# 2D Poisson equation

Dirichlet BV problem

# Dirichlet problem for the 2D Poisson equation

Let u(x) is defined for  $x = (x_1, x_2) \in G = [0, 1] \times [0, 1]$ .

$$\frac{\partial^2 u}{\partial x_1^2} + \frac{\partial^2 u}{\partial x_2^2} = f(x) \tag{1}$$

**Boundary conditions** 

$$u(x) = \mu(x)$$
 for  $x \in \Gamma$ 

# **Grids and grid functions**

Consider a 2D grid

$$\Omega_h = \{x_{kl} = (x_1^k, x_2^l)\}\$$

with

$$x_1^k = h_1 k,$$
  $k = 0, \dots, K$   $(h_1 K = 1)$   
 $x_2^l = h_1 l,$   $l = 0, \dots, L$   $(h_2 L = 1)$ 

Call  $\omega_h$  the set of internal nodes and  $\gamma_h$  the set of boundary nodes.

Define a *grid function*  $y_{kl} \equiv y(x_1^k, x_2^l)$  .

# The five-point stencil

Using the five-point stencil

Eq.(1) generates at 
$$x_{kl} \in \omega_h$$

$$\frac{y_{k+1,l} - 2y_{kl} + y_{k-1,l}}{h_1^2} + \frac{y_{k,l+1} - 2y_{kl} + y_{k,l-1}}{h_2^2} = f_{kl}$$
 (2)

and

$$y_{kl} = \mu(x_{kl})$$
 for  $x_{kl} \in \gamma_h$ 

The scheme approximates Eq. (1) with  $O(h_1^2 + h_2^2)$ .

Rewrite the FDE (2) in the matrix form.

Define the flat index

$$j = k(L+1) + l$$

$$y_j \equiv y_{kl}$$
 with  $k = 0, 1, \cdots, K$   $l = 0, 1, \cdots, L$ 

$$\frac{y_{k+1,l} - 2y_{kl} + y_{k-1,l}}{h_1^2} + \frac{y_{k,l+1} - 2y_{kl} + y_{k,l-1}}{h_2^2} = f_{kl}$$

Indices:

$$(k, l \pm 1) \Rightarrow k(L+1) + (l \pm 1) = j \pm 1$$
  
 $(k \pm 1, l) \Rightarrow (k \pm 1)(L+1) + l = j \pm (L+1)$ 

$$\frac{y_{j+L+1} - 2y_j + y_{j-(L+1)}}{h_1^2} + \frac{y_{j+1} - 2y_j + y_{j-1}}{h_2^2} = f_j$$

Taking, for simplicity,  $\mu(x) = 0$ ,

$$Ay = f$$

with  ${\bf A}$  being "tridiagonal with fringes" of  $(L-1) \times (L-1)$  blocks

- ▶ The size of the **A** matrix is  $O(L^2) = O(h^{-2})$
- ightharpoonup Each row of  ${f A}$  has at most five eq 0 elements, the sparsity is

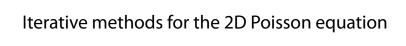
$$\sim 5/L^2 = O(h^{-2})$$

Direct solve with e.g. Gauss elimination requires

$$O((L^2)^3) = O(L^6)$$

flops.

 $\Longrightarrow$  Need a sparse solver.



#### Iterative methods, recap

The canonic form of the two-step iterative methods (a.k.a. *relaxation*)

$$\mathbf{P} \frac{\mathbf{x}^{(n+1)} - \mathbf{x}^{(n)}}{\tau} + \mathbf{A} \mathbf{x}^{(n)} = \mathbf{b}$$

$$\mathbf{A}\mathbf{x} = \mathbf{b} \longrightarrow \mathbf{A} = \mathbf{L} + \mathbf{D} + \mathbf{U}$$
 lower diagonal upper

## Iterative methods, recap

 $lackbox{Jacobi iteration: } \mathbf{D}\mathbf{x}^{(n+1)} + (\mathbf{L} + \mathbf{U})\mathbf{x}^{(n)} = \mathbf{b}$ 

ightharpoonup Seidel iteration:  $\mathbf{D}\mathbf{x}^{(n+1)} + \mathbf{L}\mathbf{x}^{(n+1)} + \mathbf{U}\mathbf{x}^{(n)} = \mathbf{b}$ 

# Jacobi method for the 2D Poisson equation

Take  $\mu(x) = 0$  for clarity.

$$\frac{y_{k+1,l}^n - 2y_{kl}^{n+1} + y_{k-1,l}^n}{h_1^2} + \frac{y_{k,l+1}^n - 2y_{kl}^{n+1} + y_{k,l-1}^n}{h_2^2} = f_{kl}, \quad x_{kl} \in \omega_h$$
$$y_{kl}^n = 0, \quad x_{kl} \in \gamma_h$$

Here  $n=0,1,\cdots$ 

The convergence rate is  $O(h^{-2})$ .

