1 Alternative hydrodynamic stellar structure equations EXPLAINED

$$\partial_r \overline{m} = -\overline{\rho} \, \overline{m} \, \overline{g}_r / \Gamma_1 \overline{P} + 4\pi r^2 \overline{\rho} \tag{1}$$

$$\partial_r \overline{P} = -\overline{\rho} \, \overline{g}_r \tag{2}$$

$$\partial_r \widetilde{L} = -4\pi r^2 \widetilde{u}_r \overline{\rho} \ \overline{g}_r / \Gamma_1 + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r \tag{3}$$

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \,\overline{\rho} \,\overline{T} \,\overline{g}_r / \Gamma_1 \overline{P} \tag{4}$$

$$\partial_t \widetilde{X}_i = \widetilde{X}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_i - \widetilde{u}_r \partial_r \widetilde{X}_i \tag{5}$$

Equations (1)-(4) can be derived by four assumptions: (i) adiabatic convection $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$, (ii) hydrostatic equilibrium $\partial_r = -\rho g_r$, (iii) $\epsilon_{nuc} << \varepsilon_i$, (iv) $\varepsilon_K << \varepsilon_i \sim \varepsilon_t$.

1.1 Continuity equation

The continuity equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$. I will avoid the _{ad} suffix for clarity reasons and at some point assume hydrostatic equilibrium $\partial_r P = -\rho g_r$.

$$\Gamma_1 = \frac{\partial \ln P}{\partial \ln \rho} \tag{6}$$

$$\Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \tag{7}$$

$$\Gamma_1 P \partial_r \rho = \rho \partial_r P \tag{8}$$

$$\Gamma_1 P \partial_r \rho = \rho (-\rho g_r) \tag{9}$$

$$\partial_r \rho = -\frac{\rho \rho g_r}{\Gamma_1 P} \quad \backslash V \text{ (volume)} = 4\pi r^3/3$$
 (10)

$$V\partial_r \rho = -\frac{V\rho \ \rho \ g_r}{\Gamma_1 P} \tag{11}$$

$$\partial_r V \rho - \rho \partial_r V = -\frac{V \rho \rho g_r}{\Gamma_1 P} \tag{12}$$

$$\partial_r m - \rho 4\pi r^2 = -\frac{m \rho g_r}{\Gamma_1 P} \tag{13}$$

$$\partial_r m = -\frac{\rho \ m \ g_r}{\Gamma_1 P} + 4\pi r^2 \rho \tag{14}$$

$$\partial_r m = -\rho \ m \ g_r / \Gamma_1 P + 4\pi r^2 \rho \tag{15}$$

1.2 Temperature equation

The temperature equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$. I will avoid the _{ad} suffix for clarity reasons and at some point assume hydrostatic equilibrium $\partial_r P = -\rho g_r$.

$$\Gamma_1 = \frac{\partial \ln P}{\partial \ln \rho} \tag{16}$$

$$\Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \quad \backslash \partial_r lnT \tag{17}$$

$$\partial_r lnT \Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \partial_r lnT \tag{18}$$

$$P \partial_r lnT \Gamma_1 = \rho \frac{\partial_r P}{\partial_r \rho} \partial_r lnT$$
(19)

$$P \partial_r lnT \Gamma_1 = \frac{\partial_r P}{\partial_r ln\rho} \partial_r lnT \tag{20}$$

$$P \partial_r lnT \Gamma_1 = \frac{-\rho g_r}{\partial_r ln\rho} \partial_r lnT \tag{21}$$

$$\frac{P \partial_r T \Gamma_1}{T} = -\frac{\rho g_r}{\partial_r ln\rho} \partial_r lnT \tag{22}$$

$$-\frac{P \partial_r T \Gamma_1}{\rho g_r T} = \frac{\partial_r lnT}{\partial_r ln\rho} \tag{23}$$

$$-\frac{P \partial_r T \Gamma_1}{\rho g_r T} = \Gamma_3 - 1 \tag{24}$$

$$\partial_r T = -\frac{(\Gamma_3 - 1)\rho g_r T}{P\Gamma_1} \tag{25}$$

$$\partial_r T = -(\Gamma_3 - 1) \rho T g_r / \Gamma_1 P \tag{26}$$

1.3 Luminosity equation

The luminosity equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$. I will avoid the _{ad} suffix for clarity reasons and at some point assume hydrostatic equilibrium $\partial_r P = -\rho g_r$, nuclear energy production and kinetic energy being negligible compared to internal energy i.e. $\epsilon_{nuc} \ll \epsilon_i$ and $\epsilon_K \ll \epsilon_i$.

