\text{bf}}^2\varrho_{\text{tf}}^4\beth$, and $x\xi_b^2\varrho_{\text{df}00}^3(\sigma^2 + Zx)^{1/2} = \beth^2\varrho_{\text{df}}\varrho_{\text{sf}}^2\varrho_{\text{ff}}^4\varrho_{\text{bf}}\varrho_{\text{tf}}^3$, and $\xi_b^2\varrho_{\text{df}00}^3x^2(\sigma^2 + Zx)^{-1/2} = \beth^2\varrho_{\text{df}}\varrho_{\text{sf}}^3\varrho_{\text{ff}}^4\varrho_{\text{bf}}^2\varrho_{\text{tf}}^3$. Adding everything together, we find

$$H_{\tilde{p}_0}\varrho_{\text{df}00} = 2\varrho_{\text{sf}}^2\varrho_{\text{ff}}^4\varrho_{\text{bf}}^2\varrho_{\text{tf}}^4\beth(1 - \varrho_{\text{df}}^2\varrho_{\text{sf}}^2\varrho_{\text{ff}}^2) - 2\varrho_{\text{df}}\varrho_{\text{sf}}^2\varrho_{\text{ff}}^4\varrho_{\text{bf}}\varrho_{\text{tf}}^3\beth^2 - Z\varrho_{\text{df}}\varrho_{\text{sf}}^3\varrho_{\text{ff}}^4\varrho_{\text{bf}}^2\varrho_{\text{tf}}^3\beth^2 \quad (461)$$

Dividing by $\varrho_{\text{df}}^{-1}x(\sigma^2 + Zx)^{1/2} = \varrho_{\text{df}}^{-1}\varrho_{\text{sf}}\varrho_{\text{ff}}^3\varrho_{\text{bf}}\varrho_{\text{tf}}^3$,

$$H_{\tilde{p}_0}^{2,0,-2}\varrho_{\text{df}00} = (2\varrho_{\text{bf}}\varrho_{\text{tf}}\beth(1 - \varrho_{\text{df}}^2\varrho_{\text{sf}}^2\varrho_{\text{ff}}^2) - 2\varrho_{\text{df}}\beth^2 - Z\varrho_{\text{df}}\varrho_{\text{sf}}\varrho_{\text{bf}}\beth^2)\varrho_{\text{df}00}, \quad (462)$$

as claimed. □

Letting β_1, β_2 be as in Proposition 5.8, $\beta_1, \beta_2 < 0$ on \mathcal{R}_0 and thus on \mathcal{R} . By Proposition 5.10, there exists an $F_1 \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ such that

$$H_{\tilde{p}}^{2,0,-2}\varrho_{\text{df}00} = F_1\varrho_{\text{df}00}, \quad (463)$$

with F_1 vanishing on \mathcal{R} . By Proposition 5.5, it is the case that for any fixed $\rho \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ equal to $\xi_{\text{sc,leC}}^2 + \eta_{\text{sc,leC}}^2$ in some neighborhood of $\text{bf} \cup \text{tf}$, there exist some symbols

$$\beta_0, F_2, F_3, F_4 \in S_{\text{cl,leC}}^{0,0,0,0}(X) \quad (464)$$

such that

$$H_{\tilde{p}}^{2,0,-2}\rho = \beta_0\rho + F_2 + F_3 + x(\sigma^2 + Zx)^{-1/2}F_4 \quad (465)$$

$$\beta_0|_{\mathcal{R}_0} < 0, \quad (466)$$

$\beta_0, F_2 \leq 0$ everywhere, and F_3 vanishes cubically at \mathcal{R} uniformly in $[0, \Sigma]$. We may choose ρ such that it is nonnegative everywhere.

It is necessary to have another weight with semidefinite sign under the Hamiltonian flow. We may use

$$\varrho_{\text{ff}} = \varrho_{\text{df}00} + \varrho_{\text{tf}00} \in S_{\text{cl,leC}}^{0,0,0,0}(X). \quad (467)$$

We have $H_{\tilde{p}}^{2,0,-2}\varrho_{\text{ff}} = \beta_3\varrho_{\text{ff}}$ for $\beta_3 = (F_1\varrho_{\text{df}00} + \beta_2\varrho_{\text{tf}00})(\varrho_{\text{df}00} + \varrho_{\text{tf}00})^{-1} \in S_{\text{cl,leC}}^{0,0,0,0}(X)$. We can write this as $\beta_3 = F_1\varrho_{\text{df}}\varrho_{\text{sf}} + \beta_2\varrho_{\text{tf}}$. Thus, $\beta_3|_{\mathcal{R}} \leq 0$. Note that β_3 vanishes at $\mathcal{R} \cap \text{tf}$. Normally, this would be problematic as far as the radial point estimate is concerned, but this fourth weight is used only to give us an extra independent order and not to manufacture positivity, so a semidefinite sign is actually acceptable.

We now have enough basic weights to construct our commutants: the basic weights are $x = \varrho_{\text{bf}}\varrho_{\text{sf}}\varrho_{\text{tf}}^2\varrho_{\text{ff}}^2$, $(\sigma^2 + \mathbf{Z}x)^{1/2} = \varrho_{\text{tf}}\varrho_{\text{ff}}$, and $\varrho_{\text{df}00} = \varrho_{\text{df}}\varrho_{\text{sf}}\varrho_{\text{ff}}$. For any $s, \varsigma, l, \ell \in \mathbb{R}$,

$$x^l(\sigma^2 + \mathbf{Z}x)^{\ell/2-l}\varrho_{\text{df}00}^{s-l}\varrho_{\text{ff}}^{\varsigma-\ell-s+l} = \varrho_{\text{df}}^{s-l}\varrho_{\text{sf}}^s\varrho_{\text{ff}}^{\varsigma}\varrho_{\text{bf}}^l\varrho_{\text{tf}}^{\ell}. \quad (468)$$

We do not care about the order at df, so the weights of the form eq. (468) (which give four independent orders) suffice. We now consider the weight (dependent on parameters $s, \varsigma, l, \ell \in \mathbb{R}$)

$$a_0 = x^{-l}(\sigma^2 + \mathbf{Z}x)^{-\ell/2+l}\varrho_{\text{df}00}^{-s+l}\varrho_{\text{ff}}^{-\varsigma+\ell+s-l} \in S_{\text{cl,leC}}^{s-l,s,\varsigma,l,\ell}(X). \quad (469)$$

Per the above and eq. (432),

$$\beta = -((s-l)F_1 + (\varsigma-\ell-s+l)\beta_3 + l\beta_1 + (\ell-2l)\beta_2) = a_0^{-1}H_{\tilde{p}}^{2,0,-2}a_0 \in S_{\text{cl,leC}}^{0,0,0,0}(X) \quad (470)$$

will have a definite sign near \mathcal{R} if l, ℓ have the same definite sign and $\varsigma - \ell - s + l$ has the same sign semidefinitely.

Hence, multiplying a_0 by an appropriate microlocal cutoff, we can arrange for a to be everywhere monotonic under the Hamiltonian flow, strictly so near \mathcal{R} . Specifically, fix $\chi \in C_c^\infty(\mathbb{R}; [0, 1])$ as in the proof of Proposition 5.9. Dilating χ if necessary, we can find $\psi \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ that is identically equal to one in some neighborhood of $\text{sf} \cup \text{ff}$ such that

$$\beta < 0 \quad (471)$$

on $\text{supp } \chi(\tilde{p}^{2,0,-2,-1,-3})\chi(\rho)\psi$ and such that

$$\chi(\tilde{p}^{2,0,-2,-1,-3})\chi_0(\rho)\psi\sqrt{-H_{\tilde{p}}^{2,0,-2}}\rho \in S_{\text{cl,leC}}^{0,0,0,0}(X) \quad (472)$$

and $\text{supp } \chi(\tilde{p}^{2,0,-2,-1,-3})\chi(\rho)\psi \cap (\mathcal{R}_+ \cup \text{df}) = \emptyset$. Now set

$$a = a_0\chi(\tilde{p}^{2,0,-2,-1,-3})^2\chi(\rho)^2\psi^2 \in S_{\text{cl,leC}}^{s-l,s,\varsigma,l,\ell}(X). \quad (473)$$

The three factors $\chi(\tilde{p}^{2,0,-2,-1,-3})$, $\chi(\rho)$, $\psi \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ together microlocalize near \mathcal{R} .

We can write $H_{\tilde{p}}^{2,0,-2}\tilde{p}^{2,0,-2,-1,-3} = \tilde{q}\tilde{p}^{2,0,-2,-1,-3}$ for $\tilde{q} \in S_{\text{cl,leC}}^{0,0,0,0}(X)$ defined by

$$\tilde{q} = \varrho_{\text{df}}^{-2}\varrho_{\text{ff}}^2\varrho_{\text{bf}}^1\varrho_{\text{tf}}^3H_{\tilde{p}}^{2,0,-2}(\varrho_{\text{df}}^2\varrho_{\text{ff}}^{-2}\varrho_{\text{bf}}^{-1}\varrho_{\text{tf}}^{-3}). \quad (474)$$

Then,

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2}a &= \psi^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\chi(\rho)^2H_{\tilde{p}}^{2,0,-2}a_0 - 2a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\psi^2H_{\tilde{p}}^{2,0,-2}\rho \\ &\quad + 2a_0\chi(\tilde{p}^{2,0,-2,-1,-3})\chi'(\tilde{p}^{2,0,-2,-1,-3})\chi(\rho)^2\psi^2\tilde{p}^{2,0,-2,-1,-3}\tilde{q} \\ &\quad + 2a_0\chi(\tilde{p}^{2,0,-2,-1,-3})^2\chi(\rho)^2\psi H_{\tilde{p}}^{2,0,-2}\psi. \end{aligned} \quad (475)$$

Observe:

(1) by eq. (471), the first term,

$$\psi^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\chi(\rho)^2H_{\tilde{p}}^{2,0,-2}a_0 = \psi^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\chi(\rho)^2\beta a_0 \quad (476)$$

will have a definite sign (the same as β , negative if $l, \ell < 0$) for appropriate χ, ψ ,

(2) the second term,

$$\begin{aligned} -a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\psi^2H_{\tilde{p}}^{2,0,-2}\rho \\ = -a_0\chi_0(\rho)^2\chi(\tilde{p}^{2,0,-2,-1,-3})^2\psi^2(\beta_0\rho + F_2 + F_3 + x(\sigma^2 + \mathbb{Z}x)^{-1/2}F_4), \end{aligned} \quad (477)$$

also has a definite sign (positive, since $\beta_0, F_2, F_3|_{\mathcal{R}} \leq 0$) for appropriate χ, ψ and is supported in an annulus around \mathcal{R}_+ , which should intersect $\text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \setminus (\mathcal{R} \cup \mathcal{R}_+)$,

(3) the third and fourth terms are supported away from $\text{Char}^{2,0,-2,-1,-3}(\tilde{P})$ and are therefore unproblematic (as that region of phase space is controlled via elliptic estimates or, in the case of $\text{bf}^\circ \cup \text{tf}^\circ$ where the fourth term might have some support, cannot be controlled by symbolic considerations anyways).

We will prove a low order radial point estimate. This means that

$$l < 0, \quad \ell < 0 \quad \varsigma \leq \ell + s - l, \quad (478)$$

so that $\beta > 0$ near \mathcal{R} . Thus, the first and second terms in eq. (475) have the opposite sign near \mathcal{R} . The second term thus contributes to the right-hand side of the radial point estimate, but this term can itself be controlled using a radial point estimate at \mathcal{R}_+ in conjunction with a propagation estimate and will therefore be unproblematic as well.

In order to “regularize,” set $\phi_\varepsilon = (1 + \varepsilon x^{-1})^{-K_1} (1 + \varepsilon \varrho_{\text{df}00}^{-1})^{-K_3}$,

$$a^{(\varepsilon)} = \phi_\varepsilon^2 a \in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, s, \varsigma, l, \ell}(X)), \quad (479)$$

for to-be-decided $K_1, K_3 \in \mathbb{R}$. We then compute that

$$\begin{aligned} H_{\tilde{p}}^{2,0,-2} a^{(\varepsilon)} &= \phi_\varepsilon^2 \left[-\chi(\tilde{p}^{2,0,-2,-1,-3})^2 \chi(\rho)^2 \psi^2 a_0 \left((\varsigma - \ell - s + l) \beta_3 \right. \right. \\ &+ F_1 \left(m - \frac{K_3 \varepsilon \varrho_{\text{df}00}^{-1}}{1 + \varepsilon \varrho_{\text{df}00}^{-1}} \right) + \beta_1 \left(l - \frac{K_1 \varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) + \beta_2 (\ell - 2l) \Big) - 2a_0 \chi_0(\rho)^2 \chi(\tilde{p}^{2,0,-2,-1,-3})^2 \psi^2 H_{\tilde{p}}^{2,0,-2} \rho \\ &\quad + 2a_0 \chi(\tilde{p}^{2,0,-2,-1,-3}) \chi'(\tilde{p}^{2,0,-2,-1,-3}) \chi(\rho)^2 \psi^2 \tilde{p}^{2,0,-2,-1,-3} \tilde{q} \\ &\quad \left. \left. + 2a_0 \chi(\tilde{p}^{2,0,-2,-1,-3})^2 \chi(\rho)^2 \psi H_{\tilde{p}}^{2,0,-2} \psi \right] \right]. \end{aligned} \quad (480)$$

In contrast to the previous radial point estimate, we can make K_1, K_3 arbitrarily large without affecting the sign of the parenthetical term, although we might need to choose ψ with smaller support and replace χ with $\chi \circ \text{dil}_\lambda$ for

$$\lambda = \lambda(m, l, \ell, K_1, K_3) > 0 \quad (481)$$

sufficiently large to ensure that the term proportional to F_1 in eq. (480) does not spoil that sign. So, for some choice of ψ, χ , we can choose $\underline{\delta} = \underline{\delta}(K_1, K_3, m, l, \ell, \chi, \psi) > 0$ sufficiently small such that there exist well-defined uniform families of leC-symbols

$$\begin{aligned} b_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{leC}}^{-\infty, (s-1)/2, (\varsigma-3)/2, (l-1)/2, (\ell-3)/2}(X)), \\ e_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, (s-1)/2, (\varsigma-3)/2, -\infty, -\infty}(X)), \\ f_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, s-1, \varsigma-1, l, \ell}(X)), \\ r_\bullet &\in L^\infty([0, 1]_\varepsilon; S_{\text{cl,leC}}^{-\infty, -\infty, -\infty, l-1, \ell-3}(X)), \end{aligned} \quad (482)$$

such that, in some neighborhood of $\{x = 0\} \subset {}^{\text{leC}}\overline{T}^*X$,

$$\begin{aligned}
b_\varepsilon &= \varrho_{\text{df}}^{-1/2} \varrho_{\text{sf}}^{1/2} \varrho_{\text{ff}}^{3/2} \varrho_{\text{bf}}^{1/2} \varrho_{\text{tf}}^{3/2} a_0^{1/2} \chi(\tilde{p}^{2,0,-2,-1,-3}) \chi(\rho) \phi_\varepsilon \psi \left[F_1 \left(m - \frac{K_3 \varepsilon \varrho_{\text{df}00}^{-1}}{1 + \varepsilon \varrho_{\text{df}00}^{-1}} \right) + \beta_1 \left(l - \frac{K_1 \varepsilon x^{-1}}{1 + \varepsilon x^{-1}} \right) \right. \\
&\quad \left. + \beta_2(\ell - 2l) + \beta_3(\varsigma - \ell - s + l) - 2\delta \phi_\varepsilon^2 \chi(\tilde{p}^{2,0,-2,-1,-3})^2 \chi(\rho)^2 \psi^2 - \varrho_{\text{df}} \varrho_{\text{sf}}^{-1} \varrho_{\text{ff}}^{-3} \varrho_{\text{bf}}^{-1} \varrho_{\text{tf}}^{-3} p_1 \right]^{1/2} \\
e_\varepsilon &= -\varrho_{\text{df}}^{-1/2} \varrho_{\text{sf}}^{1/2} \varrho_{\text{ff}}^{3/2} \varrho_{\text{bf}}^{1/2} \varrho_{\text{tf}}^{3/2} \phi_\varepsilon a_0^{1/2} \chi_0(\rho) \chi(\tilde{p}^{2,0,-2,-1,-3}) \psi \sqrt{-2H_{\tilde{p}}^{2,0,-2} \rho} \\
f_\varepsilon &= 2\varrho_{\text{df}} \varrho_{\text{sf}} \varrho_{\text{ff}} a_0 \phi_\varepsilon^2 \chi(\tilde{p}^{2,0,-2,-1,-3}) \chi'(\tilde{p}^{2,0,-2,-1,-3}) \chi(\rho)^2 \psi^2 \tilde{q} \\
r_\varepsilon &= 2\varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^3 \varrho_{\text{ff}}^3 \varrho_{\text{bf}}^3 \varrho_{\text{tf}}^3 a_0 \chi(\tilde{p}^{2,0,-2,-1,-3})^2 \chi(\rho)^2 \psi H_{\tilde{p}}^{2,0,-2} \psi.
\end{aligned} \tag{483}$$

In terms of these new symbols, we can write

$$H_{\tilde{p}} a^{(\varepsilon)} + p_1 a^{(\varepsilon)} = -2\delta \varrho_{\text{df}}^{-1} \varrho_{\text{sf}}^3 \varrho_{\text{ff}}^3 \varrho_{\text{bf}}^3 \varrho_{\text{tf}}^3 a_0^{-1} a^{(\varepsilon)2} - b_\varepsilon^2 + e_\varepsilon^2 + f_\varepsilon \tilde{p} + r_\varepsilon. \tag{484}$$

We now apply Op. Setting $A_\varepsilon = (1/2)(\text{Op}(a^{(\varepsilon)}) + \text{Op}(a^{(\varepsilon)})^*)$, $B_\varepsilon = \text{Op}(b_\varepsilon)$, $E_\varepsilon = \text{Op}(e_\varepsilon)$, $F_\varepsilon = \text{Op}(f_\varepsilon)$, we have

$$\begin{aligned}
A_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, s, \varsigma, l, \ell}(X)), \\
B_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, (s-1)/2, (\varsigma-3)/2, (l-1)/2, (\ell-3)/2}(X)), \\
E_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, (s-1)/2, (\varsigma-3)/2, -\infty, -\infty}(X)), \\
F_\bullet &\in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, s-1, \varsigma-1, l, \ell}(X)),
\end{aligned} \tag{485}$$

and

$$\begin{aligned}
-i[\Re \tilde{P}, A_\varepsilon] - \{\Im \tilde{P}, A_\varepsilon\} &= -2\delta A_\varepsilon \Lambda_{1/2, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2} A_\varepsilon - B_\varepsilon^* B_\varepsilon + E_\varepsilon^* E_\varepsilon \\
&\quad + F_\varepsilon^* \tilde{P} + R_\varepsilon
\end{aligned} \tag{486}$$

for some $R_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, s-2, \varsigma-4, l-1, \ell-3}(X))$. Moreover, we necessarily have

$$\begin{aligned}
&\text{WF}'_{L^\infty, \text{leC}}(A_\bullet), \text{WF}'_{L^\infty, \text{leC}}(B_\bullet), \text{WF}'_{L^\infty, \text{leC}}(F_\bullet), \\
&\text{WF}'_{L^\infty, \text{leC}}(E_\bullet), \text{WF}'_{L^\infty, \text{leC}}(R_\bullet) \subset \text{supp } \chi(\tilde{p}^{2,0,-2,-1,-3}) \chi(\rho) \psi,
\end{aligned} \tag{487}$$

where the last of these inclusions (the one for R_\bullet) follows from the one for A_\bullet and $\text{WF}'_{L^\infty, \text{leC}}([\tilde{P}, A_\bullet]) \subset \text{WF}'_{L^\infty, \text{leC}}(A_\bullet)$.

For each $m_0, s_0, \varsigma_0, l_0, \ell_0 \in \mathbb{R}$, there exist some $K_{1,0}, K_{3,0} > 0$ (dependent on $m_0, s_0, \varsigma_0, l_0, \ell_0$ and m, s, ς, l, ℓ) such that, given $\{u(-; \sigma)\}_{\sigma>0} \subset \mathcal{S}'(X)$ with $u(-; 0) \in H_{\text{scb}}^{m, \varsigma, \ell}(X_{1/2})$ and $u(-; \sigma) \in H_{\text{scb}}^{m, s, l}(X)$ for all $\sigma > 0$, if we take $K_1 > K_{1,0}, K_3 > K_{3,0}$ in the construction above then it is the case that (for any $\varepsilon > 0$, and for each $\sigma > 0$, implicit in the notation),

$$2\Im \langle \tilde{P}u, A_\varepsilon u \rangle_{L^2} = -\langle \{\Im \tilde{P}, A_\varepsilon\}u, u \rangle_{L^2} + i\langle [\Re \tilde{P}(\sigma), A_\varepsilon]u, u \rangle_{L^2}, \tag{488}$$

where the pairings above are well-defined distributional pairings (with the left argument of each inner product in the dual Sobolev space to a Sobolev space in which the right argument lies). Applying eq. (486) to $\{u(-; \sigma)\}_{\sigma>0}$ as above and pairing against u (and taking K_1, K_3 large enough), we have

$$\begin{aligned}
2\Im \langle \tilde{P}u, A_\varepsilon u \rangle_{L^2} &= -\|B_\varepsilon u\|_{L^2}^2 + \|E_\varepsilon u\|_{L^2}^2 + \langle \tilde{P}u, F_\varepsilon u \rangle_{L^2} + \langle R_\varepsilon u, u \rangle_{L^2} \\
&\quad - 2\delta \|x \Lambda_{1/2, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2} A_\varepsilon u\|_{L^2}^2.
\end{aligned} \tag{489}$$

Equation (489) implies

$$\begin{aligned} \|B_\varepsilon u\|_{L^2}^2 + 2\bar{\delta}\|\Lambda_{1/2, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2} A_\varepsilon u\|_{L^2}^2 \\ \leq 2|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| + |\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| + |\langle R_\varepsilon u, u \rangle_{L^2}| + \|E_\varepsilon u\|_{L^2}^2. \end{aligned} \quad (490)$$

Suppose $G \in \Psi_{\text{leC}}^{-\infty, 0, 0, 0}(X)$ is self-adjoint and such that

$$\text{WF}_{\text{leC}}^{0, 0}(1 - G) \cap \text{supp } \chi(\tilde{p}^{2, 0, -2, -1, -3})\chi(\rho)\psi = \emptyset. \quad (491)$$

We now estimate the terms in eq. (490) as in the propagation estimate, except we now must keep track of orders at bf, tf:

- Writing $\tilde{P} = (1 - G)\tilde{P} + G\tilde{P}$ we have, for each $N \in \mathbb{N}$ (and K_1, K_3 large enough),

$$|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| \leq |\langle G\tilde{P}u, A_\varepsilon u \rangle_{L^2}| + |\langle \tilde{P}u, (1 - G)A_\varepsilon u \rangle_{L^2}|. \quad (492)$$

By Lemma 2.15, for any $\bar{\delta} > 0$,

$$\begin{aligned} |\langle G\tilde{P}u, A_\varepsilon u \rangle_{L^2}| &\preceq \bar{\delta}^{-1} \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \bar{\delta} \|A_\varepsilon u\|_{\mathcal{Y}_N^*}^2 + \bar{\delta} \|u\|_{\bar{\mathcal{E}}_N}^2 \\ &\preceq \bar{\delta}^{-1} \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \bar{\delta} \|A_\varepsilon u\|_{\mathcal{Y}_{*, N}}^2 \\ &\preceq \bar{\delta}^{-1} \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \bar{\delta} \|\Lambda_{1/2, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2} A_\varepsilon u\|_{L^2}^2 + \bar{\delta} \|u\|_{\bar{\mathcal{E}}_N}^2 \end{aligned} \quad (493)$$

(where the constant does not depend on $\bar{\delta}$) for

$$\begin{aligned} \mathcal{Y}_N &= H_{\text{leC}}^{-N, (s+1)/2, (\varsigma+3)/2, (l+1)/2, (\ell+3)/2}(X), \\ \mathcal{Y}_{*, N} &= H_{\text{leC}}^{-N, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2}(X), \\ \bar{\mathcal{E}}_N &= H_{\text{leC}}^{-N, -N, -N, -N, -N}(X). \end{aligned} \quad (494)$$

On the other hand, noting that $(1 - G)A_\bullet \in L^\infty([0, 1]_\varepsilon; \Psi_{\text{leC}}^{-\infty, -\infty, -\infty, l/2, \ell/2}(X))$,

$$\begin{aligned} |\langle \tilde{P}u, (1 - G)A_\varepsilon u \rangle_{L^2}| &\preceq \|\tilde{P}u\|_{H_{\text{leC}}^{-N, -N, -N, \frac{l+1}{2}, \frac{\ell+3}{2}}}^2 + \|(1 - G)A_\varepsilon u\|_{H_{\text{leC}}^{N, N, N, -\frac{l+1}{2}, -\frac{\ell+3}{2}}}^2 \\ &\preceq \|\tilde{P}u\|_{H_{\text{leC}}^{-N, -N, -N, \frac{l+1}{2}, \frac{\ell+3}{2}}}^2 + \|(1 - G)A_\varepsilon u\|_{H_{\text{leC}}^{-N_0, -N_0, -N_0, -\frac{l+1}{2}, -\frac{\ell+3}{2}}}^2 + \|u\|_{\bar{\mathcal{E}}_{N_0}}^2 \\ &\preceq \|u\|_{H_{\text{leC}}^{-N, -N, -N, \frac{l-1}{2}, \frac{\ell-3}{2}}}^2 \end{aligned} \quad (495)$$

for N_0 sufficiently large (relative to N_0). So, $|\langle \tilde{P}u, (1 - G)A_\varepsilon u \rangle_{L^2}| \preceq \|u\|_{\bar{\mathcal{E}}_N}^2$ for $\mathcal{E}_N = H_{\text{leC}}^{-N, -N, -N, \frac{l-1}{2}, \frac{\ell-3}{2}}(X)$.

