# 1 Hydrodynamic stellar structure equations (non-local and time-dependent)

Below is a set of hydrodynamic stellar structure equations derived from RANS (viscosity explicitly neglected), where red terms are the ones used in classical approach:

$$\frac{\partial_r \overline{m}}{\partial r} = 4\pi r^2 \overline{\rho} + (4\pi r^3 / 3\widetilde{u}_r) \left[ -\nabla_r f_\rho + (f_\rho / \overline{\rho}) \partial_r \overline{\rho} - \overline{\rho} \overline{d} - \partial_t \overline{\rho} \right]$$
(1)

$$\frac{\partial_r \overline{P}}{\overline{P}} = \overline{\rho} \widetilde{g} - \overline{\rho} \partial_t \widetilde{u}_r - \nabla_r \widetilde{R}_{rr} - \overline{G}_r^M - \overline{\rho} \widetilde{u}_r \partial_r \widetilde{u}_r$$
(2)

$$\frac{\partial_r \widetilde{L}}{\partial r} = 4\pi r^2 \overline{\rho} \widetilde{\epsilon}_{nuc} + 4\pi r^2 \left[ -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{Pd} - \widetilde{R}_{ir} \partial_r \widetilde{u}_i + W_b + \overline{\rho} \widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2 - \overline{\rho} \partial_t \widetilde{\epsilon}_t \right] + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r$$
(3)

$$\frac{\partial_r \overline{T}}{T} = (1/\overline{u}_r) \left[ -\nabla_r f_T + (1 - \Gamma_3) \overline{T} \ \overline{d} + (2 - \Gamma_3) \overline{T'd'} + \epsilon_{nuc}/c_v + \nabla \cdot f_{th}/(\rho c_v) - \partial_t T \right]$$

$$\tag{4}$$

$$\frac{\partial_t \widetilde{X}_i}{\partial_t \widetilde{X}_i} = \frac{\widetilde{X}_i^{nuc}}{(1/\overline{\rho})} \nabla_r f_i - \widetilde{u}_r \partial_r \widetilde{X}_i$$
(5)

### 1.1 Continuity Equation

#### Derivation

Using full 3D hydrodynamic continuity equation, we derive its mean field counterpart in the following way:

$$\partial_{t}\rho + \nabla\rho\mathbf{u} = 0$$

$$\partial_{t}\overline{\rho} + \overline{u}_{r}\partial_{r}\overline{\rho} = -\nabla_{r}\overline{\rho'u'_{r}} - \overline{\rho}\overline{d}$$

$$\partial_{t}\overline{\rho} + \overline{u''}_{r}\partial_{r}\overline{\rho} + \widetilde{u}_{r}\partial_{r}\overline{\rho} = -\nabla_{r}\overline{\rho'u'_{r}} - \overline{\rho}\overline{d}$$

$$\partial_{t}\overline{\rho} + \widetilde{u}_{r}\partial_{r}\overline{\rho} = -\nabla_{r}\overline{\rho'u'_{r}} + (\overline{\rho'u'_{r}}/\overline{\rho})\partial_{r}\overline{\rho} - \overline{\rho}\overline{d}$$

$$\partial_{t}\overline{\rho} + \widetilde{u}_{r}\partial_{r}\overline{\rho} = -\nabla_{r}f_{\rho} + (f_{\rho}/\overline{\rho})\partial_{r}\overline{\rho} - \overline{\rho}\overline{d}$$

$$\widetilde{D}_{t}\overline{\rho} = -\nabla_{r}f_{\rho} + (f_{\rho}/\overline{\rho})\partial_{r}\overline{\rho} - \overline{\rho}\overline{d}$$

For the derivation, we used the following identities:  $\overline{\rho u_r} = \overline{\rho' u_r'} + \overline{\rho u_r}$  and  $\overline{u''}_r = \overline{u_r} - \widetilde{u_r}$  and  $\overline{\rho} \overline{u''}_r = -f_\rho$  and  $f_\rho = \overline{\rho' u_r'}$  (turbulent mass flux)

