

1. Rutherford scattering cross section  $d\sigma/d\Omega$  is calculated via the transverse impulse due to transverse electric force  $zeE_{\text{trans}}$ . When we replace the electric charge  $ze$  with the magnetic charge  $g$ ,  $d\sigma/d\Omega$  must be modified to use transverse magnetic force  $gB_{\text{trans}}$ .

(11.152) relates the transverse magnetic force to the transverse electric force

$$B_{\text{trans}} = \beta E_{\text{trans}} \quad \Rightarrow \quad gB_{\text{trans}} = g\beta E_{\text{trans}} = \frac{g\beta}{ze} (zeE_{\text{trans}}) \quad (1)$$

This allows us to write  $d\sigma/d\Omega$  due to magnetic scattering by replacing  $ze$  with  $g\beta$  in the electric case, and equation (13.1) becomes

$$\left. \frac{d\sigma}{d\Omega} \right|_{\text{mag}} = \left( \frac{\beta ge}{2pv} \right)^2 \text{cosec}^4 \left( \frac{\theta}{2} \right) \quad (2)$$

Consequently, the derivation from equation (13.1) to (13.14) can be repeated with the same replacement, and we find that the energy loss per unit length for magnetic scattering is (ignoring spin)

$$\left. \frac{dE}{dx} \right|_{\text{mag}} = 4\pi NZ \frac{g^2 e^2}{mc^2} \ln \left( \frac{2\gamma^2 \beta^2 mc^2}{\hbar \langle \omega \rangle} \right) \quad (3)$$

2. By the Dirac quantization condition (6.153) (in Gaussian units)

$$g = \frac{n\hbar c}{2e} \quad (4)$$

in order to have the same energy loss rate as the magnetic monopole with  $\beta \approx 1$ , we must have

$$ze \approx g \quad \Rightarrow \quad z \approx \frac{g}{e} = \frac{n\hbar c}{2e^2} = \frac{n}{2\alpha} \approx 68.5n \quad (5)$$

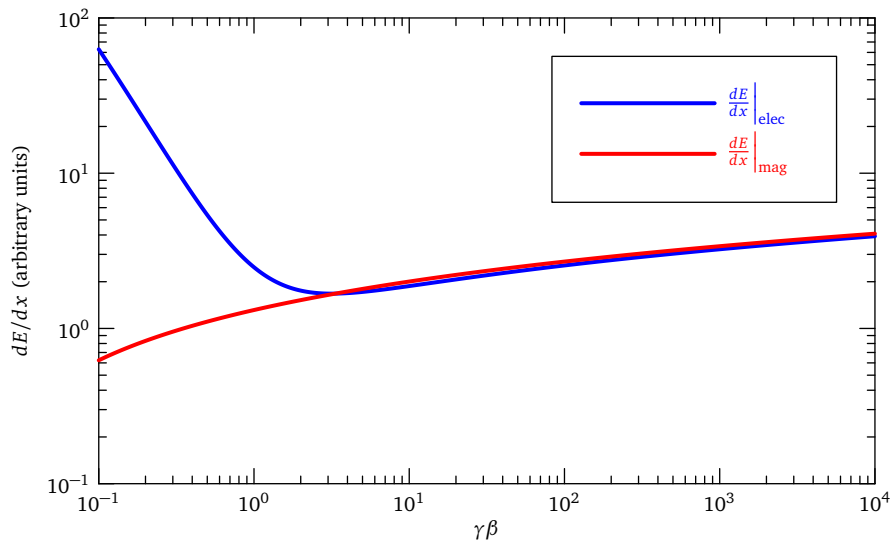
The energy loss per unit length (13.14) in the electric scattering case is

$$\begin{aligned} \left. \frac{dE}{dx} \right|_{\text{elec}} &= \overbrace{4\pi NZ \frac{z^2 e^4}{mc^2}}^{\equiv \eta} \frac{1}{\beta^2} \left[ \ln \left( \frac{2\gamma^2 \beta^2 mc^2}{\hbar \langle \omega \rangle} \right) - \beta^2 \right] \quad \text{let } x = \gamma\beta \\ &= \eta \left( 1 + \frac{1}{x^2} \right) \ln \left( \frac{2x^2 mc^2}{\hbar \langle \omega \rangle} \right) - \eta \end{aligned} \quad (6)$$

The magnetic case (3) is now written

$$\left. \frac{dE}{dx} \right|_{\text{mag}} = \overbrace{4\pi NZ \frac{g^2 e^2}{mc^2}}^{\lambda} \ln \left( \frac{2x^2 mc^2}{\hbar \langle \omega \rangle} \right) = \lambda \ln \left( \frac{2x^2 mc^2}{\hbar \langle \omega \rangle} \right) \quad (7)$$

The log-log plots of  $dE/dx \sim \gamma\beta$  are shown below (with  $\eta = \lambda = 0.15$ ,  $mc^2 = 0.511\text{MeV}$ ,  $\hbar \langle \omega \rangle = 160\text{eV}$ )



Obviously, the energy loss rate for magnetic monopole is monotonically increasing with  $\gamma\beta$ .