Thus, taking $\bar{\delta}$ sufficiently small (relative to $\bar{\delta}$),

$$2|\langle \tilde{P}u, A_\varepsilon u \rangle_{L^2}| - 2\bar{\delta}\|\Lambda_{1/2, -(s+1)/2, -(\varsigma+3)/2, -(l+1)/2, -(\ell+3)/2} A_\varepsilon u\|_{L^2}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|u\|_{\bar{\mathcal{E}}_N}^2. \quad (496)$$

- Similarly, $|\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|F_\varepsilon u\|_{\mathcal{Y}_{*, N_0}}^2 + \|u\|_{\bar{\mathcal{E}}_N}^2$, for any $N_0 \in \mathbb{N}$.

We have $\text{WF}_{L^\infty, \text{leC}}(F_\bullet) \cap \text{Char}_{\text{leC}}^{2, 0, -2, -1, -3}(\tilde{P}) = \emptyset$, so we deduce from elliptic regularity that

$$\|F_\varepsilon u\|_{\mathcal{Y}_{*, N_0}}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|u\|_{\bar{\mathcal{E}}_N}^2 \quad (497)$$

for N_0 sufficiently large (relative to N_0). Thus,

$$|\langle \tilde{P}u, F_\varepsilon u \rangle_{L^2}| \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|u\|_{\bar{\mathcal{E}}_N}^2. \quad (498)$$

- $|\langle R_\varepsilon u, u \rangle_{L^2}| \preceq \|Gu\|_{\mathcal{Z}_N}^2 + \|u\|_{\bar{\mathcal{E}}_N}^2$ for

$$\mathcal{Z}_N = H_{\text{leC}}^{-N, (s-2)/2, (\varsigma-4)/2, (l-1)/2, (\ell-3)/2}(X). \quad (499)$$

Thus,

$$\|B_\varepsilon u\|_{L^2}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|E_\varepsilon u\|_{L^2}^2 + \|u\|_{\mathcal{E}_N}^2, \quad (500)$$

I.e., letting $\mathcal{X}_N = H_{\text{leC}}^{-N, (s-1)/2, (\varsigma-3)/2, (l-1)/2, (\ell-3)/2}(X)$ and $\bar{\mathcal{X}}_N = H_{\text{leC}}^{-N, (s-1)/2, (\varsigma-3)/2, -N, -N}(X)$,

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N}^2 \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N}^2 + \|Gu\|_{\mathcal{Z}_N}^2 + \|\hat{E}_\varepsilon u\|_{\bar{\mathcal{X}}_N}^2 + \|u\|_{\mathcal{E}_N}^2, \quad (501)$$

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N} + \|Gu\|_{\mathcal{Z}_N} + \|\hat{E}_\varepsilon u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}, \quad (502)$$

where

$$\hat{B}_\varepsilon = \Lambda_{0, -(s-1)/2, -(\varsigma-3)/2, -(l-1)/2, -(\ell-3)/2} B_\varepsilon \quad (503)$$

$$\hat{E}_\varepsilon = \Lambda_{0, -(s-1)/2, -(\varsigma-3)/2, 0, 0} E_\varepsilon, \quad (504)$$

so that $\hat{B}_\bullet, \hat{E}_\bullet \in \Psi_{\text{leC}}^{0,0,0,0}(X)$.

Proposition 5.11. *Suppose that $G_0, G_1 \in \Psi_{\text{leC}}^{-\infty, 0, 0, 0}(X)$, $G_2 \in \Psi_{\text{leC}}^{-\infty, 0, 0, -\infty, -\infty}$ satisfy*

(1) $\text{WF}'_{\text{leC}}(G_2) \cap \mathcal{R}_+ = \emptyset$ and

$$\text{WF}_{\text{leC}}^{0,0}(G_1) \cap \mathcal{R}_+ = \emptyset, \quad (505)$$

(2) $\mathcal{R}, \text{WF}'_{\text{leC}}(G_1) \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_0)$,

(3) *there exist $\Theta_1, \Theta_2 \in (0, \pi)$ with $\Theta_1 < \Theta_2$ such that $\text{Ell}_{\text{leC}}^{0,0,0,0}(G_2) \supseteq \mathcal{P}[\Theta_1, \Theta_2]$ and*

$$\text{Ell}_{\text{leC}}^{0,0,0,0}(G_0) \supset \mathcal{R} \cup \bigcup_{\Theta_3 \in (\Theta_2, \pi)} \mathcal{P}[\Theta_1, \Theta_3] \quad (506)$$

Then, for any $\Sigma > 0$, $N \in \mathbb{N}$, and $m, s, \varsigma, l, \ell \in \mathbb{R}$ with $l < -1/2$, $\ell < -3/2$, and $\varsigma \leq \ell + s - l$, there exists a constant

$$C = C(\tilde{P}, G_0, G_1, G_2, G_3, \Sigma, N, m, s, \varsigma, l, \ell) > 0 \quad (507)$$

such that

$$\|G_1 u\|_{H_{\text{leC}}^{m, s, \varsigma, l, \ell}} \leq C \left[\|G_0 \tilde{P}u\|_{H_{\text{leC}}^{-N, s+1, \varsigma+3, l+1, \ell+3}} + \|G_2 u\|_{H_{\text{leC}}^{-N, s, \varsigma, -N, -N}} + \|u\|_{H_{\text{leC}}^{-N, -N, -N, l, \ell}} \right] \quad (508)$$

holds for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. ■

Proof. We can assume without loss of generality that

$$\mathcal{R} \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_1) \quad (509)$$

and

$$\text{Ell}_{\text{leC}}^{0,0,0,0}(G_0) \supseteq \text{WF}'_{\text{leC}}(G_2). \quad (510)$$

Let $s_0 = 2s + 1$, $l_0 = 2l + 1$, and $\ell_0 = 2\ell + 3$, and $\varsigma_0 = 2\varsigma + 3$. The condition $l < -1/2$ is equivalent to $l_0 < 0$, $\ell < -3/2$ is equivalent to $\ell_0 < 0$, and $\varsigma \leq \ell + s - l$ is equivalent to $\varsigma_0 \leq \ell_0 + s_0 - l_0$. We can therefore apply eq. (502) with $s_0, \varsigma_0, l_0, \ell_0$ in place of what we called “ s, ς, l, ℓ ” there. Thus, for each $\varepsilon > 0$,

$$\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G\tilde{P}u\|_{\mathcal{Y}_N} + \|Gu\|_{\mathcal{Z}_N} + \|\hat{E}_\varepsilon u\|_{\bar{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}, \quad (511)$$

where $G, \hat{B}_\bullet, \hat{E}_\bullet$ are as above and now

$$\mathcal{X}_N = H_{\text{leC}}^{-N, s, \varsigma, l, \ell}(X), \quad (512)$$

$$\mathcal{Z}_N = H_{\text{leC}}^{-N, s-1/2, \varsigma-1/2, l, \ell}(X), \quad \bar{\mathcal{X}}_N = H_{\text{leC}}^{-N, s, \varsigma, -N, -N}(X), \quad \mathcal{Z}_N = H_{\text{leC}}^{-N, s-1/2, \varsigma-1/2, l, \ell}(X),$$

$$\mathcal{Y}_N = H_{\text{leC}}^{-N, s+1, \varsigma+3, l+1, \ell+3}(X), \quad (513)$$

and $\mathcal{E}_N = H_{\text{leC}}^{-N, -N, -N, l, \ell}(X)$. If necessary, we can retroactively replace χ, ψ such that

$$\text{WF}'_{L^\infty, \text{leC}}(E_\bullet), \text{WF}_{L^\infty, \text{leC}}^{0,0}(B_\bullet) \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_0). \quad (514)$$

We can apply the propagation estimate Proposition 5.7 to bound

$$\|\hat{E}_\varepsilon u\|_{\tilde{\mathcal{X}}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (515)$$

By eq. (509), we can now retroactively choose G (for appropriate χ, ψ) such that $\text{WF}_{\text{leC}}^{0,0}(G) \subset \text{Ell}_{\text{leC}}^{0,0,0,0}(G_0) \cap \text{Ell}_{\text{leC}}^{0,0,0,0}(G_1)$, so that

$$\|G \tilde{P}u\|_{\mathcal{Y}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|u\|_{\mathcal{E}_N} \quad (516)$$

and $\|Gu\|_{\mathcal{Z}_N} \preceq \|G_1 u\|_{\mathcal{Z}_N} + \|u\|_{\mathcal{E}_N}$.

The estimate eq. (511) therefore implies $\|\hat{B}_\varepsilon u\|_{\mathcal{X}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_1 u\|_{\mathcal{Z}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}$. Via the compactness argument utilized in the proof of the propagation estimate and previous radial point estimate, $\|\hat{B}_0 u(-; \sigma)\|_{\mathcal{X}_N} \leq \liminf_{\varepsilon \rightarrow 0^+} \|\hat{B}_\varepsilon u(-; \sigma)\|_{\mathcal{X}_N}$ for each $\sigma \geq 0$, so

$$\|\hat{B}_0 u\|_{\mathcal{X}} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_1 u\|_{\mathcal{Z}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (517)$$

Unlike \hat{B}_ε for $\varepsilon > 0$, $\text{Ell}_{\text{leC}}^{0,0,0,0}(\hat{B}_0) \supset \mathcal{R}$. Thus, (for χ, ψ with sufficiently small support) $\|\hat{B}_0 u\|_{\mathcal{X}_N} \preceq \|G_1 u\|_{\mathcal{X}_N} + \|u\|_{\mathcal{E}_N}$. Substituting this into the estimate above, we get

$$\|G_1 u\|_{\mathcal{X}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_1 u\|_{\mathcal{Z}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (518)$$

Since the leC- Sobolev spaces $\mathcal{X}_N, \mathcal{Y}, \mathcal{Z}_N$ in eq. (518) get weaker as s, ς, l, ℓ decrease, we can inductively use the family of estimates eq. (518) to bound the $\|G_1 u\|_{\mathcal{Z}_N}$ term on the right-hand side: for all $N \in \mathbb{N}$,

$$\|G_1 u\|_{\mathcal{X}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_1 u\|_{\mathcal{E}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N} \quad (519)$$

$$\|G_1 u\|_{\mathcal{X}_N} \preceq \|G_0 \tilde{P}u\|_{\mathcal{Y}_N} + \|G_2 u\|_{\tilde{\mathcal{X}}_N} + \|u\|_{\mathcal{E}_N}. \quad (520)$$

This implies eq. (508). \square

5.4. Upshot. Combining the estimates above (e.g. using Lemma 2.16), we get:

Theorem 5.12. *Suppose that $G_1, G_2, G_3 \in \Psi_{\text{leC}}^{0,0,0,0}(X)$ satisfy*

- $\text{WF}_{\text{leC}}^{0,0}(G_1), \text{Char}_{\text{leC}}^{2,0,-2,-1,-3}(\tilde{P}) \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_2)$,
- $\mathcal{R}_+ \subseteq \text{Ell}_{\text{leC}}^{0,0,0,0}(G_3)$.

For every $m, s, \varsigma, l, \ell, s_0, \varsigma_0 \in \mathbb{R}$ satisfying

$$l < -1/2, \quad \ell < -3/2, \quad -1/2 < s_0 < s, \quad -3/2 < \varsigma_0 < \varsigma \leq \ell + s - l, \quad (521)$$

there exists, for each $\Sigma > 0$ and $N \in \mathbb{N}$, a constant $C = C(\tilde{P}, G_1, G_2, G_3, \Sigma, N, m, s, \varsigma, l, \ell, s_0, \varsigma_0) > 0$ such that

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C \left[\|G_2 \tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|G_3 u\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}} \right] \quad (522)$$

for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$ (in the usual strong sense that if the right-hand side is finite, then the left-hand side as well, and the stated inequality holds). \blacksquare

It will also be useful to have the following refinement of Theorem 5.12:

Proposition 5.13. *Given $G_1, G_2, m, s, \varsigma, l, \ell, s_0, \varsigma_0$ as in the setup of Theorem 5.12, if in addition $\text{Ell}_{\text{leC}}^{0,0,0,0}(G_1) \supseteq \mathcal{R}_+$, then there exists a constant*

$$c = c(\tilde{P}, G_1, G_2, \Sigma, N, m, s, \varsigma, l, \ell, s_0) > 0 \quad (523)$$

such that, for any $u \in \mathcal{S}'(X)$,

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq c \left[\|G_2 \tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}} \right] \quad (524)$$

holds for all $\sigma \in [0, \Sigma]$ such that $\|G_1(\sigma)u\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}(X)(\sigma)} < \infty$. \blacksquare

Proof. It suffices to consider the case $-N < m, l, \ell$.

Applying Theorem 5.12 with $G_3 = G_1$,

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C' (\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|G_1 u\|_{H_{\text{leC}}^{m-2,s_0,\varsigma_0,-N,-N}} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}}) \quad (525)$$

for some constant $C' > 0$. We now use the interpolation inequality Lemma 2.17, which, for each $\epsilon > 0$, allows us to bound

$$\|G_1 u\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}} \leq \epsilon \|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} + C''(\epsilon) \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}} \quad (526)$$

for some $C''(\epsilon) = C(\tilde{P}, G_1, G_2, \Sigma, N, m, s, \varsigma, l, \ell, s_0, \varsigma_0, \epsilon) > 0$. Taking $\epsilon < 1/2C'$, plugging eq. (526) into eq. (525) yields

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C' (\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + (1 + C'') \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}}) + (1/2) \|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \quad (527)$$

If $\sigma \in [0, \Sigma]$ is such that

$$\|G_1(\sigma) u(-; \sigma)\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(\sigma)} < \infty, \quad (528)$$

then we can subtract the last term on the right-hand side from both sides, getting eq. (524),

$$\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq 2C' (1 + C'') (\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}}). \quad (529)$$

If $\|G_1(\sigma) u(-; \sigma)\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}(X)(\sigma)} < \infty$, then we break into two cases:

- if one of $\|G_2 \tilde{P} u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}}$, $\|u\|_{H_{\text{b,leC}}^{-N,l,\ell}}$ is infinite, then eq. (524) holds trivially;
- if both are finite, then Theorem 5.12 implies that $\|G_1 u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} < \infty$, and therefore eq. (529) holds.

We conclude that eq. (524) holds if we take $C = 2C'(1 + C'')$. \square

Taking $G_1 = G_2 = 1$ in Proposition 5.13, we get the main claim eq. (314) at the beginning of this section:

Proposition 5.14. *For every $\Sigma > 0$, $N \in \mathbb{N}$, and $m, s, \varsigma, l, \ell, s_0, \varsigma_0 \in \mathbb{R}$ satisfying $l < -1/2 < s_0 < s$ and $\ell < -3/2 < \varsigma_0 < \varsigma \leq \ell + s - l$, there exists a constant $C = C(\tilde{P}, \Sigma, N, m, s, \varsigma, l, \ell) > 0$ such that*

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C [\|\tilde{P} u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{\text{b,leC}}^{-N,l,\ell}}] \quad (530)$$

holds for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$ such that $\|u\|_{H_{\text{leC}}^{-N,s_0,\varsigma_0,-N,-N}(X)(\sigma)} < \infty$. \blacksquare

6. PROOF OF MAIN THEOREM

Most of the results in this section are split among three subsections:

- (1) §6.1, where estimates regarding the “leC-normal operator” $N(\tilde{P})$ (see §3) are proven,
- (2) §6.2, where – using the estimates from the previous subsection – the conormality of the output of the conjugated resolvent family on $X_{\text{res}}^{\text{sp}}$ is established, along with smoothness (in terms of $E = \sigma^2$) at zf° (with the terms in the Taylor series being conormal distributions on zf) — see Proposition 6.15 — and
- (3) §6.3, where the proposition needed to upgrade the conormality established in the previous subsection to the existence of asymptotic expansions is proven.

Central to this section is the analysis of the model problem

$$2i \left[\hat{x} \partial_{\hat{x}} + \frac{\hat{x}}{1 + \hat{x}} \left(k + \frac{1}{4} \right) + l + \frac{1}{2} \right] u = f, \quad (531)$$

$u, f \in \mathcal{D}'(\mathbb{R}_{\hat{x}}^+)$, where $l, k \in \mathbb{R}$ will be required to satisfy some threshold inequalities. We record the exact solution to eq. (531): for any $l, k \in \mathbb{R}$, if $f \in C^\infty(\mathbb{R}^+)$, then $u \in \mathcal{D}'(\mathbb{R}^+)$ satisfies eq. (531) if and only if

$$u(\hat{x}) = \hat{x}^{-l-1/2}(1+\hat{x})^{-k-1/4} \left[c + \frac{i}{2} \int_{\hat{x}}^1 \hat{x}_0^{l-1/2}(1+\hat{x}_0)^{k+1/4} f(\hat{x}_0) d\hat{x}_0 \right] \quad (532)$$

for some $c \in \mathbb{C}$. The model problem above arises rewriting

$$\tilde{N}(\tilde{P}) = x^{-l-n/2}(\sigma^2 + Zx)^{-k} N(\tilde{P}) x^{l+n/2}(\sigma^2 + Zx)^k \quad (533)$$

in terms of $\hat{x} = Zx/\sigma^2$. Indeed,

$$\begin{aligned} x^{-1}(\sigma^2 + Zx)^{-1/2} \tilde{N}(\tilde{P}) &= 2i \left(x \partial_x + l + \frac{1}{2} \right) + \frac{2iZx}{\sigma^2 + Zx} \left(k + \frac{1}{4} \right) \\ &= 2i \left[\hat{x} \partial_{\hat{x}} + \frac{\hat{x}}{1+\hat{x}} \left(k + \frac{1}{4} \right) + l + \frac{1}{2} \right]. \end{aligned} \quad (534)$$

If $N(\tilde{P})u = f$ then $\tilde{N}(\tilde{P})u_0 = f_0$ for $u_0 = x^{-l-n/2}(\sigma^2 + Zx)^{-k}u$ and $f_0 = x^{-l-n/2}(\sigma^2 + Zx)^{-k}f$.

We now spell out the deduction of Theorem 1.2 (in the form of the more general Proposition 6.3) from the results in §6.2, §6.3.

Proposition 6.1. *Suppose that $g = g_0 + xg_1 + x^{\alpha_1}g_2$ for $g_1 \in C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^*X)$ and $g_2 \in S^0(X; {}^{\text{sc}}\text{Sym}^2 T^*X)$ for $\alpha_1 > 1$. Then the Laplace-Beltrami operator Δ_g has the form*

$$\Delta_g = \Delta_{g_0} + x \text{Diff}_{\text{scb}}^{2,0,-2}(X) + x^{\alpha_1} S \text{Diff}_{\text{scb}}^{2,0,-2}(X), \quad (535)$$

where g_0 is an exactly conic metric. ■

Proof. Let $h \mapsto [h]$ denote the natural bundle monomorphism ${}^{\text{sc}}\text{Sym}^2 T^*X \hookrightarrow \text{End}({}^{\text{sc}}TX, {}^{\text{sc}}T^*X)$.

The matrix identity $[g]^{-1} = [g_0]^{-1} - [g]^{-1}(x[g_1] + x^{\alpha_1}[g_2])[g_0]^{-1}$, applied inductively, yields

$$[g]^{-1} - [g_0]^{-1} = [g_0]^{-1} \sum_{k=1}^K (-1)^k ((x[g_1] + x^{\alpha_1}[g_2])[g_0]^{-1})^k + (-1)^{K+1} [g]^{-1} ((x[g_1] + x^{\alpha_1}[g_2])[g_0]^{-1})^{K+1} \quad (536)$$

for each $K \in \mathbb{N}$. Noting that $g^{-1} \in S^0(X; {}^{\text{sc}}\text{Sym}^2 TX)$, taking $K > \alpha_1$ leads to $g^{-1} - g_0^{-1} \in xC^\infty(X; {}^{\text{sc}}\text{Sym}^2 TX) + x^{\alpha_1}S^0(X; {}^{\text{sc}}\text{Sym}^2 TX)$. Also, from $[g] = [g_0](1 + x[g_0]^{-1}[g_1] + x^{\alpha_1}[g_0]^{-1}[g_2])$, $\det g \in S^0(X; {}^{\text{sc}}\Omega X)$ satisfies

$$\det g = (\det g_0)(1 + x \text{tr}([g_0]^{-1}[g_1]) + x^2 C^\infty(X) + x^{\alpha_1} S^0(X)). \quad (537)$$

Writing

$$\Delta_g = - \sum_{i,j=1}^n (g^{ij} \partial_i \partial_j + \partial_i g^{ij} \partial_j + (1/2)(\det g)^{-1} (\partial_i \det g) g^{ij} \partial_j) \quad (538)$$

in local coordinates, we conclude that eq. (535) holds locally, which suffices to show that the decomposition eq. (535) can be done globally. □

Thus:

Proposition 6.2. *Suppose that $g = g_0 + xg_1 + x^{\alpha_1}g_2$ for $g_1 \in C^\infty(X; {}^{\text{sc}}\text{Sym}^2 T^*X)$ and $g_2 \in S^0(X; {}^{\text{sc}}\text{Sym}^2 T^*X)$ for $\alpha_1 > 1$. If $P(\sigma) = \Delta_g - \sigma^2 - Zx + V$ is the spectral family of an attractive Coulomb-like Schrödinger operator with $V \in x^2 C^\infty(X) + x^{\alpha_2} S^0(X)$ for $\alpha_2 > 3/2$, then we can decompose*

$$P(\sigma) = P_0(\sigma) + P_1 + P_2 \quad (539)$$

for P_0, P_1, P_2 having the form specified in the introduction (with $a = 0$), except that

$$a_{00} = -x^4(g_1)_{00}|_{\partial X} = - \lim_{x \rightarrow 0^+} x^4 g_1(\partial_x, \partial_x) \quad (540)$$

is not necessarily constant (and the attractivity condition eq. (54) might only be satisfied for small σ). We can arrange that P_1 is fully classical and that P_2 is classical to order (β_2, β_3) with $\beta_2 = \alpha_1 - 1$ and $\beta_3 = \alpha_2 - 3/2$. \blacksquare

Proof. Let $V_0 \in C^\infty(X)$, $V_1 \in S^0(X)$ satisfy $V = x^2 V_0 + x^{\alpha_1} V_1$. It suffices to restrict attention to the boundary-collar. There, we define P_0 by

$$P_0(\sigma) = \Delta_{g_0} - x a_{00} (x^2 \partial_x)^2 - \sigma^2 - Zx. \quad (541)$$

This has the form specified in the introduction, in the sense that eq. (53) holds with $a = 0$.

By the proof of Proposition 6.1, there exists some $P_{2,0} \in \text{Diff}_{\text{scb}}^{2,-2,-3}(X) + x^{\alpha_1} S \text{Diff}_{\text{scb}}^{2,0,-2}(X)$ such that, in any local coordinate patch,

$$\Delta_g = \Delta_{g_0} + \sum_{i,j,k} x g_0^{ik} (g_1)_{kl} g_0^{\ell j} \partial_i \partial_j + P_{2,0}. \quad (542)$$

Now set $P_1 = \sum_{i,j,k} x g_0^{ik} (g_1)_{kl} g_0^{\ell j} \partial_i \partial_j - x g_0^{0k} (g_1)_{kl} g_0^{\ell 0} \partial_x^2 = \sum_{i,j,k} x g_0^{ik} (g_1)_{kl} g_0^{\ell j} \partial_i \partial_j - x^9 (g_1)_{00} \partial_x^2$. Then, P_1 has the form specified in eq. (55) and is even fully classical.

Defining a_{00} by eq. (540), $(g_1)_{00} + x^{-4} a_{00} \in x^{-3} C^\infty(X)$, so the operator $P_{2,1} = x^9 (g_1)_{00} \partial_x^2 + x a_{00} (x^2 \partial_x)^2$ is in $\text{Diff}_{\text{scb}}^{2,-2,-4}(X)$. We therefore set

$$P_2 = P_{2,0} + P_{2,1} + x^2 V_0 + x^{\alpha_2} V_1 \in \text{Diff}_{\text{sc}}^{2,-2}(X) + S \text{Diff}_{\text{scb}}^{2,-\alpha_1,-2-\alpha_1}(X) + x^{\alpha_2} S^0(X). \quad (543)$$

Then eq. (60) applies, for $\beta_2 = \alpha_1 - 1$ and $\beta_3 = \alpha_2 - 3/2$. \square

Proposition 6.3. *Given an asymptotically conic manifold (X, ι, g) of dimension $\dim X = n \geq 2$ such that g satisfies eq. (50), with $a_{00} \in \mathbb{R}$, and given $Z > 0$ and $V \in x^2 C^\infty(X) + x^{\alpha_2} S^0(X)$ for some $\alpha_2 > 3/2$, consider the Schrödinger operator*

$$P = \Delta_g - Zx + V : \mathcal{S}'(X) \rightarrow \mathcal{S}'(X). \quad (544)$$

Set

$$\Phi(x; \sigma) = \frac{1}{x} \sqrt{\sigma^2 + Zx - \sigma^2 a_{00} x} + \frac{1}{\sigma} (Z - \sigma^2 a_{00}) \operatorname{arcsinh} \left(\frac{\sigma}{x^{1/2}} \frac{1}{(Z - \sigma^2 a_{00})^{1/2}} \right) \quad (545)$$

for all $\sigma > 0$ such that $Z > \sigma^2 a_{00}$. Suppose that g is classical to α_1 th order, $\alpha_1 > 1$. Set $\delta_1 = \min\{\alpha_1 - 1, \alpha_2 - 1\}$ and $\delta_0 = \min\{\alpha_1 - 1, \alpha_2 - 3/2\}$.

Then, for any $f \in \mathcal{S}(X)$:

(I) there exists some

$$u_{0,\pm} \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > E a_{00}\}) + \mathcal{A}_{\text{loc}}^{((0,0),\delta_1),2\delta_0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > E a_{00}\}) \quad (546)$$

such that, for $E > 0$ satisfying $Z > E a_{00}$, $u_{\pm}(-; E^{1/2}) = R(E \pm i0) f \in \mathcal{S}'(X)$ can be written as

$$u_{\pm} = e^{\pm i \Phi(x; E^{1/2})} x^{(n-1)/2} (E + Zx)^{-1/4} u_{0,\pm}, \quad (547)$$

(II) $u_{\pm}(-; 0) = R(E = 0; Z \pm i0) f$ as $u_{\pm}(-; 0) = e^{\pm i \Phi(x; 0)} x^{(n-1)/2} (Zx)^{-1/4} u_{0,\pm}(-; 0)$, where $u_{0,\pm}(-; 0) \in C^\infty(X_{1/2}) + \mathcal{A}_{\text{loc}}^{2\delta_0-}(X)$ is the restriction of $u_{0,\pm}$ to $\text{zf} = \text{cl}\{\sigma = 0, x > 0\} \subset X_{\text{res}}^{\text{sp}}$.