$$\Gamma_{1} = \frac{\partial \ln P}{\partial \ln \rho} \tag{27}$$

$$\Gamma_{1} = \frac{\rho}{P} \frac{\partial_{r} P}{\partial_{r} \rho} \tag{28}$$

$$\Gamma_{1} = \frac{\rho}{P} \frac{-\rho g_{r}}{\partial_{r} \rho} = -\frac{\rho^{2}}{P} \frac{g_{r}}{\partial_{r} \rho} \tag{29}$$

$$\Gamma_{1} = -\frac{g_{r}}{P \partial_{r} \rho / \rho^{2}} \tag{30}$$

$$dq = du + P dv \quad dq = \epsilon_{nuc} << du = d\epsilon_{i} \quad v = 1/\rho \quad \text{1st thermodynamic law, e.g. Kippenhahn and Weigert, 1994, p.19} \tag{31}$$

$$d\epsilon_{i} = -P dv = -P d(1/\rho) = P d\rho / \rho^{2} \tag{32}$$

$$\Gamma_{1} = -\frac{g_{r}}{\partial_{r} \epsilon_{i}} \tag{33}$$

$$\partial_{r} \epsilon_{i} = -\frac{g_{r}}{\Gamma_{1}} \epsilon_{K} << \epsilon_{i} \text{ and } \epsilon_{i} \sim \epsilon_{t}$$

$$\partial_{r} \epsilon_{t} = -\frac{g_{r}}{\Gamma_{1}} (4\pi r^{2} \rho u_{r} g_{r}$$

$$4\pi r^{2} \rho u_{r} \partial_{r} \epsilon_{t} = -\frac{4\pi r^{2} \rho u_{r} g_{r}}{\Gamma_{1}}$$

$$(36)$$

$$\partial_r 4\pi r^2 \rho u_r \varepsilon_t - \varepsilon_t \partial_r 4\pi r^2 \rho u_r = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1}$$

$$\partial_r L - \varepsilon_t \partial_r 4\pi r^2 \rho u_r = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1}$$
(38)

$$\partial_r L = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1} + \varepsilon_t \partial_r 4\pi r^2 \rho u_r \tag{39}$$

$$\partial_r L = -4\pi r^2 u_r \rho g_r / \Gamma_1 + \varepsilon_t \partial_r 4\pi r^2 \rho u_r \tag{40}$$

Table 1: Definitions:

 ρ density

$$m = \rho V = \rho \frac{4}{3} \pi r^3$$
 mass

T temperature

P pressure

 u_r, u_θ, u_ϕ velocity components

 $\mathbf{u} = u(u_r, u_\theta, u_\phi)$ velocity

 $j_z = r \sin \theta \ u_\phi$ specific angular momentum

 $d = \nabla \cdot \mathbf{u}$ dilatation

 ϵ_I specific internal energy

h specific enthalpy

 $k = (1/2)\widetilde{u_i''u_i''}$ turbulent kinetic energy

 ϵ_k specific kinetic energy

 ϵ_t specific total energy

s specific entropy

 $v = 1/\rho$ specific volume

 X_{α} mass fraction of isotope α

 $\dot{X}_{\alpha}^{\mathrm nuc}$ rate of change of X_{α}

 A_{α} number of nucleons in isotope α

 Z_{α} charge of isotope α

A mean number of nucleons per isotope

Z mean charge per isotope

 $f_P = \overline{P'u'_r}$ acoustic flux

 g_r radial gravitational acceleration

$$M = \int \rho(r)dV = \int \rho(r)4\pi r^2 dr$$
 integrated mass

 $S = \rho \epsilon_{\text{nuc}}(q)$ nuclear energy production (cooling function)

