

```

from __future__ import annotations
from abc import ABC, abstractmethod
from dataclasses import dataclass
import sys

import numpy as np
from ase.units import Hartree

import gpaw.mpi as mpi

from gpaw.response.pw_parallelization import Blocks1D
from gpaw.response.coulomb_kernels import CoulombKernel
from gpaw.response.dyson import DysonEquation
from gpaw.response.density_kernels import DensityXCKernel
from gpaw.response.chi0 import Chi0Calculator, get_frequency_descriptor
from gpaw.response.chi0_data import Chi0Data
from gpaw.response.pair import get_gs_and_context

from typing import TYPE_CHECKING

if TYPE_CHECKING:
    from gpaw.response.groundstate import CellDescriptor
    from gpaw.response.frequencies import FrequencyDescriptor
    from gpaw.response.pair_functions import SingleQPWDescriptor

```

"""

On the notation in this module.

When calculating properties such as the dielectric function, EELS spectrum and polarizability there are many inherent subtleties relating to (ir)reducible representations and inclusion of local-field effects. For the reciprocal space representation of the Coulomb potential, we use the following notation

v or $v(q)$: The bare Coulomb interaction, $4\pi/|G+q|^2$

V or $V(q)$: The specified Coulomb interaction. Will usually be either the bare interaction or a truncated version hereof.

\bar{V} or $\bar{V}(q)$: The modified Coulomb interaction. Equal to $V(q)$ for finite reciprocal wave vectors $G > 0$, but modified to exclude long-range interactions, that is, equal to 0 for $G = 0$.

"""

```

@dataclass
class Chi0DysonEquations:
    chi0: Chi0Data
    coulomb: CoulombKernel
    xc_kernel: DensityXCKernel | None
    cd: CellDescriptor

    def __post_init__(self):
        if self.coulomb.truncation is None:
            self.bare_coulomb = self.coulomb
        else:
            self.bare_coulomb = self.coulomb.new(truncation=None)
        # When inverting the Dyson equation, we distribute frequencies globally
        blockdist = self.chi0.body.blockdist.new_distributor(nblocks='max')
        self.wblocks = Blocks1D(blockdist.blockcomm, len(self.chi0.wd))

    @staticmethod
    def _normalize(direction):
        if isinstance(direction, str):
            d_v = {'x': [1, 0, 0],
                    'y': [0, 1, 0],
                    'z': [0, 0, 1]}[direction]
        else:
            d_v = direction

```

```

d_v = np.asarray(d_v) / np.linalg.norm(d_v)
return d_v

def get_chi0_wGG(self, direction='x'):
    chi0 = self.chi0
    chi0_wGG = chi0.body.get_distributed_frequencies_array().copy()
    if chi0.qpd.optical_limit:
        # Project head and wings along the input direction
        d_v = self._normalize(direction)
        W_w = self.wblocks.myslice
        chi0_wGG[:, 0] = np.dot(d_v, chi0.chi0_WxvG[W_w, 0])
        chi0_wGG[:, :, 0] = np.dot(d_v, chi0.chi0_WxvG[W_w, 1])
        chi0_wGG[:, 0, 0] = np.dot(d_v, np.dot(chi0.chi0_Wvv[W_w], d_v).T)
    return chi0_wGG

def get_coulomb_scaled_kernel(self, modified=False, Kxc_GG=None):
    """Get the Hxc kernel rescaled by the bare Coulomb potential v(q).

    Calculates


$$\tilde{K}(q) = v^{(-1/2)}(q) K_{Hxc}(q) v^{(-1/2)}(q),$$


    where v(q) is the bare Coulomb potential and


$$K_{Hxc}(q) = V(q) + K_{xc}(q).$$


    When using the `modified` flag, the specified Coulomb kernel will be
    replaced with its modified analogue,


$$V(q) \rightarrow \tilde{V}(q)$$

    """
    qpd = self.chi0.qpd
    if self.coulomb is self.bare_coulomb:
        v_G = self.coulomb.V(qpd) # bare Coulomb interaction
        K_GG = np.eye(len(v_G), dtype=complex)
    else:
        v_G = self.bare_coulomb.V(qpd)
        V_G = self.coulomb.V(qpd)
        K_GG = np.diag(V_G / v_G)
    if modified:
        K_GG[0, 0] = 0.
    if Kxc_GG is not None:
        sqrtv_G = v_G**0.5
        K_GG += Kxc_GG / sqrtv_G / sqrtv_G[:, np.newaxis]
    return v_G, K_GG

@staticmethod
def invert_dyson_like_equation(in_wGG, K_GG, reuse_buffer=True):
    """Generalized Dyson equation inversion.

    Calculates


$$B(q, \omega) = [1 - A(q, \omega) K(q)]^{-1} A(q, \omega)$$


    while possibly storing the output B(q,ω) in the input A(q,ω) buffer.
    """
    if reuse_buffer:
        out_wGG = in_wGG
    else:
        out_wGG = np.zeros_like(in_wGG)
    for w, in_GG in enumerate(in_wGG):
        out_wGG[w] = DysonEquation(in_GG, in_GG @ K_GG).invert()
    return out_wGG

def rpa_density_response(self, direction='x', qinf_v=None):
    """Calculate the RPA susceptibility for (semi-)finite q.

    Currently this is only used by the QEH code, why we don't support a top
    level user interface.

