## 1 Alternative hydrodynamic stellar structure equations

Below is a set of alternative hydrodynamic stellar structure equations inspired by mean fields from the mean density/temperature/internal energy and pressure flux equations involving gradient term scaled by favrian Reynolds stress  $\widetilde{R}_{rr}$  and the other one turbulent dilatation flux  $\overline{u'_r d''}$ .

$$\partial_r \overline{M} \sim -\overline{\rho} \ \overline{M} \ \overline{u'_r d''} / \ \widetilde{R}_{rr} + 4\pi r^2 \overline{\rho}$$
 (1)

$$\partial_r \overline{P} \sim -\Gamma_1 \ \overline{\rho} \ \overline{P} \ \overline{u'_r d''} / \ \widetilde{R}_{rr}$$
 (2)

$$\partial_r \widetilde{L} \sim +\widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r - 4\pi r^2 \overline{\rho} \ \overline{P} \ \overline{u'_r d''} / \ \widetilde{R}_{rr}$$
 (3)

$$\partial_r \overline{T} \sim -(\Gamma_3 - 1) \, \overline{\rho} \, \overline{T} \, \overline{u'_r d''} / \, \widetilde{R}_{rr}$$
 (4)

$$\partial_t \widetilde{X}_i = \widetilde{\dot{X}}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_\alpha - \widetilde{u}_r \partial_r \widetilde{X}_\alpha \tag{5}$$

These equations could be perfectly validated by our ransX framework and are shown on the next page in Figure 1. It appears that there is a universality relation between gradient of a mean thermodynamic variable Q and dilatation flux  $\overline{u'_r d''}$ , that is:

$$\partial_r \overline{Q} \sim -const \ \overline{\rho} \ \overline{Q} \ \overline{u_r' d''} / \widetilde{R}_{rr}$$
 (6)

Moreover, we know that due to hydrostatic equilibrium,  $\partial_r \overline{P} \sim -\overline{\rho} \ \overline{g}_r$ , and based on Equation (2) we can write, that dilatation flux  $\overline{u'_r d''}$  is:

$$\overline{u_r'd''} \sim \frac{\widetilde{R}_{rr} \ \overline{g}_r}{\Gamma_1 \ \overline{P}} \tag{7}$$

Also, the expansion velocity  $\tilde{u}_r$  can be replaced by:

$$\widetilde{u}_r = -\partial_t \overline{M} / 4\pi r^2 \overline{\rho} \tag{8}$$

So the alternative hydrodynamic stellar structure equations for hydrostatic convection become a system with only **one unknown the composition flux**  $f_{\alpha}$  for which we need a proper model. See validation of these simplified alternative hydrodynamic stellar structure equation in Figure 2.

$$\partial_r \overline{M} \sim -\overline{\rho} \ \overline{M} \ \overline{g}_r / \Gamma_1 \overline{P} + 4\pi r^2 \overline{\rho}$$
 (9)

$$\partial_r \overline{P} \sim -\overline{\rho} \, \overline{g}_r$$
 (10)

$$\partial_r \widetilde{L} \sim -4\pi r^2 \overline{\rho} \ \overline{g}_r / \Gamma_1 + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r \tag{11}$$

$$\partial_r \overline{T} \sim -(\Gamma_3 - 1) \ \overline{\rho} \ \overline{T} \ \overline{g}_r / \Gamma_1 \overline{P}$$
 (12)

$$\partial_t \widetilde{X}_i = \widetilde{\dot{X}}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_\alpha - \widetilde{u}_r \partial_r \widetilde{X}_\alpha \tag{13}$$

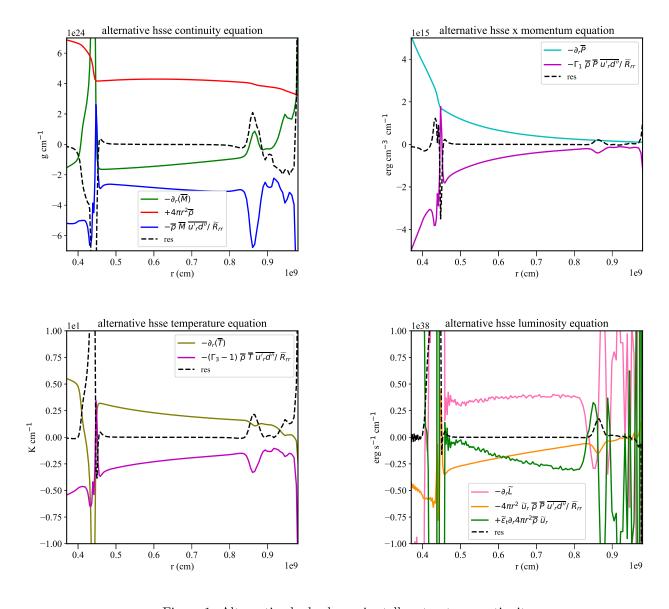


Figure 1: Alternative hydrodynamic stellar structure continuity.