 $\tau_{ij} = 2\mu S_{ij}$ viscous stress tensor (μ kinematic viscosity)

 $S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$ strain rate

 $\widetilde{R}_{ij} = \overline{\rho} \widetilde{u_i''} u_j''$ Reynolds stress tensor

 $F_T = \chi \partial_r T$ heat flux

 $\Gamma_1 = (d \ln P/d \ln \rho)|_s$

 $\Gamma_2/(\Gamma_2-1)=(d\ ln\ P/d\ ln\ T)|_s$

 $\Gamma_3 - 1 = (d \ln T/d \ln \rho)|_s$

 $\widetilde{k}^r = (1/2)\widetilde{u_r''}u_r'' = (1/2)\widetilde{R}_{rr}/\overline{\rho}$ radial turbulent kinetic energy

 $\widetilde{k}^{\theta} = (1/2)\widetilde{u_{\theta}''u_{\theta}''} = (1/2)\widetilde{R}_{\theta\theta}/\overline{\rho}$ angular turbulent kinetic energy

 $\widetilde{k}^\phi=(1/2)\widetilde{u''_\phi u''_\phi}=(1/2)\widetilde{R}_{\phi\phi}/\overline{\rho}$ angular turbulent kinetic energy

 $\widetilde{k}^h = \widetilde{k}^\theta + \widetilde{k}^\phi$ horizontal turbulent kinetic energy

 $f_k = (1/2)\overline{\rho}u_i^{\prime\prime}\overline{u_i^{\prime\prime}}u_r^{\prime\prime}$ turbulent kinetic energy flux

 $f_k^r = (1/2)\overline{\rho}u_r''\overline{u_r''}u_r''$ radial turbulent kinetic energy flux

 $f_k^{\theta} = (1/2)\overline{\rho}u_{\theta}^{"}\widetilde{u_{\theta}^{"}}u_r^{"}$ angular turbulent kinetic energy flux

 $f_k^\phi = (1/2) \overline{\rho} u_\phi'' \overline{u_\phi''} u_r''$ angular turbulent kinetic energy flux

 $f_k^h = f_k^\theta + f_k^\phi$ horizontal turbulent kinetic energy flux

 $W_p = \overline{P'd''}$ turbulent pressure dilatation

 $W_b = \overline{\rho} \overline{u_r''} \widetilde{g}_r$ buoyancy

 $f_T = -\overline{\chi \partial_r T}$ heat flux (χ thermal conductivity)

Table 2: Definitions (continued):

$$\begin{split} f_I &= \overline{\rho} \widetilde{F'_I u''_I} & \text{ internal energy flux} & f_\alpha &= \overline{\rho} \widetilde{X''_\alpha u''_I} \times X_\alpha \text{ flux} \\ f_S &= \overline{\rho} \widetilde{S'' u''_I} & \text{ entropy flux} & f_{jz} &= \overline{\rho} \widetilde{J'_2 u''_I} & \text{ angular momentum flux} \\ f_T &= \overline{u'_T T'} & \text{ turbulent heat flux} & f_A &= \overline{\rho} \overline{A'' u''_I} & A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_h &= \overline{\rho} h'' u''_I & \text{ enthalpy flux} & f_Z &= \overline{\rho} Z'' u''_I & Z & (\text{mean charge per isotope}) & \text{ flux} \\ f_T &= \overline{\rho} \overline{I''_I u''_I} & \text{ enthalpy flux} & f_Z &= \overline{\rho} Z'' u''_I & Z & (\text{mean charge per isotope}) & \text{ flux} \\ f_T &= \overline{\rho} \overline{I''_I u''_I} & V_I &= -\nabla_T f_T + \varepsilon_I & \text{ numerical effect} \\ f_T &= -\overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T + \varepsilon_I & \text{ numerical effect} \\ f_T &= -\overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\overline{f'_T u'_I u'_I} & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ radial flux of } f_I & \mathcal{N}_T &= +(\Gamma_3 - 1)\varepsilon_I & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ radial flux of } f_I & \mathcal{N}_{Rrr} &= -2\nabla_T f_T^r - 2\varepsilon_I^r & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ radial flux of } f_I & \mathcal{N}_{R\theta\theta} &= -2\nabla_T f_T^r - 2\varepsilon_I^r & \text{ numerical effect} \\ f_T &= \overline{f'_T u'_I u'_I} & \text{ radial flux of } f_A & \mathcal{N}_R &= -\nabla_T f_T - \varepsilon_I^r & \text{ numerical effect} \\ f_R &= \overline{\rho} A'' u''_I u''_I & \text{ radial flux of } f_A & \mathcal{N}_R &= -\nabla_T f_T^r - \varepsilon_I^r & \text{ numerical effect} \\ f_R &= \overline{\rho} A'' u''_I u''_I & \text{ radial flux of } f_A & \mathcal{N}_R &= -\nabla_T f_T^r - \varepsilon_I^r & \text{ numerical effect} \\ f_R &= -(1/2) \overline{G'_T r} - \overline{u''_I u''_I} & \text{ radial flux of } f_Z & \mathcal{N}_R &= -\varepsilon_R & \text{ numerical effect} \\ f_R &= -(1/2) \overline{G'_T r} - \overline{u''_I u''_I} & \text{ radial flux of } f_Z & \mathcal{N}_R &= -\varepsilon_R & \text{ numerical effect} \\ f_R &= -\varepsilon_I u''_I u''_I u''_I & \text{ radi$$