```

```

"""
# Extract  $\chi_0(q, \omega)$ 
qpd = self.chi0.qpd
chi0_wGG = self.get_chi0_wGG(direction=direction)
if qpd.optical_limit:
    # Restore the q-dependence of the head and wings in the  $q \rightarrow 0$  limit
    assert qinf_v is not None and np.linalg.norm(qinf_v) > 0.
    d_v = self._normalize(direction)
    chi0_wGG[:, 1:, 0] *= np.dot(qinf_v, d_v)
    chi0_wGG[:, 0, 1:] *= np.dot(qinf_v, d_v)
    chi0_wGG[:, 0, 0] *= np.dot(qinf_v, d_v)**2
# Invert Dyson equation,  $\chi(q, \omega) = [1 - \chi_0(q, \omega) V(q)]^{-1} \chi_0(q, \omega)$ 
V_GG = self.coulomb.kernel(qpd, q_v=qinf_v)
chi_wGG = self.invert_dyson_like_equation(chi0_wGG, V_GG)
return qpd, chi_wGG, self.wblocks

def inverse_dielectric_function(self, direction='x'):
    """Calculate  $v^{(1/2)} \chi v^{(1/2)}$ , from which  $\epsilon^{-1}(q, \omega)$  is constructed."""
    return InverseDielectricFunction.from_chi0_dyson_eqs(
        self, *self.calculate_vchi_symm(direction=direction))

def calculate_vchi_symm(self, direction='x', modified=False):
    """Calculate  $v^{(1/2)} \chi v^{(1/2)}$ .

    Starting from the TDDFT Dyson equation


$$\chi(q, \omega) = \chi_0(q, \omega) + \chi_0(q, \omega) K_{Hxc}(q, \omega) \chi(q, \omega), \quad (1)$$


the Coulomb scaled susceptibility,


$$\tilde{\chi}(q, \omega) = v^{(1/2)}(q) \chi(q, \omega) v^{(1/2)}(q)$$


can be calculated from the Dyson-like equation


$$\tilde{\chi}(q, \omega) = \tilde{\chi}_0(q, \omega) + \tilde{\chi}_0(q, \omega) \tilde{K}(q, \omega) \tilde{\chi}(q, \omega) \quad (2)$$


where


$$\tilde{K}(q, \omega) = v^{(-1/2)}(q) K_{Hxc}(q, \omega) v^{(-1/2)}(q).$$


Here  $v(q)$  refers to the bare Coulomb potential. It should be emphasized
that inversion of (2) rather than (1) is not merely a rescaling
exercise. In the optical  $q \rightarrow 0$  limit, the Coulomb kernel  $v(q)$  diverges
as  $1/|q|^2$  while the Kohn-Sham susceptibility  $\chi_0(q, \omega)$  vanishes as
 $|q|^2$ . Treating  $v^{(1/2)}(q) \chi_0(q, \omega) v^{(1/2)}(q)$  as a single variable,
the effects of this cancellation can be treated accurately within k.p
perturbation theory.
"""
chi0_wGG = self.get_chi0_wGG(direction=direction)
Kxc_GG = self.xc_kernel(self.chi0.qpd, chi0_wGG=chi0_wGG) \
    if self.xc_kernel else None
v_G, K_GG = self.get_coulomb_scaled_kernel(
    modified=modified, Kxc_GG=Kxc_GG)
# Calculate  $v^{(1/2)}(q) \chi_0(q, \omega) v^{(1/2)}(q)$ 
sqrtv_G = v_G**0.5
vchi0_symm_wGG = chi0_wGG # reuse buffer
for w, chi0_GG in enumerate(chi0_wGG):
    vchi0_symm_wGG[w] = chi0_GG * sqrtv_G * sqrtv_G[:, np.newaxis]
# Invert Dyson equation
vchi_symm_wGG = self.invert_dyson_like_equation(
    vchi0_symm_wGG, K_GG, reuse_buffer=False)
return vchi0_symm_wGG, vchi_symm_wGG

def customized_dielectric_function(self, direction='x'):
    """Calculate  $E(q, \omega) = 1 - V(q) P(q, \omega)$ ."""
    V_GG = self.coulomb.kernel(self.chi0.qpd)
    P_wGG = self.polarizability_operator(direction=direction)
    nG = len(V_GG)
    eps_wGG = P_wGG # reuse buffer

