# 1 Alternative hydrodynamic stellar structure equations EXPLAINED

$$\partial_r \overline{M} = -\overline{\rho} \ \overline{M} \ \overline{g}_r / \Gamma_1 \overline{P} + 4\pi r^2 \overline{\rho} \tag{1}$$

$$\partial_r \overline{P} = -\overline{\rho} \, \overline{g}_r \tag{2}$$

$$\partial_r \widetilde{L} = -4\pi r^2 \widetilde{u}_r \overline{\rho} \ \overline{g}_r / \Gamma_1 + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r \tag{3}$$

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \,\overline{\rho} \,\overline{T} \,\overline{g}_r / \Gamma_1 \overline{P} \tag{4}$$

$$\partial_t \widetilde{X}_i = \widetilde{X}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_i - \widetilde{u}_r \partial_r \widetilde{X}_i \tag{5}$$

Equations (1)-(4) can be derived by four assumptions: (i) adiabatic convection  $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$ , (ii) hydrostatic equilibrium  $\partial_r = -\rho g_r$ , (iii)  $\epsilon_{nuc} << \varepsilon_i$ , (iv)  $\varepsilon_K << \varepsilon_i \sim \varepsilon_t$ .

### 1.1 Continuity equation

The continuity equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of  $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$ . I will avoid the <sub>ad</sub> suffix for clarity reasons and at some point assume hydrostatic equilibrium  $\partial_r P = -\rho g_r$ .

$$\Gamma_1 = \frac{\partial \ln P}{\partial \ln \rho} \tag{6}$$

$$\Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \tag{7}$$

$$\Gamma_1 P \partial_r \rho = \rho \partial_r P \tag{8}$$

$$\Gamma_1 P \partial_r \rho = \rho \left( -\rho g_r \right) \tag{9}$$

$$\partial_r \rho = -\frac{\rho \rho g_r}{\Gamma_1 P} \quad \backslash V \text{ (volume)} = 4\pi r^3 / 3$$
 (10)

$$V\partial_r \rho = -\frac{V\rho \ \rho \ g_r}{\Gamma_1 P} \tag{11}$$

$$\partial_r V \rho - \rho \partial_r V = -\frac{V \rho \rho g_r}{\Gamma_1 P} \tag{12}$$

$$\partial_r M - \rho 4\pi r^2 = -\frac{M \rho g_r}{\Gamma_1 P} \tag{13}$$

$$\partial_r M = -\frac{\rho M g_r}{\Gamma_1 P} + 4\pi r^2 \rho \tag{14}$$

$$\partial_r M = -\rho \ M \ g_r / \Gamma_1 P + 4\pi r^2 \rho \tag{15}$$

## 1.2 Temperature equation

The temperature equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of  $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$ . I will avoid the <sub>ad</sub> suffix for clarity reasons and at some point assume hydrostatic equilibrium  $\partial_r P = -\rho g_r$ .

$$\Gamma_1 = \frac{\partial \ln P}{\partial \ln \rho} \tag{16}$$

$$\Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \quad \backslash \partial_r lnT \tag{17}$$

$$\partial_r lnT \Gamma_1 = \frac{\rho}{P} \frac{\partial_r P}{\partial_r \rho} \partial_r lnT \tag{18}$$

$$P \partial_r lnT \Gamma_1 = \rho \frac{\partial_r P}{\partial_r \rho} \partial_r lnT$$
(19)

$$P \partial_r lnT \Gamma_1 = \frac{\partial_r P}{\partial_r ln\rho} \partial_r lnT \tag{20}$$

$$P \partial_r lnT \Gamma_1 = \frac{-\rho g_r}{\partial_r ln\rho} \partial_r lnT \tag{21}$$

$$\frac{P \partial_r T \Gamma_1}{T} = -\frac{\rho g_r}{\partial_r ln\rho} \partial_r lnT \tag{22}$$

$$-\frac{P \partial_r T \Gamma_1}{\rho g_r T} = \frac{\partial_r lnT}{\partial_r ln\rho} \tag{23}$$