Table 3: Definitions (continued):

$$\begin{array}{lll} \mathcal{G}_{k}^{\theta} = -(1/2)\overline{G_{k}^{\theta}} - \overline{u_{0}''}G_{\phi}^{M} & \mathcal{N}_{b} \text{ numerical effect} \\ \mathcal{G}_{k}^{\phi} = -(1/2)\overline{G_{k}^{\theta}} - \overline{u_{0}''}G_{\phi}^{M} & \mathcal{N}_{fI} = -\nabla_{r}(c_{I}''\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} - \varepsilon_{I} \text{ numerical effect} \\ \mathcal{G}_{k}^{h} = +\mathcal{G}_{k}^{\theta} + \mathcal{G}_{k}^{\phi} & \mathcal{N}_{fh} = -\nabla_{r}(\overline{h''}\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} - \varepsilon_{I} \text{ numerical effect} \\ \mathcal{G}_{a} = +\overline{\rho'v}G_{r}^{M} & \mathcal{N}_{fs} = -\nabla_{r}(\overline{s''}\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} / T - \varepsilon_{s} \text{ numerical effect} \\ \mathcal{G}_{I} = -\overline{G_{r}^{I}} - \overline{\epsilon_{I}''}G_{r}^{M} & \mathcal{N}_{fs} = -\nabla_{r}(\overline{s''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{a} = -\overline{G_{r}^{2}} - \overline{\lambda_{u}''}G_{r}^{M} & \mathcal{N}_{fA} = -\nabla_{r}(\overline{A''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{\lambda_{u}''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{Z''}\tau_{rr}') - \varepsilon_{Z} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f3} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f3} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f3} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f3} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_$$

Table 4: Definitions (continued):

$$\begin{split} \varepsilon_k^r &= \overline{\tau_{rr}'} \overline{\partial_r u_r''} + \overline{\tau_{r\theta}'} (1/r) \partial_\theta u_\theta'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\theta &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\theta'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\phi &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\phi'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\phi'' \\ \varepsilon_k^\phi &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\phi'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\phi'' \\ \varepsilon_k^\phi &= \varepsilon_k^\phi + \varepsilon_k^\phi + \varepsilon_k^\phi \\ \varepsilon_k^\phi &= \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k^\phi &= \overline{\rho' v \nabla_r \tau_{rr}'} \\ \varepsilon_k^\phi &= \overline{\rho' v \nabla_r \tau_{rr}'} \\ \varepsilon_1^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi$$

$$\nabla(.) = \nabla_r(.) + \nabla_{\theta}(.) + \nabla_{\phi}(.) = \frac{1}{r^2} \partial_r(r^2.) + \frac{1}{r \sin \theta} \partial_{\theta}(\sin \theta.) + \frac{1}{r \sin \theta} \partial_{\phi}(.)$$