```

```

for w, P_GG in enumerate(P_wGG):
    eps_wGG[w] = np.eye(nG) - V_GG @ P_GG
return CustomizableDielectricFunction.from_chi0_dyson_eqs(
    self, eps_wGG)

def bare_dielectric_function(self, direction='x'):
    """Calculate  $v^{(1/2)} \bar{\chi} v^{(1/2)}$ , from which  $\bar{\epsilon}=1-v \bar{\chi}$  is constructed.

    The unscreened susceptibility is given by the Dyson-like equation


$$\bar{\chi}(q,\omega) = P(q,\omega) + P(q,\omega) \bar{V}(q) \bar{\chi}(q,\omega), \quad (3)$$


    In the special case of RPA, where  $P(q,\omega) = \chi_0(q,\omega)$ , one may notice that
    the Dyson-like equation (3) is exactly identical to the TDDFT Dyson
    equation (1) when replacing the Hartree-exchange-correlation kernel
    with the modified Coulomb interaction:


$$K_{Hxc}(q) \rightarrow \bar{V}(q).$$


    We may thus reuse that functionality to calculate  $v^{(1/2)} \bar{\chi} v^{(1/2)}$ .
    """
    if self.xc_kernel:
        raise NotImplementedError(
            'Calculation of the bare dielectric function within TDDFT has '
            'not yet been implemented. For TDDFT dielectric properties, '
            'please calculate the inverse dielectric function.')
    vP_symm_wGG, vchibar_symm_wGG = self.calculate_vchi_symm(
        direction=direction, modified=True)
    return BareDielectricFunction.from_chi0_dyson_eqs(
        self, vP_symm_wGG, vchibar_symm_wGG)

def polarizability_operator(self, direction='x'):
    """Calculate the polarizability operator  $P(q,\omega)$ .

    Depending on the theory (RPA, TDDFT, MBPT etc.), the polarizability
    operator is approximated in various ways see e.g.
    [Rev. Mod. Phys. 74, 601 (2002)].

    In RPA:

$$P(q,\omega) = \chi_0(q,\omega)$$


    In TDDFT:

$$P(q,\omega) = [1 - \chi_0(q,\omega) K_{xc}(q,\omega)]^{-1} \chi_0(q,\omega)$$

    """
    chi0_wGG = self.get_chi0_wGG(direction=direction)
    if not self.xc_kernel: # RPA
        return chi0_wGG
    # TDDFT (in adiabatic approximations to the kernel)
    if self.chi0.qpd.optical_limit:
        raise NotImplementedError(
            'Calculation of the TDDFT dielectric function via the '
            'polarizability operator has not been implemented for the '
            'optical limit. Please calculate the inverse dielectric '
            'function instead.')
    # The TDDFT implementation here is invalid in the optical limit
    # since the chi0_wGG variable already contains Coulomb effects,
    #  $\chi_0_wGG \sim V^{(1/2)} \chi_0 V^{(1/2)} / (4\pi)$ .
    # Furthermore, a direct evaluation of  $V(q) P(q,\omega)$  does not seem
    # sensible, since it does not account for the exact cancellation
    # of the q-dependences of the two functions.
    # In principle, one could treat the  $v(q) P(q,\omega)$  product in
    # perturbation theory, similar to the  $\chi_0(q,\omega) v(q)$  product in the
    # Dyson equation for  $\chi$ , but unless we need to calculate the TDDFT
    # polarizability using truncated kernels, this isn't really
    # necessary.
    Kxc_GG = self.xc_kernel(self.chi0.qpd, chi0_wGG=chi0_wGG)
    return self.invert_dyson_like_equation(chi0_wGG, Kxc_GG)