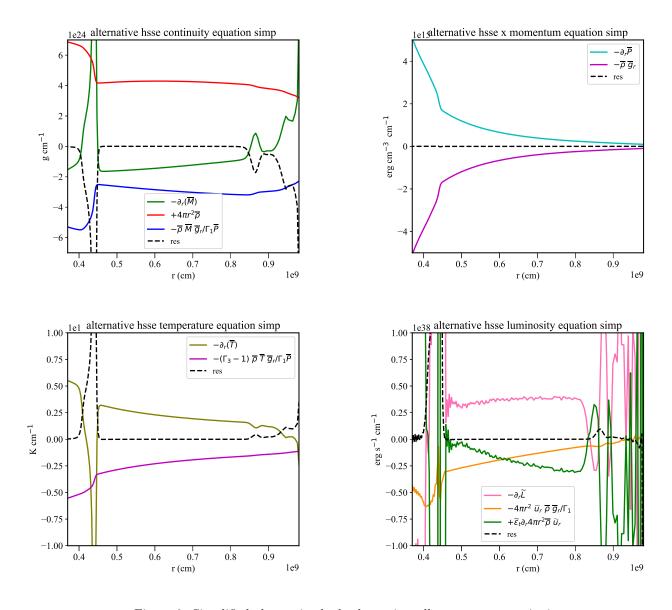


Figure 2: Simplified alternative hydrodynamic stellar structure continuity.

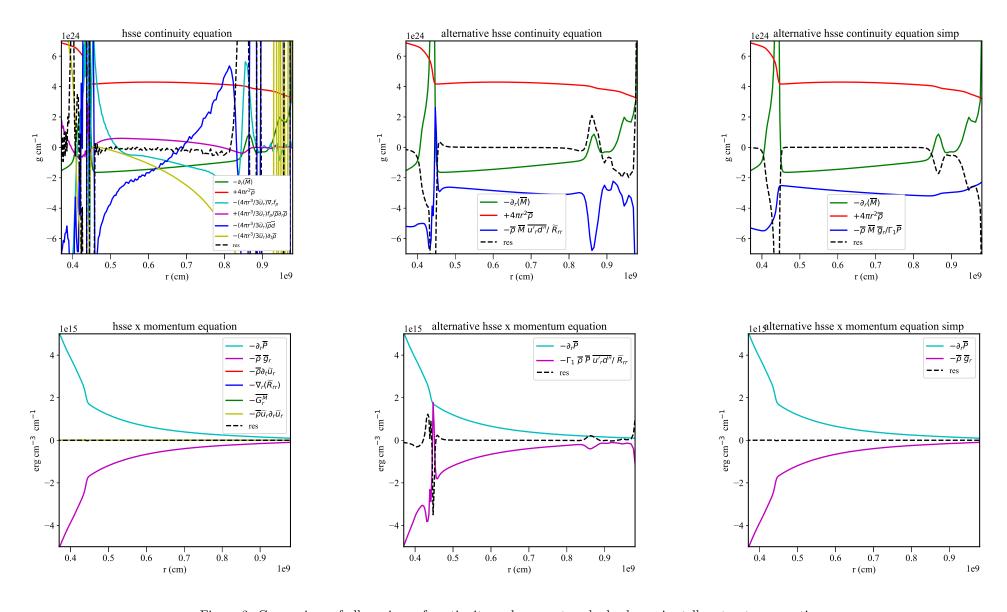


Figure 3: Comparison of all versions of continuity and momentum hydrodynamic stellar structure equations.

