

1. Griffiths 2.18

Solution: Using Euler's formula $e^{i\theta} = \cos \theta + i \sin \theta$, and noting that \cos is even and \sin is odd,

$$Ae^{ikx} + Be^{-ikx} = (A + B) \cos kx + (A - B)i \sin kx, \quad (1)$$

so

$$\begin{aligned} C &= A + B, & A &= \frac{1}{2}(C - Di), \\ D &= (A - B)i, & B &= \frac{1}{2}(C + Di). \end{aligned} \quad (2)$$

Note that since $A = \overline{B}$, we still have only two (real) degrees of freedom.

2. Gaussians (based on G1.3)

A Gaussian distribution, parametrised by μ, σ , is given by

$$\rho(x) = A \exp\left(-\frac{(x - \mu)^2}{2\sigma^2}\right). \quad (3)$$

(a) Define

$$I = \int_{-\infty}^{\infty} e^{-x^2} dx. \quad (4)$$

Show that $I = \sqrt{\pi}$.

(Hint: find I^2 by transforming to polar coordinates and then changing variables)

Solution: Substituting $u = r^2$,

$$I^2 = \int e^{-x^2-y^2} dx dy = \int e^{-r^2} r dr d\theta = 2\pi \int_0^\infty r e^{-r^2} dr = \pi \int_0^\infty e^{-u} du = \pi. \quad (5)$$

(b) Hence find the normalisation constant A .

Solution: Substituting $y = \frac{x-\mu}{\sqrt{2}\sigma}$,

$$1 = \int \rho(x) dx = A \int_{-\infty}^{\infty} e^{-(x-\mu)^2/2\sigma^2} dx = \sqrt{2}\sigma A \int_{-\infty}^{\infty} e^{-y^2} dy = \sqrt{2\pi}\sigma A, \quad (6)$$

so $A = \frac{1}{\sqrt{2\pi}\sigma}$.

(c) Find $\langle x \rangle$, $\langle x^2 \rangle$, $\sigma(x)$ of the given Gaussian.

(Hint: $\int x^2 e^{-ax^2} dx = -\frac{\partial}{\partial a} \int e^{-ax^2} dx$)

Solution:

$$\begin{aligned} \langle x \rangle &= \int x \rho(x) dx = \sqrt{2}\sigma A \int (\sqrt{2}\sigma y + \mu) e^{-y^2} dy \\ &= \sqrt{\frac{2}{\pi}}\sigma \int y e^{-y^2} dy + \frac{\mu}{\sqrt{\pi}} \int e^{-y^2} dy = \mu, \end{aligned} \quad (7)$$

since $\int_{-\infty}^{\infty} y e^{-y^2} dy = 0$ as the integrand is odd. Since

$$\int e^{-ax^2} dx = \frac{1}{\sqrt{a}} \int e^{-y^2} dy = \sqrt{\frac{\pi}{a}}, \quad (8)$$

and

$$\int x^2 e^{-ax^2} dx = \int x^2 e^{-ax^2} dx \Big|_{a=1} = -\frac{\partial}{\partial a} \sqrt{\frac{\pi}{a}} \Big|_{a=1} = \frac{1}{2} \sqrt{\frac{\pi}{a^3}} \Big|_{a=1} = \frac{\sqrt{\pi}}{2}, \quad (9)$$

$$\begin{aligned}
\langle x^2 \rangle &= \int x^2 \rho(x) dx = \frac{1}{\sqrt{\pi}} \int (\sqrt{2}\sigma y + \mu)^2 e^{-y^2} dy \\
&= \frac{2\sigma^2}{\sqrt{\pi}} \int y^2 e^{-y^2} dy + \frac{\mu^2}{\sqrt{\pi}} \int e^{-y^2} dy = \sigma^2 + \mu^2,
\end{aligned}
\tag{10}$$

$$\sigma(x) = \sqrt{\langle x^2 \rangle - \langle x \rangle^2} = \sigma. \tag{11}$$

(d) Sketch the Gaussian. Indicate μ and σ on your sketch or describe the effect of changing them.