```

```

@dataclass
class DielectricFunctionData(ABC):
    cd: CellDescriptor
    qpd: SingleQPWDescriptor
    wd: FrequencyDescriptor
    wblocks: Blocks1D
    coulomb: CoulombKernel
    bare_coulomb: CoulombKernel

    @classmethod
    def from_chi0_dyson_eqs(cls, chi0_dyson_eqs, *args, **kwargs):
        chi0 = chi0_dyson_eqs.chi0
        return cls(chi0_dyson_eqs.cd, chi0.qpd,
                   chi0.wd, chi0_dyson_eqs.wblocks,
                   chi0_dyson_eqs.coulomb, chi0_dyson_eqs.bare_coulomb,
                   *args, **kwargs)

    def _macroscopic_component(self, in_wGG):
        return self.wblocks.all_gather(in_wGG[:, 0, 0])

    @property
    def v_G(self):
        return self.bare_coulomb.V(self.qpd)

    @abstractmethod
    def macroscopic_dielectric_function(self) -> ScalarResponseFunctionSet:
        """Get the macroscopic dielectric function  $\epsilon_M(q, \omega)$ ."""

    def dielectric_constant(self):
        return self.macroscopic_dielectric_function().static_limit.real

    def eels_spectrum(self):
        """Get the macroscopic EELS spectrum.

        The spectrum is defined as


$$\text{EELS}(q, \omega) \equiv -\text{Im} \frac{1}{\epsilon_M^{-1}(q, \omega)} = -\text{Im} \frac{1}{\epsilon_M(q, \omega)}.$$


        In addition to the many-body spectrum, we also calculate the
        EELS spectrum in the relevant independent-particle approximation,
        here defined as


$$\text{EELS}_0(\omega) = -\text{Im} \frac{1}{\epsilon_{\text{IP}}^M(q, \omega)}.$$


        """
        _, eps0_W, eps_W = self.macroscopic_dielectric_function().arrays
        eels0_W = -(1. / eps0_W).imag
        eels_W = -(1. / eps_W).imag
        return ScalarResponseFunctionSet(self.wd, eels0_W, eels_W)

    def polarizability(self):
        """Get the macroscopic polarizability  $\alpha_M(q, \omega)$ .

        In the special case where dielectric properties are calculated
        solely based on the bare Coulomb interaction  $V(q) = v(q)$ , the
        macroscopic part of the electronic polarizability  $\alpha(q, \omega)$  is related
        directly to the macroscopic dielectric function  $\epsilon_M(q, \omega)$ .

        In particular, we define the macroscopic polarizability so as to
        make it independent of the nonperiodic cell vector lengths. Namely,


$$\alpha_M(q, \omega) \equiv \Lambda \alpha(q, \omega)$$


```

where Λ is the nonperiodic hypervolume of the unit cell. Thus,

$$\alpha_M(q, \omega) = \Lambda / (4\pi) (\epsilon_M(q, \omega) - 1) = \Lambda / (4\pi) (\epsilon_M(q, \omega) - 1),$$

where the latter equality only holds in the special case $V(q) = v(q)$.

In addition to $\alpha_M(q, \omega)$, we calculate also the macroscopic polarizability in the relevant independent-particle approximation by replacing ϵ_M with $\epsilon_M^{(IP)}$.

```
"""
if self.coulomb is not self.bare_coulomb:
    raise ValueError(
        'When using a truncated Coulomb kernel, the polarizability '
        'cannot be calculated based on the macroscopic dielectric '
        'function. Please calculate the bare dielectric function '
        'instead.')
_, eps0_W, eps_W = self.macroscopic_dielectric_function().arrays
return self._polarizability(eps0_W, eps_W)

def _polarizability(self, eps0_W, eps_W):
    L = self.cd.nonperiodic_hypervolume
    alpha0_W = L / (4 * np.pi) * (eps0_W - 1)
    alpha_W = L / (4 * np.pi) * (eps_W - 1)
    return ScalarResponseFunctionSet(self.wd, alpha0_W, alpha_W)
```