$$-\frac{P \partial_r T \Gamma_1}{\rho g_r T} = \Gamma_3 - 1 \tag{24}$$

$$\partial_r T = -\frac{(\Gamma_3 - 1)\rho g_r T}{P\Gamma_1} \tag{25}$$

$$\partial_r T = -(\Gamma_3 - 1) \rho T g_r / \Gamma_1 P \tag{26}$$

### 1.3 Luminosity equation

The luminosity equation can be derived from first principles for adiabatic convection in hydrostatic equilibrium. We start from the definition of  $\Gamma_1 = \partial \ln P/\partial \ln \rho|_{ad}$ . I will avoid the <sub>ad</sub> suffix for clarity reasons and at some point assume hydrostatic equilibrium  $\partial_r P = -\rho g_r$ , nuclear energy production and kinetic energy being negligible compared to internal energy i.e.  $\epsilon_{nuc} \ll \epsilon_i$  and  $\epsilon_K \ll \epsilon_i$ .

$$\Gamma_{1} = \frac{\partial \ln P}{\partial \ln \rho} \tag{27}$$

$$\Gamma_{1} = \frac{\rho}{P} \frac{\partial_{r} P}{\partial_{r} \rho} \tag{28}$$

$$\Gamma_{1} = \frac{\rho}{P} \frac{-\rho g_{r}}{\partial_{r} \rho} = -\frac{\rho^{2}}{P} \frac{g_{r}}{\partial_{r} \rho} \tag{29}$$

$$\Gamma_{1} = -\frac{g_{r}}{P \partial_{r} \rho / \rho^{2}} \tag{30}$$

$$dq = du + P dv \quad dq = \epsilon_{nuc} << du = d\epsilon_{i} \quad v = 1/\rho \quad \text{1st thermodynamic law, e.g. Kippenhahn and Weigert, 1994, p.19} \tag{31}$$

$$d\epsilon_{i} = -P dv = -P d(1/\rho) = P d\rho / \rho^{2} \tag{32}$$

$$\Gamma_{1} = -\frac{g_{r}}{\partial_{r} \epsilon_{i}} \tag{33}$$

$$\partial_{r} \epsilon_{i} = -\frac{g_{r}}{\Gamma_{1}} \epsilon_{K} << \epsilon_{i} \text{ and } \epsilon_{i} \sim \epsilon_{t}$$

$$\partial_{r} \epsilon_{t} = -\frac{g_{r}}{\Gamma_{1}} (4\pi r^{2} \rho u_{r} g_{r}$$

$$4\pi r^{2} \rho u_{r} \partial_{r} \epsilon_{t} = -\frac{4\pi r^{2} \rho u_{r} g_{r}}{\Gamma_{1}}$$

$$(36)$$

$$\partial_r 4\pi r^2 \rho u_r \varepsilon_t - \varepsilon_t \partial_r 4\pi r^2 \rho u_r = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1}$$

$$\partial_r L - \varepsilon_t \partial_r 4\pi r^2 \rho u_r = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1}$$
(38)

$$\partial_r L = -\frac{4\pi r^2 \rho u_r g_r}{\Gamma_1} + \varepsilon_t \partial_r 4\pi r^2 \rho u_r \tag{39}$$

$$\partial_r L = -4\pi r^2 u_r \rho g_r / \Gamma_1 + \varepsilon_t \partial_r 4\pi r^2 \rho u_r \tag{40}$$