Moreover, the map

$$\mathcal{S}(X) \ni f \mapsto u_{0,\pm} \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > E a_{00}\}) + \mathcal{A}_{\text{loc}}^{((0,0),\delta_1),2\delta_0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > E a_{00}\}) \quad (548)$$

is continuous. \blacksquare

Proof. It suffices to prove only the ‘+’ case of the proposition, since the ‘−’ case is similar (and follows via complex conjugation). Moreover, it suffices to construct $u_{0,\pm} = u_{0,\pm,E_0}$ over every interval of the form $[0, E_0]$ for $E_0 > 0$ satisfying $Z > E_0 a_{00}$, since then for $E \leq E_0$ the function $u_{0,\pm,E_0}(-; E^{1/2})$ does not depend on E_0 in the sense that for any E_0, E'_0 satisfying $Z > E_0 a_{00}, E'_0 a_{00}$ it is the case that u_{0,\pm,E_0} and u_{0,\pm,E'_0} agree on $[0, \min\{E_0, E'_0\}]_E$.

Given the setup of the proposition, Proposition 6.2 applies, so we can define a family $P = \{P(\sigma)\}_{\sigma \geq 0}$ satisfying the assumptions listed in §1 such that $P(\sigma) = P(0) - \sigma^2$ for $\sigma^2 \leq E_0$, with P_1 fully classical and P_2 classical to order (β_2, β_3) for $\beta_2 = \alpha_1 - 1$ and $\beta_3 = \alpha_2 - 3/2$.

Defining $\tilde{P} = \exp(-i\Phi)P\exp(+i\Phi)$, since $Pu_+ = f$, $\tilde{P}u_{0,+} = \tilde{f}$ for $u_{0,+} = \exp(-i\Phi)u_+ \in \mathcal{S}'(X)$ and $\tilde{f} = \exp(-i\Phi)f \in \mathcal{A}^{\infty,\infty,(0,0)}(X_{\text{res}}^{\text{sp}})$. Referring §6.2 for the definition of $\tilde{R}_+(\sigma)$, by [Vas21a, Theorem 1.1] it is the case that

$$u_{0,+}(-; \sigma) = \tilde{R}_+(\sigma)\tilde{f}(-; \sigma) \quad (549)$$

for each $\sigma > 0$ (since, for each $\sigma > 0$, Φ differs from Vasy's phase by a remainder the exponential of which acts as a bounded multiplication operator on b-Sobolev spaces — cf. eq. (67), with $a = 0$).

By Proposition 6.15, $u_{0,+} = x^{(n-1)/2}(\sigma^2 + Zx)^{-1/4}u_{0,+}$ for some $u_{0,+} \in \mathcal{A}_{\text{loc}}^{0-,0-, (0,0)}(X_{\text{res}}^{\text{sp}})$, depending continuously on f . Then, by Proposition 6.16, we conclude that

$$u_{0,+} \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > Ea_{00}\}) + \mathcal{A}_{\text{loc}}^{((0,0), \delta_1), 2\delta_0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{Z > Ea_{00}\}), \quad (550)$$

depending continuously on $f \in \mathcal{S}(X)$, with $\delta_1 = \min\{\beta_2, 1/2 + \beta_3\} = \{\alpha_1 - 1, \alpha_2 - 1\}$ and $\delta_0 = \min\{\beta_2, \beta_3\} = \min\{\alpha_1 - 1, \alpha_2 - 3/2\}$. This yields the first half of the proposition (as well as the continuity clause).

The second clause of this proposition then follows from the second clause of Proposition 6.15 and Proposition 4.2. Indeed, by the latter,

$$e^{-i\Phi(-;0)}u_+ = e^{-i\Phi(-;0)}R(0; Z + i0)f = \tilde{R}_+(0)\tilde{f}(-; 0). \quad (551)$$

By the second half of Proposition 6.15, the right-hand side is the restriction to zf of $u_{+,00}$. \square

6.1. Estimates.

Lemma 6.4. *For any $l, k \in \mathbb{R}$ satisfying $l < -1/2$ and $k + l \leq -3/4$, and for any $N \in \mathbb{N}$, there exists a $C = C(\bar{x}, l, k, N, Z) > 0$ such that the linear map*

$$\mathcal{N}_{l,N} = \mathcal{N}_{l,N}(k, \sigma) : H_{\text{b}}^{-N,l}[0, \bar{x}] \rightarrow H_{\text{b}}^{-N-1,l}[0, \bar{x}] \quad (552)$$

given by

$$\mathcal{N}_{l,N} = 2i\left[x\partial_x + \frac{1}{2}\right] + \frac{2iZx}{\sigma^2 + Zx}\left(k + \frac{1}{4}\right) \quad (553)$$

satisfies $\|u\|_{H_{\text{b}}^{-N,l}[0, \bar{x}]} \leq C\|\mathcal{N}_{l,N}(k, \sigma)u\|_{H_{\text{b}}^{-N-1,l}[0, \bar{x}]}$ for all $\sigma \in [0, \infty)$ and $u \in H_{\text{b}}^{-N,l}[0, \bar{x}]$.

The same estimate (possibly with a different constant) holds with $\hat{X} = [0, \bar{x}]_x \times \partial X$ in place of $[0, \bar{x}]$ if we replace $H_{\text{b}}^{-N,l}[0, \bar{x}]$ with $\hat{H}_{\text{b}}^{-N,l}(\hat{X})$, where the latter is the usual b-Sobolev space (such that $\hat{H}_{\text{b}}^{0,0}(\hat{X}) = L_{\text{b}}^2(\hat{X})$). \blacksquare

Proof. We write the proof for $[0, \bar{x}]$, and the proof for \hat{X} is similar (since angular derivatives commute with $\mathcal{N}_{l,N}$) once we interpret u as a $H^m(\partial X)$ -valued element of a b-Sobolev space on $[0, \bar{x}]$.

Via the Mellin transform, $\mathcal{N}_{l,N}(k, 0) = 2i(x\partial_x + k + 3/4) : H_{\text{b}}^{-N,l}[0, \bar{x}] \rightarrow H_{\text{b}}^{-N-1,l}[0, \bar{x}]$ is invertible for l, k as in the lemma statement, so it suffices to restrict attention to the case $\sigma > 0$, i.e. to prove that

$$\|u\|_{H_{\text{b}}^{-N,0}[0, \bar{x}]} \preceq \|x^{-l}\mathcal{N}_{l,N}(k, \sigma)x^l u\|_{H_{\text{b}}^{-N-1,0}[0, \bar{x}]} \quad (554)$$

for all $u \in H_{\text{b}}^{-N,0}[0, \bar{x}]$ and $\sigma > 0$ (the estimate required to be uniform in σ). Let $\hat{x} = Zx/\sigma^2$.

Via the dilation invariance of the b-Sobolev spaces, the estimate eq. (554) is equivalent to the following:

$$\|u\|_{H_{\text{b}}^{-N,0}[0, Z/\sigma^2]} \preceq \|\widehat{\mathcal{N}}_{l,k}u\|_{H_{\text{b}}^{-N-1,0}[0, Z/\sigma^2]} \quad (555)$$

for all $u \in H_{\text{b}}^{-N,0}[0, Z/\sigma^2]$ and $\sigma > 0$, where

$$\widehat{\mathcal{N}}_{l,k} = 2i\left[\hat{x}\partial_{\hat{x}} + l + \frac{1}{2}\right] + \frac{2i\hat{x}}{1 + \hat{x}}\left(k + \frac{1}{4}\right). \quad (556)$$

We now “radially” compactify the nonnegative real axis $[0, \infty)_{\hat{x}}$ (so that $1/(1 + \hat{x})$ becomes a bdf for the new boundary face), and we call the result $[0, \infty]$. Observe that $\widehat{\mathcal{N}}_{l,k}$ is a b-differential operator on $[0, \infty]$, with b-decay rate zero at both ends. In order to prove the estimate eq. (555), it suffices to prove that

$$\|u\|_{H_b^{-N,0,0}[0,\infty]} \preceq \|\widehat{\mathcal{N}}_{l,k}u\|_{H_b^{-N-1,0,0}[0,\infty]} \quad (557)$$

for all $u \in H_b^{-N,0,0}[0, \infty]$, where the third index is the b-decay order at $\hat{x} = \infty$.

For $l \neq -1/2$ and $l + k \neq -3/4$, $\widehat{\mathcal{N}}_{l,k}$ is (via b-ellipticity and the Taylor series expansion of $\hat{x}/(1 + \hat{x})$ around $\hat{x} = [0, \infty)$) Fredholm as an operator $H_b^{-N,0,0}[0, \infty] \rightarrow H_b^{-N-1,0,0}[0, \infty]$, and for any $N_0 \in \mathbb{N}$ we have the estimate

$$\|u\|_{H_b^{-N,0,0}[0,\infty]} \preceq \|\widehat{\mathcal{N}}_{l,k}u\|_{H_b^{-N-1,0,0}[0,\infty]} + \|u\|_{H_b^{-N_0,-1,-1}[0,\infty]} \quad (558)$$

for $u \in H_b^{-N,0,0}[0, \infty]$. Once $\widehat{\mathcal{N}}_{l,k}$ is known to be injective, a standard argument allows us to remove the last term of eq. (558), yielding the desired estimate eq. (557). It suffices to consider the case $N_0 > N$. The standard argument is as follows:

- If we could not remove the last term of eq. (558), then we would be able to find a sequence $\{u_j\}_{j \in \mathbb{N}} \subset H_b^{-N,0,0}[0, \infty]$ with $\|u_j\|_{H_b^{-N,0,0}[0,\infty]} = 1$ for all j and

$$\|\widehat{\mathcal{N}}_{l,k}u_j\|_{H_b^{-N-1,0,0}[0,\infty]} \rightarrow 0 \quad (559)$$

as $j \rightarrow \infty$. By the Banach–Alaoglu theorem, we may assume without loss of generality (by passing from $\{u_j\}_{j \in \mathbb{N}}$ to a subsequence if necessary) that there exists some $u_\infty \in H_b^{-N-1,0,0}[0, \infty]$ such that $u_j \rightarrow u_\infty$ weakly as $j \rightarrow \infty$.

Via the compactness of the inclusion $H_b^{-N,0,0}[0, \infty] \hookrightarrow H_b^{-N_0,-1,-1}[0, \infty]$ for $N_0 > N$, $u_j \rightarrow u_\infty$ strongly in the latter space. From eq. (558), we deduce that

$$\|u_\infty\|_{H_b^{-N_0,-1,-1}[0,\infty]} = \lim_{j \rightarrow \infty} \|u_j\|_{H_b^{-N_0,-1,-1}[0,\infty]} \succeq 1. \quad (560)$$

In particular, (I) u_∞ is nonzero.

Also, from the strong convergence of $u_j \rightarrow u_\infty$ in $H_b^{-N_0,-1,-1}[0, \infty]$,

$$\widehat{\mathcal{N}}_{l,k}u_j \rightarrow \widehat{\mathcal{N}}_{l,k}u_\infty \quad (561)$$

distributionally. But $\widehat{\mathcal{N}}_{l,k}u_j \rightarrow 0$ strongly in $H_b^{-N-1,0,0}[0, \infty]$, by eq. (559). Thus, (II) $\widehat{\mathcal{N}}_{l,k}u_\infty = 0$.

From (I) and (II), we conclude that $\widehat{\mathcal{N}}_{l,k}$ is not injective.

In order to show that

$$\ker_{H_b^{-N,0,0}[0,\infty]} \widehat{\mathcal{N}}_{l,k} = \{u \in H_b^{-N,0,0}[0, \infty] : \widehat{\mathcal{N}}_{l,k}u = 0\} \quad (562)$$

is trivial, we simply appeal to the solution eq. (532) of the ODE (although it is slightly simpler to integrate in the other direction). Indeed, any element u of the kernel eq. (562) must be given by $u(\hat{x}) = c\hat{x}^{-(l+1/2)}(1 + \hat{x})^{-k-1/4}$ for some $c \in \mathbb{C}$. If this is nonzero, then it is $\Omega(\hat{x}^{-(l+k+3/4)})$ as $\hat{x} \rightarrow \infty$. If $l + k + 3/4 \leq 0$, then u fails to lie in $L_b^2[0, \infty]$. By eq. (558) (applied with 0 in place of N and N in place of N_0),

$$\|u\|_{L_b^2[0,\infty]} \preceq \|\widehat{\mathcal{N}}_{l,k}u\|_{H_b^{-1,0,0}[0,\infty]} + \|u\|_{H_b^{-N,-1,-1}[0,\infty]} \preceq \|\widehat{\mathcal{N}}_{l,k}u\|_{H_b^{-1,0,0}[0,\infty]} + \|u\|_{H_b^{-N,0,0}[0,\infty]}. \quad (563)$$

Since $\|u\|_{L_b^2[0,\infty]} = \infty$ and $u \in \ker_{H_b^{-N,0,0}[0,\infty]} \widehat{\mathcal{N}}_{l,k} \Rightarrow \|\widehat{\mathcal{N}}_{l,k}u\|_{H_b^{-1,0,0}[0,\infty]} = 0$, eq. (563) implies

$$u \notin H_b^{-N,0,0}[0, \infty], \quad (564)$$

which contradicts $u \in \ker_{H_b^{-N,0,0}[0,\infty]} \widehat{\mathcal{N}}_{l,k}$. Thus, $u \in \ker_{H_b^{-N,0,0}[0,\infty]} \widehat{\mathcal{N}}_{l,k} \Rightarrow u = 0$.

This completes the proof of the lemma. \square

Proposition 6.5. *For any $l < -1/2$, $k + l \leq -3/4$, and $N \in \mathbb{N}$, there exists a constant $C = C(\tilde{P}, N, l, \ell, k) > 0$ such that, for all $\sigma \geq 0$,*

$$\|v\|_{H_b^{-N,l}(X)} \leq C \cdot (\|(\sigma^2 + Zx)^{-k} N(\tilde{P}(\sigma))((\sigma^2 + Zx)^k v)\|_{H_{b,\text{leC}}^{-N,l+1,2l+3}(X)}) \quad (565)$$

$$\|v\|_{H_{b,\text{leC}}^{-N,l,2l+2k}(X)} \leq C \cdot (\|N(\tilde{P}(\sigma))v\|_{H_{b,\text{leC}}^{-N,l+1,2l+2k+3}(X)}) \quad (566)$$

for all $v \in \mathcal{S}'(X)$ supported in $\{x < \bar{x}\} \subset X$. \blacksquare

Proof. Letting $\tilde{N}(\tilde{P}(\sigma)) = (\sigma^2 + Zx)^{-k} N(\tilde{P})(\sigma^2 + Zx)^k$, $\tilde{N}(\tilde{P}(\sigma)) = N(\tilde{P}(\sigma)) + 2Zkix^2/(\sigma^2 + Zx)^{1/2}$. Consequently, we can choose the b,leC-Sobolev and b-Sobolev norms such that

$$\|\tilde{N}(\tilde{P}(\sigma))v\|_{H_{b,\text{leC}}^{-N,l+1,2l+3}} = \|(2ix(x\partial_x - (n-1)/2) + 2Zkix^2/(\sigma^2 + Zx))v\|_{H_b^{-N,l+1}(X)}. \quad (567)$$

We can work on $\hat{X} = [0, 1]_x \times \partial X$, as

$$\|x^{-n/2}w\|_{\tilde{H}_b^{-N,l}(\hat{X})} \preceq \|w\|_{H_b^{-N,l}(X)} \preceq \|x^{-n/2}w\|_{\tilde{H}_b^{-N,l}(\hat{X})} \quad (568)$$

for all $N, l \in \mathbb{R}$ and $w \in \mathcal{S}'(X)$ supported in $\{x < \bar{x}\}$. Equation (565) is therefore equivalent to

$$\|x^{-n/2}v\|_{\tilde{H}_b^{-N,l}(\hat{X})} \preceq \|x^{-n/2}(2ix(x\partial_x - (n-1)/2) + 2Zkix^2/(\sigma^2 + Zx))v\|_{\tilde{H}_b^{-N,l+1}(X)} \quad (569)$$

for $v \in H_b^{-N,l}(X)$ supported in $\{x < \bar{x}\}$, which follows if

$$\|w\|_{\tilde{H}_b^{-N,l}(\hat{X})} \preceq \|(2i(x\partial_x + 1/2) + 2Zkix/(\sigma^2 + Zx))w\|_{\tilde{H}_b^{-N,l}(X)} \quad (570)$$

holds for $w \in \tilde{H}_b^{-N,l}(\hat{X})$, which was the conclusion of Lemma 6.4. \square

Proposition 6.6. *For each $\Sigma > 0$, $N \in \mathbb{N}$, $m, s, \varsigma, l, \ell \in \mathbb{R}$ satisfying $l < -1/2$, $\ell < -3/2$, $s > s_0 > -1/2$, $-3/2 < \varsigma_0 < \varsigma \leq \ell + s - l$, there exists a constant $C = C(\tilde{P}, \Sigma, N, m, s, \varsigma, l, \ell) > 0$ such that, for any $u \in \mathcal{S}'(X)$*

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C \cdot (\|\tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{b,\text{leC}}^{-N,l-\delta,\ell-2\delta}}) \quad (571)$$

holds for any $\sigma \in [0, \Sigma]$ such that $u(-; \sigma) \in H_{\text{scb}}^{-N,s_0,\varsigma_0}(X)$. \blacksquare

Proof. Consider $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$ as in the proposition statement. By Proposition 5.13, we have

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \preceq \|\tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{b,\text{leC}}^{-N_0-2,l,\ell}}, \quad (572)$$

where $N_0 \in \mathbb{N}$ is arbitrary. We now apply Proposition 6.5 to estimate the remainder term. Let $\chi \in C_c^\infty(X)$ be supported in $x \leq \bar{x}$ and identically equal to one in some neighborhood of $x = 0$. First of all,

$$\|u\|_{H_{b,\text{leC}}^{-N_0-2,l,\ell}} \preceq \|\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l,\ell}} + \|u\|_{H_{b,\text{leC}}^{-N_0,-N_0,-N_0}}. \quad (573)$$

Set $k = (\ell - 2l)/2$. Then, $k + l \leq -3/4$. We now apply the previous proposition with $v = \chi u$, the result being

$$\begin{aligned} \|\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l,\ell}} &\preceq \|N(\tilde{P})\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} \\ &\leq \|\tilde{P}\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} + \|E\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} \end{aligned} \quad (574)$$

for $E = N(\tilde{P}) - \tilde{P}$. By Proposition 3.8, $E \in \Psi_{b,\text{leC}}^{2,-1-\delta,-3-2\delta}(X)$ for some $\delta \in (0, 1/2)$, so

$$\|E\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} \preceq \|u\|_{H_{b,\text{leC}}^{-N_0,l-\delta,\ell-2\delta}}. \quad (575)$$

On the other hand, since χ is identically one in some neighborhood of ∂X ,

$$\begin{aligned} \|\tilde{P}\chi u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} &\preceq \|\tilde{P}u\|_{H_{b,\text{leC}}^{-N_0-2,l+1,\ell+3}} + \|u\|_{H_{b,\text{leC}}^{-N_0,-N_0,-N_0}} \\ &\preceq \|\tilde{P}u\|_{H_{b,\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} + \|u\|_{H_{b,\text{leC}}^{-N_0,-N_0,-N_0}} \end{aligned} \quad (576)$$

for sufficiently large N_0 . Combining the estimates above (for sufficiently large N_0), we get eq. (571). \square

For each $m, s, \varsigma, l, \ell \in \mathbb{R}$, we consider the families $\mathcal{X} = \mathcal{X}_{m,s,\varsigma,l,\ell} = \{\mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma)\}_{\sigma \geq 0}$ and $\mathcal{Y} = \mathcal{Y}_{m,s,\varsigma,l,\ell} = \{\mathcal{Y}_{m,s,\varsigma,l,\ell}(\sigma)\}_{\sigma \geq 0}$ given by

$$\mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma) = \{u \in H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X) : \tilde{P}u \in H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)\} \quad (577)$$

$$\mathcal{Y}_{m,s,\varsigma,l,\ell}(\sigma) = H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X), \quad (578)$$

considered as families of Banach spaces in the usual way,

$$\|u\|_{\mathcal{X}} = \|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} + \|\tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}}. \quad (579)$$

Note that $\mathcal{X}_{m,s,\varsigma,l,\ell}(0) = \mathcal{X}_{m,\varsigma,\ell}$ and $\mathcal{Y}_{m,s,\varsigma,l,\ell}(0) = \mathcal{Y}_{m,\varsigma,\ell}$, where the right-hand sides are defined eq. (307) and eq. (308). Tautologically, $\tilde{P} : \mathcal{X} \rightarrow \mathcal{Y}$ is bounded, uniformly in $\sigma \geq 0$, as $\|\tilde{P}u\|_{\mathcal{Y}} \leq \|u\|_{\mathcal{X}}$.

Proposition 6.7. *Given m, s, ς, l, ℓ satisfying the inequalities $l < -1/2$, $\ell < -3/2$, $-1/2 < s$, $-3/2 < \varsigma \leq \ell + s - l$, one of the following two alternatives holds:*

- *there exists some $\sigma \geq 0$ and nonzero $u \in \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma)$ such that $\tilde{P}(\sigma)u = 0$,*
- *there exists, for each $\Sigma > 0$, a constant $C_0 = C_0(\tilde{P}, m, s, \varsigma, l, \ell, \Sigma) > 0$ such that the estimate*

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}} \leq C_0 \|\tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} \quad (580)$$

holds for all $\sigma \in [0, \Sigma]$ and all $u \in \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma)$.

■

Proof. The following is a variant of the proof of [Hö7, Theorem 26.1.7], also used in the proof of the main theorem in [Vas21a].

Suppose that the second of the two alternatives does not hold, so that there exist $\Sigma > 0$ and sequences $\{\sigma_k\}_{k=0}^\infty \subset [0, \Sigma]$ and $\{u_k\}_{k=0}^\infty \subset \mathcal{S}'(X)$ with $\|u_k\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_k)} = 1$ and

$$\|\tilde{P}(\sigma_k)u_k\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)(\sigma_k)} < 1/k \quad (581)$$

for all $k \in \mathbb{N}$. By passing to a subsequence if necessary (and noting that eq. (581) continues to hold upon doing so), we can arrange that $\sigma_k \rightarrow \sigma_\infty$ for some $\sigma_\infty \in [0, \Sigma]$.

Even though eq. (580) might not hold, by Proposition 6.6 we at least have the bound

$$1 = \|u_k\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_k)} \leq C \cdot (\|\tilde{P}u_k\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)(\sigma_k)} + \|u_k\|_{H_{b,\text{leC}}^{-N,l-\delta,\ell-2\delta}}), \quad (582)$$

for any $N \in \mathbb{N}$, where $C = C(\tilde{P}, \Sigma, N, m, s, \varsigma, l, \ell)$ is some constant. On the other hand, for sufficiently large $N_0 > N$, we can bound

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_k)} \succeq \|\Lambda_{m,s,\varsigma,l,\ell}(\sigma_k)u\|_{L^2(X)} + \|x^{-l}(\sigma^2 + Zx)^{l-\ell/2}u\|_{H_b^{-N_0,0}(X)}, \quad (583)$$

this holding for all $u \in \mathcal{S}'(X)$ and $\sigma \in [0, \Sigma]$. Consequently, $\{\Lambda_{m,s,\varsigma,l,\ell}(\sigma_k)u_k\}_{k \in \mathbb{N}}$ is bounded in $L^2(X)$ and $\{x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\}_{k \in \mathbb{N}}$ is bounded in $H_b^{-N_0,0}(X)$. By Banach–Alaoglu – passing to a subsequence if necessary (and once again noting that eq. (581) continues to hold upon doing so) – we can assume that there exist some $v \in L^2(X)$ and $w \in H_b^{-N_0,0}(X)$ such that

$$\Lambda_{m,s,\varsigma,l,\ell}(\sigma_k)u_k \rightarrow v \quad (584)$$

as $k \rightarrow \infty$ weakly in $L^2(X)$ and $x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k \rightarrow w$ as $k \rightarrow \infty$ weakly in $H_b^{-N_0,0}(X)$.

It follows from the latter that $u_k \rightarrow x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w$ strongly in some b-Sobolev space. This has two consequences:

- First,

$$\tilde{P}(\sigma_k)u_k \rightarrow \tilde{P}(\sigma_\infty)(x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w) \quad (585)$$

in $\mathcal{S}'(X)$ as $k \rightarrow \infty$ (e.g. using Proposition 2.4 and Proposition 3.6). But, the assumption $\|\tilde{P}(\sigma_k)u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)(\sigma_k)} < 1/k$ implies that $\tilde{P}(\sigma_k)u_k \rightarrow 0$ in $\mathcal{S}'(X)$. Therefore $\tilde{P}(\sigma_\infty)(x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w) = 0$.

- Second (using Proposition 2.4), $\Lambda_{m,s,\varsigma,l,\ell}(\sigma_k)u_k \rightarrow \Lambda_{m,s,\varsigma,l,\ell}(\sigma_\infty)(x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w)$ in $\mathcal{S}'(X)$. Since $\mathcal{S}'(X)$ is Hausdorff, this implies that

$$\Lambda_{m,s,\varsigma,l,\ell}(\sigma_\infty)(x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w) = v \in L^2(X), \quad (586)$$

which in turn implies that $x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w \in H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_\infty)$ by elliptic regularity (and the fact that $w \in H_b^{-N_0,0}$).

