1 Alternative hydrodynamic stellar structure equations

Below is a set of alternative hydrodynamic stellar structure equations inspired by mean fields from the mean density/temperature/internal energy and pressure flux equations involving gradient term scaled by favrian Reynolds stress \widetilde{R}_{rr} and the other one turbulent dilatation flux $\overline{u'_r d''}$.

$$\partial_r \overline{M} = -\overline{\rho} \ \overline{M} \ \overline{u'_r d''} / \ \widetilde{R}_{rr} + 4\pi r^2 \overline{\rho}$$
 (1)

$$\partial_r \overline{P} = -\Gamma_1 \ \overline{\rho} \ \overline{P} \ \overline{u'_r d''} / \ \widetilde{R}_{rr} \tag{2}$$

$$\partial_r \widetilde{L} = +\widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r - 4\pi r^2 \overline{\rho} \ \widetilde{u}_r \ \overline{P} \ \overline{u'_r d''} / \ \widetilde{R}_{rr}$$

$$\tag{3}$$

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \,\overline{\rho} \,\overline{T} \,\overline{u'_r d''} / \,\widetilde{R}_{rr} \tag{4}$$

$$\partial_t \widetilde{X}_i = \widetilde{X}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_i - \widetilde{u}_r \partial_r \widetilde{X}_i$$
 (5)

These equations could be perfectly validated by our ransX framework and are shown on the next page in Figure 1. It appears that there is a universality relation between gradient of a mean thermodynamic variable Q and dilatation flux $\overline{u'_r d''}$, that is:

$$\partial_r \overline{Q} \sim -\overline{\rho} \ \overline{Q} \ \overline{u_r' d''} / \widetilde{R}_{rr}$$
 (6)

Moreover, we know that due to hydrostatic equilibrium, $\partial_r \overline{P} \sim -\overline{\rho} \ \overline{g}_r$, and based on Equation (2) we can write, that dilatation flux $\overline{u'_r d''}$ is:

$$\overline{u_r'd''} \sim \frac{\widetilde{R}_{rr} \ \overline{g}_r}{\Gamma_1 \ \overline{P}} \tag{7}$$

Also, the expansion velocity \tilde{u}_r can be replaced by:

$$\widetilde{u}_r = -\partial_t \overline{M} / 4\pi r^2 \overline{\rho} \tag{8}$$

So the alternative hydrodynamic stellar structure equations for hydrostatic convection become a system with only **one unknown the composition flux** f_{α} for which we need a proper model. See validation of these simplified alternative hydrodynamic stellar structure equation in Figure 2.

$$\partial_r \overline{M} = -\overline{\rho} \ \overline{M} \ \overline{g}_r / \Gamma_1 \overline{P} + 4\pi r^2 \overline{\rho} \tag{9}$$

$$\partial_r \overline{P} = -\overline{\rho} \, \overline{g}_r \tag{10}$$

$$\partial_r \widetilde{L} = -4\pi r^2 \widetilde{u}_r \ \overline{\rho} \ \overline{g}_r / \Gamma_1 + \widetilde{\epsilon}_t \partial_r 4\pi r^2 \overline{\rho} \widetilde{u}_r \tag{11}$$

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \,\overline{\rho} \,\overline{T} \,\overline{g}_r / \Gamma_1 \overline{P} \tag{12}$$

$$\partial_t \widetilde{X}_i = \widetilde{X}_i^{nuc} - (1/\overline{\rho}) \nabla_r f_i - \widetilde{u}_r \partial_r \widetilde{X}_i \tag{13}$$

1.1 Derivation of the alternative hydrodynamic stellar structure equations from flux evolution equations

From the pressure flux equation (Fig.5) we get, that:

$$+\overline{u_r'u_r''} \ \partial_r \overline{P} = -\Gamma_1 \ \overline{u_r'Pd} \ \backslash \overline{\rho}$$
 (14)

$$+\overline{\rho} \ \overline{u_r' u_r''} \ \partial_r \overline{P} = -\overline{\rho} \ \Gamma_1 \ \overline{u_r' P d} \qquad P = \overline{P} + P' \quad d = \overline{d} + d' \tag{15}$$