From there, let us now express gradient of mean density  $\partial_r \overline{\rho}$ . We get:

$$\partial_r \overline{\rho} = -(1/\widetilde{u}_r) \left( \nabla_r f_\rho + (f_\rho/\overline{\rho}) \partial_r \overline{\rho} - \overline{\rho} \overline{d} - \partial_t \overline{\rho} \right) \tag{6}$$

We want to derive hydrodynamic continuity equation into the form familiar from its classical form. Therefore, let us now find relation between  $\partial_r \overline{\rho}$  and  $\partial_r \overline{m}$  using total differentials of  $\rho$  and m. We know, that  $\rho = \rho(r, t)$  and m = m(r, t) and:

$$d\rho = \partial_r \rho \ dr + \partial_t \rho \ dt$$
$$dm = \partial_r m \ dr + \partial_t m \ dt$$

Let us now transform the  $d\rho$  equation to dm equation by multiplying it by volume  $V = 4\pi r^3/3$  and few algebraic modifications:

$$Vd\rho = V\partial_r \rho \ dr + V\partial_t \rho \ dt$$

$$dm - \rho dV = V\partial_r \rho \ dr + V\partial_t \rho \ dt$$

$$dm - 4\pi r^2 \rho dr = V\partial_r \rho \ dr + V\partial_t \rho \ dt$$

$$dm = 4\pi r^2 \rho dr + V\partial_r \rho \ dr + V\partial_t \rho \ dt$$

$$dm = (4\pi r^2 \rho + V\partial_r \rho) dr + V\partial_t \rho \ dt$$

By comparing this result to  $dm = \partial_r m \, dr + \partial_t m \, dt$  we get:

$$(4\pi r^2 \rho + V \partial_r \rho) dr + V \partial_t \rho dt = \partial_r m dr + \partial_t m dt$$

and

$$\partial_r m = V \partial_r \rho + 4\pi r^2 \rho \tag{7}$$

By space-time averaging and using equation 6 we get the desired form of the continuity equation for stellar evolution.

$$\partial_r \overline{m} = 4\pi r^2 \overline{\rho} + (4\pi r^3 / 3\widetilde{u}_r) \left( -\nabla_r f_\rho + (f_\rho / \overline{\rho}) \partial_r \overline{\rho} - \overline{\rho} \overline{d} - \partial_t \overline{\rho} \right) \tag{8}$$

$$\partial_r \overline{m} = \underbrace{4\pi r^2 \overline{\rho}}_{\text{density distribution}} + (4\pi r^3/3\widetilde{u}_r) \left( \underbrace{-\nabla_r f_\rho}_{\text{transport of turbulent density field}} + \underbrace{(f_\rho/\overline{\rho})\partial_r \overline{\rho}}_{\text{down-gradient density source/sink term}} - \underbrace{\overline{\rho} \overline{d}}_{\text{compressibility effects}} - \underbrace{\partial_t \overline{\rho}}_{\text{time-dependence}} \right)$$
 (9)

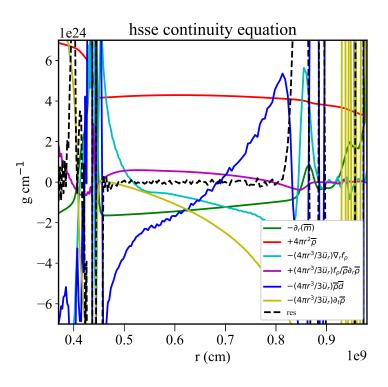


Figure 1: Hydrodynamic stellar structure continuity equation.