Let $u = x^l(\sigma_\infty^2 + Zx)^{\ell/2-l}w$. What we proved above is $\tilde{P}(\sigma_\infty)u = 0$ and $u \in H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma_\infty)$. Thus,

$$u \in \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma_\infty). \quad (587)$$

Since $N > N_0$, via the compactness of the inclusion $H_b^{-N_0,0}(X) \hookrightarrow H_b^{-N,-\delta}(X)$ it is the case that $x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k \rightarrow w$ strongly in $H_b^{-N,-\delta}(X)$, so

$$\|w\|_{H_b^{-N,-\delta}(X)} = \lim_{k \rightarrow \infty} \|x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\|_{H_b^{-N,-\delta}(X)}. \quad (588)$$

On the other hand, we can bound $\|u_k\|_{H_{b,\text{leC}}^{-N,l-\delta,\ell-2\delta}(X)(\sigma_k)} \preceq \|x^{-l}(\sigma_k^2 + Zx)^{l-\ell/2}u_k\|_{H_b^{-N,-\delta}(X)}$. So, eq. (582) yields

$$\begin{aligned} 1 &\preceq \|w\|_{H_b^{-N,-\delta}(X)} + \limsup_{k \rightarrow \infty} \|\tilde{P}u_k\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)(\sigma_k)} \\ &= \|w\|_{H_b^{-N,-\delta}(X)}. \end{aligned} \quad (589)$$

Therefore $w \neq 0$. It follows that $u \neq 0$. We have therefore succeeded in showing that the second of the two alternatives listed in the proposition holds. \square

Proposition 6.8. *Suppose that P is the spectral family of an attractive Coulomb-like Schrödinger operator for $\sigma \leq \Sigma$. Then, given $m, s, \varsigma, l, \ell \in \mathbb{R}$ satisfying $l < -1/2$, $\ell < -3/2$, $-1/2 < s$, $-3/2 < \varsigma$, it is the case that, for each $\sigma \geq \Sigma$,*

$$\tilde{P}(\sigma) : \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma) \rightarrow \mathcal{Y}_{m,s,\varsigma,l,\ell}(\sigma) \quad (590)$$

is invertible. \blacksquare

Proof. We already observed the $\sigma = 0$ case in §3. The $\sigma > 0$ case is essentially proven in [Vas21a, §4]. In order to see this, note that, for each $\sigma > 0$,

$$\mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma) = \{u \in H_{\text{scb}}^{m,s,l}(X) : \tilde{P}(\sigma)u \in H_{\text{scb}}^{m-2,s+1,l+1}(X)\} \quad (591)$$

$$\mathcal{Y}_{m,s,\varsigma,l,\ell}(\sigma) = H_{\text{scb}}^{m-2,s+1,l+1}(X) \quad (592)$$

at the level of sets, where

$$\tilde{P}(\sigma) = e^{-i\Phi_0 + i\Phi} \tilde{P} e^{i\Phi_0 - i\Phi} \quad (593)$$

is Vasy's conjugated operator, Φ_0 simply being defined by $\Phi_0 = \sigma x^{-1}$. Since the leC-Sobolev spaces are just scb-Sobolev spaces for $\sigma > 0$, the crux of the previous claim is that the \tilde{P} on the

right-hand side of eq. (577) can be replaced by Vasy's eq. (593). Indeed, $\Phi_0(x; \sigma) - \Phi(x; \sigma) \in \log x C^\infty([0, \bar{x})_x) + C^\infty([0, \bar{x})_x)$ for each $\sigma > 0$, so

$$\tilde{\tilde{P}}(\sigma) = \tilde{P}(\sigma) + T(\sigma) \quad (594)$$

for $T(\sigma) = e^{-i\Phi_0(-; \sigma) + i\Phi(-; \sigma)}[\tilde{P}(\sigma), e^{+i\Phi_0(-; \sigma) - i\Phi(-; \sigma)}] \in \text{Diff}_{\text{sc}}^{1, -1, -1}(X)$. Thus,

$$T(\sigma) : H_{\text{scb}}^{m, s, l}(X) \rightarrow H_{\text{scb}}^{m-2, s+1, l+1}(X). \quad (595)$$

So, for $u \in H_{\text{scb}}^{m, s, l}(X)$, $\tilde{\tilde{P}}(\sigma)u \in H_{\text{scb}}^{m-2, s+1, l+1}(X)$ if and only if $\tilde{P}(\sigma)u \in H_{\text{scb}}^{m-2, s+1, l+1}(X)$. The operators we consider have slightly more general σ -dependence than the ones in Vasy (as some additional assumptions are needed for the $\sigma \notin \mathbb{R}$ case of [Vas21a, Theorem 1.1]), but since we are only considering real σ his proof of the real case of [Vas21a, Theorem 1.1] goes through in this slightly greater generality *mutatis mutandis*.

Alternatively, the $\sigma > 0$ case of Proposition 6.6 suffices as a replacement for [Vas21a, Prop. 4.16] in his proof of [Vas21a, Theorem 1.1], the rest of which is identical. (For the purpose of the proof above, we do not need to know that the estimate in Proposition 6.6 is uniform down to $\sigma = 0$, so the $\varsigma \leq \ell + s - l$ hypothesis there is not relevant here.) \square

6.2. Smoothness at zf, Conormality elsewhere. For this subsection, we suppose that $P(\sigma)$ is the spectral family of an attractive Coulomb-like Schrödinger operator for σ in some neighborhood of $[0, \Sigma]$, $\Sigma > 0$. By analogy with the terminology in [Mel94], we might say that $u \in \mathcal{S}'(X)$ “satisfies the conjugated Sommerfeld radiation condition” for some given $\sigma \geq 0$ if $u \in \mathcal{X}_{m, s, \varsigma, l, \ell}(X)(\sigma)$ for some $m, s, \varsigma, l, \ell \in \mathbb{R}$ satisfying $l < -1/2$, $\ell < -3/2$, $s > -1/2$, $\ell + s - l \geq \varsigma > -3/2$. One of the main tasks of this subsection is to show that the limiting resolvent output converges as $\sigma \rightarrow 0^+$ to something satisfying the zero energy version of the Sommerfeld radiation condition.

For each $\sigma \in [0, \Sigma]$, let $\tilde{R}_+(\sigma) : \mathcal{Y}_{m, s, \varsigma, l, \ell}(\sigma) \rightarrow \mathcal{X}_{m, s, \varsigma, l, \ell}(\sigma)$ denote the set-theoretic inverse to eq. (590) (which, of course, must actually be an isomorphism of Banach spaces e.g. by the closed graph theorem). (The ‘+’ subscript of $\tilde{R}_+ = \{\tilde{R}_+(\sigma)\}_{\sigma \geq 0}$ refers to the choice of sign in defining the conjugation.) This extends the definition of the operator $\tilde{R}_+(0)$ introduced at the end of §3 to the $\sigma > 0$ case. For each $\sigma \in (0, \Sigma]$,

$$\tilde{R}_+(\sigma) : H_{\text{scb}}^{m-2, s+1, l+1}(X) \rightarrow H_{\text{scb}}^{m, s, l}(X) \quad (596)$$

is bounded (but not uniformly in σ). Considering the case $s = m + l$, $\tilde{R}_+(\sigma) : H_{\text{b}}^{m, l+1}(X) \rightarrow H_{\text{b}}^{m, l}(X)$ if $l < -1/2 < m + l$. Note that the mapping properties of the resolvent with respect to the b-Sobolev spaces are slightly lossy, in the sense that we can no longer keep track of the fact that the map eq. (596) smooths by two orders. As the notation in eq. (596) indicates, the operator $\tilde{R}_+(\sigma)$ makes sense as a map

$$\bigcup_{\substack{m, s, l \in \mathbb{R} \\ s > -1/2 > l}} H_{\text{scb}}^{m-2, s+1, l+1}(X) \rightarrow \bigcup_{\substack{m, s, l \in \mathbb{R} \\ s > -1/2 > l}} H_{\text{scb}}^{m, s, l}(X), \quad (597)$$

hence we can just write “ $\tilde{R}_+(\sigma)$ ” without specifying m, s, l . A similar statement holds for $\sigma = 0$.

Proposition 6.9. *Given $m, s, \varsigma, l, \ell \in \mathbb{R}$ satisfying $l < -1/2$, $\ell < -3/2$, $s > -1/2$, $\ell + s - l \geq \varsigma > -3/2$, there exists some constant $C = C(m, s, \varsigma, l, \ell, \Sigma) > 0$ such that*

$$\|\tilde{R}_+(\sigma)\tilde{f}\|_{H_{\text{leC}}^{m, s, \varsigma, l, \ell}(X)(\sigma)} \leq C\|\tilde{f}\|_{H_{\text{leC}}^{m-2, s+1, \varsigma+3, l+1, \ell+3}(X)(\sigma)} \quad (598)$$

for all $\sigma \in [0, \Sigma]$ and $\tilde{f} \in H_{\text{leC}}^{m-2, s+1, \varsigma+3, l+1, \ell+3}(X)(\sigma)$. Moreover, for $m, l \in \mathbb{R}$ with $l < -1/2$ and $-1 < m + 2l$:

(I) for any $\tilde{f} \in H_{\text{b}}^{m, l+5/4}(X)$, we have $\|(\sigma^2 + \mathbf{Z}x)^{1/4}\tilde{R}_+(\sigma)\tilde{f}\|_{H_{\text{b}}^{m, l}} \leq C_0\|\tilde{f}\|_{H_{\text{b}}^{m, l+5/4}}$ for some constant $C_0 = C_0(\tilde{P}, m, l, \Sigma) > 0$

(II) for any $\tilde{f} \in H_b^{m,l+1}(X)$ we have $\|(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)(\sigma^2 + Zx)^{1/4} \tilde{f}\|_{H_b^{m,l}} \leq C_1 \|\tilde{f}\|_{H_b^{m,l+1}}$ for some constant $C_1 = C_1(\tilde{P}, m, l, \Sigma) > 0$ for all $\sigma \in [0, \Sigma]$. ■

Proof. As a corollary of Proposition 6.7 and Proposition 6.8, we get that for m, s, ς, l, ℓ as above and P the spectral family of an attractive Coulomb-like Schrödinger operator,

$$\|u\|_{H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma)} \leq C(m, s, \varsigma, l, \ell, \Sigma) \|\tilde{P}u\|_{H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}} \quad (599)$$

holds for all $\sigma \in [0, \Sigma]$ and all $u \in \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma)$. We also have

$$\tilde{R}_+(\sigma) : H_{\text{leC}}^{m-2,s+1,\varsigma+3,l+1,\ell+3}(X)(\sigma) \rightarrow \mathcal{X}_{m,s,\varsigma,l,\ell}(\sigma) \subseteq H_{\text{leC}}^{m,s,\varsigma,l,\ell}(X)(\sigma). \quad (600)$$

Taking $u = \tilde{R}_+(\sigma)\tilde{f}$ in eq. (599) yields Equation (598).

Suppose now that $m, l \in \mathbb{R}$ satisfy $l < -1/2$ and $-1 < m + 2l$ (in which case $-1/2 < m + l$ holds as well). First suppose that $\tilde{f} \in H_b^{m,l+5/4}(X)$. Applying eq. (598) (observing that the required inequalities $l < -1/2 < m + l$ and $-3/2 < m + 2l - 1/2$ hold),

$$\begin{aligned} \|(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)\tilde{f}\|_{H_b^{m,l}} &\leq \|(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)\tilde{f}\|_{H_{\text{leC}}^{m,m+l,m+2l,l,2l}(\sigma)} \\ &\leq \|\tilde{R}_+(\sigma)\tilde{f}\|_{H_{\text{leC}}^{m,m+l,m+2l-1/2,l,2l-1/2}(\sigma)} \\ &\leq \|\tilde{f}\|_{H_{\text{leC}}^{m-2,m+l+1,m+2l+5/2,l+1,2l+5/2}(\sigma)} \\ &\leq \|\tilde{f}\|_{H_b^{m,l+5/4}}. \end{aligned} \quad (601)$$

Now supposing that $\tilde{f} \in H_b^{m,l+1}(X)$,

$$\begin{aligned} \|(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)(\sigma^2 + Zx)^{1/4} \tilde{f}\|_{H_b^{m,l}} &\leq \|(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)(\sigma^2 + Zx)^{1/4} \tilde{f}\|_{H_{\text{leC}}^{m,m+l,m+2l,l,2l}(\sigma)} \\ &\leq \|\tilde{R}_+(\sigma)(\sigma^2 + Zx)^{1/4} \tilde{f}\|_{H_{\text{leC}}^{m,m+l,m+2l-1/2,l,2l-1/2}(\sigma)} \\ &\leq \|(\sigma^2 + Zx)^{1/4} \tilde{f}\|_{H_{\text{leC}}^{m-2,m+l+1,m+2l+5/2,l+1,2l+5/2}(\sigma)} \\ &\leq \|\tilde{f}\|_{H_{\text{leC}}^{m-2,m+l+1,m+2l+2,l+1,2l+2}(\sigma)} \\ &\leq \|\tilde{f}\|_{H_b^{m,l+1}}. \end{aligned} \quad (602)$$

□

Proposition 6.10. For any $m, l \in \mathbb{R}$ with $l < -1/2$ and $-1 < m + 2l$, for any $f \in H_b^{m,l+5/4}(X)$,

$$(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)f \rightarrow Z^{1/4}x^{1/4} \tilde{R}_+(0)f \quad (603)$$

weakly in $H_b^{m,l}(X)$ as $\sigma \rightarrow 0^+$. In fact, the map $(\sigma, f) \mapsto (\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)f$ defines a jointly continuous map

$$[0, \Sigma] \times H_b^{m,l+5/4}(X) \rightarrow H_b^{m-\epsilon,l-\epsilon}(X), \quad (604)$$

for any $\epsilon > 0$, where we are using the strong topologies on $H_b^{m,l+5/4}(X)$ and $H_b^{m-\epsilon,l-\epsilon}(X)$.

This (applied for slightly smaller m, l) implies that

$$\{(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)\}_{\sigma \geq 0} \subset \mathcal{L}(H_b^{m,l+5/4}(X), H_b^{m-\epsilon,l-\epsilon}(X)) \quad (605)$$

is continuous with respect to the uniform operator topology. ■

Proof. First consider the claim of joint continuity. By the metrizable of $[0, \Sigma] \times H_b^{m,l+5/4}(X)$ and $H_b^{m-\epsilon,l-\epsilon}(X)$, joint continuity follows from the claim that

- given $f \in H_b^{m,l+5/4}(X)$ and $\sigma_\infty \in [0, \Sigma)$, for any and sequences $\{\sigma_k\}_{k \in \mathbb{N}} \subset [0, \Sigma)$ and $\{f_k\}_{k \in \mathbb{N}} \subset H_b^{m,l+5/4}(X)$ with $\sigma_k \rightarrow \sigma_\infty$ and $f_k \rightarrow f$ as $k \rightarrow \infty$,

$$(\sigma_k^2 + Zx)^{1/4} \tilde{R}_+(\sigma_k) f_k \rightarrow (\sigma_\infty^2 + Zx)^{1/4} \tilde{R}_+(\sigma_\infty) f \quad (606)$$

strongly in $H_b^{m-\epsilon, l-\epsilon}(X)$.

Since a sequence of elements of a metric space converges to some element if and only if every subsequence thereof contains a further subsequence converging to that same element, it suffices to show that

- given any $\{\sigma_k\}_{k \in \mathbb{N}} \subset [0, \Sigma)$, $\{f_k\}_{k \in \mathbb{N}} \subset H_b^{m,l+5/4}(X)$ with $\sigma_k \rightarrow \sigma_\infty$ and $f_k \rightarrow f$ as $k \rightarrow \infty$, there exists a subsequence $\{k_\kappa\}_{\kappa \in \mathbb{N}} \subset \{k\}_{k \in \mathbb{N}}$ such that

$$(\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa} \rightarrow (\sigma_\infty^2 + Zx)^{1/4} \tilde{R}_+(\sigma_\infty) f \quad (607)$$

strongly in $H_b^{m-\epsilon, l-\epsilon}(X)$.

We now handle the case of $\sigma_\infty = 0$. The case $\sigma_\infty > 0$ follows by a similar but even easier argument.

By Proposition 6.9 and Banach–Alaoglu, we can find a subsequence $\{k_\kappa\}_{\kappa \in \mathbb{N}} \subset \{k\}_{k \in \mathbb{N}}$ such that $(\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa}$ converges weakly in $H_b^{m,l}(X)$. Let $w \in H_b^{m,l}(X)$ denote the weak limit. We first want to show that $w = Z^{1/4} x^{1/4} \tilde{R}_+(0) f$.

- We first check that w solves the PDE $\tilde{P}(0)(Z^{-1/4} x^{-1/4} w) = f$. Indeed, for any m_0, l_0 ,

$$[0, \infty)_\sigma \times H_b^{m_0, l_0}(X) \ni (\sigma, u) \mapsto \tilde{P}(\sigma) u \in \mathcal{S}'(X) \quad (608)$$

is jointly continuous with respect to the strong topology on $H_b^{m_0, l_0}(X)$. (Besides being clear from the explicit formulas for \tilde{P} in §3, this follows from Proposition 2.4 and Proposition 3.6.) Since

$$Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa} \rightarrow Z^{-1/4} x^{-1/4} w \quad (609)$$

weakly in $H_b^{m, l-1/4}(X)$, this convergence occurs strongly in $H_b^{m-\epsilon, l-1/4-\epsilon}(X)$ for any $\epsilon > 0$, so (by the aforementioned joint continuity)

$$\tilde{P}(\sigma_{k_\kappa})(Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa}) \rightarrow \tilde{P}(0)(Z^{-1/4} x^{-1/4} w) \quad (610)$$

$\mathcal{S}'(X)$.

Moreover, it is not difficult to see that $\tilde{P}(\sigma_{k_\kappa})(Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa}) \rightarrow f$ in $\mathcal{S}'(X)$: indeed

$$\begin{aligned} \tilde{P}(\sigma_{k_\kappa})(Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa}) \\ = [\tilde{P}(\sigma_{k_\kappa}), Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4}] \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa} + f_{k_\kappa}, \end{aligned} \quad (611)$$

and $[\tilde{P}(\sigma), Z^{-1/4} x^{-1/4} (\sigma^2 + Zx)^{1/4}] \in S \text{Diff}_{\text{leC}}^{1, -5/4, -3, -5/4, -3}(X)$ satisfies

$$[\tilde{P}(\sigma), Z^{-1/4} x^{-1/4} (\sigma^2 + Zx)^{1/4}]|_{\sigma=0} = [\tilde{P}(0), 1] = 0, \quad (612)$$

so the boundedness of $\tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa}$ in $H_b^{m,l}(X)$ in some b-Sobolev space (as given by Proposition 6.9) and Proposition 2.4 show that $[\tilde{P}(\sigma_{k_\kappa}), Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4}] \tilde{R}_+(\sigma_{k_\kappa}) f_{k_\kappa} \rightarrow 0$ in $\mathcal{S}'(X)$.

Since $\mathcal{S}'(X)$ is Hausdorff, it follows that $\tilde{P}(0)(Z^{-1/4} x^{-1/4} w) = f$.

- Thus, we have

$$\tilde{P}(0)(Z^{-1/4} x^{-1/4} w) = f. \quad (613)$$

But, also,

$$\tilde{P}(0)(\tilde{R}_+(0) f) = f. \quad (614)$$

Set $l_0 = l - 1/4$, so that $f \in H_b^{m, l+5/4}(X) = H_b^{l_0+3/2}(X)$. In terms of $\ell_0 = 2l_0$, the inequalities $l < -1/2$ and $-1 < m + 2l$ become $\ell_0 < -3/2$ and $-3/2 < m + \ell$, so Proposition 4.3 applies:

$$\tilde{P}(0) : \{u \in H_b^{m, l-1/4}(X) : \tilde{P}(0)u \in H_b^{m, l+5/4}(X)\} \rightarrow H_b^{m, l+5/4}(X) \quad (615)$$

is invertible, and the inverse is $\tilde{R}_+(0)$. Thus, $u = \tilde{R}_+(0)f$ is the unique solution to $\tilde{P}(0)u = f$ in the domain of eq. (615). But $Z^{-1/4}x^{-1/4}w$ is in the codomain, and as we saw in eq. (614) solves this PDE. We conclude that

$$Z^{-1/4}x^{-1/4}w = \tilde{R}_+(0)f. \quad (616)$$

Via the compactness of the inclusion $H_b^{m, l} \hookrightarrow H_b^{m-\epsilon, l-\epsilon}$, we conclude that

$$(\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa})f_{k_\kappa} \rightarrow Z^{1/4}x^{1/4} \tilde{R}_+(0)f \quad (617)$$

strongly in $H_b^{m-\epsilon, l-\epsilon}$. This completes the proof of joint continuity.

By Proposition 6.9, for each $\Sigma > 0$ the set $\{(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)f\}_{\sigma \in [0, \Sigma]}$ is bounded in $H_b^{m, l}(X)$, and the result above shows that eq. (603) holds in the topology generated by Schwartz test functions (and, in fact, even the strong topologies $H_b^{m-\epsilon, l-\epsilon}$). It follows from the conjunction of these observations that eq. (603) holds with respect to the weak topology of $H_b^{m, l}(X)$ (since $\mathcal{S}(X)$ is dense in $H_b^{m, l}(X)$).

We now deduce the uniform continuity statement from the joint continuity statement. Suppose, to the contrary, that we have some $m, l \in \mathbb{R}$ with $l < -1/2$ and $-1 < m + 2l$ and some $\epsilon > 0$ such that $\{(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)\}_{\sigma \geq 0}$ is *not* continuous with respect to the uniform operator topology, generated by the norm

$$\|-\|_{\mathcal{L}(H_b^{m, l+5/4}(X), H_b^{m-\epsilon, l-\epsilon}(X))}. \quad (618)$$

We handle the case of a discontinuity at $\sigma = 0$, and the case of $\sigma > 0$ follows via a similar, easier argument. The discontinuity statement means that there exists some $\varepsilon > 0$ such that there exist sequences $\{\sigma_k\}_{k \in \mathbb{N}}$ with $\sigma_k \rightarrow 0^+$ and $\{f_k\}_{k \in \mathbb{N}} \subset H_b^{m, l+5/4}$ with $\|f_k\|_{H_b^{m, l+5/4}} \leq 1$ such that

$$\|(\sigma_k^2 + Zx)^{1/4} \tilde{R}_+(\sigma_k)f_k - (Zx)^{1/4} \tilde{R}_+(0)f_k\|_{H_b^{m-\epsilon, l-\epsilon}} \geq \varepsilon \quad (619)$$

for all k . By Banach–Alaoglu, we can choose these sequences such that there exists some

$$f_\infty \in H_b^{m, l+5/4} \quad (620)$$

such that $f_k \rightarrow f_\infty$ weakly, which implies strong convergence in $H_b^{m-\epsilon', l+5/4-\epsilon'}$ for any $\epsilon' > 0$. But then, by the joint continuity statement already proven, as long as ϵ' is sufficiently small such that $-1 < m + 2l - 3\epsilon'$, we have

$$\|(\sigma_k^2 + Zx)^{1/4} \tilde{R}_+(\sigma_k)f_k - (Zx)^{1/4} \tilde{R}_+(0)f_k\|_{H_b^{m-\epsilon'-\epsilon'', l-\epsilon'-\epsilon''}} \rightarrow 0 \quad (621)$$

as $k \rightarrow \infty$, for any $\epsilon'' > 0$. But we can take ϵ', ϵ'' sufficiently small such that $\epsilon' + \epsilon'' < \epsilon$, in which case eq. (621) contradicts eq. (619). \square

Via Sobolev embedding, the previous proposition already yields the following corollary on the continuity of the resolvent output at zero energy:

- for any $f \in H_b^{\infty, l+5/4}(X)$, $l < -1/2$, and for any $\chi \in C_c^\infty(X^\circ)$, $\chi \tilde{R}_+(\sigma)f \rightarrow \chi \tilde{R}_+(0)f$ in $C_c^\infty(X^\circ)$ as $\sigma \rightarrow 0^+$.

We will need to strengthen this result to apply to ∂_σ derivatives of the resolvent output. In order to handle the compositions that arise, we will use the following variant of the preceding proposition.

Proposition 6.11. *For any $m, l \in \mathbb{R}$ with $l < -1/2$ and $-1 < m + 2l$, the map $(\sigma, f) \mapsto (\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) f (\sigma^2 + Zx)^{1/4}$ defines a jointly continuous map*

$$[0, \Sigma) \times H_b^{m, l+1}(X) \rightarrow H_b^{m-\epsilon, l-\epsilon}(X), \quad (622)$$

for any $\epsilon > 0$ (with respect to the strong topologies on $H_b^{m, l+1}(X)$ and $H_b^{m-\epsilon, l-\epsilon}(X)$). Consequently, $\{(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) (\sigma^2 + Zx)^{1/4}\}_{\sigma \geq 0} \subset \mathcal{L}(H_b^{m, l+1}(X), H_b^{m-\epsilon, l-\epsilon}(X))$ is continuous with respect to the uniform operator topology. \blacksquare

Proof. We mimic the proof of Proposition 6.10. By the metrizable of $[0, \Sigma) \times H_b^{m, l+1}(X)$ and $H_b^{m-\epsilon, l-\epsilon}(X)$, joint continuity follows from the claim that

- given $f \in H_b^{m, l+1}(X)$ and $\sigma_\infty \in [0, \Sigma)$, for any and sequences $\{\sigma_k\}_{k \in \mathbb{N}}, \{f_k\}_{k \in \mathbb{N}} \subset H_b^{m, l+1}(X)$ with $\sigma_k \rightarrow \sigma_\infty$ and $f_k \rightarrow f$ as $k \rightarrow \infty$,
$$(\sigma_k^2 + Zx)^{1/4} \tilde{R}_+(\sigma_k) (\sigma_k^2 + Zx)^{1/4} f_k \rightarrow (\sigma_\infty^2 + Zx)^{1/4} \tilde{R}_+(\sigma_\infty) (\sigma_\infty^2 + Zx)^{1/4} f \quad (623)$$
strongly in $H_b^{m-\epsilon, l-\epsilon}(X)$.

It suffices to show that

- given any $\{\sigma_k\}_{k \in \mathbb{N}} \subset [0, \Sigma)$, $\{f_k\}_{k \in \mathbb{N}} \subset H_b^{m, l+1}(X)$ with $\sigma_k \rightarrow \sigma_\infty$ and $f_k \rightarrow f$ as $k \rightarrow \infty$, there exists a subsequence $\{k_\kappa\}_{\kappa \in \mathbb{N}} \subset \{k\}_{k \in \mathbb{N}}$ such that
$$(\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) (\sigma_{k_\kappa}^2 + Zx)^{1/4} f_{k_\kappa} \rightarrow (\sigma_\infty^2 + Zx)^{1/4} \tilde{R}_+(\sigma_\infty) (\sigma_\infty^2 + Zx)^{1/4} f \quad (624)$$
strongly in $H_b^{m-\epsilon, l-\epsilon}(X)$.

