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Implementing Density Functional Theory

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Preface

Importance of density functional theory

Implementation of density functional theory in various program packages (free and commercial)

Several books about density functional theories

The problem: not yet giving necessary details

This book is our humble attempt to demystifying several aspects of practical density functional theory to beginners in the field.

Objective of this book: show the reader how to implement a density functional theory for simple system containing only model potential (such as harmonic potential) and to non-local pseudopotentials which are usually used in typical DFT calculations for molecular and crystalline systems.

Outline of the book: 1d, 2d, 3d, Schrodinger equation, Poisson equation, Kohn-Sham equation for local (pseudo)potentials, Kohn-Sham equation for nonlocal pseudopotentials.

This is for acknowledgments.

Bandung, month year Fadjar Fathurrahman Hermawan Kresno Dipojono

Contents

1	An introduction to density functional theory	1
	References	1
2	Schroedinger equation in 1d	3
	2.1 Approximating second derivative	4
	2.2 Harmonic potential	6
	2.3 Higher order finite difference	8
	2.4 Exercises	9
3	Schroedinger equation in 2d	11
	3.1 Finite difference grid in 2d	11
	3.2 Laplacian operator	12
	3.3 Iterative methods for eigenvalue problem	13
4	Schroedinger equation in 3d	15
5	Numerical solution of Poisson equation	17
6	Kohn-Sham equation part I	19
7	Numerical solution of Kohn-Sham equation (part II)	21
A	Introduction to Julia programming language	23
	A.1 Installation	23

viii	Contents

	A.2	Using Julia	23
		A.2.1 Using Julia REPL	23
		A.2.2 Julia script file	24
	A.3	Basic programming construct	25
	A.4	Mathematical operators	25
В	Intr	oduction to Octopus DFT code	27
Ind	ex		29

An introduction to density functional theory

Intro to DFT [1].

Kohn-Sham equation

$$\nabla^2 + V(\mathbf{r})\psi_i \mathbf{r} = E\psi_i(\mathbf{r}) \tag{1.1}$$

References

1. Test Kohn Sham

Schroedinger equation in 1d

The equation we want to solve is (in Hartree atomic unit):

$$\left[-\frac{1}{2} \frac{\mathrm{d}^2}{\mathrm{d}x^2} + V(x) \right] \psi(x) = E \psi(x) \tag{2.1}$$

with the boundary conditions:

$$\lim_{x \to +\infty} = 0 \tag{2.2}$$

We will discretize the wave function, potentials, and various spatial quantities using regular grid within the finite computational interval $[x_{\min}, x_{\max}]$. This computational interval should be choosen such that the boundary condition 2.2 approximately satisfied.

We will choose the grid points x_i , i = 1, 2, ... as:

$$x_i = x_{\min} + (i - 1)h \tag{2.3}$$

where N is the number of grid points and h is the spacing between the grid points is:

$$h = \frac{x_{\text{max}} - x_{\text{min}}}{N - 1} \tag{2.4}$$

The following code can be used to initialize the grid points:

```
function init_FDld_grid( x_min::Float64, x_max::Float64, N::Int64 )
    L = x_max - x_min
    h = L/(N-1)
    x = zeros(Float64,N)
    for i = 1:N
        x[i] = x_min + (i-1)*h
    end
    return x, h
```

end

```
init_FDld_grid( X, N ) = init_FDld_grid( X[1], X[2], N )
```

2.1 Approximating second derivative

With the following notation: $\psi_i = \psi(x_i)$, we can use 3-point finite difference to approximate second derivative of $\psi(x)$:

$$\frac{d^2}{dx^2}\psi_i = \frac{\psi_{i+1} - 2\psi_i + \psi_{i-1}}{h^2}$$
 (2.5)

Take $\{\psi_i\}$ as (column) vector, we can represent the second derivative operation as matrix multiplication:

$$" = \mathbb{D}^{(2)} \tag{2.6}$$

where $\mathbb{D}^{(2)}$ is the second derivative matrix operator

$$\mathbb{D}^{(2)} = \frac{1}{h^2} \begin{bmatrix} -2 & 1 & 0 & 0 & 0 & \cdots & 0 \\ 1 & -2 & 1 & 0 & 0 & \cdots & 0 \\ 0 & 1 & -2 & 1 & 0 & \cdots & 0 \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & 0 & 1 & -2 & 1 & 0 \\ 0 & \cdots & \cdots & 0 & 1 & -2 & 1 \\ 0 & \cdots & \cdots & \cdots & 0 & 1 & -2 \end{bmatrix}$$
 (2.7)

An example implementation can be found in file build_D2_matrix_3pt.jl.

