Please note, all Griffiths problems come from our class text, the second edition.

## 1. Bloch Sphere

A particle with magnetic moment  $\hat{\mu}$  is placed in a magnetic field **B**. As a result, it has energy

$$\hat{H} = -\hat{\mu} \cdot \mathbf{B}.$$

The magnetic moment of an electron (which is spin  $\frac{1}{2}$ ) is related to its spin by its gyromagnetic ratio  $\gamma$  (which has units of charge over mass):

$$\hat{\mu} = \gamma \hat{\mathbf{S}}.$$

We will perform this problem in the  $S_z$  eigenbasis:  $|+z\rangle$  has eigenvalue  $+\hbar/2$  and  $|-z\rangle$  has eigenvalue  $-\hbar/2$ .

- (a) If we orient the magnetic field such that  $\mathbf{B} = B_0 \hat{\mathbf{z}}$ , (here  $\hat{\mathbf{z}}$  is a unit vector not an operator) what are the energy eigenstates and eigenvalues?
- (b) If we orient the magnetic field in another direction, but we do not change its magnitude, argue why the energy eigenvalues are the same.
- (c) Let us orient the magnetic field such that  $\mathbf{B} = B_0 \hat{\mathbf{n}}$ , where

$$\hat{\mathbf{n}} = \sin \theta \cos \phi \hat{\mathbf{x}} + \sin \theta \sin \phi \hat{\mathbf{y}} + \cos \theta \hat{\mathbf{z}}.$$

Again,  $\hat{\mathbf{n}}$ ,  $\hat{\mathbf{x}}$ ,  $\hat{\mathbf{y}}$ , and  $\hat{\mathbf{z}}$  are unit vectors and not operators. Show that the energy eigenstates are given as

$$\left|+n\right\rangle = \cos\left(\theta/2\right)\left|+z\right\rangle + e^{i\phi}\sin\left(\theta/2\right)\left|-z\right\rangle$$

 $|-n\rangle = \sin(\theta/2) |+z\rangle - e^{i\phi}\cos(\theta/2) |-z\rangle$ .

(d) The state  $|+n\rangle$  gives us a way to map from the surface of a sphere, parametrised by  $(\theta, \phi)$ , to the space of physical states of a two-level quantum mechanical system,

$$|\psi(\theta,\phi)\rangle = \cos(\theta/2) |+z\rangle + e^{i\phi} \sin(\theta/2) |-z\rangle$$
.

Show that this is a bijection between the sphere and space of physical states. That is, show that any normalized state of a two-level system can be uniquely determined by a point on a sphere with coordinates  $(\theta, \phi)$ . [Hint: Begin by showing that any two-level state is equivalent to a  $|+n\rangle$ . Then show that all coordinate pairs  $(\theta, \phi)$  that refer to the same point on a sphere also refer to the same quantum state]. We call the sphere made of these points the Bloch sphere.

(e) Draw the Bloch sphere and place points indicating where the eigenstates of  $S_x$ ,  $S_y$ , and  $S_z$  are mapped to. [Hint: the state

$$\cos\frac{\theta}{2}\left|+z\right\rangle + e^{i\phi}\sin\frac{\theta}{2}\left|-z\right\rangle$$

is at the coordinate  $(\theta, \phi)$ .]

(f) Calculate  $\langle S_x \rangle$ ,  $\langle S_y \rangle$ , and  $\langle S_z \rangle$  for the state  $|\psi(\theta, \phi)\rangle$ .

## 2. Rotation Matrices

Do this problem after problem 1. It will make more sense if you do it that way

- (a) If we rotate the  $|+z\rangle$  state around the  $\hat{\mathbf{x}}$ -axis by angle  $\varphi$ , what will it become? What about  $|-z\rangle$ ? [Hint: Make use of the physical connection from the previous problem. If a state was in the  $|+z\rangle$  state, it would have a specific energy when a magnetic field is pointing in the  $\hat{\mathbf{z}}$  direction. If we rotate out (unphysical) coordinates so that the magnetic field points in the  $\hat{\mathbf{n}}$  direction, the state must still have the same energy, and the state that has that energy is  $|+n\rangle$ . So if we rotate  $\hat{\mathbf{z}}$  to  $\hat{\mathbf{n}}$ , then  $|+z\rangle$  becomes  $|+n\rangle$ .]
- (b) Use the above result to find the matrix representation (in the  $S_z$  eigenbasis) of the operator  $\hat{R}(\varphi \hat{\mathbf{x}})$  that rotates a spin state around the  $\hat{\mathbf{x}}$ -axis by an angle  $\varphi$ .

## 3. Griffiths 4.2

\*Problem 4.2 Use separation of variables in *cartesian* coordinates to solve the infinite *cubical* well (or "particle in a box"):

$$V(x, y, z) = \begin{cases} 0, & \text{if } x, y, z \text{ are all between 0 and } a; \\ \infty, & \text{otherwise.} \end{cases}$$

- (a) Find the stationary states, and the corresponding energies.
- (b) Call the distinct energies  $E_1, E_2, E_3, \ldots$ , in order of increasing energy. Find  $E_1, E_2, E_3, E_4, E_5$ , and  $E_6$ . Determine their degeneracies (that is, the number of different states that share the same energy). Comment: In one dimension degenerate bound states do not occur (see Problem 2.45), but in three dimensions they are very common.
- (c) What is the degeneracy of  $E_{14}$ , and why is this case interesting?

## 4. Griffiths 4.38

\*Problem 4.38 Consider the three-dimensional harmonic oscillator, for which the potential is

$$V(r) = \frac{1}{2}m\omega^2 r^2.$$
 [4.188]

(a) Show that separation of variables in cartesian coordinates turns this into three one-dimensional oscillators, and exploit your knowledge of the latter to determine the allowed energies. *Answer:* 

$$E_n = (n + 3/2)\hbar\omega.$$
 [4.189]

(b) Determine the degeneracy d(n) of  $E_n$ .