As before, we only consider the case of $\sigma_\infty = 0$, since the case $\sigma_\infty > 0$ follows by a similar but even easier argument.

By Proposition 6.9 and Banach–Alaoglu, we can find a subsequence $\{k_\kappa\}_{\kappa \in \mathbb{N}} \subset \{k\}_{k \in \mathbb{N}}$ such that $(\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) (\sigma_{k_\kappa}^2 + Zx)^{1/4} f_{k_\kappa}$ converges weakly in $H_b^{m, l}(X)$. Let $w \in H_b^{m, l}(X)$ denote the weak limit. We now show that $w = Z^{1/2} x^{1/4} \tilde{R}_+(0) x^{1/4} f$.

- We first check that w solves the PDE $\tilde{P}(0)(Z^{-1/4} x^{-1/4} w) = Z^{1/4} x^{1/4} f$. As above, we use that, for any m_0, l_0 ,

$$[0, \infty)_\sigma \times H_b^{m_0, l_0}(X) \ni (\sigma, u) \mapsto \tilde{P}(\sigma)u \in \mathcal{S}'(X) \quad (625)$$

is jointly continuous with respect to the strong topology on $H_b^{m_0, l_0}(X)$. Since

$$Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) (\sigma_{k_\kappa}^2 + Zx)^{1/4} f_{k_\kappa} \rightarrow Z^{-1/4} x^{-1/4} w \quad (626)$$

weakly in $H_b^{m, l-1/2}(X)$,

$$\tilde{P}(\sigma_{k_\kappa})(Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) (\sigma_{k_\kappa}^2 + Zx)^{1/4} f_{k_\kappa}) \rightarrow \tilde{P}(0)(Z^{-1/4} x^{-1/4} w) \quad (627)$$

$\mathcal{S}'(X)$. Moreover, it is not difficult to see that, as before, $\tilde{P}(\sigma_{k_\kappa})(Z^{-1/4} x^{-1/4} (\sigma_{k_\kappa}^2 + Zx)^{1/4} \tilde{R}_+(\sigma_{k_\kappa}) (\sigma_{k_\kappa}^2 + Zx)^{1/4} f_{k_\kappa}) \rightarrow Z^{1/4} x^{1/4} f$ in $\mathcal{S}'(X)$.

Since $\mathcal{S}'(X)$ is Hausdorff, it follows that $\tilde{P}(0)(Z^{-1/4} x^{-1/4} w) = f$.

- Thus, we have

$$\begin{aligned} \tilde{P}(0)(Z^{-1/4} x^{-1/4} w) &= Z^{1/4} x^{1/4} f \\ \tilde{P}(0)(\tilde{R}_+(0)(Z^{1/4} x^{1/4} f)) &= Z^{1/4} x^{1/4} f. \end{aligned} \quad (628)$$

Set $l_0 = l - 1/4$, so that $x^{1/4} f \in H_b^{m, l+5/4}(X) = H_b^{l_0+3/2}(X)$. By Proposition 4.3,

$$\begin{aligned} \tilde{P}(0) : \{u \in H_b^{m, l_0}(X) : \tilde{P}(0)u \in H_b^{m, l_0+3/2}(X)\} &\rightarrow H_b^{m, l_0+3/2}(X) \\ : \{u \in H_b^{m, l-1/4}(X) : \tilde{P}(0)u \in H_b^{m, l+5/4}(X)\} &\rightarrow H_b^{m, l+5/4}(X) \end{aligned} \quad (629)$$

is invertible, and the inverse is $\tilde{R}_+(0)$. Thus, $u = \tilde{R}_+(0)Z^{1/4}x^{1/4}f$ is the unique solution to $\tilde{P}(0)u = Z^{1/4}x^{1/4}f$ in the domain of eq. (629). But $Z^{-1/4}x^{-1/4}w$ is in the codomain, and as we saw in eq. (628) solves this PDE. We conclude that

$$Z^{-1/4}x^{-1/4}w = \tilde{R}_+(0)Z^{1/4}x^{1/4}f. \quad (630)$$

Via the compactness of the inclusion $H_b^{m,l} \hookrightarrow H_b^{m-\epsilon,l-\epsilon}$, we conclude that

$$(\sigma_{k_\kappa}^2 + Zx)^{1/4}\tilde{R}_+(\sigma_{k_\kappa})(\sigma_{k_\kappa}^2 + Zx)^{1/4}f_{k_\kappa} \rightarrow Z^{1/4}x^{1/4}\tilde{R}_+(0)Z^{1/4}x^{1/4}f \quad (631)$$

strongly in $H_b^{m-\epsilon,l-\epsilon}$.

Uniform continuity follows as in the proof of Proposition 6.10. \square

Proposition 6.12. *Fix $\psi \in C_c^\infty(\mathbb{R})$ supported sufficiently close to $(-\infty, \Sigma^2]$ and $\chi \in C_c^\infty(X_{\text{res}}^{\text{sp}})$ supported away from bf. Then, for each $k, K \in \mathbb{N}$ with $k + K > 0$,*

$$\{\psi(E)(E\partial_E)^k(\chi x\partial_E)^K \tilde{P}(E^{1/2})\}_{E \geq 0} \in S\text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X) \subseteq S\text{Diff}_{\text{b,leC}}^{1,-1,-3}(X). \quad (632)$$

■

Proof. For $\sigma \in [0, \Sigma]$, $\tilde{P}(\sigma) = K(\sigma) + C$ for $K(\sigma) \in S\text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X)$ containing the σ -dependent part of $\tilde{P}(\sigma)$ and $C \in S\text{Diff}_{\text{scb}}^{2,0,-2}(X)$ constant in σ . Thus, for $k, K \in \mathbb{N}$ which are not both zero,

$$\{\psi(E)(E\partial_E)^k(\chi x\partial_E)^K \tilde{P}(\sigma)\}_{\sigma \geq 0} \in (E\partial_E)^k(\chi x\partial_E)^K S\text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X). \quad (633)$$

Since $E\partial_E$ lifts to a b-vector field on $X_{\text{res}}^{\text{sp}}$ and $\chi x\partial_E$ lifts to a smooth vector field which is b- at tf (and identically zero near bf),

$$(\sigma\partial_\sigma)^k(\chi x\partial_E)^K S\text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X) \subseteq S\text{Diff}_{\text{leC}}^{1,0,-2,-1,-3}(X), \quad (634)$$

so eq. (632) follows. \square

Proposition 6.13. *For each $m, l \in \mathbb{R}$ with $m + l > -1/2 > l$ and for each $f \in H_b^{m,l+1}(X)$, the map*

$$\tilde{R}_+(\bullet)f : (0, \Sigma) \ni \sigma \mapsto \tilde{R}_+(\sigma)f \in \mathcal{S}'(X) \quad (635)$$

is smooth as a map $(0, \Sigma) \rightarrow \mathcal{S}'(X)$. Thus, $\partial_\sigma^k(\tilde{R}_+(\sigma)\tilde{f}(-; \sigma)) : \mathbb{R}_\sigma^+ \rightarrow \mathcal{S}'(X)$ is well-defined for $k \in \mathbb{N}$ and $\tilde{f} \in C^\infty((0, \Sigma); H_b^{m,l+1}(X))$ and is given by

$$\left[\sum_{\kappa=0}^k \binom{k}{\kappa} \partial_\sigma^\kappa \left(\tilde{R}_+(\sigma) (\partial_{\sigma_0}^{k-\kappa} \tilde{f}(-; \sigma_0)) \right) \right] \Big|_{\sigma=\sigma_0}. \quad (636)$$

Moreover, for $k \in \mathbb{N}$, it is the case that, for each $\sigma > 0$, the map

$$\partial_\sigma^k \tilde{R}_+(\sigma) = \partial_{\sigma_0}^k \tilde{R}_+(\sigma_0)|_{\sigma_0=\sigma} : \bigcup_{\substack{m, l \in \mathbb{R} \\ k-m-1/2 < l < -1/2}} H_b^{m,l+1}(X) \rightarrow \mathcal{S}'(X) \quad (637)$$

$$: f \mapsto \partial_\sigma^k(\tilde{R}_+(\sigma)f) \quad (638)$$

satisfies $\partial_\sigma^k \tilde{R}_+(\sigma) \in \mathcal{L}(H_b^{m,l+1}(X), H_b^{m-k,l}(X))$. For each $k \in \mathbb{N}^+$, the identity

$$\partial_\sigma^k \tilde{R}_+(\sigma) = \left[\sum_{\{k_i\}_{i=1}^I \in \mathcal{I}_k} c_{\{k_i\}_{i=1}^I} \prod_{i=1}^I (\tilde{R}_+(\sigma) \partial_\sigma^{k_i} \tilde{P}(\sigma)) \right] \tilde{R}_+(\sigma) \quad (639)$$

holds for some $c_{\{k_i\}_{i=1}^I} \in \mathbb{Z}$, where \mathcal{I}_k is the set of finite sequences of positive integers summing to k . \square

Proof. The proposition holds for the usual conjugated resolvent family

$$\tilde{R}_+ = e^{-i\sigma/x+i\Phi} \tilde{R}_+ e^{+i\sigma/x-i\Phi} \quad (640)$$

in place of \tilde{R}_+ and $\tilde{P} = e^{-i\sigma/x} P e^{+i\sigma/x}$ in place of \tilde{P} — see [Hin22, Lemma 2.10], [HHV21, Proposition 2.10]. Since multiplication by $\partial_\sigma^\kappa \exp(\pm i(\Phi_0 - \Phi))$ acts boundedly (but not uniformly so as $\sigma \rightarrow 0^+$) on the b-Sobolev spaces for each $\kappa \in \mathbb{N}$ — see eq. (67) — the mapping properties asserted in the proposition follow from those for the usual conjugated resolvent family. The specific formula eq. (639) follows via inductively applying $\partial_\sigma \tilde{R}_+(\sigma) f = -\tilde{R}_+(\sigma) (\partial_\sigma \tilde{P}(\sigma)) \tilde{R}_+(\sigma) f$. \square

Remark 9. By the conjunction of

- (1) $\tilde{R}_+(\sigma) : H_b^{m,l+1}(X) \rightarrow H_b^{m,l}(X)$ holding for all $m, l \in \mathbb{R}$ with $l < -1/2 < m + l$ and $\sigma > 0$,
and
- (2) $\partial_\sigma^k \tilde{P}(\sigma) : H_b^{m,l}(X) \rightarrow H_b^{m-1,l+1}(X)$ holding for all $m, l \in \mathbb{R}$ and $\sigma > 0$,

it is the case that

$$\tilde{R}_+(\sigma) \partial_\sigma^k \tilde{P}(\sigma) : H_b^{m,l}(X) \rightarrow H_b^{m-1,l}(X) \quad (641)$$

for each $\sigma > 0$, $k \in \mathbb{N}$, and $m, l \in \mathbb{R}$ satisfying $l < -1/2 < m - 1 + l$.

Consequently, for $\{i_i\}_{i=1}^I \in \mathcal{I}_k$,

$$\prod_{i=1}^I [\tilde{R}_+(\sigma) \partial_\sigma^{k_i} \tilde{P}(\sigma)] : H_b^{m,l}(X) \rightarrow H_b^{m-I,l}(X) \quad (642)$$

is a well-defined composition whenever $l < -1/2 < m + l - I$. The right-hand side of eq. (639) is therefore a well-defined map $H_b^{m,l+1}(X) \rightarrow H_b^{m-k,l}(X)$ whenever $l < -1/2 < m + l - k$.

The identity eq. (639) should therefore be read as stating that both sides agree as maps $H_b^{m,l+1}(X) \rightarrow H_b^{m-k,l}(X)$. \blacksquare

Proposition 6.14. *For each $k, K \in \mathbb{N}$ and $l < -1/2$, there exists an $m_0(l, k, K) \in \mathbb{R}$ such that for all $m > m_0$ and $\chi \in C_c^\infty(X_{\text{res}}^{\text{sp}})$ supported away from bf,*

$$\begin{aligned} & \{ (E \partial_E)^k (\chi x \partial_E)^K ((E + Zx)^{1/4} \tilde{R}_+(E^{1/2})) \}_{E \in (0, \Sigma^2)} \\ & \in L^\infty \cap C^0((0, \Sigma); \mathcal{L}(H_b^{m,l+5/4}(X), H_b^{m-k-K-\epsilon, l-\epsilon}(X))) \end{aligned} \quad (643)$$

for all $\epsilon > 0$. \blacksquare

Proof. We explicitly consider the $K = 0$ case, and the $K \in \mathbb{N}^+$ case is similar (if a bit messier).

So, we want to prove that

$$\{ (\sigma \partial_\sigma)^k ((\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)) \}_{\sigma \in (0, \Sigma)} \in L^\infty \cap C^0((0, \Sigma); \mathcal{L}(H_b^{m,l+5/4}(X), H_b^{m-k-\epsilon, l-\epsilon}(X))). \quad (644)$$

By eq. (639),

$$\begin{aligned} (\sigma \partial_\sigma)^k ((\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)) &= \sum_{j=0}^k \binom{k}{j} (\sigma^2 + Zx)^{-1/4} ((\sigma \partial_\sigma)^j (\sigma^2 + Zx)^{1/4}) \\ & \left[\sum_{\{k_i\}_{i=1}^I \in \mathcal{I}_{k-j}} c_{\{k_i\}_{i=1}^I} \prod_{i=1}^I (\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) (\sigma \partial_\sigma)^{k_i} \tilde{P}(\sigma) (\sigma^2 + Zx)^{-1/4} \right] [(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma)] \end{aligned} \quad (645)$$

for all $\sigma > 0$. By Proposition 6.12, for any $k \in \mathbb{N}^+$ we can write

$$(\sigma^2 + Zx)^{-1/4} ((\sigma \partial_\sigma)^k \tilde{P}(\sigma)) (\sigma^2 + Zx)^{-1/4} \in L^\infty \cap C^0([0, \Sigma]; S \text{Diff}_b^{1, -1+\epsilon}(X)) \quad (646)$$

for any $\epsilon > 0$. On the other hand, by Proposition 6.11, for all $\epsilon > 0$,

$$(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) (\sigma^2 + Zx)^{1/4} \in C^0([0, \Sigma]; \mathcal{L}(H_b^{m-1, l+1-\epsilon/2}(X), H_b^{m-1-\epsilon, l-\epsilon}(X))) \quad (647)$$

for any $m, l \in \mathbb{R}$ satisfying $l < -1/2$ and $m > -2l$.

Combining these two observations, for any $k \in \mathbb{N}^+$ and $\epsilon > 0$ we have that

$$\begin{aligned} (\sigma^2 + Zx)^{\frac{1}{4}} \tilde{R}_+(\sigma) ((\sigma \partial_\sigma)^k \tilde{P}(\sigma)) (\sigma^2 + Zx)^{-\frac{1}{4}} \\ = (\sigma^2 + Zx)^{\frac{1}{4}} \tilde{R}_+(\sigma) (\sigma^2 + Zx)^{\frac{1}{4}} (\sigma^2 + Zx)^{-\frac{1}{4}} ((\sigma \partial_\sigma)^k \tilde{P}(\sigma)) (\sigma^2 + Zx)^{-\frac{1}{4}} \end{aligned} \quad (648)$$

lies in

$$C^0([0, \Sigma]; \mathcal{L}(H_b^{m-1, l+1-\frac{\epsilon}{2}}, H_b^{m-1-\epsilon, l-\epsilon}) S \text{Diff}_b^{1, -1+\frac{\epsilon}{2}}(X)) \subseteq C^0([0, \Sigma]; \mathcal{L}(H_b^{m, l}, H_b^{m-1-\epsilon, l-\epsilon})) \quad (649)$$

for any $m, l \in \mathbb{R}$ satisfying $l < -1/2$ and $m > -2l$. Therefore, by Proposition 6.10,

$$\begin{aligned} \left[\sum_{\{k_i\}_{i=1}^I \in \mathcal{I}_{k-j}} c_{\{k_i\}_{i=1}^I} \prod_{i=1}^I (\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) (\sigma \partial_\sigma)^{k_i} \tilde{P}(\sigma) (\sigma^2 + Zx)^{-1/4} \right] \left[(\sigma^2 + Zx)^{1/4} \tilde{R}_+(\sigma) \right] \\ \in C^0([0, \Sigma]_\sigma; \mathcal{L}(H_b^{m, l+5/4}, H_b^{m-k-\epsilon, l-\epsilon})) \end{aligned} \quad (650)$$

for any $\epsilon > 0$ and for any $m, l \in \mathbb{R}$ satisfying $l < -1/2$ and $-2l + k - 1 < m$.

Since $(\sigma^2 + Zx)^{-1/4} (\sigma \partial_\sigma)^j (\sigma^2 + Zx)^{1/4} \in C^0([0, \infty)_\sigma; S_b^\epsilon(X))$ (as checked in Lemma 2.1), we conclude eq. (644), eq. (643). \square

The preceding results amount to:

Proposition 6.15. *Given $\tilde{f} \in C^\infty(X_{\text{res}}^{\text{sp}})$ vanishing rapidly at tf, bf (i.e. $\tilde{f} \in \mathcal{A}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}})$), set $u_{0,+}(\sigma) = \tilde{R}_+(\sigma) \tilde{f}(-; \sigma)$ for all $\sigma \in [0, \Sigma]$. Defining $u_{0,+}(\sigma) = x^{-(n-1)/2} (\sigma^2 + Zx)^{1/4} u_{00,+}(\sigma)$, $u_{0,+} = \{u_{0,+}(-; \sigma)\}_{\sigma \in [0, \Sigma]} \in \mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\})$.*

Moreover, the mapping $\mathcal{A}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}}) \ni \tilde{f} \mapsto u_{0,+} \in \mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\})$ is continuous. \blacksquare

Proof. Fix $\chi \in C_c^\infty(X_{\text{res}}^{\text{sp}})$ supported away from bf and nonvanishing near zf.

- We first show that $u_{0,+} = \{u_{0,+}(-; \sigma)\}_{\sigma \in (0, \Sigma)} \in \mathcal{A}_{\text{loc}}^{0-, 0-, 0}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\})$.

For any $k, K \in \mathbb{N}$, $(E \partial_E)^k (\chi x \partial_E)^K \tilde{f} \in \mathcal{A}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}})$. Thus, Proposition 6.13 and Proposition 6.14 imply that, for all $m \in \mathbb{R}$, $l < -1/2$, $K \in \mathbb{N}$,

$$\begin{aligned} \{(\chi x \partial_E)^K u_{0,+}(-; E^{1/2})\}_{E \in (0, \Sigma^2)} \\ = \left\{ \left(\chi x \frac{\partial}{\partial E} \right)^K \left[x^{-(n-1)/2} (E + Zx)^{1/4} \tilde{R}_+(E^{1/2}) \tilde{f}(-; E^{1/2}) \right] \right\}_{E \in (0, \Sigma^2)} \end{aligned} \quad (651)$$

is in $\mathcal{A}_{\text{loc}}^0([0, \Sigma^2]_E; H_b^{m, l-(n-1)/2}(X))$.

By Proposition 2.19, we deduce that $(\chi x \partial_E)^K u_{0,+} \in \mathcal{A}_{\text{loc}}^{0-, 0-, 0}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\})$. Moreover, Proposition 6.14 shows that the map

$$\mathcal{A}^{\infty, \infty, (0,0)}(X) \ni \tilde{f} \mapsto (\chi x \partial_E)^K u_{0,+}(-; E^{1/2}) \in \mathcal{A}_{\text{loc}}^0([0, \Sigma^2]_E; H_b^{m, l-(n-1)/2}(X)) \quad (652)$$

is continuous, for each $K \in \mathbb{N}$, and the inclusion in Proposition 2.19 is continuous, so the map $\mathcal{A}^{\infty, \infty, (0,0)}(X) \ni \tilde{f} \mapsto (\chi x \partial_E)^K u_{0,+}(-; E^{1/2}) \in \mathcal{A}_{\text{loc}}^{0-, 0-, 0}(X_{\text{res}}^{\text{sp}})$ is continuous as well.

- It now follows from Proposition 2.20 that $\{u_{0,+}(-; \sigma)\}_{\sigma > 0}$ defines an element of the space $\mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\})$. By the remark following Proposition 2.20,

$$\mathcal{A}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}}) \ni \tilde{f} \mapsto u_{0,+} \in \mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{res}}^{\text{sp}} \cap \{\sigma < \Sigma\}) \quad (653)$$

is continuous.

- Finally, we observe from Proposition 6.12 that $\{u_{0,+}(-; \sigma)\}_{\sigma > 0}$, considered as an element of $\mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{res}}^{\text{sp}})$, restricts to $x^{-(n-1)/2} (Zx)^{1/4} \tilde{R}_+(0) \tilde{f}(-; 0)$ at zf.

\square

6.3. Asymptotics at bf, tf. For this subsection, we do not assume that P is the spectral family of an attractive Coulomb-like Schrödinger operator, only that P satisfies the hypotheses in §1. Let \mathcal{E} denote the index set

$$\mathcal{E} = \{(k, \kappa) \in \mathbb{N} \times \mathbb{N} : \kappa \leq \lfloor k/2 \rfloor\}. \quad (654)$$

Suppose that P_1 is classical to order $\beta_1 > 0$ and P_2 is classical to order (β_2, β_3) . Then, by Proposition 3.8,

$$N(\tilde{P}) - \tilde{P} \in \text{Diff}_{b, \text{leC}}^{2, -2, -4}(X) + S \text{Diff}_{b, \text{leC}}^{2, -1 - \delta_1, -3 - 2\delta_0}(X) \quad (655)$$

holds for $\delta_1 = \min\{1 + \beta_1, \beta_2, 1/2 + \beta_3\}$ and $\delta_0 = \min\{1 + \beta_1, \beta_2, \beta_3\}$.

Proposition 6.16. *Suppose that P_1 is classical to β_1 th order and P_2 is classical to order (β_2, β_3) , and suppose that $\delta_0 \notin \mathbb{N}$. If $u_0 \in \mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)}(X_{\text{loc}}^{\text{sp}})$ and $\tilde{f} \in \mathcal{A}_{\text{loc}}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}})$, then, setting*

$$u_{00}(-; \sigma) = x^{(n-1)/2}(\sigma^2 + Zx)^{-1/4}u_0(-; \sigma), \quad (656)$$

if $\tilde{P}u_{00} = \tilde{f}$, then

$$u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0), \delta_1), 2\delta_0-, (0,0)}(X_{\text{res}}^{\text{sp}}), \quad (657)$$

and $(\mathcal{A}_{\text{loc}}^{0-, 0-, (0,0)} \times \mathcal{A}_{\text{loc}}^{\infty, \infty, (0,0)})|_{\tilde{P}u_{00}=\tilde{f}} \ni (u_0, \tilde{f}) \mapsto u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)} + \mathcal{A}_{\text{loc}}^{((0,0), \delta_1), 2\delta_0-, (0,0)}$ is continuous. In particular, if P is fully classical, then

$$u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}). \quad (658)$$

■

Proof. We will prove via induction on α, β that, for all $\alpha \in (-\infty, \delta_1]$ and $\beta \in (-\infty, \delta_0)$,

$$u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0), \alpha), 2\beta, (0,0)}(X_{\text{loc}}^{\text{sp}}). \quad (659)$$

The $\alpha, \beta < 0$ case of eq. (659) is the hypothesis of the proposition.

Consider α, β such that eq. (659) holds. From $\tilde{P}u_{00} = \tilde{f}$, writing $\tilde{P} = N(\tilde{P}) + (\tilde{P} - N(\tilde{P}))$, we get

$$\begin{aligned} N(\tilde{P})u_{00} &= \tilde{f} - (\tilde{P} - N(\tilde{P}))u_{00} \in \mathcal{A}_{\text{loc}}^{\infty, \infty, (0,0)}(X_{\text{res}}^{\text{sp}}) \\ &+ (\text{Diff}_{b, \text{leC}}^{2, -2, -4}(X) + S \text{Diff}_{b, \text{leC}}^{2, -1 - \delta_1, -3 - 2\delta_0}(X))x^{(n-1)/2}\mathcal{A}_{\text{loc}}^{((0,0), \alpha), 2\beta - 1/2, (0,0)}(X_{\text{res}}^{\text{sp}}) \\ &+ (\text{Diff}_{b, \text{leC}}^{2, -2, -4}(X) + S \text{Diff}_{b, \text{leC}}^{2, -1 - \delta_1, -3 - 2\delta_0}(X))x^{(n-1)/2}\mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E} - 1/2, (0,0)}(X_{\text{res}}^{\text{sp}}). \end{aligned} \quad (660)$$

The set on the right-hand side is $x^{(n-1)/2}$ times

$$\begin{aligned} &\mathcal{A}_{\text{loc}}^{(2,0), \mathcal{E} + 7/2, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((2,0), 2 + \alpha), 2\beta + 7/2, (0,0)}(X_{\text{res}}^{\text{sp}}) \\ &+ \mathcal{A}_{\text{loc}}^{\min\{1 + \delta_1, 1 + \alpha + \delta_1\}, \min\{5/2 + 2\delta_0, 5/2 + 2\delta_0 + 2\alpha\}, (0,0)}(X_{\text{res}}^{\text{sp}}). \end{aligned} \quad (661)$$

For α sufficiently close to zero or positive, we can apply Proposition 6.17 below to conclude that

$$\begin{aligned} u_0 &\in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)} + \mathcal{A}_{\text{loc}}^{((0,0), 1 + \alpha), 1 + 2\beta, (0,0)} + \mathcal{A}_{\text{loc}}^{((0,0), \min\{\delta_1, \alpha + \delta_1\}), \min\{2\delta_0, 2\delta_0 + 2\beta\} -, (0,0)} \\ &\subseteq \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0), \min\{1 + \alpha, \delta_1, \alpha + \delta_1\}), \min\{1 + 2\beta, 2\delta_0, 2\delta_0 + 2\beta\} -, (0,0)}(X_{\text{res}}^{\text{sp}}). \end{aligned} \quad (662)$$

If $\delta_1 \leq \alpha + 1$, $\alpha + \delta_1$ and $\delta_0 \leq 2\beta + 1$, $2\beta + 2\delta_0$, then

$$\mathcal{A}_{\text{loc}}^{((0,0), \min\{\alpha + 1, \delta_1, \alpha + \delta_1\}), \min\{1 + 2\beta, 2\delta_0, 2\delta_0 + 2\beta\} -, (0,0)}(X_{\text{res}}^{\text{sp}}) = \mathcal{A}_{\text{loc}}^{((0,0), \delta_1), 2\delta_0 -, (0,0)}(X_{\text{res}}^{\text{sp}}), \quad (663)$$

so we have concluded that eq. (657) holds. Otherwise,

$$u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0), \alpha + \varepsilon_1), 2\beta + 2\varepsilon_0 -, (0,0)}(X_{\text{res}}^{\text{sp}}) \quad (664)$$

for $\varepsilon_1 = \min\{1, \delta_1\}$ and $\varepsilon_0 = \min\{1/2, \delta_0\}$. The claim therefore follows by induction.