. . .

```
Build second derivative matrix using 3-points centered finite difference approximation.
```

```
# Arguments
- `N::Int64`: number of grid points
- `h::Float64`: spacing between grid points
"""

function build_D2_matrix_3pt( N::Int64, h::Float64 )
    mat = zeros(Float64,N,N)
    for i = 1:N-1
        mat[i,i] = -2.0
        mat[i,i+1] = 1.0
        mat[i+1,i] = mat[i,i+1]
    end
```

Test with Gaussian function:

$$\psi(x) = e^{-\alpha x^2} \tag{2.8}$$

which second derivative can be calculated as

$$\psi''(x) = (-2\alpha + 4\alpha^2 x^2) e^{-\alpha x^2}$$
 (2.9)

They are implemented in the following code

```
function my_gaussian(x; \alpha=1.0) return exp(-\alpha*x^2) end function d2_my_gaussian(x; \alpha=1.0) return (-2*\alpha + 4*\alpha^2 * x^2) * exp(-\alpha*x^2) end
```

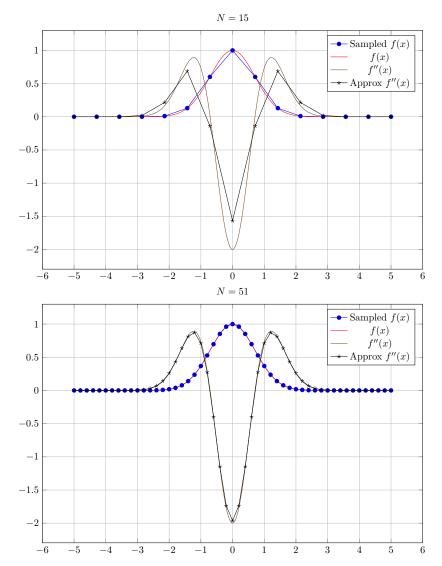


Fig. 2.1 Finite difference approximation to a Gaussian function and its second derivative

2.2 Harmonic potential

We will start with a simple potential with known exact solution, namely the harmonic potential:

$$V(x) = \frac{1}{2}\omega^2 x^2 {(2.10)}$$

The Hamiltonian in finite difference representation:

$$\mathbb{H} = -\frac{1}{2}\mathbb{D}^{(2)} + \mathbb{V} \tag{2.11}$$

where \mathbb{V} is a diagonal matrix whose elements are:

$$V_{ij} = V(x_i)\delta_{ij} \tag{2.12}$$

Code to solve harmonic oscillator:

```
using Printf
using LinearAlgebra
using PGFPlotsX
using LaTeXStrings
include("init_FD1d_grid.jl")
include("build_D2_matrix_3pt.jl")
function pot_harmonic( x; \omega = 1.0 )
   return 0.5 * \omega^2 * x^2
end
function main()
   # Initialize the grid points
   xmin = -5.0
   xmax = 5.0
   N = 51
   x, h = init_FDld_grid(xmin, xmax, N)
   # Build 2nd derivative matrix
   D2 = build_D2_matrix_3pt(N, h)
   # Potential
   Vpot = pot_harmonic.(x)
   # Hamiltonian
   Ham = -0.5*D2 + diagm( 0 => Vpot )
   # Solve the eigenproblem
   evals, evecs = eigen( Ham )
   # We will show the 5 lowest eigenvalues
   Nstates = 5
    @printf("Eigenvalues\n")
    for i in 1:Nstates
       @printf("%5d %18.10f\n", i, evals[i])
   # normalize the first three eigenstates
    for i in 1:3
        ss = dot(evecs[:,i], evecs[:,i])*h
        evecs[:,i] = evecs[:,i]/sqrt(ss)
   # Plot up to 3rd eigenstate
```

Compare with analytical solution.

Plot of eigenfunctions:

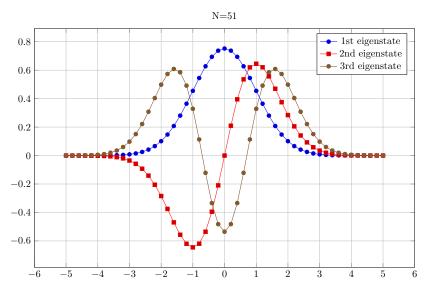


Fig. 2.2 Eigenstates of harmonic oscillator

2.3 Higher order finite difference

To obtain higher accuracy

Implementing higher order finite difference.