#### Table 1: Definitions:

 $\rho$  density  $g_r$  radial gravitational acceleration  $S = \rho \epsilon_{\text{nuc}}(q)$  nuclear energy production (cooling function) T temperature  $\tau_{ij} = 2\mu S_{ij}$  viscous stress tensor ( $\mu$  kinematic viscosity) P pressure  $S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$  strain rate  $u_r, u_\theta, u_\phi$  velocity components  $\widetilde{R}_{ij} = \overline{\rho} \widetilde{u_{i'}' u_{i'}'}$  Reynolds stress tensor  $\mathbf{u} = u(u_r, u_\theta, u_\phi)$  velocity  $j_z = r \sin \theta \ u_\phi$  specific angular momentum  $F_T = \chi \partial_r T$  heat flux  $d = \nabla \cdot \mathbf{u}$  dilatation  $\Gamma_1 = (d \ln P/d \ln \rho)|_{\mathfrak{s}}$  $\Gamma_2/(\Gamma_2-1)=(d \ln P/d \ln T)|_s$  $\epsilon_I$  specific internal energy  $\Gamma_3 - 1 = (d \ln T/d \ln \rho)|_s$ h specific enthalpy  $k = (1/2)\widetilde{u_i''u_i''}$  turbulent kinetic energy  $\widetilde{k}^r = (1/2)\widetilde{u_r''}u_r'' = (1/2)\widetilde{R}_{rr}/\overline{\rho}$  radial turbulent kinetic energy  $\widetilde{k}^{\theta} = (1/2)\widetilde{u_{\theta}''u_{\theta}''} = (1/2)\widetilde{R}_{\theta\theta}/\overline{\rho}$  angular turbulent kinetic energy  $\epsilon_k$  specific kinetic energy  $\widetilde{k}^{\phi} = (1/2)\widetilde{u_{\phi}''u_{\phi}''} = (1/2)\widetilde{R}_{\phi\phi}/\overline{\rho}$  angular turbulent kinetic energy  $\epsilon_t$  specific total energy  $\widetilde{k}^h = \widetilde{k}^\theta + \widetilde{k}^\phi$  horizontal turbulent kinetic energy s specific entropy  $f_k = (1/2)\overline{\rho}u_i^{\prime\prime}\overline{u_i^{\prime\prime}}u_r^{\prime\prime}$  turbulent kinetic energy flux  $v = 1/\rho$  specific volume  $f_k^r = (1/2)\overline{\rho}u_r^{\prime\prime}\overline{u_r^{\prime\prime}}u_r^{\prime\prime}$  radial turbulent kinetic energy flux  $X_{\alpha}$  mass fraction of isotope  $\alpha$  $\dot{X}_{\alpha}^{\mathrm nuc}$  rate of change of  $X_{\alpha}$  $f_k^{\theta} = (1/2)\overline{\rho}u_{\theta}^{\prime\prime}u_{\theta}^{\prime\prime}u_r^{\prime\prime}$  angular turbulent kinetic energy flux  $f_k^{\phi} = (1/2) \overline{\rho} u_{\phi}^{\prime\prime} u_{\phi}^{\prime\prime} u_r^{\prime\prime}$  angular turbulent kinetic energy flux  $A_{\alpha}$  number of nucleons in isotope  $\alpha$  $f_k^h = f_k^\theta + f_k^\phi$  horizontal turbulent kinetic energy flux  $Z_{\alpha}$  charge of isotope  $\alpha$  $W_p = \overline{P'd''}$  turbulent pressure dilatation A mean number of nucleons per isotope  $W_b = \overline{\rho} \overline{u_r''} \widetilde{g}_r$  buoyancy Z mean charge per isotope  $f_P = \overline{P'u'_r}$  acoustic flux  $f_T = -\overline{\chi \partial_r T}$  heat flux ( $\chi$  thermal conductivity)

Table 2: Definitions (continued):

$$\begin{split} f_I &= \overline{\rho} \widetilde{F'_I u''_I} & \text{ internal energy flux} & f_\alpha &= \overline{\rho} \widetilde{X''_\alpha u''_I} \times X_\alpha \text{ flux} \\ f_S &= \overline{\rho} \widetilde{S'' u''_I} & \text{ entropy flux} & f_{jz} &= \overline{\rho} \widetilde{J'_2 u''_I} & \text{ angular momentum flux} \\ f_T &= \overline{u'_T T'} & \text{ turbulent heat flux} & f_A &= \overline{\rho} \overline{A'' u''_I} & A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_h &= \overline{\rho} h'' u''_I & \text{ enthalpy flux} & f_Z &= \overline{\rho} Z'' u''_I & Z & (\text{mean charge per isotope}) & \text{ flux} \\ f_T &= \overline{\rho} \overline{I''_I u''_I} & \text{ enthalpy flux} & f_Z &= \overline{\rho} Z'' u''_I & Z & (\text{mean charge per isotope}) & \text{ flux} \\ f_T &= \overline{\rho} \overline{I''_I u''_I} & V_I &= -\nabla_T f_T + \varepsilon_I & \text{ numerical effect} \\ f_T &= -\overline{f'_T u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T + \varepsilon_I & \text{ numerical effect} \\ f_T^{\rho} &= -\overline{f'_{\theta T} u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T & \text{ numerical effect} \\ f_T^{\rho} &= -\overline{f'_{\theta T} u'_I u'_I} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_T f_T & \text{ numerical effect} \\ f_T^{\rho} &= f_T^{\rho} + f_T^{\rho} & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\overline{V_T f_T} & -\overline{V_$$