# 1.2 Momentum Equation

### Derivation

We start from the RANS equation for X momentum and modify it into form familiar from classical stellar evolution theory.

$$\overline{\rho}\widetilde{D}_{t}\widetilde{u}_{r} = -\nabla_{r}\widetilde{R}_{rr} - \overline{G_{r}^{M}} - \partial_{r}\overline{P} + \overline{\rho}\widetilde{g}_{r}$$

$$\tag{10}$$

$$\overline{\rho}\partial_t \widetilde{u}_r + \overline{\rho}\widetilde{u}_r \partial_r \widetilde{u}_r = -\nabla_r \widetilde{R}_{rr} - \overline{G_r^M} - \partial_r \overline{P} + \overline{\rho}\widetilde{g}_r$$

$$\tag{11}$$

$$\partial_r \overline{P} = \overline{\rho} \widetilde{g}_r - \overline{\rho} \partial_t \widetilde{u}_r - \nabla_r \widetilde{R}_{rr} - \overline{G}_r^M - \overline{\rho} \widetilde{u}_r \partial_r \widetilde{u}_r$$
(12)

$$\partial_r \overline{P} = \underbrace{\overline{\rho} \widetilde{g}_r}_{\text{gravity}} - \underbrace{\overline{\rho} \partial_t \widetilde{u}_r}_{\text{acceleration due to expansion}} - \underbrace{\nabla_r \widetilde{R}_{rr}}_{\text{transport of turbulent velocity field}} - \underbrace{\overline{G}_r^M}_{\text{centrifugal forces}} - \underbrace{\overline{\rho} \widetilde{u}_r \partial_r \widetilde{u}_r}_{\text{advection due to expansion}}$$

$$(13)$$

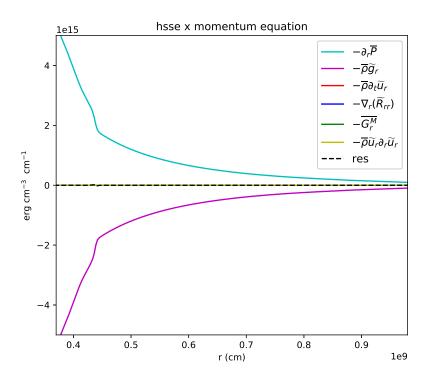


Figure 2: Hydrodynamic stellar structure momentum equation.

### 1.3 Luminosity Equation

$$\overline{\rho}\widetilde{D}_{t}\widetilde{\epsilon}_{t} = -\nabla_{r}(f_{i} + f_{th} + f_{K} + f_{p}) - \overline{P} \ \overline{d} - \widetilde{R}_{ir}\partial_{r}\widetilde{u}_{r} + W_{b} + \overline{\rho}\widetilde{\epsilon}_{nuc} + \overline{\rho}\widetilde{D}_{t}\widetilde{u}_{i}\widetilde{u}_{i}/2$$

$$\tag{14}$$

$$\overline{\rho}\partial_t \widetilde{\epsilon}_t + \overline{\rho}\widetilde{u}_r \partial_r \widetilde{\epsilon}_t = -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{P} \ \overline{d} - \widetilde{R}_{ir} \partial_r \widetilde{u}_r + W_b + \overline{\rho}\widetilde{\epsilon}_{nuc} + \overline{\rho}\widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2$$

$$\tag{15}$$

$$4\pi r^2 \overline{\rho} \partial_t \widetilde{\epsilon}_t + 4\pi r^2 \overline{\rho} \widetilde{u}_r \partial_r \widetilde{\epsilon}_t = 4\pi r^2 \left[ -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{P} \ \overline{d} - \widetilde{R}_{ir} \partial_r \widetilde{u}_r + W_b + \overline{\rho} \widetilde{\epsilon}_{nuc} + \overline{\rho} \widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2 \right]$$

$$\tag{16}$$

$$4\pi r^2 \overline{\rho} \partial_t \widetilde{\epsilon}_t + \underbrace{\partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r \widetilde{\epsilon}_t}_{\widetilde{\epsilon}_t} - \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r = 4\pi r^2 \left[ -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{P} \ \overline{d} - \widetilde{R}_{ir} \partial_r \widetilde{u}_r + W_b + \overline{\rho} \widetilde{\epsilon}_{nuc} + \overline{\rho} \widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2 \right]$$

$$(17)$$

$$\partial_r \widetilde{L} = 4\pi r^2 \left[ -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{P} \ \overline{d} - \widetilde{R}_{ir} \partial_r \widetilde{u}_r + W_b + \overline{\rho} \widetilde{\epsilon}_{nuc} + \overline{\rho} \widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2 - \overline{\rho} \partial_t \widetilde{\epsilon}_t \right] + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r$$
(18)