The continuity clause of the proposition can be proven using the same argument, keeping track of topologies. \square

Proposition 6.17. *Suppose that P satisfies the minimal hypotheses of §1.*

Suppose further that $u \in \mathcal{A}_{\text{loc}}^{-\infty, -\infty, (0,0)}(X_{\text{res}}^{\text{sp}}) = \mathcal{A}_{\text{loc}}^{-\infty, (0,0)}([0, \infty)_E \times X)$ satisfies $\tilde{N}(\tilde{P})u = f$ for some

$$f \in x^{(n-1)/2} \mathcal{A}_{\text{loc}}^{(2,0), \mathcal{E}+7/2, (0,0)}(X_{\text{res}}^{\text{sp}}) + x^{(n-1)/2} \mathcal{A}_{\text{loc}}^{((2,0), \alpha+1), \beta+5/2, (0,0)}(X_{\text{res}}^{\text{sp}}) \quad (665)$$

for $\alpha, \beta \in \mathbb{R}^+$, $\beta \notin 2\mathbb{N}$. Then, setting $u_0 = x^{-(n-1)/2}(\sigma^2 + Zx)^{1/4}u$,

$$u_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0), \alpha), \beta, (0,0)}(X_{\text{res}}^{\text{sp}}) \quad (666)$$

holds. ■

Proof. Now letting $\tilde{N}(\tilde{P}) = x^{-(n-1)/2}(\sigma^2 + Zx)^{1/4} \tilde{N}(\tilde{P}) x^{(n-1)/2}(\sigma^2 + Zx)^{-1/4}$, $\tilde{N}(\tilde{P})u_0 = f_0$ for $u_0 = x^{-(n-1)/2}(\sigma^2 + Zx)^{1/4}u$ and $f_0 = x^{-(n-1)/2}(\sigma^2 + Zx)^{1/4}f$. Below, it will be slightly more convenient to work with $f_1 = x^{-1}(\sigma^2 + Zx)^{-1/2}f_0$. Equation (665) yields

$$f_1 \in x(\sigma^2 + Zx)^{-1/2} \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((1,0), \alpha), \beta, (0,0)}(X_{\text{res}}^{\text{sp}}). \quad (667)$$

In order to prove the proposition, it suffices to restrict attention to $\hat{X} = [0, \bar{x}]_x \times \partial X_y$. By eq. (534) (with $k = -1/4$ and $l = -1/2$), we have

$$\tilde{N}(\tilde{P}) = 2ix^2(\sigma^2 + Zx)^{1/2}\partial_x. \quad (668)$$

Thus, integrating $\tilde{N}(\tilde{P})u_0 = f_0$, we get

$$u_0(x, y; \sigma) = c(y; \sigma) + \frac{i}{2} \int_x^{\bar{x}} x_0^{-1} f_1(x_0, y; \sigma) dx_0 \quad (669)$$

for some $c(y; \sigma) \in \mathbb{C}$, for each $\sigma \geq 0$. Since $u_0 \in \mathcal{A}_{\text{loc}}^{-\infty, (0,0)}([0, \infty)_E \times \hat{X})$, and since the same applies to the integral in eq. (669) (cut-off near ∂X), we deduce that $c \in \mathcal{A}_{\text{loc}}^{-\infty, (0,0)}([0, \infty)_E \times \hat{X})$. Since $c(y, \sigma)$ does not depend on x , this implies that $c \in C^\infty([0, \infty)_E \times \partial X)$.

Equation (666) then follows from the mapping properties of the integral in eq. (669), which we record in Corollary 6.19 and Proposition 6.20 below. □

For the following proposition, we use $\hat{X}_{\text{res}}^{\text{sp}}$, defined as $X_{\text{res}}^{\text{sp}}$ with $\hat{X} = [0, \bar{x}] \times \partial X$ in place of X .

Proposition 6.18. *Let $\mathcal{F} \subset \mathbb{N} \times \mathbb{N}$ be some index set containing $(0, 0)$.*

If $g \in x(\sigma^2 + Zx)^{-1/2} \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{F}, (0,0)}(\hat{X}_{\text{res}}^{\text{sp}}) = \mathcal{A}_{\text{loc}}^{(1,0), \mathcal{F}+1, (0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$, then

$$I = \int_x^{\bar{x}} x_0^{-1} g(x_0, y; \sigma) dx_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{F}_+, (0,0)}(\hat{X}_{\text{res}}^{\text{sp}}), \quad (670)$$

where $\mathcal{F}_+ = \mathcal{F} \cup \{(k+1, \kappa+1), (k+2, \kappa+1) : (k, \kappa) \in \mathcal{F}, k \equiv 1 \pmod{2}\}$. ■

Proof. It suffices to consider the case when g is supported in $[0, \bar{x}/2]_x$. Let us write $G(x, y; E) = x^{-1}(\sigma^2 + Zx)^{1/2}g$, so

$$I = \int_x^{\bar{x}} (E + Zx_0)^{-1/2} G(x_0, y; E) dx_0 = 2 \int_{x^{1/2}}^{\bar{x}^{1/2}} \left(\frac{\rho_0^2}{E + Z\rho_0^2} \right)^{1/2} G(\rho_0^2, y; E) d\rho_0. \quad (671)$$

This is evidently smooth away from tf.

By the polyhomogeneity of G at $\text{zf} \cap \text{tf}$, there exists a $G_0 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{F}}([0, \infty)_{\hat{E}} \times [0, \infty)_\rho \times \partial X)$ such that $G_0(\rho, y; \hat{E}) = G(x, y; E)$ when $x = \rho^2$ and $E = \hat{E}x$. That is, $G(x, y; E) = G_0(x^{1/2}, y; Ex^{-1})$.

Then, as a function of ρ, ρ_0, y , and $\hat{E} = E/x$, $G(\rho_0^2, y; E) = G_0(\rho_0, y; E\rho_0^{-2}) = G_0(\rho_0, y; \hat{E}\rho^2\rho_0^{-2})$. We can therefore write

$$I = I(\rho, y; \hat{E}) = 2 \int_\rho^{\bar{x}^{1/2}} \frac{1}{(\hat{E} + Z)^{1/2}} G_0\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 = \int_\rho^{\bar{x}^{1/2}} G_1\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 \quad (672)$$

for $G_1 \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{F}}([0, \infty)_{\hat{E}} \times [0, \infty)_\rho \times \partial X)$ defined by $G_1(\rho, y; \hat{E}) = 2(\hat{E} + Z)^{-1/2}G_0(\rho, y; \hat{E})$.

We now expand $G_1(\rho, y, ; \hat{E})$ polyhomogeneously around $\rho = 0$: there exist $G_1^{(k, \kappa)} \in C^\infty([0, \infty)_{\hat{E}} \times \partial X)$ and $F_1^{(k)} \in \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{F}|_{\geq k-k}}([0, \infty)_{\hat{E}} \times [0, \infty)_\rho \times \partial X)$ for $k \in \mathbb{N}$ such that

$$G_1\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) = \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \rho_0^k \log^\varkappa(\rho_0) G_1^{(k, \varkappa)}\left(y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) + \rho_0^{K+1} F_1^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right), \quad (673)$$

where $\mathcal{F}|_{\geq k-k} = \{(k' - k, \kappa) \in \mathcal{E} : \Re k' \geq k\}$, $\varkappa_k = \max\{\varkappa : (k, \varkappa) \in \mathcal{F}\}$, and $G_1^{(k, \varkappa)}$ can be nonzero only if $(k, \varkappa) \in \mathcal{F}$.

Expand $G_1^{(k, \varkappa)}$ in the second slot: there exist $G_1^{(j, k, \varkappa)} \in C^\infty(\partial X)$ and $F_2^{(j, k, \varkappa)} \in C^\infty([0, \infty)_{\hat{E}} \times \partial X)$ such that

$$G_1^{(k, \varkappa)}\left(y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) = \sum_{j=0}^J \rho^{2j} \rho_0^{-2j} \hat{E}^j G_1^{(j, k, \varkappa)}(y) + \rho^{2J+2} \rho_0^{-2J-2} \hat{E}^{J+1} F_2^{(J+1, k, \varkappa)}\left(y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right). \quad (674)$$

Let $F_{k, \varkappa} = \int \rho^k \log^\varkappa(\rho) d\rho$ for $k \in \mathbb{Z}$ and $\varkappa \in \mathbb{N}$, with the additive constant chosen for convenience. Then, if $k \neq -1$, we can take

$$F_{k, \varkappa}(\rho) = \rho^{k+1} \sum_{\kappa=0}^{\varkappa} c_{k, \varkappa, \kappa} \log^\kappa(\rho) \quad (675)$$

for some $c_{k, \varkappa, \kappa} \in \mathbb{R}$. If instead $k = -1$, we have $F_{-1, \varkappa}(\rho) = (\varkappa + 1)^{-1} \log^{\varkappa+1}(\rho)$. Integrating G_1 , we write

$$\begin{aligned} I(\rho, y; \hat{E}) &= \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{j=0}^J \rho^{2j} \hat{E}^j G_1^{(j, k, \varkappa)}(y) \int_\rho^{\bar{x}^{1/2}} \rho_0^{k-2j} \log^\varkappa(\rho_0) d\rho_0 \\ &\quad + \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \rho^{2J+2} \hat{E}^{J+1} \int_\rho^{\bar{x}^{1/2}} \rho_0^{k-2J-2} \log^\varkappa(\rho_0) F_2^{(J+1, k, \varkappa)}\left(y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 \\ &\quad + \int_\rho^{\bar{x}^{1/2}} \rho_0^{K+1} F_1^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0. \end{aligned} \quad (676)$$

We decompose this as $I = I_1 + I_2 + I_3$ for

$$\begin{aligned} I_1 &= \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{j=0}^J \rho^{2j} \hat{E}^j G_1^{(j, k, \varkappa)}(y) \int_\rho^{\bar{x}^{1/2}} \rho_0^{k-2j} \log^\varkappa(\rho_0) d\rho_0 \\ &= \tilde{I}_1 + \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{\substack{j=0 \\ k-2j \neq -1}}^J \rho^{2j} \hat{E}^j G_1^{(j, k, \varkappa)}(y) [F_{k-2j, \varkappa}(\bar{x}^{1/2}) - F_{k-2j, \varkappa}(\rho)] \\ &= \tilde{I}_1 + \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{\substack{j=0 \\ k-2j \neq -1}}^J \hat{E}^j G_1^{(j, k, \varkappa)}(y) \sum_{\kappa=0}^{\varkappa} c_{k-2j, \varkappa, \kappa} [\rho^{2j} \bar{x}^{k-2j+1} \log^\kappa(\bar{x}) - \rho^{k+1} \log^\kappa(\rho)], \end{aligned} \quad (677)$$

where

$$\begin{aligned} \tilde{I}_1 &= \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{\substack{j=0 \\ 2j=k+1}}^J \rho^{2j} \hat{E}^j G_1^{(j, k, \varkappa)}(y) [F_{-1, \varkappa}(\bar{x}^{1/2}) - F_{-1, \varkappa}(\rho)] \\ &= \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \sum_{\substack{j=0 \\ 2j=k+1}}^J \rho^{k+1} \hat{E}^j G_1^{(j, k, \varkappa)}(y) \frac{1}{\varkappa+1} [\log^{\varkappa+1}(\bar{x}^{1/2}) - \log^{\varkappa+1}(\rho)], \end{aligned} \quad (678)$$

and

$$I_2 = \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \rho^{2J+2} \hat{E}^{J+1} \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{k-2J-2} \log^{\varkappa}(\rho_0) F_2^{(J+1,k,\varkappa)}\left(y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 \quad (679)$$

$$I_3 = \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{K+1} F_1^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0. \quad (680)$$

Since $(0, 0) \in \mathcal{F}$,

$$\sum_{k=0}^K \sum_{\substack{j=0 \\ k-2j \neq -1}}^J \hat{E}^j G_1^{(j,k,\varkappa)}(y) \sum_{\kappa=0}^{\varkappa} c_{k-2j,\varkappa,\kappa} \rho^{2j} \bar{x}^{k-2j+1} \log^{\varkappa}(\bar{x}) \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X). \quad (681)$$

Also,

$$\sum_{k=0}^K \sum_{\substack{j=0 \\ k-2j \neq -1}}^J \hat{E}^j G_1^{(j,k,\varkappa)}(y) \sum_{\kappa=0}^{\varkappa} c_{k-2j,\varkappa,\kappa} \rho^{k+1} \log^{\varkappa}(\rho) \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}+1}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X). \quad (682)$$

Thus, $I_1 - \tilde{I}_1 \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F},(0,0)}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X)$. Moreover, $\tilde{I}_1 \in \hat{E} \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}+,(0,0)}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X)$.

On the other hand,

$$I_3 \in \mathcal{A}_{\text{loc}}^{(0,0),((0,0),K+1)}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X), \quad (683)$$

and we can write

$$I_2 = \sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \rho^k \hat{E}^{J+1} \left[\log^{\varkappa}(\rho) \int_{\rho/\bar{x}^{1/2}}^1 t^{2J-k} F_2^{(k,\varkappa)}(y; \hat{E}t^2) dt - \int_{\rho/\bar{x}^{1/2}}^1 t^{2J-k} \log^{\varkappa}(t) F_2^{(J+1,k,\varkappa)}(y; \hat{E}t^2) dt \right]. \quad (684)$$

Choosing J such that $2J \geq K$, we have $\int_{\rho/\bar{x}^{1/2}}^1 t^{2J-k} F_2^{(J+1,k,\varkappa)}(y; \hat{E}t^2) dt \in C^\infty([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X)$, so

$$\sum_{k=0}^K \sum_{\varkappa=0}^{\varkappa_k} \rho^k \hat{E}^{J+1} \log^{\varkappa}(\rho) \int_{\rho/\bar{x}^{1/2}}^1 t^{2J-k} F_2^{(k,\varkappa)}(y; \hat{E}t^2) dt \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X). \quad (685)$$

On the other hand, $\int_{\rho/\bar{x}^{1/2}}^1 t^{2J-k} \log^{\varkappa}(t) F_2^{(J+1,k,\varkappa)}(y; \hat{E}t^2) dt \in \mathcal{A}_{\text{loc}}^{(0,0),((0,0),2J-K+1)}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X)$. All in all,

$$I_2 \in \mathcal{A}_{\text{loc}}^{(0,0),(\mathcal{F},2J-K+1)}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X). \quad (686)$$

Combining all of the parts above, we have

$$I \in \mathcal{A}_{\text{loc}}^{(0,0),(\mathcal{F}+,\min\{K+1,2J-K+1\})}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X) \quad (687)$$

By taking J to be very large, $K+1$ and $2J-K+1$ can be made arbitrarily large simultaneously, so we conclude that I defines an element of

$$\bigcap_{L=1}^{\infty} \mathcal{A}_{\text{loc}}^{(0,0),(\mathcal{F}+,(L))}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X) = \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}+}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X). \quad (688)$$

This shows – in conjunction with the smoothness of I away from tf – that I is locally in $\mathcal{A}^{(0,0),\mathcal{F}+,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$ everywhere except possibly the corner $\text{bf} \cap \text{tf}$, to which we now turn.

We consider the following two cases:

- Suppose that G is supported in some set of the form $\{x/E > C\}$. Let $\varrho = x/E$. Near $\text{bf} \cap \text{tf}$,

$$I(x, y, E) = I(\varrho E, y, E) = I(CE, y, E). \quad (689)$$

Defining $\rho = C^{1/2}\sigma$ and $\hat{E} = C^{-1/2}$, $I(CE, y, E) = I(\rho^2, y, \hat{E}\rho^2)$. Since we already know that $I(\rho^2, y, \hat{E}\rho^2)$ depends polyhomogeneously on ρ, \hat{E}, y , with the desired index set \mathcal{F}_+ , we conclude that I depends polyhomogeneously on σ, y alone near $\text{bf} \cap \text{tf}$.

- On the other extreme, suppose that G is supported in some set of the form $\{x/E < c\}$, for $c > 0$. In order to study I near $\text{bf} \cap \text{tf}$, we work with the coordinate $\varrho = x/E$. By the polyhomogeneity of G , there exists a $G_2 \in \mathcal{A}_{\text{loc}}^{\mathcal{F},(0,0)}([0, \infty)_\sigma \times [0, \infty)_\varrho \times \partial X)$ such that $G_2(\varrho, y; \sigma) = G(x, y; E)$ whenever $\varrho = x/E$ and $\sigma^2 = E$. We can then write

$$I = \sigma \int_{\varrho}^{\min\{\bar{x}/\sigma^2, c\}} (1 + Z\varrho_0)^{-1/2} G_2(\varrho_0, y; \sigma) d\varrho_0 = \sigma \int_{\varrho}^{\min\{\bar{x}/\sigma^2, c\}} G_3(\varrho_0, y; \sigma) d\varrho_0 \quad (690)$$

for $G_3 \in \mathcal{A}_{\text{loc}}^{\mathcal{F},(0,0)}([0, \infty)_\sigma \times [0, \infty)_\varrho \times \partial X)$ defined by $G_3(\varrho, y; \sigma) = (1 + Z\varrho)^{-1/2} G_2(\varrho, y; \sigma)$.

The right-hand side of eq. (690) is in $\sigma \mathcal{A}^{\mathcal{F},(0,0)}([0, \infty)_\sigma \times [0, \infty)_\varrho \times \partial X)$ for $\sigma < \bar{x}^{1/2}/c^{1/2}$.

Since any $G \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F},(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$ can be decomposed $G = G_1 + G_2$ into $G_1, G_2 \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F},(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$ with G_1 supported on $\{x/E > C\}$ and G_2 supported on $\{x/E < c\}$ for some $c, C > 0$, and since $I = I[G]$ depends linearly on G , we can conclude that

$$I \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{F}_+,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}). \quad (691)$$

□

A simple example worth keeping mind is $g = x\sigma^2(\sigma^2 + Zx)^{-1}$, for which

$$\int_x^{\bar{x}} x_0^{-1} g(x_0; \sigma) dx_0 = -\frac{\sigma^2}{Z} \log\left(\frac{\sigma^2 + Zx}{\sigma^2 + Z\bar{x}}\right). \quad (692)$$

This shows that, even when $g \in xC^\infty(\hat{X}_{\text{res}}^{\text{sp}})$, solving $x\partial_x u = g$ for u may produce logarithms at tf . To compare with Proposition 6.18,

$$(\sigma^2 + Zx)^{1/2} x^{-1} g \in \mathcal{A}_{\text{loc}}^{(0,0),(1,0),(1,0)}(\hat{X}_{\text{res}}^{\text{sp}}) \subset \mathcal{A}_{\text{loc}}^{(0,0),(0,0),(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}), \quad (693)$$

while

$$-\frac{\sigma^2}{Z} \log\left(\frac{\sigma^2 + Zx}{\sigma^2 + Z\bar{x}}\right) \in \mathcal{A}_{\text{loc}}^{(0,0),(2,1),(1,0)}(\hat{X}_{\text{res}}^{\text{sp}}) \subset \mathcal{A}_{\text{loc}}^{(0,0),(2,1),(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}). \quad (694)$$

Setting \mathcal{F} to be the index set generated by $(0, 0)$, \mathcal{F}^+ is generated by $(0, 0)$ and $(2, 1)$, so this is in accordance with Proposition 6.18.

Since the index set $\mathcal{E} \subset \mathbb{N} \times \mathbb{N}$ contains $(0, 0)$ and satisfies $\mathcal{E} = \mathcal{E}_+$:

Corollary 6.19. *If $g \in x(\sigma^2 + Zx)^{-1/2} \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$, then*

$$I = \int_x^{\bar{x}} x_0^{-1} g(x_0, y; \sigma) dx_0 \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}). \quad (695)$$

■□

Proposition 6.20. *If $g \in \mathcal{A}_{\text{loc}}^{((\alpha_0,0),\alpha),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$ for $\alpha_0 \in \mathbb{N}^+$, $\alpha, \beta > 0$, $\beta \notin 2\mathbb{N}$, then*

$$\int_x^{\bar{x}} x_0^{-1} g(x_0, y; \sigma) dx_0 \in C^\infty(\hat{X}_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{((0,0),\alpha),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}) \quad (696)$$

holds. ■

Proof. Let $I = \int_x^{\bar{x}} x_0^{-1} g(x_0, y; \sigma) dx_0$. It suffices to prove the following two claims:

- (I) if g is supported in $\{x/E > C\}$, then $I \in C^\infty(\hat{X}_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{(0,0),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$,

(II) if g is supported in $\{x/E < c\}$, then $I \in \mathcal{A}_{\text{loc}}^{((0,0),\alpha),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$.

In order to prove (II), we write I in terms of $\varrho = x/E$:

$$I = \int_{\varrho}^{\min\{\bar{x}/\sigma^2, c\}} \varrho_0^{-1} G(\varrho_0, y; \sigma) d\varrho_0, \quad (697)$$

where $G \in \mathcal{A}_{\text{loc}}^{\beta,((\alpha_0,0),\alpha)}([0, \infty)_{\sigma} \times [0, \infty)_{\varrho} \times \partial X)$. Since $\alpha_0, \alpha > 0$, we can write

$$I = + \int_0^c \varrho_0^{-1} G(\varrho_0, y; \sigma) d\varrho_0 - \int_0^{\varrho} \varrho_0^{-1} G(\varrho_0, y; \sigma) d\varrho_0 \quad (698)$$

for σ sufficiently small. The first term is in $\mathcal{A}_{\text{loc}}^{\beta}([0, \infty)_{\sigma} \times \partial X) \subset \mathcal{A}_{\text{loc}}^{\beta,(0,0)}([0, \infty)_{\sigma} \times [0, \infty)_{\varrho} \times \partial X)$, while the second is in $\mathcal{A}_{\text{loc}}^{\beta,((\alpha_0,0),\alpha)}([0, \infty)_{\sigma} \times [0, \infty)_{\varrho} \times \partial X)$. Thus, it is also the case that

$$I \in \mathcal{A}_{\text{loc}}^{((0,0),\alpha),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}}) \quad (699)$$

away from zf. Since I is identically zero in the set $\{Ec < x\}$ and therefore smooth in some neighborhood of zf, this suffices to show that, globally, $I \in \mathcal{A}_{\text{loc}}^{((0,0),\alpha),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$.

To prove (I), we write I in terms of $\rho = x^{1/2}$ and $\hat{E} = E/x$:

$$I = \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{-1} G\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0, \quad (700)$$

where $G(\rho, y; \hat{E}) \in C^{\infty}([0, \infty)_{\hat{E}}; \mathcal{A}_{\text{loc}}^{\beta}([0, \infty)_{\rho} \times \partial X))$ is supported in $\hat{E} < C^{-1}$.

We expand $G(\rho, y; \hat{E})$ in Taylor series around $\hat{E} = 0$: there exist $G^{(k)} \in \mathcal{A}_{\text{loc}}^{\beta}([0, \infty)_{\rho} \times \partial X)$ and $F^{(k)} \in C^{\infty}([0, \infty)_{\hat{E}}; \mathcal{A}_{\text{loc}}^{\beta}([0, \infty)_{\rho} \times \partial X))$ such that

$$G(\rho, y; \hat{E}) = \sum_{k=0}^K \hat{E}^k G^{(k)}(\rho, y) + \hat{E}^{K+1} F^{(K+1)}(\rho, y; \hat{E}). \quad (701)$$

Substituting this into eq. (700),

$$I = \sum_{k=0}^K \hat{E}^k \rho^{2k} \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 + \hat{E}^{K+1} \rho^{2K+2} \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{-3-2K} F^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 \quad (702)$$

for each $K \in \mathbb{N}$. We take $K = \lfloor \beta/2 \rfloor$, in which case we can write

$$\int_{\rho}^{\bar{x}^{1/2}} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 = \int_0^{\bar{x}^{1/2}} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 - \int_0^{\rho} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0, \quad (703)$$

where the integrals on the right-hand side are well-defined because $\rho^{-1-2k} G^{(k)} \in \mathcal{A}_{\text{loc}}^{0+}$ for $k = 0, \dots, K$ (owing to $\beta \notin 2\mathbb{N}$). We split $I = I_0 + I_1 + I_2$, where

$$\begin{aligned} I_0 &= \sum_{k=0}^K \hat{E}^k \rho^{2k} \int_0^{\bar{x}^{1/2}} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 \in C^{\infty}([0, \infty)_{\hat{E}} \times [0, \infty)_{\rho} \times \partial X) \\ I_1 &= - \sum_{k=0}^K \hat{E}^k \rho^{2k} \int_0^{\rho} \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 \\ I_2 &= \hat{E}^{K+1} \rho^{2K+2} \int_{\rho}^{\bar{x}^{1/2}} \rho_0^{-3-2K} F^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0. \end{aligned} \quad (704)$$

Since $\rho^{2k} \int_0^\rho \rho_0^{-1-2k} G^{(k)}(\rho_0, y) d\rho_0 \in \mathcal{A}_{\text{loc}}^\beta([0, \infty)_\rho \times \partial X)$, $I_1 \in C^\infty([0, \infty)_{\hat{E}}; \mathcal{A}_{\text{loc}}^\beta([0, \infty)_\rho \times \partial X))$. The same holds for I_2 . The L^∞ case of this estimate is

$$\begin{aligned} \left| \rho^{2K+2} \int_\rho^{\bar{x}^{1/2}} \rho_0^{-3-2K} F^{(K+1)}\left(\rho_0, y; \hat{E}\left(\frac{\rho}{\rho_0}\right)^2\right) d\rho_0 \right| &\leq \rho^{2K+2} \int_\rho^{\bar{x}^{1/2}} \rho_0^{-3-2K+\beta} d\rho_0 \\ &\leq \rho^{2K+2} \left[(\bar{x}^{1/2})^{-2-2K+\beta} + \rho^{-2-2K+\beta} \right] \\ &\leq \rho^{2[\beta/2]+2} + \rho^\beta = O(\rho^\beta), \end{aligned} \quad (705)$$

and $\partial_{\hat{E}}^\kappa(\rho \partial_\rho)^\kappa I_2$ is estimated similarly.