@dataclass

```
class InverseDielectricFunction(DielectricFunctionData):
    """Data class for the inverse dielectric function  $\epsilon^{-1}(q, \omega)$ .
```

The inverse dielectric function characterizes the longitudinal response

$$V_{\text{tot}}(q, \omega) = \epsilon^{-1}(q, \omega) V_{\text{ext}}(q, \omega),$$

where the induced potential due to the electronic system is given by $v\chi$,

$$\epsilon^{-1}(q, \omega) = 1 + v(q) \chi(q, \omega).$$

In this data class, ϵ^{-1} is cast in terms of its symmetrized representation

$$\tilde{\epsilon}^{-1}(q, \omega) = v^{(-1/2)}(q) \epsilon^{-1}(q, \omega) v^{(1/2)}(q),$$

that is, in terms of $v^{(1/2)}(q) \chi(q, \omega) v^{(1/2)}(q)$.

Please remark that $v(q)$ here refers to the bare Coulomb potential regardless of whether $\chi(q, \omega)$ was determined using a truncated analogue.

```
"""
vchi0_symm_wGG: np.ndarray # v^{(1/2)}(q) \chi_0(q, \omega) v^{(1/2)}(q)
vchi_symm_wGG: np.ndarray
```

```
def macroscopic_components(self):
    vchi0_W = self._macroscopic_component(self.vchi0_symm_wGG)
    vchi_W = self._macroscopic_component(self.vchi_symm_wGG)
    return vchi0_W, vchi_W
```

```
def macroscopic_dielectric_function(self):
    """Get the macroscopic dielectric function  $\epsilon_M(q, \omega)$ .
```

Calculates

$$\epsilon_M(q, \omega) = \frac{1}{\epsilon_{00}^{-1}(q, \omega)}$$

along with the macroscopic dielectric function in the independent-particle random-phase approximation [Rev. Mod. Phys. 74, 601 (2002)],

$$\epsilon_M(q, \omega) = 1 - \frac{v(q)}{\omega^2} \chi^0(q, \omega)$$

that is, neglecting local-field and exchange-correlation effects.

```

vchi0_W, vchi_W = self.macroscopic_components()
eps0_W = 1 - vchi0_W
eps_W = 1 / (1 + vchi_W)
return ScalarResponseFunctionSet(self.wd, eps0_W, eps_W)

```

```

def dynamic_susceptibility(self):
    """Get the macroscopic components of  $\chi(q, \omega)$  and  $\chi^0(q, \omega)$ ."""
    vchi0_W, vchi_W = self.macroscopic_components()
    v0 = self.v_G[0] # Macroscopic Coulomb potential ( $4\pi/q^2$ )
    return ScalarResponseFunctionSet(self.wd, vchi0_W / v0, vchi_W / v0)

```

@dataclass

```

class CustomizableDielectricFunction(DielectricFunctionData):
    """Data class for customized dielectric functions  $E(q, \omega)$ .

```

$E(q, \omega)$ is customizable in the sense that bare Coulomb interaction $v(q)$ is replaced by an arbitrary interaction $V(q)$ in the formula for $\epsilon(q, \omega)$,

$$E(q, \omega) = 1 - V(q) P(q, \omega),$$

where P is the polarizability operator [Rev. Mod. Phys. 74, 601 (2002)]. Thus, for any truncated or otherwise customized interaction $V(q) \neq v(q)$, $E(q, \omega) \neq \epsilon(q, \omega)$ and $E^{-1}(q, \omega) \neq \epsilon^{-1}(q, \omega)$.

eps_wGG: np.ndarray

```

def macroscopic_customized_dielectric_function(self):
    """Get the macroscopic customized dielectric function  $E_M(q, \omega)$ .

    We define the macroscopic customized dielectric function as