$$+\widetilde{R}_{rr} \partial_r \overline{P} = -\overline{\rho} \Gamma_1 \left(\overline{Pu'_r} \overline{d} + \overline{P} \overline{u'_r d'} + \overline{P'u'_r d'} \right) \qquad \overline{P} \overline{u'_r d'} >> \overline{Pu'_r} \overline{d} + \overline{P'u'_r d'} \text{ and } \overline{a'b'} = \overline{a'b''} = \overline{a''b'}$$

$$(16)$$

$$+\widetilde{R}_{rr} \partial_r \overline{P} = -\overline{\rho} \Gamma_1 \overline{P} \overline{u'_r d''}$$

$$\tag{17}$$

$$\partial_r \overline{P} = -\overline{\rho} \Gamma_1 \overline{P} \overline{u'_r d''} / \widetilde{R}_{rr}$$
(18)

$$\partial_r \overline{P} = -\overline{\rho} \widetilde{g}_r$$
 hydrostatic equilibrium (19)

$$\overline{u'_r d''} = + \widetilde{R}_{rr} \widetilde{g}_r / \Gamma_1 \overline{P}$$
 model for dilatation flux used later (20)

From the turbulent mass flux equation (Fig.5) we get, that:

$$+\widetilde{R}_{rr}/\overline{\rho} \ \partial_r \overline{\rho} = -\overline{\rho} \ \overline{u_r' d''} \tag{21}$$

$$\partial_r \overline{\rho} = -\overline{\rho} \ \overline{\rho} \ \overline{u_r' d''} / \widetilde{R}_{rr} \quad \backslash V \text{ (volume)} = 4\pi r^3 / 3$$
 (22)

$$V\partial_r \overline{\rho} = -\overline{\rho} V \overline{\rho} \overline{u'_r d''} / \widetilde{R}_{rr}$$
(23)

$$\partial_r \overline{\rho} V - \rho \partial_r V = -\overline{\rho} V \overline{\rho} \overline{u'_r d''} / \widetilde{R}_{rr}$$
(24)

$$\partial_r \overline{M} - \rho 4\pi r^2 = -\overline{M} \,\overline{\rho} \,\overline{u_r' d''} / \widetilde{R}_{rr} \tag{25}$$

$$\partial_r \overline{M} = + \rho 4\pi r^2 - \overline{M} \, \overline{\rho} \, \overline{u'_r d''} / \widetilde{R}_{rr} \tag{26}$$

$$\partial_r \overline{M} = + \rho 4\pi r^2 - \overline{\rho} \ \overline{M} \ \widetilde{g}_r / \Gamma_1 \overline{P} \tag{27}$$

$$\partial_r \overline{M} = -\overline{\rho} \ \overline{M} \ \widetilde{g}_r / \Gamma_1 \overline{P} + \rho 4\pi r^2 \tag{28}$$

From the temperature flux equation (Fig.5) we get, that:

$$-\overline{u_r'u_r''}\ \partial_r \overline{T} = +(\Gamma_3 - 1)\ \overline{T}\ \overline{u_r'd''} \ \backslash \overline{\rho}$$
 (29)

$$-\overline{\rho} \ \overline{u'_r u''_r} \ \partial_r \overline{T} = + (\Gamma_3 - 1) \ \overline{\rho} \ \overline{T} \ \overline{u'_r d''}$$

$$(30)$$

$$-\widetilde{R}_{rr} \ \partial_r \overline{T} = + (\Gamma_3 - 1) \ \overline{\rho} \ \overline{T} \ \overline{u'_r d''}$$
(31)

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \ \overline{\rho} \ \overline{T} \ \overline{u_r' d''} / \widetilde{R}_{rr}$$
(32)

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \,\overline{\rho} \,\overline{T} \,(\widetilde{R}_{rr} \widetilde{g}_r / \Gamma_1 \overline{P}) / \widetilde{R}_{rr} \tag{33}$$

$$\partial_r \overline{T} = -(\Gamma_3 - 1) \ \overline{\rho} \ \overline{T} \ \widetilde{g}_r / \Gamma_1 \overline{P}$$
(34)