Thus, $I \in C^\infty([0, \infty)_{\hat{E}} \times [0, \infty)_\rho \times \partial X) + C^\infty([0, \infty)_{\hat{E}}; \mathcal{A}_{\text{loc}}^\beta([0, \infty)_\rho \times \partial X))$. Using eq. (689) as in the proof of the previous proposition, this suffices to show that $I \in C^\infty(\hat{X}_{\text{res}}^{\text{sp}}) + \mathcal{A}_{\text{loc}}^{(0,0),\beta,(0,0)}(\hat{X}_{\text{res}}^{\text{sp}})$. \square

APPENDIX A. THE MODEL ODE

We record in this appendix some computations regarding the model ODE, eq. (62), now allowing a Schwarzschild-like subleading term and nonzero forcing:

$$(1 + \mathbf{a}x) \left(x^2 \frac{\partial}{\partial x} \right)^2 u + \sigma^2 u + \mathbf{Z}xu = -f, \quad (706)$$

where $\mathbf{a} \neq 0$, $f \in C_c^\infty(\mathbb{R}_x^+)$, and $u \in C^\infty(\mathbb{R}^+)$. The case $\mathbf{a} < 0$ is computationally similar to the case $\mathbf{a} \geq 0$, but in the former it is necessary to restrict x to $(0, 1/|\mathbf{a}|)$ to avoid the second order term in eq. (706) vanishing at $x = 1/|\mathbf{a}|$. We will therefore only consider the case $\mathbf{a} \geq 0$. If we allow \mathbf{Z} to depend on σ , the $\mathbf{a} > 0$ case can be reduced to the $\mathbf{a} = 0$ case via a simple change of variables: let $x_0 = (\mathbf{a} + x^{-1})^{-1} = x/(1 + \mathbf{a}x)$, so that $r_0 = 1/x_0$ is given by $r_0 = \mathbf{a} + r$, where $r = 1/x$. Then $\partial_r = \partial_{r_0}$, and the ODE eq. (706) is equivalent to

$$\left(x_0^2 \frac{\partial}{\partial x_0} \right)^2 u + \sigma^2 u + (\mathbf{Z} - \sigma^2 \mathbf{a}) x_0 u = -f_0(x_0) \quad (707)$$

for $f_0(x_0) = (1 - \mathbf{a}x_0)f(x)$. The interval $[0, \infty)_x$ becomes $[0, 1/\mathbf{a})_{x_0}$, but we can analyze eq. (707) on the larger region $[0, \infty)_{x_0}$, so that the analysis of eq. (706) is reduced to the $\mathbf{a} = 0$ case, with $\mathbf{Z}(\sigma) = \mathbf{Z} - \sigma^2 \mathbf{a}$. Since the results in this section mainly serve to illustrate the general features of the problem observed in the body of the paper, proofs are either sketched or omitted entirely when elementary. References for many of the elementary statements can be found in [Bat53][Sla60][AS64][Olv97]. We will mostly cite [AS64] when an explicit reference is desired.

The (nonzero) solutions to the homogeneous ODE cannot be written in terms of elementary functions for any value of $\sigma \geq 0$. In fact, for $\sigma \neq 0$, the homogeneous ODE is essentially a special case of *Whittaker's ODE*

$$\frac{d^2 W}{dz^2} + \left(-\frac{1}{4} + \frac{\kappa}{z} + \frac{1/4 - \mu^2}{z^2} \right) W = 0, \quad (708)$$

where $\kappa \in \mathbb{C}$, $\mu \in \mathbb{C}$ are parameters and $W \in \mathcal{D}'(\mathbb{R}_z^+)$. When $\mu \notin -2^{-1}\mathbb{N}^+$, there exist two named solutions to eq. (708), and these extend from the nonnegative real axis to analytic functions

$$\text{WhittM}_{\kappa,\mu}, \text{WhittW}_{\kappa,\mu} : \mathbb{C}_z \setminus (-\infty, 0] \rightarrow \mathbb{C} \quad (709)$$

which solve Whittaker's ODE in the complex analytic sense. These functions are Whittaker's M- and W- functions. They can be written in terms of Kummer's and Tricomi's confluent hypergeometric functions [AS64, §13.1.32, §13.1.33]. For certain values of μ , the branch cuts in eq. (709) can be removed.

Proposition A.1. *For $\sigma > 0$, the set of $u \in \mathcal{D}'(\mathbb{R}_r^+)$ solving the original homogeneous ODE $\partial_r^2 u + \sigma^2 u + \mathbf{Z}r^{-1}u = 0$ is given by*

$$\mathcal{W}_\sigma = \{c_1 \text{WhittM}_{\kappa,1/2}(2i\sigma r) + c_2 \text{WhittW}_{\kappa,1/2}(2i\sigma r) : c_1, c_2 \in \mathbb{C}\} \quad (710)$$

for $\kappa = -iZ/2\sigma$. ■□

The asymptotic expansions of the Whittaker M- and W-functions at large imaginary argument are due originally to Whittaker. For fixed $\sigma > 0$, we have

$$\begin{aligned} \text{Whitt}M_{\kappa,1/2}(2i\sigma r) = & \left[-\frac{(-2i\sigma)^{-iZ/2\sigma}}{\Gamma(1+\kappa)} e^{-i\sigma r - (iZ/2\sigma) \log r} \right. \\ & \left. + \frac{(+2i\sigma)^{iZ/2\sigma}}{\Gamma(1-\kappa)} e^{+i\sigma r + (iZ/2\sigma) \log r} \right] \left(1 + O_{\kappa,\sigma}\left(\frac{1}{r}\right) \right) \end{aligned} \quad (711)$$

$$\text{Whitt}W_{\kappa,1/2}(2i\sigma r) = (2i\sigma)^{-iZ/2\sigma} e^{-i\sigma r - (iZ/2\sigma) \log r} \left(1 + O_{\kappa,\sigma}\left(\frac{1}{r}\right) \right) \quad (712)$$

as $r \rightarrow \infty$. Here, and below, we leave the Z dependence of the estimates implicit. We are using the principal branch of the logarithm in making sense of $(2i\sigma)^{-iZ/2\sigma}$. As we are concerned with $r \rightarrow \infty$ behavior, the Whittaker M- function (which is singled out of the space of all solutions to Whittaker's ODE by its behavior at small argument) is not one of the solutions of direct interest. For $\sigma > 0$,

$$\mathcal{W}_{\sigma,-} = \text{span}_{\mathbb{C}}\{\text{Whitt}W_{\kappa,1/2}(2i\sigma r)\} \quad (713)$$

$$\begin{aligned} \mathcal{W}_{\sigma,+} = \text{span}_{\mathbb{C}} \left\{ \left[(2i\sigma)^{-iZ/2\sigma} \text{Whitt}M_{\kappa,1/2}(2i\sigma r) \right. \right. \\ \left. \left. + (-2i\sigma)^{-(iZ/2\sigma)} \Gamma(1+\kappa)^{-1} \text{Whitt}W_{\kappa,1/2}(2i\sigma r) \right] \right\} \end{aligned} \quad (714)$$

are the spaces of “incoming” or “outgoing” solutions to the ODE, and $\mathcal{W}_{\sigma} = \mathcal{W}_{\sigma,-} \oplus \mathcal{W}_{\sigma,+}$. It is the spaces $\mathcal{W}_{\sigma,-}, \mathcal{W}_{\sigma,+}$ that concern us.

Set $w_{-}(r; \sigma) = (2i\sigma)^{iZ/2\sigma} \text{Whitt}W_{\kappa,1/2}(2i\sigma r)$ and

$$w_{+}(r; \sigma) = \left[(2i\sigma)^{-iZ/2\sigma} \Gamma(1-\kappa) \text{Whitt}M_{\kappa,1/2}(2i\sigma r) + (-2i\sigma)^{-iZ/2\sigma} \frac{\Gamma(1-\kappa)}{\Gamma(1+\kappa)} \text{Whitt}W_{\kappa,1/2}(2i\sigma r) \right]. \quad (715)$$

Thus, $w_{-}(-; \sigma) \in \mathcal{W}_{\sigma,-}$ and $w_{+}(-; \sigma) \in \mathcal{W}_{\sigma,+}$, and $w_{\pm}(r; \sigma) \in C^{\infty}(\mathbb{R}_{\sigma}^{+} \times \mathbb{R}_{r}^{+})$. These have normalized oscillatory behavior $\exp(\pm i\sigma r)$ as $r \rightarrow \infty$,

$$w_{\pm}(r; \sigma) = e^{\pm i\sigma r \pm (iZ/2\sigma) \log r} \left(1 + O_{\sigma}\left(\frac{1}{r}\right) \right). \quad (716)$$

For each $\sigma > 0$, w_{\pm} are the unique solutions to the ODE satisfying eq. (716). It follows from this and the fact that the ODE has real coefficients that $w_{-}(r; \sigma) = w_{+}(r; \sigma)^{*}$. Expanding to higher order [AS64, §13.5]: for each $\sigma > 0$ and $K \in \mathbb{N}^{+}$,

$$w_{\pm}(r; \sigma) = e^{\pm i\sigma r} r^{\pm iZ/2\sigma} \left[1 + Z \sum_{k=1}^{K-1} \frac{(\pm i)^k}{8^k k! \sigma^{3k} r^k} (Z \pm 2ki\sigma) \prod_{j=1}^{k-1} (Z \pm 2ji\sigma)^2 + O_{\sigma}\left(\frac{1}{r^K}\right) \right] \quad (717)$$

as $r \rightarrow \infty$.

For $\sigma = 0$, the set of $u \in \mathcal{D}'(\mathbb{R}_x^{+})$ solving the ODE $(x^2 \partial_x)^2 u + Zxu = 0$ is

$$\{c_1 r^{1/2} J_1(2Z^{1/2} r^{1/2}) + c_2 r^{1/2} Y_1(2Z^{1/2} r^{1/2}) : c_1, c_2 \in \mathbb{C}\}, \quad (718)$$

where J_1, Y_1 denote the Bessel J- and Y- functions of order one. As $r \rightarrow \infty$,

$$r^{1/2} J_1(2Z^{1/2} r^{1/2}) = \frac{r^{1/4}}{\pi^{1/2} Z^{1/4}} \left[-\cos\left(2\sqrt{Z}r + \frac{\pi}{4}\right) + \frac{3}{16} \frac{1}{\sqrt{Z}r} \sin\left(2\sqrt{Z}r + \frac{\pi}{4}\right) \right] \left(1 + O\left(\frac{1}{r}\right) \right) \quad (719)$$

$$r^{1/2} Y_1(2Z^{1/2} r^{1/2}) = \frac{r^{1/4}}{\pi^{1/2} Z^{1/4}} \left[-\sin\left(2\sqrt{Z}r + \frac{\pi}{4}\right) - \frac{3}{16} \frac{1}{\sqrt{Z}r} \cos\left(2\sqrt{Z}r + \frac{\pi}{4}\right) \right] \left(1 + O\left(\frac{1}{r}\right) \right), \quad (720)$$

and we have a full expansion in powers of $r^{-1/2}$. For $\sigma = 0$, set

$$v_-(r; 0) = r^{1/2} J_1(2Z^{1/2} r^{1/2}) - i r^{1/2} Y_1(2Z^{1/2} r^{1/2}) = r^{1/2} H_1^{(2)}(2Z^{1/2} r^{1/2}) \quad (721)$$

$$v_+(r; 0) = r^{1/2} J_1(2Z^{1/2} r^{1/2}) + i r^{1/2} Y_1(2Z^{1/2} r^{1/2}) = r^{1/2} H_1^{(1)}(2Z^{1/2} r^{1/2}), \quad (722)$$

where $H_1^{(1)}, H_1^{(2)}$ denote the Hankel functions of order one. These have the asymptotics

$$v_{\pm}(r; \sigma) = e^{\pm 2iZ^{1/2} r^{1/2} \mp 3\pi i/4} \frac{r^{1/4}}{\pi^{1/2} Z^{1/4}} \left(1 + O\left(\frac{1}{r^{1/2}}\right) \right) \quad (723)$$

in the $r \rightarrow \infty$ limit. Expanding eq. (723) to higher order [AS64, §9.2]: for each $K \in \mathbb{N}^+$,

$$v_{\pm}(r; 0) = e^{\pm 2iZ^{1/2} r^{1/2} \mp 3\pi i/4} \frac{r^{1/4}}{\pi^{1/2} Z^{1/4}} \sum_{k=0}^{K-1} (\pm i)^k Z^{-k/2} \frac{(2k+1)!(2k)!}{64^k (k!)^3} \frac{1}{r^{k/2}} + O\left(\frac{1}{r^{K/2-1/4}}\right) \quad (724)$$

as $r \rightarrow \infty$.

Proposition A.2. v_{\pm} are the unique solutions to the $\sigma = 0$ ODE satisfying the asymptotic eq. (723). \blacksquare

Proposition A.3. *Setting*

$$C_{\pm}(\sigma) = -\frac{Z^{1/2}}{2\pi\sigma} (\mp 2i\sigma)^{\pm iZ/(2\sigma)} \Gamma\left(\mp \frac{iZ}{2\sigma}\right) \quad (725)$$

and $v_{\pm}(r; \sigma) = C_{\pm}(\sigma) w_{\pm}(r; \sigma)$, the functions $v_{\pm}(r; E^{1/2}) : [0, \infty)_E \times (0, \infty)_r \rightarrow \mathbb{C}$ are both smooth all the way down to $E = 0$. Consequently, if

$$C_{\pm,0}(\sigma) = e^{\pm \pi i/4} \pi^{-1/2} \sigma^{-1/2} \exp\left[\pm \frac{Zi}{\sigma} \left(\log\left(\frac{2\sigma}{Z^{1/2}}\right) + \frac{1}{2}\right)\right], \quad (726)$$

then $C_{\pm,0}(\sigma) w_{\pm}(r; \sigma) : [0, \infty)_{\sigma} \times (0, \infty)_r \rightarrow \mathbb{C}$ are smooth all the way down to $\sigma = 0$, with restriction $v_{\pm}(-; 0)$ to $\sigma = 0$. \blacksquare

Remark. The C^0 case of this proposition is similar to [AS64, §13.3.4, §13.3.5], except that our κ is purely imaginary rather than purely real. See also [Tay39][Bat53, §6.13.3, Eq. (21) - (24)]. Of course, if we were to multiply $C_{\pm}(\sigma)$ by any element of $C^{\infty}[0, \infty)_E$, the resultant functions would also satisfy the proposition above, and likewise with $C_{\pm,0}(\sigma)$ in place of $C_{\pm}(\sigma)$ and $C^{\infty}[0, \infty)_{\sigma}$ in place of $C^{\infty}[0, \infty)_E$.

Remark. Once we know that C_{\pm} satisfies the first clause of the conclusion of Proposition A.3, the second clause of the proposition follows from the large argument asymptotics of $\log \Gamma$, which can be found in [AS64, §6.1.40]. Observe, using [AS64, §6.1.40], that the ratio $C_{\pm}/C_{\pm,0} : [0, \infty)_{\sigma} \rightarrow \mathbb{C}$ is a smooth function of σ , not $E = \sigma^2$. This is related to the fact that, for $u_{0,\pm}$ as in Theorem 1.2, $u_{0,\pm}|_{\text{bf}} : [0, \infty)_{\sigma} \times \partial X \rightarrow \mathbb{C}$ is only smooth with respect to σ , in contrast to $u_{0,\pm}|_{\{x=\varepsilon\}} : [0, \infty)_{\sigma} \times \partial X \rightarrow \mathbb{C}$ for $\varepsilon \in (0, \bar{x})$, which is smooth with respect to E .

Up to the $\exp(\pm \pi i/4) \pi^{-1/2}$ in eq. (726), eq. (726) is suggested by Theorem 1.2, as, for each $\sigma > 0$,

$$\Phi(x; \sigma) = \frac{\sigma}{x} - \frac{Z}{2\sigma} \log x + \frac{Z}{2\sigma} + \frac{Z}{\sigma} \log\left(\frac{2\sigma}{Z^{1/2}}\right) + O_{\sigma}(x) \quad (727)$$

as $x \rightarrow 0^+$, while the phase of the bracketed term in eq. (715) consists only of the first two terms of eq. (727). The $\sigma^{-1/2}$ term in eq. (726) is needed to match the $r^{1/2}$ in eq. (721), eq. (722). This heuristic should be taken with a grain of salt, as the $O_{\sigma}(x)$ term blows up as $\sigma \rightarrow 0^+$. Assuming that there exist some C_{\pm} satisfying the conclusion of Proposition A.3, the sufficiency of the formula above can be seen from the $r \rightarrow 0^+$ asymptotics of the Whittaker functions. Indeed, the ODE can be written in the form

$$\left(r \frac{\partial}{\partial r}\right)^2 u - \left(r \frac{\partial}{\partial r}\right) u + (\sigma^2 r^2 + Zr) u = 0, \quad (728)$$

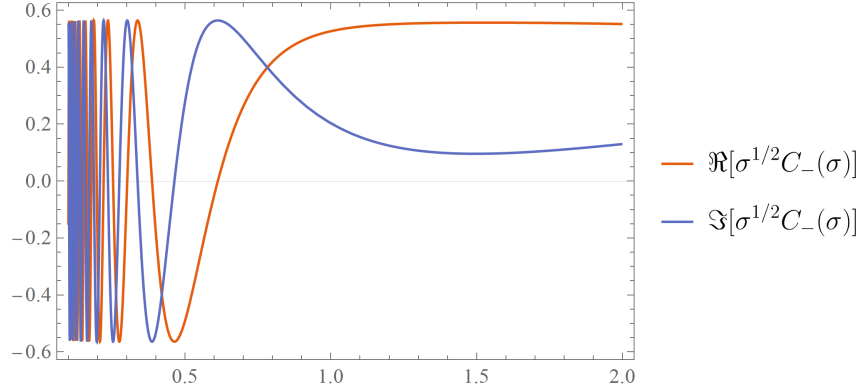


FIGURE 8. The real and imaginary parts of the function $C_-(\sigma)$ defined by eq. (725), rescaled so that the oscillations have constant amplitude.

which has a regular singular point at $r = 0$. The “normal operator” of $(r\partial_r)^2 - (r\partial_r) + \sigma^2 r^2 + Zr$ at $r = 0$ is $(r\partial_r)^2 - r\partial_r$, the indicial roots of which are 0 and 1. We can deduce that a family $\{u(-; \sigma)\}_{\sigma \geq 0} \subset \mathcal{D}'(\mathbb{R}^+)$ of solutions to eq. (728) can be smooth at $\{r > 0, \sigma = 0\}$ only if it extends continuously to $[0, 1]_\sigma \times [0, 1]_r$ (we will not have smoothness at $r = 0$ because of the presence of a $r \log r$ term), in which case $\lim_{\sigma \rightarrow 0^+} \lim_{r \rightarrow 0^+} u = \lim_{r \rightarrow 0^+} \lim_{\sigma \rightarrow 0^+} u$ exists. Now observe that, as $r \rightarrow 0^+$,

$$w_\pm(r; \sigma) = \pm \frac{2i}{Z} (\mp 2i\sigma)^{\mp iZ/2\sigma} \frac{\sigma}{\Gamma(\mp iZ/2\sigma)} + O_\sigma(r \log r), \quad (729)$$

$$\lim_{r \rightarrow 0^+} w_\pm(r; \sigma) = \pm \frac{2i}{Z} (\mp 2i\sigma)^{\mp iZ/2\sigma} \frac{\sigma}{\Gamma(\mp iZ/2\sigma)} \quad (730)$$

for each $\sigma > 0$. We now compare eq. (730) with

$$\lim_{r \rightarrow 0^+} v_\pm(r; 0) = \mp \frac{i}{\pi Z^{1/2}} + O(r \log r) \quad (731)$$

[AS64, §9.1.9, §9.1.11]. This yields $C_\pm(\sigma) = -(2\pi\sigma)^{-1} Z^{1/2} (\mp 2i\sigma)^{\pm iZ/2\sigma} \Gamma(\mp iZ/2\sigma)$, which is eq. (725). A more thorough analysis of the ODE near $r = 0$ suffices to prove Proposition A.3 properly.

Letting $\chi \in C_c^\infty[0, \infty)$ be identically one in some neighborhood of the origin, $v_\pm^\circ = \chi(1/r)v_\pm(r; \sigma)$ satisfies

$$\frac{\partial^2 v_\pm^\circ}{\partial r^2} + \left(\sigma^2 + \frac{Z}{r}\right) v_\pm^\circ = f_\pm \quad (732)$$

for some $f_\pm \in \cap_{E_0 > 0} C_c^\infty([0, E_0]_E \times \mathbb{R}^+)$. We can deduce (e.g. by appealing to Theorem 1.2 plus Remark 1 in the case of spherical symmetry) that v_\pm° has the form

$$v_\pm^\circ = \exp(\pm i\Phi(r^{-1}; E^{1/2}))(E + Zr^{-1})^{-1/4} v_{0,\pm}^\circ \quad (733)$$

for some $v_{0,\pm} \in \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}})$, where $X = [0, \infty)_x$ and $X_{\text{res}}^{\text{sp}} = [[0, \infty)_E \times X; \{0\} \times \{0\}; 1/2]$. Thus, $w_\pm(r, \sigma)$ satisfy

$$w_\pm(r; \sigma) \in C_\pm^{-1}(\sigma) e^{\pm i\Phi(r^{-1}; E^{1/2})} (E + Zx)^{-1/4} \mathcal{A}_{\text{loc}}^{(0,0),\mathcal{E},(0,0)}(X_{\text{res}}^{\text{sp}}). \quad (734)$$

That is, in the ‘−’ case:

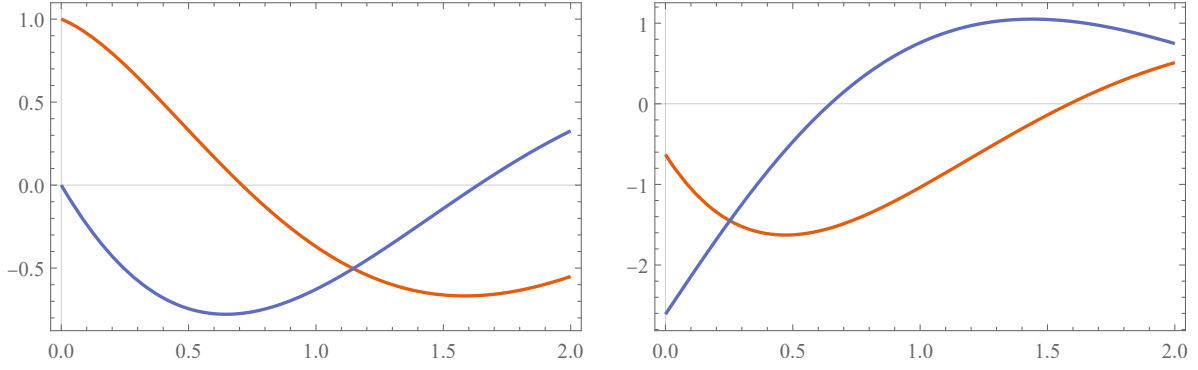


FIGURE 9. The function $U(r; E) = v_-(r; E^{1/2})/v_-(r; 0)$ (left) and its first derivative $\partial_E U(r; E)$ (right) evaluated at $r = 5$, plotted against $E \in [0, 2]$. The real parts are plotted in orange and the imaginary parts are in blue. From the figures it appears that $U \in C^2[0, 2)_E$ and $U(0) = 1$, in accord with Proposition A.3 (and with Corollary 1.4).

Proposition A.4. *The Whittaker W-function $\text{Whitt}W_{\kappa, 1/2}$ satisfies*

$$\begin{aligned} \text{Whitt}W_{-iZ/2\sigma, 1/2}(2i\sigma r) &\in \sigma \Gamma\left(\frac{iZ}{2\sigma}\right)^{-1} \left(\sigma^2 + \frac{Z}{r}\right)^{-1/4} \\ &\exp\left(-ir\sqrt{\sigma^2 + \frac{Z}{r}} - \frac{iZ}{\sigma} \log\left(\frac{\sigma r^{1/2}}{Z^{1/2}} + \left(1 + \frac{\sigma^2 r}{Z}\right)^{1/2}\right)\right) \mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}}). \end{aligned} \quad (735)$$

■□

Remark. It seems that there are no logarithmic terms in the expansion of the polyhomogeneous function in eq. (735). In order to prove this, it should be possible to combine the previous proposition with a WKB type expansion at tf, with σ being the semiclassical parameter. The point here is that the Whittaker W-function has simple asymptotics at bf, and the asymptotic expansion at tf without log terms can be concluded from this. The alternative argument above is indirect, utilizing the asymptotics at tf, and from this perspective it is somewhat miraculous that the asymptotic expansion at tf is one-step polyhomogeneous. For general forcing $f \in C_c^\infty(\mathbb{R}^+)$, we should not expect the outgoing solution to the forced ODE to have a one-step polyhomogeneous expansion at tf.

Proposition A.4 can be strengthened by specifying the behavior of $\text{Whitt}W_{-iZ/2\sigma, 1/2}(2i\sigma r)$ at small r . Namely, if we replace X by $\bar{X} = [0, \infty]_x$, then, replacing r by $\langle r \rangle$ in eq. (735), the same statement holds, with $\mathcal{A}_{\text{loc}}^{(0,0), \mathcal{E}, (0,0)}(X_{\text{res}}^{\text{sp}})$ replaced by the set of polyhomogeneous functions on $\bar{X}_{\text{res}}^{\text{sp}}$ with index set \mathbb{N} at zf, tf, bf and with some other index set (which can be specified) at the new face formed by the lift of $[0, \infty)_E \times \{\infty\}_x$ to $\bar{X}_{\text{res}}^{\text{sp}}$.

