

MMAE-502 ENGINEERING ANALYSIS II FULL DERIVATION FOR 2D-DIFFUSION EQUATION

EYOB GHEBREIUSUS

BASICS

2D DIFFUSION EQUATION

2D – Determines the statistical randomness of moving particles across a medium.

Earliest known theory was in 1883 by Adolf Fick.

$$\partial_t - k \nabla_{2D}^2 \quad \left| \quad \Theta(t) \left(\frac{1}{4\pi kt} \right) e^{-\rho^2/4kt} \right.$$

Problem Definition

$$\frac{\partial u}{\partial t} = D \nabla^2 u + f(\vec{r} + t)$$

or

$$\nabla^2 P - \frac{1}{D} \frac{\partial P}{\partial t} = 0 \quad \text{in cartesian coordinates}$$

$D \mapsto$ diffusion coefficient constant

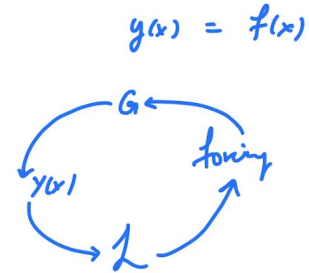
\mapsto sometimes referred as κ

$f(\vec{r} + t) \rightarrow$ forcing function

back to the original problem

$$\nabla^2 P - \frac{1}{D} \frac{\partial P}{\partial t} = 0$$

$$\frac{\partial^2 P}{\partial x^2} + \frac{\partial^2 P}{\partial y^2} - \frac{1}{D} \frac{\partial P}{\partial t} = 0$$



$$\mathcal{L} = \partial_t - \kappa \nabla_{2D}^2$$

↳ linear PDE

↳ x and y do not mix so it's separable

We begin with the 2D-FT in Cartesian solution.

Let \hat{P} is the FT of our original function P

∴ then

$$\hat{P}(k_x, k_y, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x x + k_y y)} \cdot P(x, y, t) dx dy$$

$$P(x, y, t) = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{i2\pi(k_x x + k_y y)} \cdot \hat{P}(k_x, k_y, t) dk_x dk_y$$

- Symmetry in the Fourier Space
- Thus introduce integration by Parts

$$\int_a^b u dv = uv \Big|_a^b - \int_a^b v du$$

$$\text{let } u = e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \quad \text{and} \quad v = \frac{\partial P}{\partial x}$$

$$du = \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} dx \quad \text{and} \quad dv = \frac{\partial^2 P}{\partial x^2} dx$$

by parts:

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial^2 P}{\partial x^2} dx dy =$$

B.C. $\frac{\partial P}{\partial x(x,y)} = 0$
as $|x| = |y| \rightarrow \infty$

$$= e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \frac{\partial P}{\partial x} \Big|_{-\infty(x,y)}^{\infty} - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy$$

$$= e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \frac{\partial P}{\partial x} \Big|_{-\infty(x,y)}^{\infty} - \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy$$

$$= \underbrace{e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \frac{\partial P}{\partial x} \Big|_{-\infty(x,y)}^{\infty}}_0 + 2i\pi k_x \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy$$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial^2 P}{\partial x^2} dx dy = 2i\pi k_x \underbrace{\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy}_{\text{FT of } \frac{dP}{dx} \text{ therefore I follow the same manner}}$$

$$\begin{aligned} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial^2 P}{\partial x^2} dx dy &= 2i\pi k_x \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \frac{\partial}{\partial x} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy \\ &= \underbrace{P \cdot e^{-i2\pi(k_x \cdot x + k_y \cdot y)}}_0 \Big|_{-\infty(x,y)}^{\infty} + 2i\pi k_x \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} P \cdot e^{-i2\pi(k_x \cdot x + k_y \cdot y)} dx dy \end{aligned}$$

P = 0 as $|x|, |y| \rightarrow \infty$

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial P}{\partial x} dx dy = 2i\pi k_x \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} P \cdot e^{-i2\pi(k_x \cdot x + k_y \cdot y)} dx dy$$

therefore

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial^2 P}{\partial x^2} dx dy = (\partial_{x,i} k_x)^2 \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} P \cdot e^{-i2\pi(k_x \cdot x + k_y \cdot y)} dx dy$$
$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-i2\pi(k_x \cdot x + k_y \cdot y)} \cdot \frac{\partial^2 P}{\partial x^2} dx dy = (2\pi i k_x)^2 \hat{P}$$

Typically for n in general we would have:

$$\frac{\partial^n P}{\partial x^n} = (2\pi i k_x)^n \hat{P}$$

- since we used a spatial FT function the derivative in time does not change.
- we can write the 2D-Diffusion Equation after FT

$$(2\pi)^2 \hat{P} \cdot (k_x^2 + k_y^2) + \frac{1}{D} \frac{\partial \hat{P}}{\partial t} = 0$$

$$\frac{\partial \hat{P}}{\partial t} + D(2\pi)^2 (k_x^2 + k_y^2) \cdot \hat{P} = 0$$