From the internal energy flux equation (Fig.5) we get, that:

$$\begin{split} -\widetilde{R}_{rr} \ \partial_{r}\widetilde{\varepsilon}_{i} &= + \overline{u_{r}^{\prime\prime}Pd} \quad P = \overline{P} + P' \quad d = \overline{d} + d' \\ -\widetilde{R}_{rr} \ \partial_{r}\widetilde{\varepsilon}_{i} &= + \overline{P}u_{r}^{\prime\prime\prime} \ \overline{d} + \overline{P} \ \overline{u_{r}^{\prime\prime}d^{\prime}} + \overline{P'u_{r}^{\prime\prime}d^{\prime}} \quad \overline{P} \ \overline{u_{r}^{\prime\prime\prime}d^{\prime}} >> \overline{Pu_{r}^{\prime\prime}} \ \overline{d} + \overline{P'u_{r}^{\prime\prime\prime}d^{\prime}} \quad \text{and} \quad \overline{a'b'} = \overline{a'b''} = \overline{a''b'} \end{split}$$

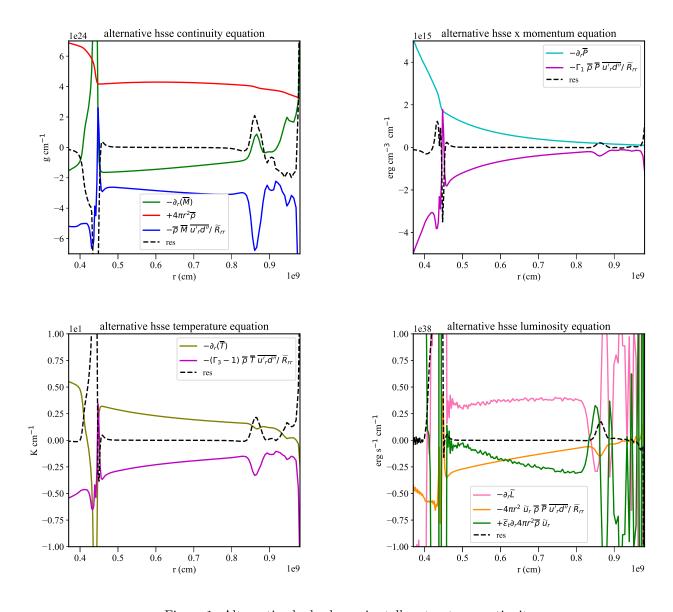


Figure 1: Alternative hydrodynamic stellar structure continuity.

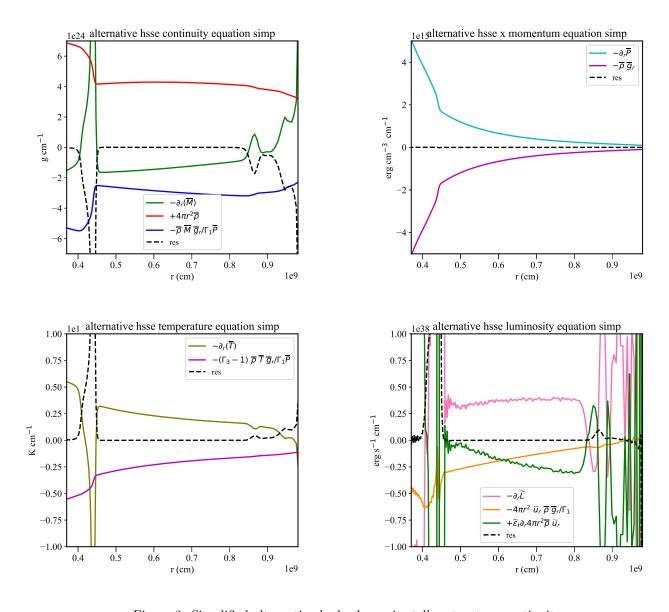


Figure 2: Simplified alternative hydrodynamic stellar structure continuity.

