

For non-relativistic motion, we can use (14.20), (14.21) for the approximate angular distribution of radiated power

$$\frac{dP}{d\Omega} = \frac{e^2}{4\pi c} |\mathbf{n} \times (\mathbf{n} \times \dot{\boldsymbol{\beta}})|^2 = \frac{e^2}{4\pi c^3} |\dot{\mathbf{v}}|^2 \sin^2 \Theta \quad \text{where } \Theta \text{ is the angle between } \dot{\mathbf{v}} \text{ and } \mathbf{n} \quad (1)$$

Note that the RHS is evaluated at the retarded time, and the angular distribution is in the sense as described by the second paragraph of section 14.3, i.e., *the energy per unit area per unit time detected at an observation point at time t of radiation emitted by the charge at time $t' = t - R(t')/c$* . I.e., the angular distribution from the perspective of the moving charge, not of the observer at rest.

With this in mind, we can write

$$\mathbf{n} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}} \quad (2)$$

Thus

- For harmonic motion along z -axis, $\dot{\mathbf{v}} = -\omega_0^2 a \cos(\omega_0 t) \hat{\mathbf{z}}$, $\Theta = \theta$, giving

$$\frac{dP}{d\Omega} = \frac{e^2}{4\pi c^3} \omega_0^4 a^2 \cos^2(\omega_0 t) \sin^2 \theta \quad \Rightarrow \quad \left\langle \frac{dP}{d\Omega} \right\rangle = \frac{e^2}{8\pi c^3} \omega_0^4 a^2 \sin^2 \theta \quad (3)$$

- For circular motion around z -axis, we have

$$\dot{\mathbf{v}}(t) = -\omega_0^2 a [\cos(\omega_0 t) \hat{\mathbf{x}} + \sin(\omega_0 t) \hat{\mathbf{y}}] \quad (4)$$

therefore

$$\cos \Theta = -\sin \theta \cos(\phi - \omega_0 t) \quad \Rightarrow \quad \sin^2 \Theta = 1 - \sin^2 \theta \cos^2(\phi - \omega_0 t) \quad (5)$$

giving

$$\frac{dP}{d\Omega} = \frac{e^2}{4\pi c^3} \omega_0^4 a^2 [1 - \sin^2 \theta \cos^2(\phi - \omega_0 t)] \quad \Rightarrow \quad \left\langle \frac{dP}{d\Omega} \right\rangle = \frac{e^2}{4\pi c^3} \omega_0^4 a^2 \left(1 - \frac{\sin^2 \theta}{2}\right) \quad (6)$$

The plots are shown below.