## Seidel iteration in components

$$\mathbf{A} = \begin{pmatrix} a_{11} & a_{12} & a_{13} & \cdots & a_{1m} \\ a_{21} & a_{22} & a_{23} & \cdots & a_{2m} \\ a_{31} & a_{32} & a_{33} & \cdots & a_{3m} \\ & & & \ddots & & \end{pmatrix}$$

$$a_{11}x_1^{(n+1)} + a_{12}x_2^{(n)} + a_{13}x_3^{(n)} + \dots + a_{1m}x_m^{(n)} = b_1$$

$$a_{22}x_2^{(n+1)} + a_{21}x_1^{(n+1)} + a_{23}x_3^{(n)} + \dots + a_{2m}x_m^{(n)} = b_2$$

$$\dots$$

At each iteration, sweep down the system of equations; at each step use  $\mathbf{x}^{(n+1)}$ -s from previous steps: for  $x_i^{(n+1)}$  use  $x_1^{(n+1)} \cdots x_{i-1}^{(n+1)}$  (Jacobi iteration uses  $x_1^{(n)} \cdots x_{i-1}^{(n)}$ )

# Seidel method for the 2D Poisson equation

Take  $\mu(x) = 0$  for clarity.

$$\frac{y_{k+1,l}^n - 2y_{kl}^{n+1} + y_{k-1,l}^{n+1}}{h_1^2} + \frac{y_{k,l+1}^n - 2y_{kl}^{n+1} + y_{k,l-1}^{n+1}}{h_2^2} = f_{kl}, \quad x_{kl} \in \omega_h$$
$$y_{kl}^{n+1} = 0, \quad x_{kl} \in \gamma_h$$

Here  $n=0,1,\cdots$ 

The convergence rate is  $O(h^{-2})$ .

#### **Relaxation methods**

- ► Seidel iteration converges faster then Jacobi's (symmetric **A**)
- ▶ The convergence rate is still  $O(h^{-2})$
- ▶ S.O.R. and/or Chebyshev iterations: can push to  $O(h^{-1})$
- ▶ More specialized methods: can push to  $O(h^{-1/2})$

see, e.g. Part III Chap 5 in S.G.

## Spectral approach to the 2D Poisson FDE

1D fast Fourier transform + tridiagonal solve

# **Dirichlet problem for 2D Poisson equation**

Consider

$$\frac{y_{k+1,l}-2y_{kl}+y_{k-1,l}}{h_1^2}+\frac{y_{k,l+1}-2y_{kl}+y_{k,l-1}}{h_2^2}=f_{kl} \qquad \text{(3)}$$
 
$$y_{kl}=0 \qquad \text{for } x_{kl}\in\gamma_h$$

Here

$$\Omega_h = \{ x_{kl} = (x_1^k, x_2^l) : x_1^k = h_1 k, x_2^l = h_2 l$$

$$k = 0, 1, \dots, N_1, \quad l = 0, 1, \dots, N_2$$

$$N_\alpha h_\alpha = a_\alpha, \alpha = 1, 2 \}$$

#### 1D eigenvalue problem

Consider the 1D eigenvalue problem

$$\frac{g(l+1) - 2g(l) + g(l-1)}{h_2^2} + \lambda g(l) = 0$$
  
$$g(0) = g(N_2) = 0$$

 $l=1,2,\cdots,N_2-1$  and  $l_2N_2=a_2$ . Here  $g(l)\equiv g(x_2^l)$  for brevity.

## 1D eigenvalue problem

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The eigenvalues and eigenfunctions are

$$g_s(l) = \sqrt{\frac{2}{a_2}} \sin \frac{\pi s x_2^l}{a_2}$$
$$\lambda_s = \frac{4}{h_2^2} \sin^2 \frac{\pi s h_2}{2a_2}, \qquad s = 1, \dots N_2 - 1$$

## Fourier expansion over the $x_2$ -eigenfunctions

At a fixed value of k,  $0 < k < N_1$ , expand

$$y_{kl} = \sum_{s=1}^{N_2 - 1} c_s(k) g_s(l)$$
$$f_{kl} = \sum_{s=1}^{N_2 - 1} \widehat{f}_s(k) g_s(l)$$

NB: for each value of k, there are  $N_2-1$  coefficients  $c_s(k)$ .

The Fourier coefficients of 
$$f_{kl}$$
 are  $\hat{f}_s(k) = h_2 \sum_{l=1}^{N_2-1} f_{kl} g_s(l)$ 

## Fourier expansion over the $x_2$ -eigenfunctions

Substitute the expansions into the FDE, use the completeness of the eigensystem:

$$\frac{c_s(k+1) - 2c_s(k) + c_s(k-1)}{h_1^2} - \lambda_s c_s(k) = \hat{f}_s(k)$$
$$c_s(0) = c_s(N_1) = 0$$

$$k = 1, 2, \cdots, N_1 - 1.$$

This is a tridiagonal system at each value of s.

## The spectral algorithm

- 1. Compute the Fourier coefficients of the r.h.s.,  $\widehat{f}_s(k)$ . For each  $k,s=1,\cdots,N_2-1$ . Using the discrete FFT, the complexity is  $O(N_2\ln N_2)$  per k value; the total complexity is  $O(N_1N_2\ln N_2)$
- 2. Solve for  $c_s(k)$ . For each  $s,s=1,\cdots,N_1-1$ , use a tridiagonal solver. The total complexity is  $O(N_1N_2)$
- 3. Restore  $y_{kl}$  via

$$y_{kl} = \sum_{s=1}^{N_2 - 1} c_s(k) g_s(l)$$

With FFT, the complexity is  $O(N_1N_2 \ln N_2)$