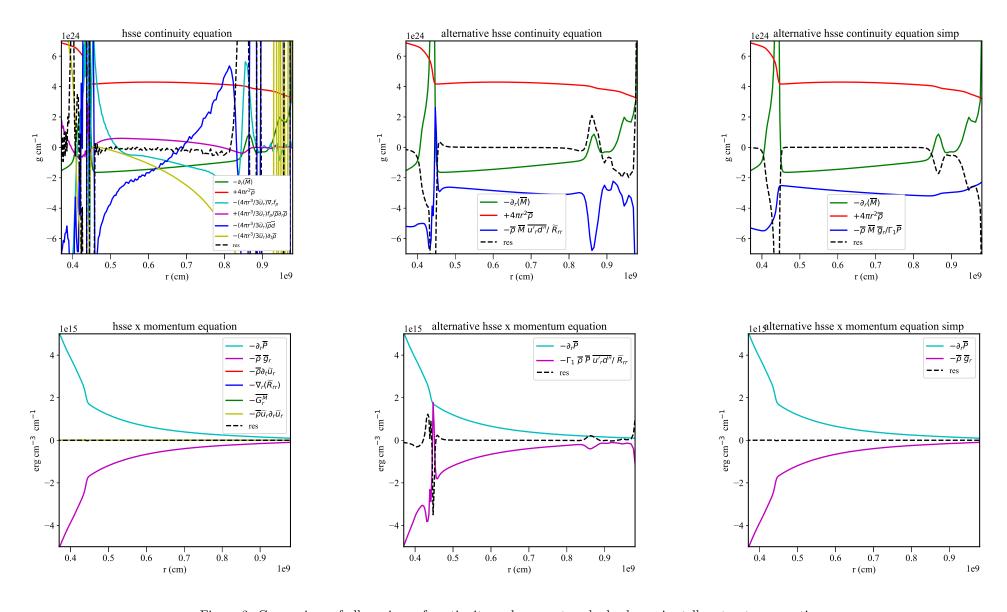


Figure 3: Comparison of all versions of continuity and momentum hydrodynamic stellar structure equations.

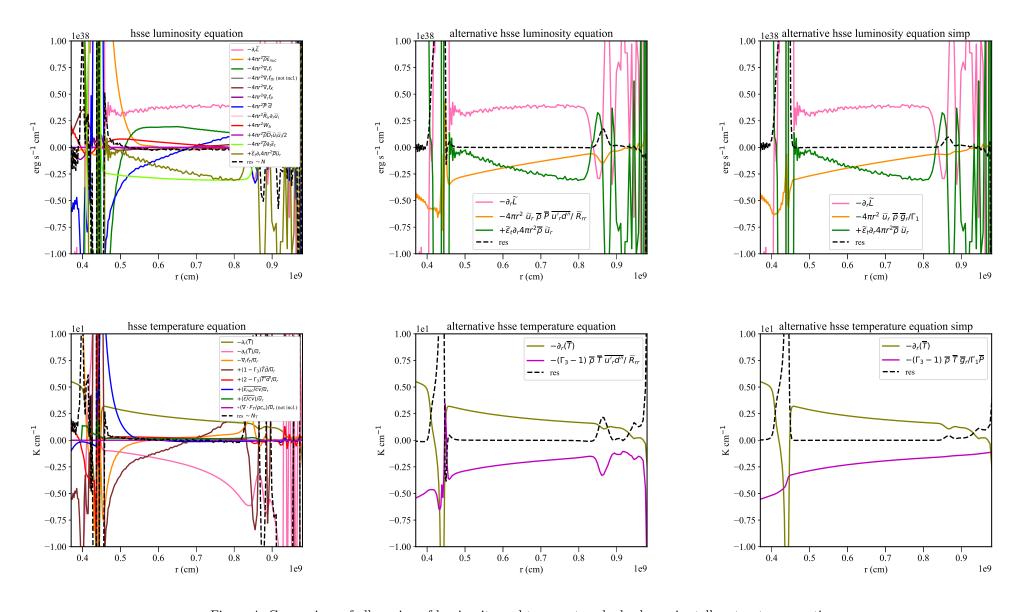


Figure 4: Comparison of all version of luminosity and temperature hydrodynamic stellar structure equations.