```

$$E_M(q, \omega) = \frac{1}{E^{-1}(q, \omega)},$$

and calculate also the macroscopic dielectric function in the customizable independent-particle approximation:

$$\epsilon_M^{\text{cIP}}(q, \omega) = 1 - \frac{V(q)}{\omega^2} = \frac{E(q, \omega)}{\omega^2}.$$

```

eps0_W = self._macroscopic_component(self.eps_wGG)
# Invert  $E(q, \omega)$  one frequency at a time to compute  $E_M(q, \omega)$ 
eps_w = np.zeros((self.wblocks.nlocal,), complex)
for w, eps_GG in enumerate(self.eps_wGG):
    eps_w[w] = 1 / np.linalg.inv(eps_GG)[0, 0]
eps_W = self.wblocks.all_gather(eps_w)
return ScalarResponseFunctionSet(self.wd, eps0_W, eps_W)

```

```

def macroscopic_dielectric_function(self):
    """Get the macroscopic dielectric function  $\epsilon_M(q, \omega)$ .

```

In the special case $V(q) = v(q)$, $E_M(q, \omega) = \epsilon_M(q, \omega)$.

```

if self.coulomb is not self.bare_coulomb:
    raise ValueError(
        'The macroscopic dielectric function is defined in terms of '
        'the bare Coulomb interaction. To truncate the Hartree '
        'electron-electron correlations, please calculate the inverse '

```

```

        'dielectric function instead.')
    return self.macroscopic_customized_dielectric_function()

```

```
@dataclass
```

```
class BareDielectricFunction(DielectricFunctionData):
```

```
    """Data class for the bare (unscreened) dielectric function.
```

```

    The bare dielectric function is defined in terms of the unscreened
    susceptibility,

```

$$\bar{\epsilon}(q, \omega) = 1 - v(q) \bar{\chi}(q, \omega),$$

```

    here represented in terms of  $v^{(1/2)}(q) \bar{\chi}(q, \omega) v^{(1/2)}(q)$ .
    """

```

```

    vP_symm_wGG: np.ndarray #  $v^{(1/2)} P v^{(1/2)}$ 
    vchibar_symm_wGG: np.ndarray

```

```
def macroscopic_components(self):
```

```
    vP_W = self._macroscopic_component(self.vP_symm_wGG)
```

```
    vchibar_W = self._macroscopic_component(self.vchibar_symm_wGG)
```

```
    return vP_W, vchibar_W
```

```
def macroscopic_bare_dielectric_function(self):
```

```
    """Get the macroscopic bare dielectric function  $\epsilon_M(q, \omega)$ .
```

```

    Calculates

```

$$\epsilon_M(q, \omega) = \bar{\epsilon}(q, \omega)$$

```

    along with the macroscopic dielectric function in the independent-
    particle approximation:

```

$$\epsilon_M(q, \omega) = 1 - vP(q, \omega).$$

```

    """
    vP_W, vchibar_W = self.macroscopic_components()

```

```
    eps0_W = 1. - vP_W
```

```
    eps_W = 1. - vchibar_W
```

```
    return ScalarResponseFunctionSet(self.wd, eps0_W, eps_W)
```

```
def macroscopic_dielectric_function(self):
```

```
    """Get the macroscopic dielectric function  $\epsilon_M(q, \omega)$ .
```

```

    In the special case where the unscreened susceptibility is calculated
    using the bare Coulomb interaction,

```

$$\bar{\chi}(q, \omega) = P(q, \omega) + P(q, \omega) \bar{v}(q) \bar{\chi}(q, \omega),$$

```

    it holds identically that [Rev. Mod. Phys. 74, 601 (2002)]:

```

$$\epsilon_M(q, \omega) = \epsilon_M(q, \omega).$$

```

    """
    if self.coulomb is not self.bare_coulomb:

```

```
        raise ValueError(
```

```
            'The macroscopic dielectric function cannot be obtained from '
```

```
            'the bare dielectric function calculated based on a truncated '
```

```
            'Coulomb interaction. Please calculate the inverse dielectric '
```

```
            'function instead')
```

```
    return self.macroscopic_bare_dielectric_function()
```

```
def polarizability(self):
```

```
    """Get the macroscopic polarizability  $\alpha_M(q, \omega)$ .
```

```

    In the most general case, the electronic polarizability can be defined
    in terms of the unscreened susceptibility

```

-

$$\alpha(q, \omega) \equiv -1/(4\pi) \ v(q) \ \chi(q, \omega) = 1/(4\pi) \ (\epsilon(q, \omega) - 1).$$