2.4 Exercises 9

2.4 Exercises

Gaussian potential

Schroedinger equation in 2d

Schrodinger equation in 2d:

$$\left[-\frac{1}{2} \nabla^2 + V(x, y) \right] \psi(x, y) = E \psi(x, y)$$
(3.1)

where ∇^2 is the Laplacian operator:

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \tag{3.2}$$

3.1 Finite difference grid in 2d

```
struct FD2dGrid
   Npoints::Int64
   Nx:: Int64
   Ny:: Int64
   hx::Float64
   hy::Float64
   dA::Float64
   x::Array{Float64,1}
   y::Array{Float64,1}
   r::Array{Float64,2}
   idx_ip2xy::Array{Int64,2}
    idx_xy2ip::Array{Int64,2}
function FD2dGrid( x_domain, Nx, y_domain, Ny )
   x, hx = init_FD1d_grid(x_domain, Nx)
   y, hy = init_FD1d_grid(y_domain, Ny)
   dA = hx*hy
   Npoints = Nx*Ny
   r = zeros(2,Npoints)
```

```
ip = 0
idx_ip2xy = zeros(Int64,2,Npoints)
idx_xy2ip = zeros(Int64,Nx,Ny)
for j in 1:Ny
    for i in 1:Nx
        ip = ip + 1
        r[1,ip] = x[i]
        r[2,ip] = y[j]
        idx_ip2xy[1,ip] = i
        idx_ip2xy[2,ip] = j
        idx_xy2ip[i,j] = ip
    end
end
return FD2dGrid(Npoints, Nx, Ny, hx, hy, dA, x, y, r, idx_ip2xy,
        idx_xy2ip)
end
```

3.2 Laplacian operator

Given second derivative matrix in x, $\mathbb{D}_x^{(2)}$, y direction, $\mathbb{D}_x^{(2)}$, we can construct finite difference representation of the Laplacian operator \mathbb{L} by using

$$\mathbb{L} = \mathbb{D}_{x}^{(2)} \otimes \mathbb{I}_{y} + \mathbb{I}_{x} \otimes \mathbb{D}_{y}^{(2)} \tag{3.3}$$

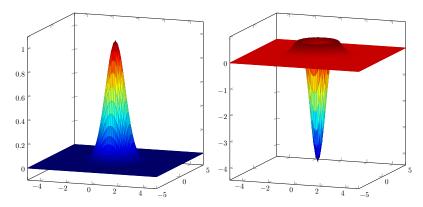
where \otimes is Kronecker product. In Julia, we can use the function kron to form the Kronecker product between two matrices A and B as kron(A,B).

```
function build_nabla2_matrix( fdgrid::FD2dGrid; func_1d=build_D2_matrix_3pt )
    Nx = fdgrid.Nx
    hx = fdgrid.hx
    Ny = fdgrid.Ny
    hy = fdgrid.hy

    D2x = func_1d(Nx, hx)
    D2y = func_1d(Ny, hy)

    ∇2 = kron(D2x, speye(Ny)) + kron(speye(Nx), D2y)
    return ∇2
end
```

Example to the approximation of 2nd derivative of 2d Gaussian function



 $\textbf{Fig. 3.1} \ \, \textbf{Two-dimensional Gaussian function and its finite difference approximation of second derivative}$

3.3 Iterative methods for eigenvalue problem

The Hamiltonian matrix:

```
\nabla 2 = build\_nabla2\_matrix( fdgrid, func_1d=build_D2\_matrix_9pt )  

Ham = -0.5*\nabla 2 + spdiagm( 0 => Vpot )
```

The Hamiltonian matrix size is large. The use eigen method to solve this eigenvalue problem is not practical. We also do not need to solve for all eigenvalues. We must resort to the so called iterative methods.

Schroedinger equation in 3d

Extension to 3d [?]