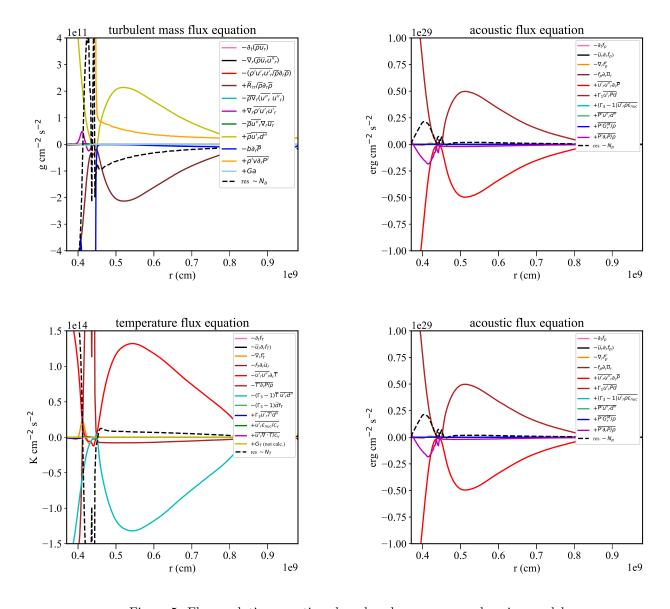


Figure 5: Flux evolution equations based on low-res oxygen burning model.

Table 1: Definitions:

 ρ density g_r radial gravitational acceleration $S = \rho \epsilon_{\text{nuc}}(q)$ nuclear energy production (cooling function) T temperature $\tau_{ij} = 2\mu S_{ij}$ viscous stress tensor (μ kinematic viscosity) P pressure $S_{ij} = (1/2)(\partial_i u_j + \partial_j u_i)$ strain rate u_r, u_θ, u_ϕ velocity components $\widetilde{R}_{ij} = \overline{\rho} \widetilde{u_{i'}' u_{i'}'}$ Reynolds stress tensor $\mathbf{u} = u(u_r, u_\theta, u_\phi)$ velocity $j_z = r \sin \theta \ u_\phi$ specific angular momentum $F_T = \chi \partial_r T$ heat flux $d = \nabla \cdot \mathbf{u}$ dilatation $\Gamma_1 = (d \ln P/d \ln \rho)|_{\mathfrak{s}}$ $\Gamma_2/(\Gamma_2-1)=(d \ln P/d \ln T)|_s$ ϵ_I specific internal energy $\Gamma_3 - 1 = (d \ln T/d \ln \rho)|_s$ h specific enthalpy $k = (1/2)\widetilde{u_i''u_i''}$ turbulent kinetic energy $\widetilde{k}^r = (1/2)\widetilde{u_r''}u_r'' = (1/2)\widetilde{R}_{rr}/\overline{\rho}$ radial turbulent kinetic energy $\widetilde{k}^{\theta} = (1/2)\widetilde{u_{\theta}''u_{\theta}''} = (1/2)\widetilde{R}_{\theta\theta}/\overline{\rho}$ angular turbulent kinetic energy ϵ_k specific kinetic energy $\widetilde{k}^{\phi} = (1/2)\widetilde{u_{\phi}''u_{\phi}''} = (1/2)\widetilde{R}_{\phi\phi}/\overline{\rho}$ angular turbulent kinetic energy ϵ_t specific total energy $\widetilde{k}^h = \widetilde{k}^\theta + \widetilde{k}^\phi$ horizontal turbulent kinetic energy s specific entropy $f_k = (1/2)\overline{\rho}u_i^{\prime\prime}\overline{u_i^{\prime\prime}}u_r^{\prime\prime}$ turbulent kinetic energy flux $v = 1/\rho$ specific volume $f_k^r = (1/2)\overline{\rho}u_r^{\prime\prime}\overline{u_r^{\prime\prime}}u_r^{\prime\prime}$ radial turbulent kinetic energy flux X_{α} mass fraction of isotope α $\dot{X}_{\alpha}^{\mathrm nuc}$ rate of change of X_{α} $f_k^{\theta} = (1/2)\overline{\rho}u_{\theta}^{\prime\prime}u_{\theta}^{\prime\prime}u_r^{\prime\prime}$ angular turbulent kinetic energy flux $f_k^{\phi} = (1/2) \overline{\rho} u_{\phi}^{\prime\prime} u_{\phi}^{\prime\prime} u_r^{\prime\prime}$ angular turbulent kinetic energy flux A_{α} number of nucleons in isotope α $f_k^h = f_k^\theta + f_k^\phi$ horizontal turbulent kinetic energy flux Z_{α} charge of isotope α $W_p = \overline{P'd''}$ turbulent pressure dilatation A mean number of nucleons per isotope $W_b = \overline{\rho} \overline{u_r''} \widetilde{g}_r$ buoyancy Z mean charge per isotope $f_P = \overline{P'u'_r}$ acoustic flux $f_T = -\overline{\chi \partial_r T}$ heat flux (χ thermal conductivity)