Or

$$\partial_r \widetilde{L} = 4\pi r^2 \overline{\rho} \widetilde{\epsilon}_{nuc} + 4\pi r^2 \left[ -\nabla_r (f_i + f_{th} + f_K + f_p) - \overline{P} \ \overline{d} - \widetilde{R}_{ir} \partial_r \widetilde{u}_r + W_b + \overline{\rho} \widetilde{D}_t \widetilde{u}_i \widetilde{u}_i / 2 - \overline{\rho} \partial_t \widetilde{\epsilon}_t \right] + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r$$

$$\tag{19}$$

Some term description:

$$\partial_{r}\widetilde{L} = \underbrace{4\pi r^{2}\overline{\rho}\widetilde{\epsilon}_{nuc}}_{\text{nuclear}} + 4\pi r^{2} \left[ \underbrace{-\nabla_{r}(f_{i} + f_{K} + f_{p} + f_{th})}_{\text{transport of internal energy, kinetic energy, pressure and heat due to conduction and radiation} - \underbrace{\overline{P}\ \overline{d}}_{\text{compressibility}} - \underbrace{\widetilde{R}_{ir}\partial_{r}\widetilde{u}_{r}}_{\text{down-gradient source/sink term}} + \underbrace{W_{b}}_{\text{buoyancy work}} + \underbrace{\overline{\rho}\widetilde{D}_{t}\widetilde{u}_{i}\widetilde{u}_{i}/2 - \overline{\rho}\partial_{t}\widetilde{\epsilon}_{t}}_{\text{time-dependence}} + \widetilde{\epsilon}_{t}\partial_{r}4\pi r^{2}\overline{\rho}\widetilde{u}_{r}}_{\text{compressibility}} \right]$$

$$(20)$$

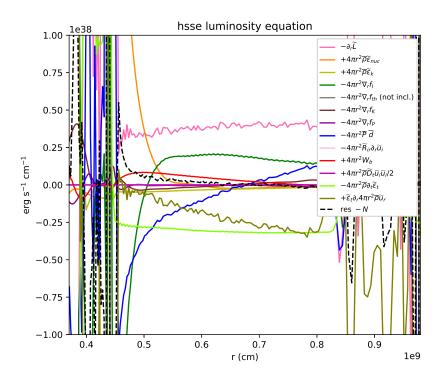


Figure 3: Hydrodynamic stellar structure luminosity equation.

## 1.4 Temperature Equation

#### Derivation

We start from the RANS equation for temperature evolution. At the end, there will be no resemblance to the temperature equation of the classical stellar evolution theory.

$$\overline{D}_{t}\overline{T} = -\nabla_{r}f_{T} + (1 - \Gamma_{3})\overline{T} \ \overline{d} + (2 - \Gamma_{3})\overline{T'd'} + \overline{(\nabla \cdot f_{th})/\rho c_{v}} + \overline{\epsilon_{\text{nuc}}/c_{v}}$$

$$\partial_{t}\overline{T} + \overline{u}_{r}\partial_{r}\overline{T} = -\nabla_{r}f_{T} + (1 - \Gamma_{3})\overline{T} \ \overline{d} + (2 - \Gamma_{3})\overline{T'd'} + \overline{(\nabla \cdot f_{th})/\rho c_{v}} + \overline{\epsilon_{\text{nuc}}/c_{v}}$$

$$\overline{u}_{r}\partial_{r}\overline{T} = -\nabla_{r}f_{T} + (1 - \Gamma_{3})\overline{T} \ \overline{d} + (2 - \Gamma_{3})\overline{T'd'} + \overline{(\nabla \cdot f_{th})/\rho c_{v}} + \overline{\epsilon_{\text{nuc}}/c_{v}} - \partial_{t}\overline{T}$$