→ simple first order ODE with time

- We can Apply method of separation of variables
- assuming $\hat{P}(x, y, t) = T(t)X(x)Y(y) \rightarrow$ product of functions

$$\frac{1}{T} \frac{\partial T}{\partial t} + D(2\pi)^2(k_x^2 + k_y^2) = 0 \rightarrow \text{sub \& dividing by } \hat{P}$$

$$\frac{1}{T} \frac{dT}{dt} = -D(2\pi)^2(k_x^2 + k_y^2) = \text{constant} = \lambda \rightarrow \text{or } A \text{ (Norm)}$$

$$\left(\frac{1}{D} \frac{1}{T} \frac{dT}{dt} = -\lambda \right) \quad \text{and} \quad (2\pi)^2 D(k_x^2 + k_y^2) X \cdot Y = -\lambda T \cdot X \cdot Y$$

sub \& divide

$$\rightarrow T(t) = C e^{-\lambda D t}$$

$$(2\pi)^2(k_x^2 + k_y^2) = -\lambda$$

eigen value problem

$$X(t) = A \sin(k_x \cdot x)$$

$$Y(t) = B \sin(k_y \cdot y)$$

$$\hat{P}(x, y, t) = \sum_{n, m} C_{n, m} \sin(k_x \cdot n x) \sin(k_y \cdot m y) e^{-D(2\pi)^2(k_x^2 + k_y^2)t}$$

or

$$\hat{P} = \lambda e^{-D(2\pi)^2(k_x^2 + k_y^2)t} \rightarrow \text{Norm } \lambda \rightarrow \text{or } A \text{ used for conservation of momentum}$$

- Now back to the function

$$P = A \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-D(2\pi)^2(k_x^2 + k_y^2)t} \cdot e^{i2\pi(k_x \cdot x + k_y \cdot y)} \cdot dk_x dk_y$$

↳ entire integral separation of variables

so:

$$P = A \left[\int_{-\infty}^{\infty} e^{i2\pi k_x \cdot x - D(2\pi k_x)^2 t} dk_x \right] \left[\int_{-\infty}^{\infty} e^{i2\pi k_y \cdot y - D(2\pi k_y)^2 t} dk_y \right]$$

One of the features of the original diffusion equation comes into play here; the entire integral is *separable* by spatial variable:

$$P = A \left(\int_{-\infty}^{\infty} e^{i2\pi k_x \cdot x - D(2\pi k_x)^2 t} dk_x \right) \cdot \left(\int_{-\infty}^{\infty} e^{i2\pi k_y \cdot y - D(2\pi k_y)^2 t} dk_y \right)$$

In general, the number of these separable integrals is directly related to the dimensionality of the diffusion equation; in three dimensions there would be three such integrals. This integral is not trivial...it requires completing the square of the exponent, rescaling the integration variable, and changing to polar coordinates. I will not cover it here, but it is just an integral, so Maple[®] it! Using Maple[®] v8.0 I got:

$$P(x, y, t) = A \frac{e^{\left(\frac{-(x^2 + y^2)}{4D \cdot t} \right)}}{4\pi D \cdot t}$$

We're almost there! What about this normalization constant A ? If we are calling this function P a probability distribution, then it makes sense that the particle must exist somewhere in the x-y plane from x and y to plus and

I WILL DERIVE THIS BY HAND AGAIN FULLY IN THE NEXT SLIDES

back to our original function

$$p = \lambda \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-D(2\pi)^2(k_x^2 + k_y^2)t} \cdot e^{i2\pi(k_x \cdot x + k_y \cdot y)} \cdot dk_x dk_y$$

↳ entire integral separation of variables

so:

$$p = \lambda \left[\int_{-\infty}^{\infty} e^{\underbrace{i2\pi k_x \cdot x - D(2\pi k_x)^2 t}_{1}} dk_x \right] \left[\int_{-\infty}^{\infty} e^{\underbrace{i2\pi k_y \cdot y - D(2\pi k_y)^2 t}_{2}} dk_y \right]$$

completing
the square
integral
& changing
to polar

$$i2\pi k_x \cdot x - D(2\pi k_x)^2 t = -4\pi^2 D t \left(k_x - \frac{ix}{(2\pi D t)} \right)^2 - \frac{x^2}{4Dt}$$

↳ let $u_x = k_x - \frac{ix}{2\pi D t}$

$$-4\pi^2 D t u_x^2 - \frac{x^2}{4Dt}$$

polar $x^2 + y^2 = r^2$
for isotropic
"independent
of direction"

