

# INF5620 - Numerical methods for partial differential equations

## Problem set 3

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- For the Github repository containing programs and results, follow this link: <https://github.com/evenmn/master/tree/master/FYSKJM4480/Project1>

## 1 Introduction

The aim of this project is to study Hartree-Fock restricted and unrestricted. Many-particle system, the system is defined by the total Slater determinant  $|\Phi\rangle = |\phi_1\phi_2\cdots\phi_N\rangle$  where  $\phi_i$  is a single particle function (SPF). Electrons

## 2 Theory

For both RHF and UHF we can split up the wave function in a spin part and a position part:

$$\phi_{p,\sigma}(x) = \varphi_p(\vec{r})\chi_\sigma(s) \quad (1)$$

where each orbit  $\varphi$  is double occupied.

$$|\Phi_{RHF}\rangle = |\phi_{1\uparrow}\phi_{1\downarrow}\cdots\phi_{N\uparrow}\phi_{N\downarrow}\rangle = |\varphi_1\cdots\varphi_N\bar{\varphi}_1\cdots\bar{\varphi}_N\rangle \quad (2)$$

$$\phi_{p,\sigma}(x) = \varphi_p^\sigma(\vec{r})\chi_\sigma(s) \quad (3)$$

Total Slater determinant:

$$|\Phi_{RHF}\rangle = |\varphi_1^\uparrow\cdots\varphi_{N\uparrow}^\uparrow\bar{\varphi}_1^\downarrow\cdots\bar{\varphi}_{N\downarrow}^\downarrow\rangle \quad (4)$$

Write down some general formulas here, Wave functions for restricted and unrestricted HF ex.

## 2.1 Energy expression

The Hamiltonian of the system is assumed to on the form

$$\hat{H} = \hat{h}(i) + \hat{w}(i, j) \quad (5)$$

where  $\hat{h}$  is the kinetic energy operator and  $\hat{w}$  is the potential energy. We want to find the total energy of the system, which is given by  $\langle \Phi | \hat{H} | \Phi \rangle$ . If we consider the kinetic energy, we can easily see that the kinetic energy of particle  $i$  is given by  $T_i = \langle \phi_i | \hat{h} | \phi_i \rangle$ . For the potential it gets more complicated, since we have interaction between all the particles. We also need to make sure that the bracket is anti-symmetric as we look at a system of fermions (electrons). For one electron pair  $i$  and  $j$  we therefore get the potential

$$V_{i,j} = \frac{1}{2} \langle \phi_i \phi_j | \hat{w} | \phi_i \phi_j - \phi_j \phi_i \rangle$$

with negative sign in the ket to get an anti-symmetric bracket. Consider now a system of  $N + 1$  electrons, the total energy is given by

$$E = \langle \Phi | (\hat{h}(i) + \hat{w}(i, j)) | \Phi \rangle = \sum_{i=0}^N \langle \phi_i | \hat{h} | \phi_i \rangle + \frac{1}{2} \sum_{i=0}^N \sum_{\substack{j=0 \\ j \neq i}}^N \langle \phi_i \phi_j | \hat{w} | \phi_i \phi_j - \phi_j \phi_i \rangle.$$

We can apply this formula to find an expression for the UHF-energy. By plugging in the general UHF wave function from above, we obtain

$$\begin{aligned} E_{UHF} &= \langle \Phi_{UHF} | \hat{H} | \Phi_{UHF} \rangle \\ &= \sum_{i,\sigma} \langle \varphi_i^\sigma | \hat{h} | \varphi_i^\sigma \rangle + \frac{1}{2} \sum_{i,j,\sigma,\tau} \langle \varphi_i^\sigma \varphi_j^\tau | \hat{w} | \varphi_i^\sigma \varphi_j^\tau - \varphi_j^\tau \varphi_i^\sigma \rangle \\ &= \sum_{i,\sigma} \langle \varphi_i^\sigma | \hat{h} | \varphi_i^\sigma \rangle + \frac{1}{2} \sum_{i,j,\sigma,\tau} \langle \varphi_i^\sigma \varphi_j^\tau | \hat{w} | \varphi_i^\sigma \varphi_j^\tau \rangle - \frac{1}{2} \sum_{i,j,\sigma} \langle \varphi_i^\sigma \varphi_j^\sigma | \hat{w} | \varphi_j^\sigma \varphi_i^\sigma \rangle. \quad (6) \end{aligned}$$

Note that we do not have a sum over  $\tau$  in the last term, but both  $\tau$ s are replaced by  $\sigma$ s. This is because the last term represents the exchange energy, and only particles with the same spin can exchange energy. Here we can also see why the anti-symmetric bra needs to have negative sign: fermions cause repulsive "exchange force", and with the standard positive force direction the exchange energy should be negative for fermions.

## 2.2 UHF equations

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Recall the UHF energy from equation (6) and write out  $\delta E_{UHF}$ :

$$\begin{aligned} \delta E_{UHF} = & \sum_{i,\sigma} \langle \delta \varphi_i^\sigma | \hat{h} | \varphi_i^\sigma \rangle + \frac{1}{2} \sum_{i,j,\sigma,\tau} \langle \delta \varphi_i^\sigma \varphi_j^\tau | \hat{w} | \varphi_i^\sigma \varphi_j^\tau \rangle \\ & - \frac{1}{2} \sum_{i,j,\sigma} \langle \delta \varphi_i^\sigma \varphi_j^\sigma | \hat{w} | \varphi_j^\sigma \varphi_i^\sigma \rangle. \end{aligned} \quad (7)$$

## 2.3 Roothan-Hall equations

## 2.4 Restricted Hartree-Fock

With the restrictions  $U = U^\uparrow = U^\downarrow$  and  $N = N^\uparrow = N^\downarrow$  the UHF equations are still valid. The Roothan-Hall equation for RHF is

$$F(D)U = SU\epsilon \quad (8)$$

where the Fock-operator is  $F = h + J(D) - 1/2K(D)$  and the density matrix reads  $D = 2U_{occ}U_{occ}^H$ . Similarly for UHF we have the Roothan-Hall equation

$$F^\sigma(D^\uparrow, D^\downarrow)U^\sigma = SU^\sigma\epsilon^\sigma \quad (9)$$

with  $F = h + J(D^\uparrow + D^\downarrow) - K(D^\sigma)$  and  $D = U_{occ}^\uparrow(U_{occ}^\uparrow)^H + U_{occ}^\downarrow(U_{occ}^\downarrow)^H$ . As we can see, RHF is like a more general version of UHF with  $D^\uparrow = D^\downarrow = D$  and  $U_{occ}^\uparrow = U_{occ}^\downarrow = U$ , and a solution of RHF Roothan-Hall equations will therefore also be a solution of the UHF Roothan-Hall equations.