$$\partial_r \overline{T} = -\left(1/\overline{u}_r\right) \left(\nabla_r f_T + (1-\Gamma_3)\overline{T}\ \overline{d} + (2-\Gamma_3)\overline{T'd'} + \overline{(\nabla \cdot f_{th})/\rho c_v} + \overline{\epsilon_{\text{nuc}}/c_v} - \partial_t \overline{T}\right)$$

$$\partial_{r}\overline{T} = -(1/\overline{u}_{r})(\underbrace{\nabla_{r}f_{T}}_{\text{transport of turbulent temperature field}} + \underbrace{(1-\Gamma_{3})\overline{T}\ \overline{d} + (2-\Gamma_{3})\overline{T'd'}}_{\text{compressibility effects}} + \underbrace{(\overline{\nabla \cdot f_{th}})/\rho c_{v}}_{\text{source/sink term due to thermal transport}} + \underbrace{\overline{\epsilon_{\text{nuc}}/c_{v}}}_{\text{source/sink due to nuclear burning}} - \underbrace{\partial_{t}\overline{T}}_{\text{time-dependence}}$$

$$(21)$$

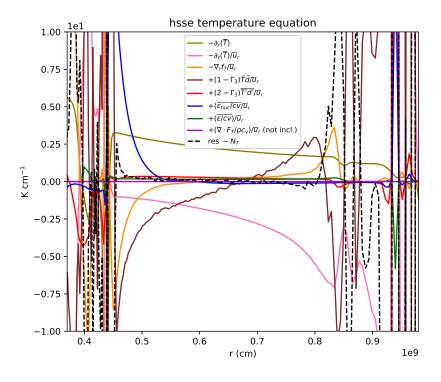


Figure 4: Hydrodynamic stellar structure temperature equation.

## 1.5 Composition Equation

### Derivation

We start from the RANS equation for compostion and modify it into form familiar from classical stellar evolution theory.

$$\begin{split} \overline{\rho}\widetilde{D}_{t}\widetilde{X}_{\alpha} &= -\nabla_{r}f_{\alpha} + \overline{\rho}\widetilde{\dot{X}_{\alpha}^{nuc}} \\ \overline{\rho}\partial_{t}\widetilde{X}_{\alpha} + \overline{\rho}\widetilde{u}_{r}\partial_{r}\widetilde{X}_{\alpha} &= -\nabla_{r}f_{\alpha} + \overline{\rho}\widetilde{\dot{X}_{\alpha}^{nuc}} \\ \overline{\rho}\partial_{t}\widetilde{X}_{\alpha} &= -\nabla_{r}f_{\alpha} + \overline{\rho}\widetilde{\dot{X}_{\alpha}^{nuc}} - \overline{\rho}\widetilde{u}_{r}\partial_{r}\widetilde{X}_{\alpha} \\ \partial_{t}\widetilde{X}_{\alpha} &= \widetilde{\dot{X}}_{\alpha}^{nuc} - (1/\overline{\rho})\nabla_{r}f_{\alpha} - \widetilde{u}_{r}\partial_{r}\widetilde{X}_{\alpha} \end{split}$$

$$\partial_t \widetilde{X}_{\alpha} = \underbrace{\widetilde{X}_{\alpha}^{nuc}}_{\text{nuclear burning transport of turbulent composition field}} - \underbrace{\widetilde{u}_r \partial_r \widetilde{X}_{\alpha}}_{\text{advection due to expansion}}$$
(22)

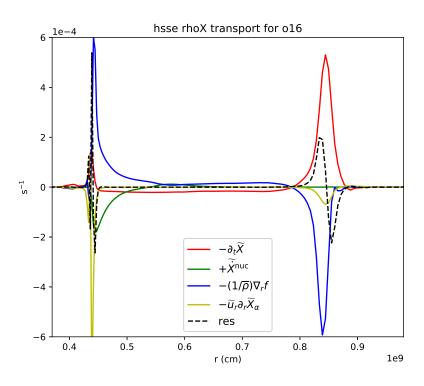


Figure 5: Hydrodynamic stellar structure composition transport equation.

