Recall: 
$$Kn - \lambda \cdot n = f$$

$$\Leftrightarrow \int_{a}^{b} \kappa(x, y) \cdot n(y) \cdot dy - \lambda \cdot n(x) = f(x).$$

What about if 
$$K(x,y) = \sum_{j=1}^{n} \alpha_{j}(x) \cdot \beta_{j}(y)$$
?

$$= \int_{a}^{b} \left( \sum_{j=1}^{n} \alpha_{j}(x) \cdot \beta_{j}(y) \right) \cdot u(y) \cdot dy - \lambda \cdot u(x) = f(x)$$

$$= \sum_{J=1}^{n} \alpha_{J}(x) \int_{a}^{b} \beta_{J}(y) \cdot u(y) dy - \lambda \cdot u(x) = f(x)$$

$$C_{J}$$



Consider that xxx) & By(y) are independent.

$$\sum_{\sigma=0}^{N} \alpha_{\sigma}(x) \cdot C_{\sigma} - \lambda \cdot u(x) = S(x) \qquad (\beta_{i,\sigma})$$

$$\sum_{\sigma=1}^{h} (\alpha_{\sigma}, \beta_{i}) \cdot C_{\sigma} - \lambda \cdot (\beta_{i}, u) = (\beta_{i}, \beta_{i}) = \beta_{i}$$

$$C_{i}$$

$$= \sum (A - \lambda I) \cdot \vec{z} = \vec{\beta}$$

$$\vec{z} = (A - \lambda \cdot I)^{-1} \cdot \vec{\beta}$$

How do we get u(x)?

From  $\sum_{J=1}^{n} \alpha_{J}(x) \cdot \int_{a}^{b} \beta_{J}(y) \cdot u(y) dy - \lambda \cdot u(y) = f(x)$   $C_{J} \leq k_{nown} now.$   $u(x) = \frac{1}{\lambda} \left( \sum_{J=1}^{n} \alpha_{J}(x) \cdot C_{J} - f(x) \right)$ 

Th: Consider the integral equation  $Ku - \lambda \cdot u = S$ , where the Kernel is given by  $K(x,y) = \sum_{s=1}^{n} \alpha_s(x) \cdot \beta_s(y)$  &  $\lambda \neq 0$ . Let A be the matrix  $A = ((\beta_i, \lambda_s)) \cdot If$   $\lambda$  is not an eigenvalue of A, then  $Ku - \lambda \cdot u = S$  has a unique solution given by  $u(x) = \frac{1}{\lambda} \left( \sum_{s=1}^{n} \alpha_s(x) \cdot C_s - S(x) \right)$ .

Th: Consider the equation Ku = f with separable Kernel  $K(x,y) = \sum_{j=1}^{n} \alpha_{j}(x) \cdot \beta_{j}(y)$ . If f is not a linear combination of the  $\alpha_{j}$  then there is no solution; if f is a linear combination of the  $\alpha_{j}$ , then there are infinitely many solutions.

The integral operators K with separable

Th: The integral operator K with separable Kernel  $K(x,y) = \sum_{j=1}^{n} \alpha_{\sigma}(x) \cdot \beta_{\sigma}(y)$  has finitely many nonzero eigenvalues of finite multiplicity, and zero is an eigenvalue of infinite multiplicity

Th: Consider the integral operator  $Ku = \int_a^b K(x,y) \cdot u(y) \cdot dy$  where the Kernel K(x,y) is real, continuous, symmetric d not degenerated. Then K has infinitely many eigenvalues  $\lambda_1, \lambda_2, \ldots,$  each with finite multiplicity an they can be ordered as:

 $0 < \dots \leqslant |\lambda_n| \leqslant \dots \leqslant |\lambda_2| \leqslant |\lambda_i|$ 

With  $\lim_{n\to\infty} \lambda_n = 0$ . Moreover, any square integrable function S(x) can be expanded in terms of the set of orthonormal eigenfunctions  $\beta_K(x)$  associated with the eigenvalues as:

 $f(x) = \sum_{k=1}^{\infty} f_k \cdot \rho_k(x),$ 

where the series converges to S(x) in L'[a,b], and the coefficients are given by

 $S_{k} = (S, \phi_{k}), \quad k = 1, 2, ...$ 

Ex: Let K.M-M.M=5, where K is not degenerated, real, continuous k symmetric.