Schrodinger equation in 3d:

$$\left[-\frac{1}{2}\nabla^{2} + V(\mathbf{r})\right]\psi(\mathbf{r}) = E\psi(\mathbf{r})$$
(4.1)

where **r** is abbreviation to (x, y, z). where ∇^2 is the Laplacian operator:

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$
 (4.2)

```
const ⊗ = kron
function build_nabla2_matrix( fdgrid::FD3dGrid; func_1d=build_D2_matrix_3pt )
    D2x = func_1d(fdgrid.Nx, fdgrid.hx)
    D2y = func_1d(fdgrid.Ny, fdgrid.hy)
    D2z = func_1d(fdgrid.Nz, fdgrid.hz)
    IIx = speye(fdgrid.Nx)
    IIy = speye(fdgrid.Ny)
    IIz = speye(fdgrid.Nz)
    ∇2 = D2x⊗IIy⊗IIz + IIx⊗D2y⊗IIz + IIx⊗IIy⊗D2z
    return ∇2
end
```

Numerical solution of Poisson equation

We will turn our attention to the Poisson equation:

$$\nabla^2 V_H(\mathbf{r}) = 4\pi \,\rho(\mathbf{r}) \tag{5.1}$$

Introduction to conjugate gradient problem 3d dimensional problem

Kohn-Sham equation part I

Using local potential only

Numerical solution of Kohn-Sham equation (part II)

Using nonlocal potential (pseudopotential)

Appendix A

Introduction to Julia programming language

This chapter is intended to as an introduction to the Julia programming language.

This chapter assumes familiarity with command line interface.

A.1 Installation

Go to https://julialang.org/downloads/ and download the suitable file for your platform. For example, on 64 bit Linux OS, we can download the file julia-1.x.x-linux-x86_64.tar.gz where 1.x.x referring to the version of Julia. After you have downloaded the tarball you can unpack it.

```
tar xvf julia-1.x.x-linux-x86_64.tar.gz
```

After unpacking the tarball, there should be a new folder called julia-1.x.x. You might want to put this directory under your home directory (or another directory of your preference).

A.2 Using Julia

A.2.1 Using Julia REPL

Let's assume that you have put the Julia distribution under your home directory. You can start the Julia interpreter by typing:

/home/username/julia-1.x.x/bin/julia

You should see something like this in your terminal:

\$ julia



This is called the Julia REPL (read-eval-print loop) or the Julia command prompt. You can type the Julia program and see the output. This is useful for interactive exploration or debugging the program.

The Julia code can be typed after the julia> prompt. In this way, we can write Julia code interactively.

Example Julia session

```
julia> 1.2 + 3.4
4.6

julia> sin(2*pi)
-2.4492935982947064e-16

julia> sin(2*pi)^2 + cos(2*pi)^2
1.0

Using Unicode:
julia> α = 1234;
julia> β = 3456;
julia> α * β
4264704

To exit type
julia> exit()
```

A.2.2 Julia script file

In a text file with .jl extension.

You can experiment with Julia REPL by typing julia at terminal:

We also can put the code in a text file with .jl extension and execute it with the command:

```
julia filename.jl
The following code
function say_hello(name)
    println("Hello: ", name)
end
say_hello("efefer")
```

A.3 Basic programming construct

Julia has similarities with several popular programming languages such as Julia, MATLAB, and R, to name a few.

A.4 Mathematical operators

```
if a >= 1
  println("a is larger or equal to 1")
end
Example code 3
using PGFPlotsX
using LaTeXStrings
include("init_FD1d_grid.jl")
function my_gaussian(x::Float64; \alpha=1.0)
  return exp( -\alpha*x^2 )
function main()
 A = -5.0
 B = 5.0
 Npoints = 8
 x, h = init_FD1d_grid( A, B, Npoints )
  NptsPlot = 200
  x_dense = range(A, stop=5, length=NptsPlot)
  f = @pgf(
    Axis( \{\text{height} = "6cm", \text{width} = "10cm" \},
```

```
PlotInc( {mark="none"}, Coordinates(x_dense, my_gaussian.(x_dense)) ),
    LegendEntry(L"f(x)"),
    PlotInc( Coordinates(x, my_gaussian.(x)) ),
    LegendEntry(L"Sampled $f(x)$"),
    )
    )
    pgfsave("TEMP_gaussian_ld.pdf", f)
end
main()
```

Appendix B

Introduction to Octopus DFT code

Prerequisites:

- Autotools and GNU Make
- C, C++ and Fortran compilers
- Libxc

autoreconf --install
./configure --prefix=path_to_install
make
make install

Index

Kohn-Sham equation, 1