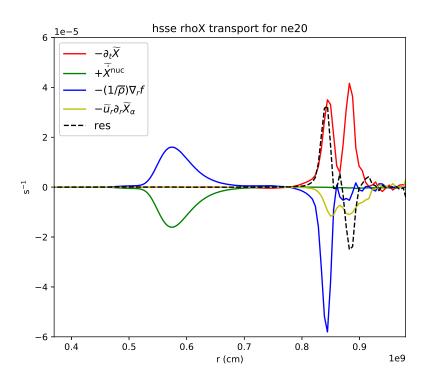


Figure 6: Hydrodynamic stellar structure composition transport equation.

Table 1: Definitions:

$\rho$ density
$m = \rho V = \rho \frac{4}{3}\pi r^3$ mass
T temperature
P pressure
$u_r, u_\theta, u_\phi$ velocity components
$\mathbf{u} = u(u_r, u_\theta, u_\phi)$ velocity
$j_z = r \sin \theta \ u_\phi$ specific angular momentum
$d = \nabla \cdot \mathbf{u}$ dilatation
$\epsilon_I$ specific internal energy
h specific enthalpy
$k = (1/2)\widetilde{u_i''u_i''}$ turbulent kinetic energy
$\epsilon_k$ specific kinetic energy
$\epsilon_t$ specific total energy
s specific entropy
$v = 1/\rho$ specific volume
$X_{\alpha}$ mass fraction of isotope $\alpha$
$\dot{X}_{lpha}^{\mathrm nuc}$ rate of change of $X_{lpha}$
$A_{\alpha}$ number of nucleons in isotope $\alpha$
$Z_{\alpha}$ charge of isotope $\alpha$
A mean number of nucleons per isotope
Z mean charge per isotope

 $f_P = \overline{P'u'_n}$  acoustic flux

Table 2: Definitions (continued):

$$\begin{split} f_I &= \overline{\rho_L^{\prime\prime}u_l^{\prime\prime}} & \text{ internal energy flux} & f_\alpha &= \overline{\rho X_\alpha^{\prime\prime}u_l^{\prime\prime}} X_\alpha \text{ flux} \\ f_s &= \overline{\rho s^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ entropy flux} & f_{jz} &= \overline{\rho j_z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ angular momentum flux} \\ f_T &= \overline{u_r^{\prime\prime}T^{\prime\prime}} & \text{ turbulent heat flux} & f_A &= \overline{\rho A^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_h &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ enthalpy flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ enthalpy flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ enthalpy flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ A (mean number of nucleons per isotope) flux} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ Numerical effect} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho Z^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ numerical effect} \\ f_{\tau} &= \overline{\rho h^{\prime\prime\prime}u_l^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ viscous flux} & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{\rho L^{\prime\prime\prime}u_l^{\prime\prime}u_l^{\prime\prime}} & \text{ radial flux of } f_Z & f_Z &= \overline{$$