Find U(x):

Let 
$$u(x) = \sum_{k=1}^{\infty} u_k \cdot \rho_k(x)$$
,  $f(x) = \sum_{k=1}^{\infty} f_k \cdot \rho_k(x)$   
 $u_k = (u, \phi_k)$ ,  $f(x) = (f, \rho_k)$   
 $K \cdot u = f$   
 $K \cdot u = f$ 

$$\sum_{\mathbf{K}} u_{\mathbf{K}} \cdot \lambda_{\mathbf{K}} \cdot p_{\mathbf{k}} - \lambda \cdot \sum_{\mathbf{K}} u_{\mathbf{K}} \cdot p_{\mathbf{k}} = \sum_{\mathbf{K}} S_{\mathbf{K}} \cdot p_{\mathbf{k}} / (p_{\mathbf{i}}, \bullet)$$

$$u_{i} \cdot \lambda_{i} - u_{i} \cdot \lambda = S_{i}$$

$$u_{j} = \frac{S_{i}}{\lambda_{i} - \lambda}$$

$$u(x) = \sum_{k=1}^{\infty} \frac{S_{a}^{b} S(x) \cdot \beta_{k}(x) dx}{\lambda_{k} - \lambda} \cdot \beta_{k}(x)$$

but, 
$$u(x) = \sum_{k=1}^{\infty} \int_{a}^{b} S(x) \cdot \phi_{k}(x) dy$$

Can be written ar

$$u(x) = \int_{a}^{b} \left( \sum_{k=1}^{\infty} \frac{p_{k}(x) \cdot p_{k}(x)}{\lambda_{k} - \lambda} \right) \cdot f(x) dy$$
Green's function

Recall: 
$$K \cdot u - \lambda \cdot u = f$$
  
 $(K - \lambda \cdot I) \cdot u = f$   
 $u = (K - \lambda \cdot I)^{-1} \cdot f$   
 $u = (K - \lambda \cdot I)^{-1} \cdot f$   
 $u(x) = \int_{a}^{b} G(x, y) \cdot f(y) dy$ 

Also, if 
$$\lambda = \lambda m$$
 &  $f_m = (f, f_m) = 0$   

$$= \lambda m = (f, f_m) = 0$$

## Green's Function

What is it?

(1) Mathematically:
$$u(x) = \int_{a}^{b} K(x,y) \cdot S(y) dy$$
the Kernel

Physically:

It is the response of a system when a unit point source is applied to a system.

$$\frac{Ex:}{D} \quad Kn - \lambda \cdot n = 5$$

$$0 \quad n(x) = \int_{a}^{b} G(x, y) \cdot f(y) \, dy$$

$$\begin{array}{c|c}
\hline
\text{(2)} & \overline{K} \cdot G - \lambda \cdot G = \delta(x, y) \\
G = \int_{a}^{b} G(x, y) \cdot \delta(x, y) \, dy \\
= G(x, y)
\end{array}$$

$$\begin{array}{c|c}
S(x, y) = 0, & x \neq y \\
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\end{array}$$

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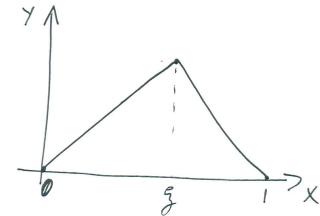
Rule of thumb: Green's function & Inverse of differential operator

$$\underline{\epsilon} \times : -u'(x) = \xi(x), \quad o(x < 1, u(0) = u(1) = 0$$

$$-\widetilde{g}' = S(x, \widetilde{g}) = \widetilde{a} \cdot x + \widetilde{b}$$

$$g(x) = a \cdot x, \quad g(0) = 0, \quad x < g$$

$$g_2(x) = b \cdot (1-x), g(1) = 0, x > 5$$



$$a \cdot \mathbf{j} = b \cdot (1 - \mathbf{j})$$

$$= \sum_{y=\xi}^{g+\xi} \int_{y=\xi}^{g+\xi} \int_{y=\xi}^{g+\xi$$

$$-\left(u'(\xi+\varepsilon)-u'(\xi-\varepsilon)\right)=1$$

$$\lim_{\xi \to 0} -(u'(\xi+\xi)-u'(\xi-\xi)) = -u'(\xi^{\dagger}) + u'(\xi^{-})$$

