Implementing Density Functional Theory using Finite Difference Method

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1 Introduction

Kohn-Sham equations:

$$\hat{H}_{KS}\,\psi_i(\mathbf{r}) = \epsilon_i\,\psi_i(\mathbf{r})\tag{1}$$

2 Schroedinger equation in 1d

We are interested in finding bound states solution to 1d time-independent Schroedinger equation:

$$\left[-\frac{1}{2} \frac{\mathrm{d}^2}{\mathrm{d}x^2} + V(x) \right] \psi(x) = E \psi(x) \tag{2}$$

with the boundary conditions:

$$\lim_{x \to +\infty} \psi(x) = 0 \tag{3}$$

First we need to define a spatial domain $[x_{\min}, x_{\max}]$ where x_{\min}, x_{\max} chosen such that the boundary condition 3 is approximately satisfied. The next step is to divide the spatial domain x using equally-spaced grid points which we will denote as $\{x_1, x_2, \ldots, x_N\}$ where N is number of grid points. Various spatial quantities such as wave function and potential will be discretized on these grid points. The grid points x_i , $i = 1, 2, \ldots$ are chosen as:

$$x_i = x_{\min} + (i-1)h \tag{4}$$

where h is the spacing between the grid points:

$$h = \frac{x_{\text{max}} - x_{\text{min}}}{N - 1} \tag{5}$$

The following code can be used to initialize the grid points:

Approximating second derivative

Our next task is to find an approximation to the second derivative operator present in the Equation (2). One simple approximation that we can use is the 3-point (central) finite difference:

$$\frac{\mathrm{d}^2}{\mathrm{d}x^2}\psi_i = \frac{\psi_{i+1} - 2\psi_i + \psi_{i-1}}{h^2} \tag{6}$$

where we have the following notation have been used: $\psi_i = \psi(x_i)$. By taking $\{\psi_i\}$ as a column vector, the second derivative operation can be expressed as matrix multiplication:

$$\vec{\psi''} = \mathbb{D}^{(2)}\vec{\psi} \tag{7}$$

where $\mathbb{D}^{(2)}$ is the second derivative matrix operator:

$$\mathbb{D}^{(2)} = \frac{1}{h^2} \begin{bmatrix} -2 & 1 & 0 & 0 & 0 & \cdots & 0 \\ 1 & -2 & 1 & 0 & 0 & \cdots & 0 \\ 0 & 1 & -2 & 1 & 0 & \cdots & 0 \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & 0 & 1 & -2 & 1 & 0 \\ 0 & \cdots & \cdots & 0 & 1 & -2 & 1 \\ 0 & \cdots & \cdots & 0 & 1 & -2 \end{bmatrix}$$
 (8)

An example implementation can be found in the following function.

```
function build_D2_matrix_3pt( N::Int64, h::Float64 )
   mat = zeros(Float64,N,N)
   for i = 1:N-1
        mat[i,i] = -2.0
        mat[i,i+1] = 1.0
        mat[i+1,i] = mat[i,i+1]
   end
   mat[N,N] = -2.0
   return mat/h^2
```

Before use this function to solve Schroedinger equation we will to test the operation in Equation (8) for a simple function which second derivative can be calculated analytically.

$$\psi(x) = e^{-\alpha x^2} \tag{9}$$

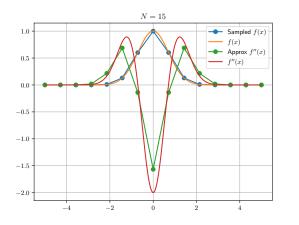
which second derivative can be calculated as

$$\psi''(x) = \left(-2\alpha + 4\alpha^2 x^2\right) e^{-\alpha x^2} \tag{10}$$

They are implemented in the following code

```
function my_gaussian(x; a=1.0)
    return exp(-a*x^2)
end

function d2_my_gaussian(x; a=1.0)
    return (-2*a + 4*a^2 * x^2) * exp(-a*x^2)
end
```



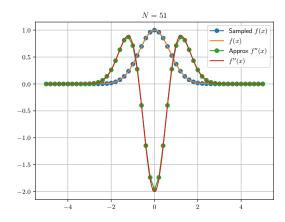


Figure 1: Finite difference approximation to a Gaussian function and its second derivative

Harmonic potential

We will start with a simple potential with known exact solution, namely the harmonic potential:

$$V(x) = \frac{1}{2}\omega^2 x^2 \tag{11}$$

The Hamiltonian in finite difference representation:

$$\mathbb{H} = -\frac{1}{2}\mathbb{D}^{(2)} + \mathbb{V} \tag{12}$$

where $\mathbb V$ is a diagonal matrix whose elements are:

$$V_{ij} = V(x_i)\delta_{ij} \tag{13}$$

Code to solve harmonic oscillator:

```
using Printf
using LinearAlgebra
using LaTeXStrings
import PyPlot
const plt = PyPlot
plt.rc("text", usetex=true)
include("init_FD1d_grid.jl")
include("build_D2_matrix_3pt.jl")
function pot_harmonic( x; \omega=1.0 )
    return 0.5 * \omega^2 * x^2
end
function main()
    # Initialize the grid points
    xmin = -5.0
    xmax = 5.0
    N = 51
    x, h = init_FD1d_grid(xmin, xmax, N)
    # Build 2nd derivative matrix
    D2 = build_D2_matrix_3pt(N, h)
    # Potential
    Vpot = pot_harmonic.(x)
    # Hamiltonian
    Ham = -0.5*D2 + diagm(0 => Vpot)
    # Solve the eigenproblem
    evals, evecs = eigen( Ham )
    # We will show the 5 lowest eigenvalues
    Nstates = 5
    @printf("Eigenvalues\n")
    for i in 1:Nstates
        @printf("%5d %18.10f\n", i, evals[i])
    end
    # normalize the first three eigenstates
    for i in 1:3
        ss = dot(evecs[:,i], evecs[:,i])*h
        evecs[:,i] = evecs[:,i]/sqrt(ss)
    end
    # Plot up to 3rd eigenstate
    plot_title = "N="*string(N)
    plt.plot(x, evecs[:,1], label="1st eigenstate", marker="o")
    plt.plot(x, evecs[:,2], label="2nd eigenstate", marker="o")
    plt.plot(x, evecs[:,3], label="3rd eigenstate", marker="o")
    plt.legend()
    plt.tight_layout()
    plt.savefig("IMG_main_harmonic_01_"*string(N) *".pdf")
end
main()
```

Compare with analytical solution.

Plot of eigenfunctions:

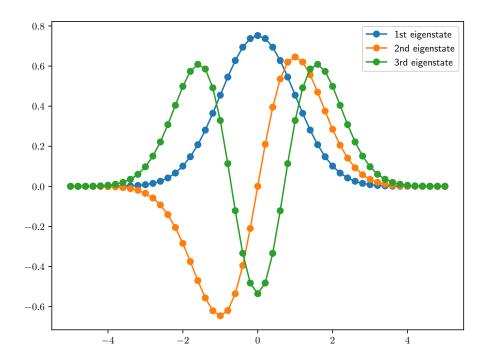


Figure 2: Eigenstates of harmonic oscillator

2.1 Higher order finite difference

An alternative to using more grid points To obtain higher accuracy Implementing higher order finite difference.

A Alternative solutions

```
function my_func()
    println("OK ...")
end
```