#### Table 3: Definitions (continued):

$$\begin{array}{lll} \mathcal{G}_{k}^{\theta} = -(1/2)\overline{G_{k}^{\theta}} - \overline{u_{0}''}G_{\phi}^{M} & \mathcal{N}_{b} \text{ numerical effect} \\ \mathcal{G}_{k}^{\phi} = -(1/2)\overline{G_{k}^{\theta}} - \overline{u_{0}''}G_{\phi}^{M} & \mathcal{N}_{fI} = -\nabla_{r}(c_{I}''\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} - \varepsilon_{I} \text{ numerical effect} \\ \mathcal{G}_{k}^{h} = +\mathcal{G}_{k}^{\theta} + \mathcal{G}_{k}^{\phi} & \mathcal{N}_{fh} = -\nabla_{r}(\overline{h''}\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} - \varepsilon_{I} \text{ numerical effect} \\ \mathcal{G}_{a} = +\overline{\rho'v}G_{r}^{M} & \mathcal{N}_{fs} = -\nabla_{r}(\overline{s''}\tau_{rr}') + \overline{u_{r}''}\tau_{ij}\partial_{i}u_{j} / T - \varepsilon_{s} \text{ numerical effect} \\ \mathcal{G}_{I} = -\overline{G_{r}^{I}} - \overline{\epsilon_{I}''}G_{r}^{M} & \mathcal{N}_{fs} = -\nabla_{r}(\overline{s''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{a} = -\overline{G_{r}^{2}} - \overline{\lambda_{u}''}G_{r}^{M} & \mathcal{N}_{fA} = -\nabla_{r}(\overline{A''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{\lambda_{u}''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{Z''}\tau_{rr}') - \varepsilon_{Z} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{A} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{A} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f2} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{A''}G_{r}^{M} & \mathcal{N}_{f3} = -\nabla_{r}(\overline{a''}\tau_{rr}') - \varepsilon_{B} \text{ numerical effect} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f3} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f4} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f4} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_{f4} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} \\ \mathcal{G}_{B} = -\overline{G_{r}^{2}} - \overline{J''}G_{r}^{M} & \mathcal{N}_$$

Table 4: Definitions (continued):

$$\begin{split} \varepsilon_k^r &= \overline{\tau_{rr}'} \overline{\partial_r u_r''} + \overline{\tau_{r\theta}'} (1/r) \partial_\theta u_\theta'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\theta &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\theta'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\phi &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\phi'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\phi'' \\ \varepsilon_k^\phi &= \overline{\tau_{dr}'} \partial_r u_\theta'' + \overline{\tau_{d\theta}'} (1/r) \partial_\theta u_\phi'' + \tau_{d\phi}' (1/r \sin \theta) \partial_\phi u_\phi'' \\ \varepsilon_k^\phi &= \varepsilon_k^\phi + \varepsilon_k^\phi + \varepsilon_k^\phi \\ \varepsilon_k^\phi &= \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k^\phi &= \overline{\rho' v \nabla_r \tau_{rr}'} \\ \varepsilon_k^\phi &= \overline{\rho' v \nabla_r \tau_{rr}'} \\ \varepsilon_1^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi &= \overline{\tau_{rr}'} \partial_r \varepsilon_1'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi \varepsilon_1'' \\ \varepsilon_2^\phi$$

$$\nabla(.) = \nabla_r(.) + \nabla_{\theta}(.) + \nabla_{\phi}(.) = \frac{1}{r^2} \partial_r(r^2.) + \frac{1}{r \sin \theta} \partial_{\theta}(\sin \theta.) + \frac{1}{r \sin \theta} \partial_{\phi}(.)$$