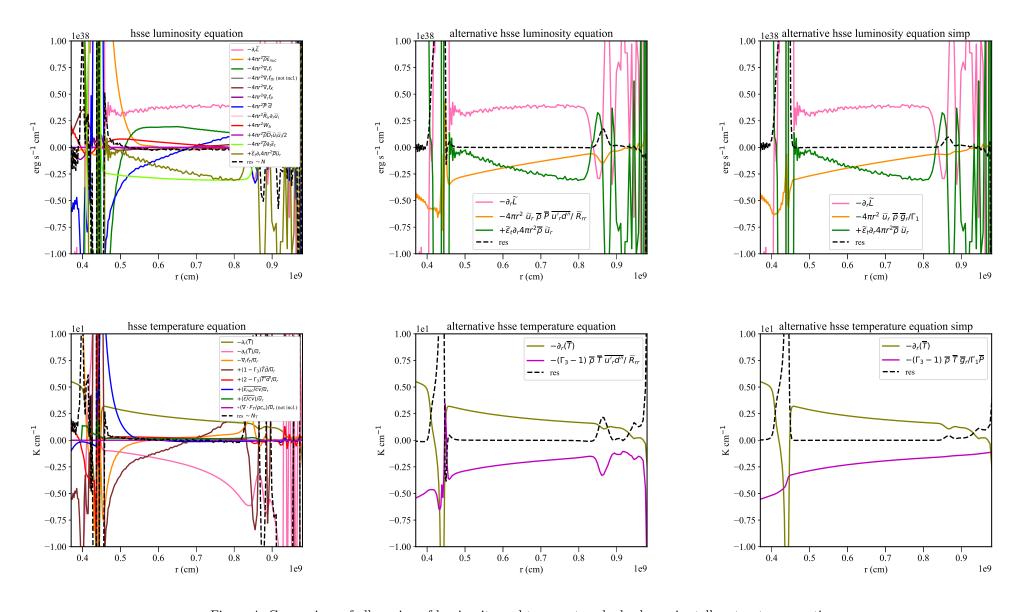


Figure 4: Comparison of all version of luminosity and temperature hydrodynamic stellar structure equations.

## Table 1: Definitions:

 $\rho$  density  $g_r$  radial gravitational acceleration  $S = \rho \epsilon_{\text{nuc}}(q)$  nuclear energy production (cooling function) T temperature  $\tau_{ij} = 2\mu S_{ij}$  viscous stress tensor ( $\mu$  kinematic viscosity) P pressure  $S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$  strain rate  $u_r, u_\theta, u_\phi$  velocity components  $\widetilde{R}_{ij} = \overline{\rho} \widetilde{u_{i'}' u_{i'}'}$  Reynolds stress tensor  $\mathbf{u} = u(u_r, u_\theta, u_\phi)$  velocity  $j_z = r \sin \theta \ u_\phi$  specific angular momentum  $F_T = \chi \partial_r T$  heat flux  $d = \nabla \cdot \mathbf{u}$  dilatation  $\Gamma_1 = (d \ln P/d \ln \rho)|_{\mathfrak{s}}$  $\Gamma_2/(\Gamma_2-1)=(d \ln P/d \ln T)|_s$  $\epsilon_I$  specific internal energy  $\Gamma_3 - 1 = (d \ln T/d \ln \rho)|_s$ h specific enthalpy  $k = (1/2)\widetilde{u_i''u_i''}$  turbulent kinetic energy  $\widetilde{k}^r = (1/2)\widetilde{u_r''}u_r'' = (1/2)\widetilde{R}_{rr}/\overline{\rho}$  radial turbulent kinetic energy  $\widetilde{k}^{\theta} = (1/2)\widetilde{u_{\theta}''u_{\theta}''} = (1/2)\widetilde{R}_{\theta\theta}/\overline{\rho}$  angular turbulent kinetic energy  $\epsilon_k$  specific kinetic energy  $\widetilde{k}^{\phi} = (1/2)\widetilde{u_{\phi}''u_{\phi}''} = (1/2)\widetilde{R}_{\phi\phi}/\overline{\rho}$  angular turbulent kinetic energy  $\epsilon_t$  specific total energy  $\widetilde{k}^h = \widetilde{k}^\theta + \widetilde{k}^\phi$  horizontal turbulent kinetic energy s specific entropy  $f_k = (1/2)\overline{\rho}u_i^{\prime\prime}\overline{u_i^{\prime\prime}}u_r^{\prime\prime}$  turbulent kinetic energy flux  $v = 1/\rho$  specific volume  $f_k^r = (1/2)\overline{\rho}u_r^{\prime\prime}\overline{u_r^{\prime\prime}}u_r^{\prime\prime}$  radial turbulent kinetic energy flux  $X_{\alpha}$  mass fraction of isotope  $\alpha$  $\dot{X}_{\alpha}^{\mathrm nuc}$  rate of change of  $X_{\alpha}$  $f_k^{\theta} = (1/2)\overline{\rho}u_{\theta}^{\prime\prime}u_{\theta}^{\prime\prime}u_r^{\prime\prime}$  angular turbulent kinetic energy flux  $f_k^{\phi} = (1/2) \overline{\rho} u_{\phi}^{\prime\prime} u_{\phi}^{\prime\prime} u_r^{\prime\prime}$  angular turbulent kinetic energy flux  $A_{\alpha}$  number of nucleons in isotope  $\alpha$  $f_k^h = f_k^\theta + f_k^\phi$  horizontal turbulent kinetic energy flux  $Z_{\alpha}$  charge of isotope  $\alpha$  $W_p = \overline{P'd''}$  turbulent pressure dilatation A mean number of nucleons per isotope  $W_b = \overline{\rho} \overline{u_r''} \widetilde{g}_r$  buoyancy Z mean charge per isotope  $f_P = \overline{P'u'_r}$  acoustic flux  $f_T = -\overline{\chi \partial_r T}$  heat flux ( $\chi$  thermal conductivity)