Similarly apply the same process to 2nd part
and substitute u_y

$$\int_{-\infty}^{\infty} e^{(i2\pi k_x \cdot x - D(2\pi k_x)^2 \cdot t)} dt = e^{\frac{-x^2}{4Dt}} \cdot \int_{-\infty}^{\infty} e^{-4\pi^2 Dt \cdot u_x^2} \cdot du_x$$

say $v = 2\pi\sqrt{Dt} u$

$$\int_{-\infty}^{\infty} e^{-\frac{v^2}{4}} dv$$

non elementary
integral
the Louville
-Ritt

$$\left[\int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} dx \right]^2$$

$$= \left[\int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} dx \right] \left[\int_{-\infty}^{\infty} e^{-\frac{y^2}{2}} dy \right]$$

$$= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{(x^2+y^2)}{2}} dx dy$$

$$= \int_0^{2\pi} \int_0^{\infty} e^{-\frac{r^2}{2}} r dr d\theta$$

$$\int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} dx = \sqrt{2\pi}$$

here

$$\left. \begin{aligned} x &= r \cos \theta \\ y &= r \sin \theta \\ x^2 + y^2 &= r^2 \\ dx dy &= r dr d\theta \end{aligned} \right\} \text{polar}$$

make a sub $u = r^2/2$

$$= 2\pi [e^{-u}]_0^\infty$$

$$= 2\pi$$

$$\left[\int_{-\infty}^{\infty} e^{-\frac{x^2}{2}} dx \right]^2 = \int_0^{2\pi} \int_0^{2\pi} e^{-\frac{r^2}{2}} r dr d\theta = 2\pi$$

$$\int_{-\infty}^{\infty} e^{-x^2} dx = \sqrt{2\pi}$$

Thus

$$\int_{-\infty}^{\infty} e^{-4\pi^2 D \cdot t \cdot u^2} du = \int_{-\infty}^{\infty} e^{-\frac{v^2}{4}} dv = 2\sqrt{\pi}$$

sub back for v and multiply by $e^{-\frac{x^2}{4Dt}}$.

$$= \frac{\sqrt{\pi}}{4Dt} e^{-\frac{x^2}{4Dt}}$$

$$\int_{-\infty}^{\infty} e^{(i2\pi k_x \cdot x - D(2\pi k_x)^2 \cdot t)} dt = e^{\frac{-x^2}{4Dt}} \cdot \int_{-\infty}^{\infty} e^{-4\pi^2 Dt \cdot u_x^2} \cdot du_x$$

$$= \frac{\sqrt{\pi}}{4Dt} e^{\frac{-x^2}{4Dt}}$$

$$\int_{-\infty}^{\infty} e^{(i2\pi k_y \cdot y - D(2\pi k_y)^2 \cdot t)} dt = e^{\frac{-y^2}{4Dt}} \int_{-\infty}^{\infty} e^{-4\pi^2 Dt \cdot u_y^2} \cdot du_y$$

$$= \frac{\sqrt{4}}{4Dt} \cdot e^{\frac{-y^2}{4Dt}}$$

→ non-elementary
integral
the Liouville
-Ritt

Combine all

$$p(x, y, t) = \frac{A}{4\pi Dt} \left[e^{\frac{-(x^2 + y^2)}{4Dt}} \right]$$

$x^2 + y^2 = r^2$ in polar coordinates
but what is A ?

A \rightarrow how many non-interacting particles we have in the system.

Impose a condition by making the total probability sum to 1. i.e.

$$\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} P(x, y, t) dx dy = 1 \quad \rightarrow \text{unity function}$$

Then for

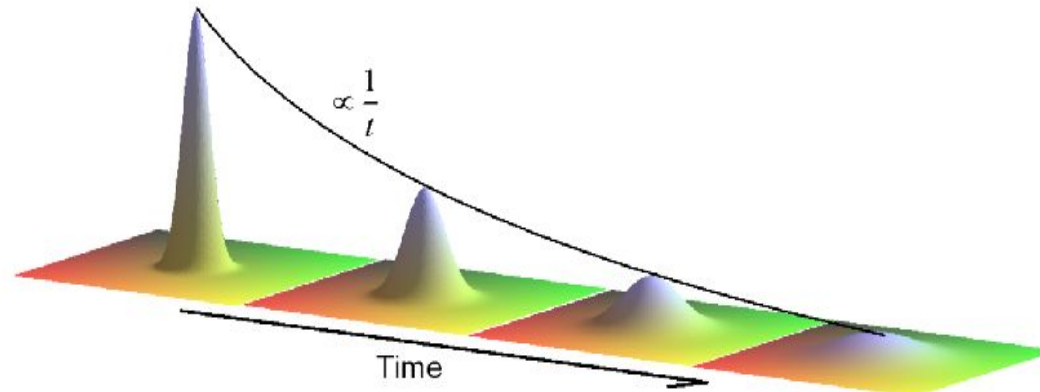
$$P(r, t) = \frac{e^{-\frac{r^2}{4Dt}}}{4\pi Dt}$$

- But what is diffusion coefficient D ?
- has units of $\frac{L^2}{t}$
- relates to the rate of increase the mean distance of the particles with time.

$$D \propto \frac{k_B T}{\eta \cdot R} \Rightarrow D = \frac{k_B T}{6\pi\eta \cdot R} \quad \nearrow \text{spherical}$$

- Helps determine the medium viscosity.

This is called a normalized Gaussian function. The first plot below shows diffusion at early times; notice how high the peak is, meaning the particles are localized around the area where they were first introduced. The remaining plots demonstrate how the particles spread out from their original entry point with time:



Source: Prof. Urser Caltech

QUESTIONS