Table 2: Definitions (continued):

$$\begin{split} f_I &= \overline{\rho} \widetilde{F'_I u''_I} & \text{ internal energy flux} & f_\alpha &= \overline{\rho} \widetilde{X''_\alpha u''_I} \times X_\alpha \text{ flux} \\ f_S &= \overline{\rho} \widetilde{S'' u''_I} & \text{ entropy flux} & f_{jz} &= \overline{\rho} \widetilde{J'_2 u''_I} & \text{ angular momentum flux} \\ f_T &= \overline{u'_T T'} & \text{ turbulent heat flux} & f_A &= \overline{\rho} \overline{A'' u''_I} & A & (\text{mean number of nucleons per isotope}) & \text{flux} \\ f_h &= \overline{\rho} \widetilde{h'' u''_I} & \text{ enthalpy flux} & f_Z &= \overline{\rho} \overline{Z'' u''_I} & Z & (\text{mean charge per isotope}) & \text{flux} \\ b &= \overline{v'\rho'} & \text{ density-specific volume covariance} & \mathcal{N}_\rho, \mathcal{N}_{ur}, \mathcal{N}_{u\theta}, \mathcal{N}_{u\phi}, \mathcal{N}_{Jz}, \mathcal{N}_\alpha, \mathcal{N}_A, \mathcal{N}_Z & \text{ numerical effect} \\ f_\tau &= f_\tau^\tau + f_\tau^\theta + f_\tau^\phi & \text{ viscous flux} & \mathcal{N}_{\epsilon I} &= -\nabla_r f_\tau + \varepsilon_k & \text{ numerical effect} \\ f_\tau^\theta &= -\overline{v_{pr}} u'_{\theta} & \text{ viscous flux} & \mathcal{N}_{\epsilon k} &= -\varepsilon_k & \text{ numerical effect} \\ f_\tau^\theta &= -\overline{v_{pr}} u'_{\theta} & \text{ viscous flux} & \mathcal{N}_{k} &= -\overline{v_{k}} / T & \text{ numerical effect} \\ f_\tau^\theta &= f_\tau^\theta + f_\tau^\phi & \text{ viscous flux} & \mathcal{N}_h &= -\nabla_r f_\tau & \text{ numerical effect} \\ f_\tau^\theta &= \overline{\rho} \varepsilon''_1 u''_n u''_n & \text{ radial flux of } f_I & \mathcal{N}_P &= + (\Gamma_3 - 1)\varepsilon_k & \text{ numerical effect} \\ f_S^\theta &= \overline{\rho} \widetilde{v''_1} u'''_n & \text{ radial flux of } f_s & \mathcal{N}_T &= + (\overline{v_{ij}} \partial_j u_i) / (c_v \rho) & \text{ numerical effect} \\ f_T^\theta &= \overline{\rho} \widetilde{v''_1} u'''_n & \text{ radial flux of } f_T & \mathcal{N}_{R\theta\theta} &= -2\nabla_r f_\tau^\theta - 2\varepsilon_k^\theta & \text{ numerical effect} \\ f_S^\theta &= \overline{\rho} \widetilde{v''_1} u'''_n & \text{ radial flux of } f_2 & \mathcal{N}_R \phi \phi &= -2\nabla_r f_\tau^\theta - 2\varepsilon_k^\theta & \text{ numerical effect} \\ f_R^\phi &= \overline{\rho} \widetilde{v''_1} u'''_n u'''_n & \text{ radial flux of } f_A & \mathcal{N}_k &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= \overline{\rho} \widetilde{v''_1} u'''_n u'''_n & \text{ radial flux of } f_A & \mathcal{N}_k &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon_k^\theta & \text{ numerical effect} \\ f_N^\theta &= -\nabla_r f_\tau^\theta - \varepsilon$$