$$u'(s^{+}) - u'(s^{-}) = -1$$

$$-b - a = -1 = a = 1 - b$$

$$a \cdot 3 = (1 - b) \cdot 3 = b \cdot (1 - 3)$$

$$a = 1 - 3$$

$$b = 3$$

=> 
$$g = (1-3) \cdot x$$
,  $x < 3$   
 $g = (1-x) \cdot g$ ,  $x > 5$ 

$$= g(x,g) = (1-g) \cdot x \cdot H(g-x) + g \cdot (1-x) \cdot H(x-g)$$

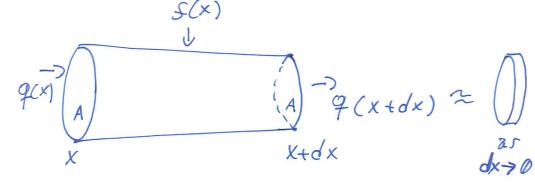
where
$$H(x) = \begin{cases} 1, & x \ge 0 \\ 0, & \text{otherwise} \end{cases}$$

$$= \sum_{x \in X} u(x) = \int_{0}^{1} g(x, y) \cdot S(y) dy$$

Noticethols(x,y) is not a function, it is a distribution but this goes beyond the scape of this course.

## Physical Interpretation

"steady state heat flow problem"



U(x)= "temperature at cross section x" (10 problem) q(x) = "energy flux across the face at x" [ Energy Area. time]

S(x) = "distributed heat source over the length of the bar" [ Energy | Volume · time] q(x)>0 = "flux moves to the right"

by Conservation

$$A \cdot q(x) - A \cdot q(x+dx) + S(x) \cdot A \cdot dx = 0$$

A \( \frac{1}{1000} \)

Outflow

outflow

inside

state

$$\frac{q(x+dx)-q(x)}{dx}=S(x)$$

$$\frac{q(x+dx)-q(x)}{dx}=S(x)$$

$$g'(x) = f(x)$$

Assuming Fourier's heat conduction law (the flux is proportional to the negative temperature gradient)
$$q(x) = -k \cdot u(x)$$

$$= -k \cdot u'(x) = f(x)$$

Note: If this were not a steady state we get the heat equation with a source

$$U_t(x,t) = k \cdot U_{xx}(x,t) + f(x)$$

A

diffusion source
term

50, if u(x,t) only depends on x=> u(x,t) = u(x)=>  $-k \cdot u_{xx}(x) = -k \cdot u'(x) = f(x)$ 

50, let 
$$k=1$$
 d
$$-u'' = S(x), \quad 0 < x < 1$$

$$u(0) = u(1) = 0$$
What about if  $S(x)$  is a "heat source" that only acts at point  $\S$ ?

Bonus: 
$$\int_{\Omega} u \, dv = uv \Big|_{\partial\Omega} - \int_{\Omega} v \cdot du$$
(i) 
$$\int_{\Omega} u = -u^{2} = \int_{\Omega} v \cdot du$$
(2) 
$$\int_{\Omega} G = -G^{2} = \delta \quad delan = 0$$
(1) 
$$\int_{\Omega} u^{2} \cdot G \, dx = \int_{\Omega} \int_{\Omega} G \, dx$$

$$-G \cdot u^{2} \Big|_{\partial\Omega} + \int_{\Omega} u^{2} \cdot G \, dx = \int_{\Omega} \int_{\Omega} G \, dx$$

$$\int_{\Omega} u^{2} \cdot G \, dx = \int_{\Omega} \int_{\Omega} G \, dx = \int_{\Omega} \int_{\Omega} G \, dx$$

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 $u(x) = \int_{\mathcal{L}} f(y) \cdot G(y, x) \, dy$ 

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Exercises:

M(1) = M(e) = 0.

Discuss the solvability of the boundary value problem:  $n' + n^2 \cdot n = f(x)$ , n(0) = n(1) = 0o(x < 1)

(ii) Consider the BVP  $u' - 2 \cdot x \cdot u' = f(x)$ , o(x) = u'(1) = 0Find Green's Function or explain why there isn't one.

iii) Find the inverse of the differential operator  $Lu = -(x^2 \cdot u)$  on 1 < x < e subject to

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