Table 2: Definitions (continued):

$$\begin{split} f_I &= \overline{\rho_I} \widetilde{I} u_I'' & \text{ internal energy flux} & f_\alpha &= \overline{\rho} \overline{\lambda_\alpha''} u_I'' \times \lambda_\alpha \text{ flux} \\ f_s &= \overline{\rho} \overline{s''} u_I'' & \text{ entropy flux} & f_{jz} &= \overline{\rho} \overline{j}_z'' u_I'' \text{ angular momentum flux} \\ f_T &= \overline{u_r'} \overline{V}' & \text{ turbulent heat flux} & f_A &= \overline{\rho} \overline{A''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_h &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{h''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{J''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{J''} u_I'' & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' \times A & (\text{mean number of nucleons per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{J''} u_I'' u_I'' & \text{ viscous flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' & X_{\eta} \times A & (\text{mean charge per isotope}) & \text{ flux} \\ f_{\tau} &= \overline{\rho} \overline{J''} u_I'' u_I'' \text{ viscous flux} & f_Z &= \overline{\rho} \overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{\rho} \overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{\rho} \overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{\rho} \overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{J''} u_I'' u_I'' \text{ radial flux of } f_Z & f_Z -\overline{J''} u_I'' u_I'' \text{ radial flux} \text{ of } f_Z & f_Z -\overline{J''} u_I'' u_I''$$

## Table 3: Definitions (continued):

$$\begin{array}{lll} \mathcal{G}_k^\theta = -(1/2)\overline{G_{Q_0}^R} & \mathcal{N}_b \text{ numerical effect} \\ \mathcal{G}_k^\phi = -(1/2)\overline{G_{Q_0}^R} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_k^\phi = -(1/2)\overline{G_{Q_0}^R} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_k^h = +\mathcal{G}_k^\theta + \mathcal{G}_k^\phi & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_k^h = +\mathcal{G}_k^\theta + \mathcal{G}_k^\phi & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_a = +\overline{\rho'vG_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_a = +\overline{\rho'vG_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_f = -\overline{G_r^T} - \overline{\epsilon_f''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_f \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_g \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_{g_r} \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_{g_r} \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_{g_r} \text{ numerical effect} \\ \mathcal{G}_g = -\overline{G_r^T} - \overline{\lambda_a''G_r^M} & \mathcal{N}_{g_r} \text{ numerical effect} \\ \mathcal{N}_{g_r} = +2\overline{\lambda_a''G_r^M} \text{ numerical effect} \\ \mathcal{N}_{g_r} = +2\overline{\lambda_a''G_r^M} \text{ numerical effect} \\ \mathcal{N}_{g_s} = +2\overline{\lambda_a''G_r^M} \text{ numerical effect} \\ \mathcal{N}_{g_s} = -2\overline{\lambda_a''G_r^M} \text{ numerical effect} \\ \mathcal{N}_{g_s} = -2\overline{\lambda_a''G_r^M}$$

Table 4: Definitions (continued):

$$\begin{split} \varepsilon_k^r &= \overline{\tau_{rr}'} \overline{\partial_r u_r''} + \overline{\tau_{r\theta}'} (1/r) \partial_\theta u_r'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\theta &= \overline{\tau_{r\theta}'} \partial_r u_\theta'' + \overline{\tau_{\theta\theta}'} (1/r) \partial_\theta u_\theta'' + \tau_{\theta\phi}' (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\phi &= \overline{\tau_{r\theta}'} \partial_r u_\theta'' + \overline{\tau_{\theta\theta}'} (1/r) \partial_\theta u_\phi'' + \overline{\tau_{\phi\phi}'} (1/r \sin \theta) \partial_\phi u_\phi'' \\ \varepsilon_k &= \varepsilon_{r\theta}' + \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k &= \varepsilon_k^r + \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k^\theta &= \overline{\rho' v \nabla_r \tau_r'} \\ \varepsilon_k &= \varepsilon_k^\theta + \varepsilon_k^\phi \\ \varepsilon_k &= \overline{\rho' v \nabla_r \tau_r'} \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r) \partial_\theta \varepsilon_l'' + \overline{\tau_{r\phi}'} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\tau_{r\theta}'} \partial_r \varepsilon_l'' + \overline{\tau_{r\theta}'} (1/r \sin \theta) \partial_\phi$$

$$\nabla(.) = \nabla_r(.) + \nabla_{\theta}(.) + \nabla_{\phi}(.) = \frac{1}{r^2} \partial_r(r^2.) + \frac{1}{r \sin \theta} \partial_{\theta}(\sin \theta.) + \frac{1}{r \sin \theta} \partial_{\phi}(.)$$