#### Table 3: Definitions (continued):

$$\begin{array}{lll} \mathcal{G}_k^{\theta} = -(1/2)\overline{G_{Q_\theta}^{\theta}} & \overline{u_\theta''G_\phi^{\theta}} & \mathcal{N}_b \text{ numerical effect} \\ \mathcal{G}_k^{\phi} = -(1/2)\overline{G_{Q_\phi}^{\theta}} & \overline{u_\phi''G_\phi^{\theta}} & \mathcal{N}_{f1} = -\nabla_r(\overline{e_1''\tau_r'}) + \overline{u_r''\tau_{ij}\partial_i u_j} - \varepsilon_I \text{ numerical effect} \\ \mathcal{G}_k^{h} = +\mathcal{G}_k^{\theta} + \mathcal{G}_k^{\phi} & \mathcal{N}_{fh} = -\nabla_r(\overline{h''\tau_{rr}'}) + \overline{u_r''\tau_{ij}\partial_i u_j} - \overline{u_r''\nabla_i u_i\tau_{ji}} - \varepsilon_h \text{ numerical effect} \\ \mathcal{G}_a = +\overline{\rho'vG_\rho^{M}} & \mathcal{N}_{fs} = -\nabla_r(\overline{s''\tau_{rr}'}) + \overline{u_r''\tau_{ij}\partial_i u_j} - \overline{u_r''\nabla_i u_i\tau_{ji}} - \varepsilon_h \text{ numerical effect} \\ \mathcal{G}_I = -\overline{G_I^r} - \overline{\epsilon_I''G_\rho^{M}} & \mathcal{N}_{fs} = -\nabla_r(\overline{s''\tau_{rr}'}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_A = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{Z''\tau_{rr}'}) - \varepsilon_Z \text{ numerical effect} \\ \mathcal{G}_A = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{a_r''\tau_{rr}'}) - \varepsilon_Z \text{ numerical effect} \\ \mathcal{G}_Z = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{a_r''\tau_{rr}'}) - \varepsilon_{\alpha} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\alpha} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f1} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f1} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_I^r} - \overline{A_r''G_\rho^{M}} & \mathcal{N}_{f2} = -\nabla_r(\overline{g_1''\tau_{rr}'}) - \varepsilon_{\beta} \text{ n$$

Table 4: Definitions (continued):

$$\begin{split} \varepsilon_k^r &= \overline{\tau_{rr}'} \partial_r u_l''' + \overline{\tau_{r\theta}'(1/r) \partial_\theta u_l''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi u_l''} \\ \varepsilon_k^\theta &= \overline{\tau_{lr}'} \partial_r u_\theta''' + \overline{\tau_{l\theta}'(1/r) \partial_\theta u_\theta''} + \overline{\tau_{l\theta}'(1/r \sin \theta) \partial_\phi u_\theta''} \\ \varepsilon_k^\theta &= \overline{\tau_{lr}'} \partial_r u_\theta'' + \overline{\tau_{l\theta}'(1/r) \partial_\theta u_\theta''} + \overline{\tau_{l\theta}'(1/r \sin \theta) \partial_\phi u_\theta''} \\ \varepsilon_k^\theta &= \overline{\tau_{lr}'} \partial_r u_\theta'' + \overline{\tau_{l\theta}'(1/r) \partial_\theta u_\theta''} + \overline{\tau_{l\theta}'(1/r \sin \theta) \partial_\phi u_\theta''} \\ \varepsilon_k &= \varepsilon_k^r + \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k^\theta &= \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k &= \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k &= \overline{\varepsilon_l''} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'(1/r) \partial_\theta \varepsilon_l''} + \overline{\tau_{r\theta}'(1/r \sin \theta) \partial_\phi \varepsilon_l''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'(1/r) \partial_\theta \varepsilon_l''} + \overline{\tau_{r\theta}'(1/r \sin \theta) \partial_\phi \varepsilon_l''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{X_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta X_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi X_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi X_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi X_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi X_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi X_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\theta}'(1/r) \partial_\theta Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\phi}'(1/r \sin \theta) \partial_\phi Z_\theta''} \\ \varepsilon_k &= \overline{\tau_{rr}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\phi}'} \partial_r \overline{Z_\theta''} + \overline{\tau_{r\phi}'} \partial_r \overline{Z_\theta''} - \overline{Z_\theta''} \overline{Z_\theta''} \overline{Z_\theta''} - \overline{Z_\theta''} \overline{Z_\theta''} \overline{Z_\theta''}$$

$$\nabla(.) = \nabla_r(.) + \nabla_{\theta}(.) + \nabla_{\phi}(.) = \frac{1}{r^2} \partial_r(r^2.) + \frac{1}{r \sin \theta} \partial_{\theta}(\sin \theta.) + \frac{1}{r \sin \theta} \partial_{\phi}(.)$$