Table 3: Definitions (continued):

$$\begin{array}{lll} \mathcal{G}_k^\theta = -(1/2)\overline{G_{Q_0}^\theta} & \mathcal{N}_b \text{ numerical effect} \\ \mathcal{G}_k^\phi = -(1/2)\overline{G_{Q_0}^\phi} & \mathcal{N}_g \mathcal{G}_d^M \\ \mathcal{G}_k^\phi = -(1/2)\overline{G_{Q_0}^\phi} & \mathcal{N}_{fI} = -\nabla_r(\overline{e_I''\tau_r'}) + \overline{u_r''\tau_{ij}\partial_i u_j} - \varepsilon_I \text{ numerical effect} \\ \mathcal{G}_k^h = +\mathcal{G}_k^\theta + \mathcal{G}_k^\phi & \mathcal{N}_{fh} = -\nabla_r(\overline{h''\tau_r'}) + u_r''(\Gamma_3 - 1)\tau_{ij}\partial_i u_j - \overline{u_r''\nabla_i u_i\tau_{ji}} - \varepsilon_h \text{ numerical effect} \\ \mathcal{G}_a = +\overline{\rho'vG_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{s''\tau_{rr}}) + u_r''(\Gamma_3 - 1)\tau_{ij}\partial_i u_j / T - \varepsilon_s \text{ numerical effect} \\ \mathcal{G}_I = -\overline{G_r^I} - \overline{\epsilon_I''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{s''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_A = -\overline{G_r^\alpha} - \overline{X_a''G_r^M} & \mathcal{N}_{fz} = -\nabla_r(\overline{Z''\tau_{rr}}) - \varepsilon_Z \text{ numerical effect} \\ \mathcal{G}_A = -\overline{G_r^A} - A''G_r^M & \mathcal{N}_{fa} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_Z = -\overline{G_r^Z} - \overline{Z''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_A \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_B \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_B \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_B \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_B \text{ numerical effect} \\ \mathcal{G}_B = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{fs} = -\nabla_r(\overline{a''\tau_{rr}}) - \varepsilon_B \text{ numerical effect} \\ \mathcal{N}_{g_1} = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{g_2} = -\overline{G_r^A} - \overline{h''G_r^M} \\ \mathcal{N}_{g_1} = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{g_2} = -\overline{G_r^A} - \overline{h''G_r^M} \\ \mathcal{N}_{g_1} = -\overline{G_r^A} - \overline{h''G_r^M} & \mathcal{N}_{g_2} = -\overline{G_r^A} - \overline{h''G_r^M} \\ \mathcal{N}_{g_1} = -\overline{G_r^A} - \overline{h''G_r^M} \\ \mathcal{N}_{g_2} = -\overline{G_r^A} - \overline{h''G_r^M}$$

Table 4: Definitions (continued):

$$\begin{split} \varepsilon_k^r &= \overline{\gamma_{r'}} \partial_r u_l''' + \overline{\gamma_{r'}} (1/r) \partial_\theta u_l''' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi u_l'' \\ \varepsilon_k^\theta &= \overline{\gamma_{r'}} \partial_r u_\theta''' + \overline{\gamma_{r'}} (1/r) \partial_\theta u_\theta'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\theta &= \overline{\gamma_{r'}} \partial_r u_\theta'' + \overline{\gamma_{r'}} (1/r) \partial_\theta u_\theta'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k^\theta &= \overline{\gamma_{r'}} \partial_r u_\theta'' + \overline{\gamma_{r'}} (1/r) \partial_\theta u_\theta'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi u_\theta'' \\ \varepsilon_k &= \varepsilon_k^r + \varepsilon_k^\theta + \varepsilon_k^\theta \\ \varepsilon_k^\theta &= \varepsilon_k^\theta + \varepsilon_k^\theta \\ \varepsilon_k &= \varepsilon_k^\theta + \varepsilon_k^\theta \\ \varepsilon_k &= \varepsilon_k^\theta + \varepsilon_k^\theta \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r \varepsilon_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta \varepsilon_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta \varepsilon_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi \varepsilon_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r) \partial_\theta v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} (1/r \sin \theta) \partial_\phi v_l'' \\ \varepsilon_k &= \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} \partial_r v_l'' + \overline{\gamma_{r'}} \partial_r v_$$

$$\nabla(.) = \nabla_r(.) + \nabla_{\theta}(.) + \nabla_{\phi}(.) = \frac{1}{r^2} \partial_r(r^2.) + \frac{1}{r \sin \theta} \partial_{\theta}(\sin \theta.) + \frac{1}{r \sin \theta} \partial_{\phi}(.)$$