The most interesting regime is tf, which we can probe using the coordinates $\varsigma = \sigma r^{1/2}$ and $r^{-1/2}$. As a corollary of Proposition A.4, analogous to Corollary 1.6, we deduce that

$$\begin{aligned} \text{Whitt}W_{-iZr^{1/2}/2\varsigma, 1/2}(2i\varsigma r^{1/2}) &\in \varsigma^{1/2} e^{\frac{\pi Z}{4\varsigma} r^{1/2}} \exp\left(-ir^{1/2}\sqrt{\varsigma^2 + Z} + \frac{iZ}{2\varsigma} r^{1/2}\right) \\ &\exp\left[-\frac{iZ}{\varsigma} r^{1/2} \left[\frac{1}{2} \log\left(\frac{Zr^{1/2}}{2\varsigma}\right) + \log\left(\frac{\varsigma}{Z^{1/2}} + \left(1 + \frac{\varsigma^2}{Z}\right)^{1/2}\right)\right]\right] \mathcal{A}^{(0,0), \mathcal{F}}([0, \infty)_\varsigma \times [0, \infty)_{r^{-1/2}}) \end{aligned} \quad (736)$$

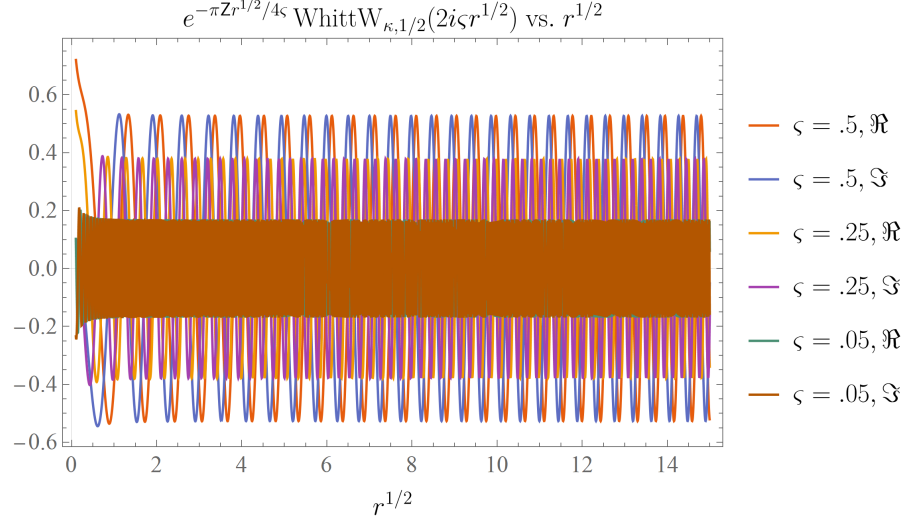


FIGURE 10. The real (in red, orange, and green) and imaginary (in blue, purple, and brown) parts of the function $\exp(-\pi Z r^{1/2}/4\varsigma) \text{Whitt}_{\kappa,1/2}(2i\varsigma r^{1/2})$ as a function of $r^{1/2}$, for $\varsigma \in \{.05, .25, .5\}$ and $Z = 3$ fixed, where $\kappa = -iZr^{1/2}/2\varsigma$.

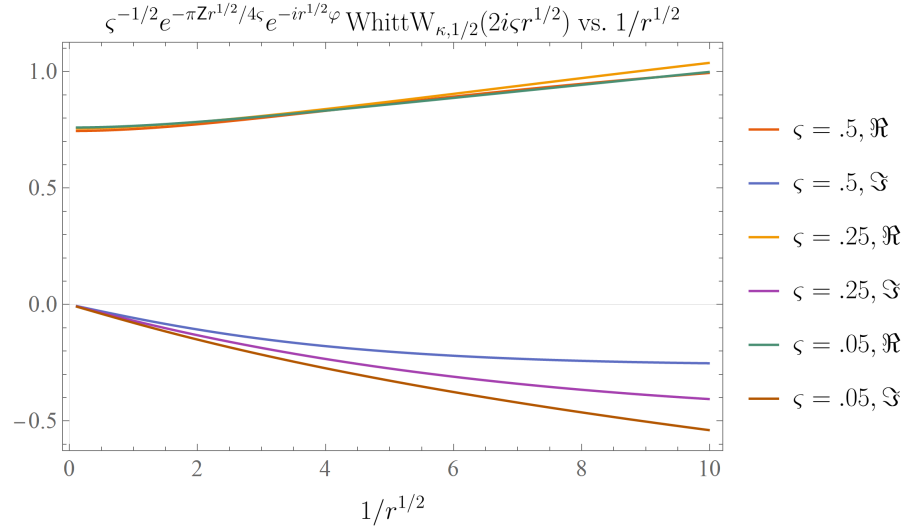


FIGURE 11. The real and imaginary parts of the function $\varsigma^{-1/2} \exp(-\pi Z r^{1/2}/4\varsigma) \exp(-ir^{1/2}\varphi) \text{Whitt}_{\kappa,1/2}(2i\varsigma r^{1/2})$ (with the same color scheme as in Figure 10), now plotted as a function of $\rho = 1/r^{1/2}$, for $\varsigma \in \{.05, .25, .5\}$ and $Z = 3$.

for some index set $\mathcal{F} \subset \mathbb{N} \times \mathbb{N}$. Here we used the large argument expansion of the Γ -function, i.e. Stirling's formula [AS64, §6.1.37], which implies

$$\Gamma(iZr^{1/2}/2\varsigma) \in r^{-1/4} \varsigma^{1/2} e^{-\pi Z r^{1/2}/4\varsigma} e^{-iZr^{1/2}/2\varsigma} (Zr^{1/2}/2\varsigma)^{iZr^{1/2}/2\varsigma} C^\infty([0, \infty)_\varsigma \times [0, \infty)_{r^{-1/2}}). \quad (737)$$

Thus, if we let

$$\varphi = -(\varsigma^2 + Z)^{1/2} + \frac{Z}{2\varsigma} - \frac{Z}{\varsigma} \left[\frac{1}{2} \log \left(\frac{Zr^{1/2}}{2\varsigma} \right) + \log \left(\frac{\varsigma}{Z^{1/2}} + \left(1 + \frac{\varsigma^2}{Z} \right)^{1/2} \right) \right], \quad (738)$$

then, for each $\varsigma > 0$, $e^{-\pi Z r^{1/2}/4\varsigma} e^{-ir^{1/2}\varphi} \text{Whitt}_{\kappa,1/2}(2i\varsigma r^{1/2})$ is a polyhomogeneous function of $r^{-1/2}$, all the way down to $r^{-1/2} = 0$. The convergence aspect of this result, in particular the fact that the multiplication by $\exp(-ir^{1/2}\varphi)$ kills off the oscillations of $\exp(-\pi Z r^{1/2}/4\varsigma) \text{Whitt}_{\kappa,1/2}(2i\varsigma r^{1/2})$, is depicted in Figure 11 (and the contrast with Figure 10, where the $\exp(-ir^{1/2}\varphi)$ factor is missing). The C^0 case of this is similar to [Bat53, §6.13.3, Eq. (21) - (24)] (though we did not compute out the leading order term in the asymptotic expansion).

ACKNOWLEDGEMENTS

The author thanks Peter Hintz for guidance and comments on the manuscript, and in addition Nick Lohr for comments on the introduction. This work was supported by a Hertz fellowship.

REFERENCES

- [AS64] M. Abramowitz and I. A. Stegun. *Handbook of Mathematical Functions with Formulas, Graphs, and Mathematical Tables*. National Bureau of Standards Applied Mathematics Series **55**. U. S. Government Printing Office, 1964 (cit. on pp. 100–103, 105).
- [Bar+12] J. Barranco et al. “Schwarzschild black holes can wear scalar wigs”. *Phys. Rev. Lett.* **109** (8 2012). DOI: [10.1103/PhysRevLett.109.081102](https://doi.org/10.1103/PhysRevLett.109.081102) (cit. on p. 15).
- [Bar+14] ———. “Schwarzschild scalar wigs: Spectral analysis and late time behavior”. *Phys. Rev. D* **89** (8 2014). DOI: [10.1103/PhysRevD.89.083006](https://doi.org/10.1103/PhysRevD.89.083006) (cit. on p. 15).
- [Bat53] H. Bateman. *Higher Transcendental Functions*. Vol. I-III. McGraw-Hill, 1953 (cit. on pp. 100, 102, 106).
- [BK04] L. M. Burko and G. Khanna. “Universality of massive scalar field late-time tails in black-hole spacetimes”. *Phys. Rev. D* **70** (4 2004). DOI: [10.1103/PhysRevD.70.044018](https://doi.org/10.1103/PhysRevD.70.044018) (cit. on p. 15).
- [Bou11] J.-M. Bouclet. “Low frequency estimates and local energy decay for asymptotically Euclidean Laplacians”. *Comm. Partial Differential Equations* **36** (2011), 1239–1286. DOI: [10.1080/03605302.2011.558553](https://doi.org/10.1080/03605302.2011.558553) (cit. on p. 6).
- [BOW67] P. G. Burke, S. Ormonde, and W. Whitaker. “Low-energy electron scattering by atomic hydrogen: I. The close-coupling approximation”. *Proceedings of the Physical Society* **92.2** (1967), 319–335. DOI: [10.1088/0370-1328/92/2/308](https://doi.org/10.1088/0370-1328/92/2/308) (cit. on p. 15).
- [BT67] P. G. Burke and A. J. Taylor. “Low-energy electron scattering by atomic hydrogen: II. The correlation method”. *Proceedings of the Physical Society* **92.2** (1967), 336–344. DOI: [10.1088/0370-1328/92/2/309](https://doi.org/10.1088/0370-1328/92/2/309) (cit. on p. 15).
- [BTO67] P. G. Burke, A. J. Taylor, and S. Ormonde. “Low-energy electron scattering by atomic hydrogen: III. Comparison of theory and experiment for the electron-induced excitation of hydrogen to the $n = 2$ level”. *Proceedings of the Physical Society* **92.2** (1967), 345–350. DOI: [10.1088/0370-1328/92/2/310](https://doi.org/10.1088/0370-1328/92/2/310) (cit. on p. 15).
- [Det80] S. Detweiler. “Klein-Gordon equation and rotating black holes”. *Phys. Rev. D* **22** (10 1980), 2323–2326. DOI: [10.1103/PhysRevD.22.2323](https://doi.org/10.1103/PhysRevD.22.2323) (cit. on p. 15).
- [Dol64] J. D. Dollard. “Asymptotic convergence and the Coulomb interaction”. *J. Mathematical Phys.* **5** (1964). DOI: [10.1063/1.1704171](https://doi.org/10.1063/1.1704171) (cit. on pp. 1, 2).
- [DS09a] J. Dereziński and E. Skibsted. “Quantum scattering at low energies”. *J. Funct. Anal.* **257** (2009), 1828–1920. DOI: [10.1016/j.jfa.2009.05.026](https://doi.org/10.1016/j.jfa.2009.05.026) (cit. on p. 6).
- [DS09b] ———. “Scattering at zero energy for attractive homogeneous potentials”. *Ann. Henri Poincaré* **10** (2009), 549–571. DOI: [10.1007/s00023-009-0408-x](https://doi.org/10.1007/s00023-009-0408-x) (cit. on p. 6).
- [DSS11] R. Donninger, W. Schlag, and A. Soffer. “A proof of Price’s law on Schwarzschild black hole manifolds for all angular momenta”. *Adv. Math.* **226.1** (2011), 484–540. DOI: [10.1016/j.aim.2010.06.026](https://doi.org/10.1016/j.aim.2010.06.026) (cit. on p. 15).
- [DSS12] ———. “On pointwise decay of linear waves on a Schwarzschild black hole background”. *Comm. Math. Phys.* **309.1** (2012), 51–86. DOI: [10.1007/s00220-011-1393-8](https://doi.org/10.1007/s00220-011-1393-8) (cit. on p. 15).
- [FS04] S. Fournais and E. Skibsted. “Zero energy asymptotics of the resolvent for a class of slowly decaying potentials”. *Math. Z.* **248.3** (2004), 593–633. DOI: [10.1007/s00209-004-0673-9](https://doi.org/10.1007/s00209-004-0673-9) (cit. on pp. 6, 10, 14).
- [GH08] C. Guillarmou and A. Hassell. “Resolvent at low energy and Riesz transform for Schrödinger operators on asymptotically conic manifolds. I”. *Math. Ann.* **341.4** (2008), 859–896. DOI: [10.1007/s00208-008-0216-5](https://doi.org/10.1007/s00208-008-0216-5) (cit. on pp. 2, 5, 7).
- [GH09] ———. “Resolvent at low energy and Riesz transform for Schrödinger operators on asymptotically conic manifolds. II”. *Ann. Inst. Fourier (Cie)* **59.4** (2009), 1553–1610. URL: http://aif.cedram.org/item?id=AIF_2009__59_4_1553_0 (cit. on pp. 2, 5).

- [GHS13] C. Guillarmou, A. Hassell, and A. Sikora. “Resolvent at low energy III: The spectral measure”. *Trans. Amer. Math. Soc.* **365**.11 (2013), 6103–6148. DOI: [10.1090/S0002-9947-2013-05849-7](https://doi.org/10.1090/S0002-9947-2013-05849-7) (cit. on pp. [2](#), [5](#), [7](#), [14](#)).
- [GR+20] J. Gell-Redman et al. “Existence and asymptotics of nonlinear Helmholtz eigenfunctions”. *SIAM J. Math. Anal.* **52**.6 (2020), 6180–6221. DOI: [10.1137/19M1307238](https://doi.org/10.1137/19M1307238) (cit. on pp. [6](#), [7](#)).
- [Hö7] L. Hörmander. *The Analysis of Linear Partial Differential Operators*. Vol. I-IV. Classics in Mathematics. Springer, 2007 (cit. on pp. [28](#), [84](#)).
- [HHV21] D. Häfner, P. Hintz, and A. Vasy. “Linear stability of slowly rotating Kerr black holes”. *Invent. Math.* **223**.3 (2021), 1227–1406. DOI: [10.1007/s00222-020-01002-4](https://doi.org/10.1007/s00222-020-01002-4) (cit. on p. [92](#)).
- [Hin22] P. Hintz. “A sharp version of Price’s law for wave decay on asymptotically flat spacetimes”. *Comm. Math. Phys.* **389** (2022), 491–542. DOI: [10.1007/s00220-021-04276-8](https://doi.org/10.1007/s00220-021-04276-8) (cit. on pp. [3](#), [5](#), [11](#), [13–15](#), [19](#), [38](#), [92](#)).
- [HMY04] A. Hassell, R. Melrose, and A. Vasy. “Spectral and scattering theory for symbolic potentials of order zero”. *Adv. Math.* **181**.1 (2004), 1–87. DOI: [10.1016/S0001-8708\(03\)00020-3](https://doi.org/10.1016/S0001-8708(03)00020-3) (cit. on p. [6](#)).
- [HP98] S. Hod and T. Piran. “Late-time tails in gravitational collapse of a self-interacting (massive) scalar-field and decay of a self-interacting scalar hair”. *Phys. Rev. D* **58** (4 1998). DOI: [10.1103/PhysRevD.58.044018](https://doi.org/10.1103/PhysRevD.58.044018) (cit. on p. [15](#)).
- [HV99] A. Hassell and A. Vasy. “The spectral projections and the resolvent for scattering metrics”. *J. Anal. Math.* **79** (1999), 241–298. DOI: [10.1007/BF02788243](https://doi.org/10.1007/BF02788243) (cit. on p. [5](#)).
- [HW05] A. Hassell and J. Wunsch. “The Schrödinger propagator for scattering metrics”. *Ann. of Math. (2)* **162**.1 (2005), 487–523. DOI: [10.4007/annals.2005.162.487](https://doi.org/10.4007/annals.2005.162.487) (cit. on pp. [16](#), [20](#)).
- [Ivr19] V. Ivrii. *Microlocal Analysis, Sharp Spectral Asymptotics and Applications*. Vol. II. Springer, 2019 (cit. on p. [7](#)).
- [JSB99] M. S. Joshi and A. Sá Barreto. “Recovering asymptotics of metrics from fixed energy scattering data”. *Invent. Math.* **137**.1 (1999), 127–143. DOI: [10.1007/s002220050326](https://doi.org/10.1007/s002220050326) (cit. on p. [16](#)).
- [KT01] H. Koyama and A. Tomimatsu. “Asymptotic tails of massive scalar fields in a Schwarzschild background”. *Phys. Rev. D* **64** (4 2001). DOI: [10.1103/PhysRevD.64.044014](https://doi.org/10.1103/PhysRevD.64.044014) (cit. on p. [15](#)).
- [KT02] ———. “Slowly decaying tails of massive scalar fields in spherically symmetric spacetimes”. *Phys. Rev. D* **65** (8 2002). DOI: [10.1103/PhysRevD.65.084031](https://doi.org/10.1103/PhysRevD.65.084031) (cit. on p. [15](#)).
- [KZM07] R. A. Konoplya, A. Zhidenko, and C. Molina. “Late time tails of the massive vector field in a black hole background”. *Phys. Rev. D* **75** (8 2007). DOI: [10.1103/PhysRevD.75.084004](https://doi.org/10.1103/PhysRevD.75.084004) (cit. on p. [15](#)).
- [Loo21] S.-Z. Looi. *Pointwise decay for the wave equation on nonstationary spacetimes*. 2021. arXiv: [2105.02865](https://arxiv.org/abs/2105.02865) [[math.AP](#)] (cit. on p. [15](#)).
- [Mac67] J. Macek. “Low-energy electron scattering by atomic hydrogen: V. Photoionization of H^- in the vacuum ultra-violet”. *Proceedings of the Physical Society* **92**.2 (1967), 365–369. DOI: [10.1088/0370-1328/92/2/312](https://doi.org/10.1088/0370-1328/92/2/312) (cit. on p. [15](#)).
- [MB67] J. Macek and P. G. Burke. “Low-energy electron scattering by atomic hydrogen: IV. The effect of resonances on excitation in e-H scattering”. *Proceedings of the Physical Society* **92**.2 (1967), 351–364. DOI: [10.1088/0370-1328/92/2/311](https://doi.org/10.1088/0370-1328/92/2/311) (cit. on p. [15](#)).
- [Mel92] R. B. Melrose. “Calculus of conormal distributions on manifolds with corners”. *Internat. Math. Res. Notices* **3** (1992), 51–61. DOI: [10.1155/S1073792892000060](https://doi.org/10.1155/S1073792892000060) (cit. on pp. [7](#), [13](#), [38](#)).
- [Mel93] ———. *The Atiyah-Patodi-Singer Index Theorem*. Research Notes in Mathematics **4**. CRC Press, 1993. URL: <https://www.maths.ed.ac.uk/~v1ranick/papers/melrose.pdf> (cit. on pp. [12](#), [21](#), [38](#)).
- [Mel94] ———. “Spectral and Scattering Theory for the Laplacian on Asymptotically Euclidian Spaces”. *Spectral and Scattering Theory*. Ed. by M. Ikawa. CRC Press, 1994. URL: <https://klein.mit.edu/~rbm/papers/sslaes/sslaes1.pdf> (cit. on pp. [1](#), [2](#), [5](#), [7](#), [9](#), [11](#), [13](#), [14](#), [16](#), [50](#), [54–56](#), [86](#)).
- [Mel95] ———. *Geometric Scattering Theory*. Stanford Lectures. Cambridge University Press, 1995 (cit. on pp. [1](#), [2](#), [16](#)).
- [Mor19] K. Morgan. *Wave decay in the asymptotically flat stationary setting*. Thesis (Ph.D.)—The University of North Carolina at Chapel Hill. 2019. URL: http://gateway.proquest.com/openurl?url_ver=Z39.88-2004&rft_val_fmt=info:ofi/fmt:kev:mtx:dissertation&res_dat=xri:pqm&rft_dat=xri:pqdiss:22587071 (cit. on p. [15](#)).
- [MTT12] J. Metcalfe, D. Tataru, and M. Tohaneanu. “Price’s law on nonstationary space-times”. *Adv. Math.* **230**.3 (2012), 995–1028. DOI: [10.1016/j.aim.2012.03.010](https://doi.org/10.1016/j.aim.2012.03.010) (cit. on p. [15](#)).
- [MW21] K. Morgan and J. Wunsch. *Generalized Price’s law on fractional-order asymptotically flat stationary spacetimes*. 2021. arXiv: [2105.02305](https://arxiv.org/abs/2105.02305) [[math.AP](#)] (cit. on p. [15](#)).
- [Nak94] S. Nakamura. “Low energy asymptotics for Schrödinger operators with slowly decreasing potentials”. *Comm. Math. Phys.* **161**.1 (1994), 63–76. URL: <http://projecteuclid.org/euclid.cmp/1104269792> (cit. on pp. [6](#), [7](#), [10](#)).
- [Olv97] F. W. J. Olver. *Asymptotics and Special Functions*. AKP Classics. A K Peters, Ltd., 1997 (cit. on p. [100](#)).

- [PSRM23] F. Pasqualotto, Y. Shlapentokh-Rothman, and M. Van de Moortel. *The asymptotics of massive fields on stationary spherically symmetric black holes for all angular momenta*. 2023. arXiv: [2303.17767 \[gr-qc\]](#) (cit. on p. [15](#)).
- [RS79] M. Reed and B. Simon. *Methods of modern mathematical physics. III. Scattering theory*. Academic Press [Harcourt Brace Jovanovich, Publishers], New York-London, 1979 (cit. on p. [1](#)).
- [Ski13] E. Skibsted. “Sommerfeld radiation condition at threshold”. *Comm. Partial Differential Equations* **38.9** (2013), 1601–1625. DOI: [10.1080/03605302.2013.804551](#) (cit. on pp. [6](#), [14](#)).
- [Sla60] L. J. Slater. *Confluent Hypergeometric Functions*. Cambridge University Press, 1960 (cit. on p. [100](#)).
- [SN17] J. J. Sakurai and J. Napolitano. *Modern Quantum Mechanics*. 2nd ed. Cambridge University Press, 2017. DOI: [10.1017/9781108499996](#) (cit. on p. [7](#)).
- [SR14] Y. Shlapentokh-Rothman. “Exponentially growing finite energy solutions for the Klein-Gordon equation on sub-extremal Kerr spacetimes”. *Comm. Math. Phys.* **329.3** (2014), 859–891. DOI: [10.1007/s00220-014-2033-x](#) (cit. on p. [15](#)).
- [SRM23] Y. Shlapentokh-Rothman and M. Van de Moortel. *Polynomial time decay for solutions of the Klein-Gordon equation on a subextremal Reissner-Nordström black hole*. 2023. arXiv: [2401.00048 \[gr-qc\]](#) (cit. on p. [15](#)).
- [Tat13] D. Tataru. “Local decay of waves on asymptotically flat stationary space-times”. *Amer. J. Math.* **135.2** (2013), 361–401. DOI: [10.1353/ajm.2013.0012](#). URL: <https://doi.org/10.1353/ajm.2013.0012> (cit. on p. [15](#)).
- [Tay39] W. C. Taylor. “A complete set of asymptotic formulas for the Whittaker function and the Laguerre polynomials”. *Journal of Mathematics and Physics* **18.1-4** (1939), 34–49. DOI: <https://doi.org/10.1002/sapm193918134> (cit. on p. [102](#)).
- [Vas13] A. Vasy. “Microlocal analysis of asymptotically hyperbolic and Kerr-de Sitter spaces (with an appendix by Semyon Dyatlov)”. *Invent. Math.* **194.2** (2013), 381–513. DOI: [10.1007/s00222-012-0446-8](#) (cit. on p. [7](#)).
- [Vas18] ——. “A minicourse on microlocal analysis for wave propagation”. *Asymptotic Analysis in General Relativity*. London Math. Soc. Lecture Note Ser. **443**. Cambridge Univ. Press, 2018, 219–374 (cit. on pp. [6](#), [7](#), [12](#), [20](#), [26](#), [28](#), [38](#), [68](#), [69](#)).
- [Vas21a] ——. “Limiting absorption principle on Riemannian scattering (asymptotically conic) spaces, a Lagrangian approach”. *Comm. Partial Differential Equations* **46.5** (2021), 780–822. DOI: [10.1080/03605302.2020.1857400](#) (cit. on pp. [1](#), [5–7](#), [12–14](#), [17](#), [19–21](#), [28](#), [32](#), [39](#), [40](#), [49–52](#), [55](#), [56](#), [68](#), [81](#), [84–86](#)).
- [Vas21b] ——. “Resolvent near zero energy on Riemannian scattering (asymptotically conic) spaces”. *Pure Appl. Anal.* **3** (2021), 1–74. DOI: [10.2140/paa.2021.3.1](#) (cit. on pp. [2](#), [5](#), [6](#), [21](#)).
- [Vas21c] ——. “Resolvent near zero energy on Riemannian scattering (asymptotically conic) spaces, a Lagrangian approach”. *Comm. Partial Differential Equations* **46.5** (2021), 823–863. DOI: [10.1080/03605302.2020.1857401](#) (cit. on pp. [3](#), [5–7](#), [12–14](#), [19–21](#), [23–28](#)).
- [VZ00] A. Vasy and M. Zworski. “Semiclassical estimates in asymptotically Euclidean scattering”. *Comm. Math. Phys.* **212** (2000), 205–217. DOI: [10.1007/s002200000207](#) (cit. on p. [5](#)).
- [Wan06] X. P. Wang. “Asymptotic expansion in time of the Schrödinger group on conical manifolds”. *Annales de l’Institut Fourier* **56.6** (2006), 1903–1945. DOI: [10.5802/aif.2230](#) (cit. on p. [7](#)).
- [Wun99] J. Wunsch. “Propagation of singularities and growth for Schrödinger operators”. *Duke Math. J.* **98.1** (1999), 137–186. DOI: [10.1215/S0012-7094-99-09804-6](#) (cit. on pp. [12](#), [20](#)).
- [Yaf82] D. R. Yafaev. “The low energy scattering for slowly decreasing potentials”. *Comm. Math. Phys.* **85.2** (1982), 177–196. URL: <http://projecteuclid.org/euclid.cmp/1103921410> (cit. on pp. [1](#), [6](#), [7](#), [14](#)).

Email address: ethanws@mit.edu

DEPARTMENT OF MATHEMATICS, MASSACHUSETTS INSTITUTE OF TECHNOLOGY, MASSACHUSETTS 02139-4307, USA