3. G2.20

Solution:

(a) We are essentially applying the analysis of G2.18 to each term in the sum:

$$\begin{aligned}
f(x) &= \sum_{n=0}^{\infty} (a_n \sin(n\pi x/a) + b_n \cos(n\pi x/a)) \\
&= \sum_{n=0}^{\infty} \frac{1}{2} ((b_n - ia_n)e^{in\pi x/a} + (b_n + ia_n)e^{-in\pi x/a}),
\end{aligned}
\tag{12}$$

$$c_n = \begin{cases} \frac{1}{2}(b_n - ia_n) & n \geq 0, \\ \frac{1}{2}(b_n + ia_n) & n \leq 0. \end{cases} \tag{13}$$

Note that $c_n = \overline{c_{-n}}$, for real input and $c_0 = b_0$ is the constant part of the function (we can always choose $a_0 = 0$ as convention).

(b)

$$\begin{aligned}
\frac{1}{2a} \int_{-a}^a f(x) e^{-in\pi x/a} dx &= \frac{1}{2a} \sum_{m=-\infty}^{\infty} c_m \int_{-a}^a e^{i(m-n)\pi x/a} dx \\
&= \frac{1}{2\pi} \sum_m c_m \int_{-\pi}^{\pi} e^{i(m-n)\theta} d\theta \\
&= \sum_m c_m \delta_{m,n} = c_n,
\end{aligned}
\tag{14}$$

since for any integer $n \neq 0$, $\int e^{in\theta} d\theta$ over a range of 2π cancels out, and for $n = 0$, is simply 2π . Also

$$\delta_{a,b} = \begin{cases} 1 & a = b, \\ 0 & \text{else,} \end{cases} \tag{15}$$

is the Kronecker delta function.

(c)

$$\begin{aligned}
f(x) &= \sum_{n=-\infty}^{\infty} c_n e^{in\pi x/a} = \frac{1}{a} \sqrt{\frac{\pi}{2}} \sum_{n=-\infty}^{\infty} F(k) e^{ikx} = \frac{1}{\sqrt{2\pi}} \sum_n \frac{\pi}{a} F(k) e^{ikx} \\
&= \frac{1}{\sqrt{2\pi}} \sum_n \Delta k F(k) e^{ikx}, \\
F(k) &= \sqrt{\frac{2}{\pi}} a c_n = \frac{1}{\sqrt{2\pi}} \int_{-a}^a f(x) e^{-in\pi x/a} dx = \frac{1}{\sqrt{2\pi}} \int_{-a}^a f(x) e^{-ikx} dx.
\end{aligned}
\tag{16}$$

- (d) As $a \rightarrow \infty$, k becomes continuous rather than discrete (it can take on any value), so we replace $\sum_n \Delta k \rightarrow \int dk$, to get

$$\begin{aligned} f(x) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(k) e^{ikx} dk, \\ F(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx. \end{aligned} \quad (17)$$

4. G2.21

Solution:

(a)

$$\int_{-\infty}^{\infty} \Psi^*(x, 0) \Psi(x, 0) dx = A^2 \int_{-\infty}^{\infty} e^{-2a|x|} dx = 2A^2 \int_0^{\infty} e^{-2ax} dx = \frac{A^2}{a} = 1, \quad (18)$$

so $A = \sqrt{a}$.

(b)

$$\begin{aligned} \phi(k) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \Psi(x, 0) e^{-ikx} dx = \sqrt{\frac{a}{2\pi}} \int_{-\infty}^{\infty} e^{-a|x|-ikx} dx \\ &= \sqrt{\frac{a}{2\pi}} \int_0^{\infty} (e^{-(a+ik)x} + e^{-(a-ik)x}) dx = \sqrt{\frac{a}{2\pi}} \left(\frac{1}{a+ik} + \frac{1}{a-ik} \right) \\ &= \sqrt{\frac{a}{2\pi}} \frac{2a}{a^2 + k^2} = \sqrt{\frac{2}{\pi a}} \left(1 + \left(\frac{k}{a} \right)^2 \right)^{-1}. \end{aligned} \quad (19)$$

(c) Defining $E_k = \hbar^2 k^2 / 2m$,

$$\Psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} e^{-iE_k t / \hbar} dk = \frac{1}{\pi \sqrt{a}} \int_{-\infty}^{\infty} \frac{e^{ikx} e^{-iE_k t / \hbar}}{1 + \left(\frac{k}{a} \right)^2} dk. \quad (20)$$

- (d) Since a sets the scale for k and $\frac{1}{x}$, for small a , the wavefunction is localised in momentum and spread out in position, while for large a it is localised in position and spread out in momentum.