This is especially useful in systems of reduced dimensionality, where $\chi(q, \omega)$ needs to be calculated using a truncated Coulomb kernel in order to achieve convergence in a feasible way.

```
_, eps0_W, eps_W = self.macroscopic_bare_dielectric_function().arrays
return self._polarizability(eps0_W, eps_W)
```

```
class DielectricFunctionCalculator:
```

```
    def __init__(self, chi0calc: Chi0Calculator):
```

```
        self.chi0calc = chi0calc
```

```
        self.gs = chi0calc.gs
```

```
        self.context = chi0calc.context
```

```
        self._chi0cache: dict[tuple[str, ...], Chi0Data] = {}
```

```
    def get_chi0(self, q_c: list | np.ndarray) -> Chi0Data:
```

```
        """Get the Kohn-Sham susceptibility  $\chi_0(q, \omega)$  for input wave vector q.
```

```
        Keeps a cache of  $\chi_0$  for the latest calculated wave vector, thus
        allowing for investigation of multiple dielectric properties,
        Coulomb truncations, xc kernels etc. without recalculating  $\chi_0$ .
        """
```

```
        # As cache key, we round off and use a string representation.
```

```
        # Not very elegant, but it should work almost always.
```

```
        q_key = [f'{q:.10f}' for q in q_c]
```

```
        key = tuple(q_key)
```

```
        if key not in self._chi0cache:
```

```
            self._chi0cache.clear()
```

```
            self._chi0cache[key] = self.chi0calc.calculate(q_c)
```

```
            self.context.write_timer()
```

```
        return self._chi0cache[key]
```

```
    def get_chi0_dyson_eqs(self,
```

```
        q_c: list | np.ndarray = [0, 0, 0],
```

```
        truncation: str | None = None,
```

```
        xc: str = 'RPA',
```

```
        **kwargs
```

```
    ) -> Chi0DysonEquations:
```

```
        """Set up the Dyson equation for  $\chi(q, \omega)$  at given wave vector q.
```

```
        Parameters
```

```
        -----
```

```
        truncation : str or None
```

```
            Truncation of the Hartree kernel.
```

```
        xc : str
```

```
            Exchange-correlation kernel for LR-TDDFT calculations.
```

```
            If xc == 'RPA', the dielectric response is treated in the random
            phase approximation.
```

```
        **kwargs
```

```
            Additional parameters for the chosen xc kernel.
```

```
        """
```

```
        chi0 = self.get_chi0(q_c)
```

```
        coulomb = CoulombKernel.from_gs(self.gs, truncation=truncation)
```

```
        if xc == 'RPA':
```

```
            xc_kernel = None
```

```
        else:
```

```
            xc_kernel = DensityXCKernel.from_functional(
```

```
                self.gs, self.context, functional=xc, **kwargs)
```

```
        return Chi0DysonEquations(chi0, coulomb, xc_kernel, self.gs.cd)
```

```
    def get_bare_dielectric_function(self, direction='x', **kwargs
```

```
    ) -> BareDielectricFunction:
```

```
        """Calculate the bare dielectric function  $\bar{\epsilon}(q, \omega) = 1 - v(q) \bar{\chi}(q, \omega)$ .
```

```
        Here v(q) is the bare Coulomb potential while  $\bar{\chi}$  is the unscreened
        susceptibility calculated based on the modified (and possibly
```

```

truncated) Coulomb potential.
"""
return self.get_chi0_dyson_eqs(
    **kwargs).bare_dielectric_function(direction=direction)

def get_literal_dielectric_function(self, direction='x', **kwargs
    ) -> CustomizableDielectricFunction:
    """Calculate the dielectric function  $\epsilon(q,\omega) = 1 - v(q) P(q,\omega)$ ."""
    return self.get_chi0_dyson_eqs(
        truncation=None, **kwargs).customized_dielectric_function(
            direction=direction)

def get_customized_dielectric_function(self, direction='x', *,
    truncation: str, **kwargs
    ) -> CustomizableDielectricFunction:
    """Calculate the customized dielectric function  $E(q,\omega) = 1 - V(q)P(q,\omega)$ .

    In comparison with the literal dielectric function, the bare Coulomb
    interaction has here been replaced with a truncated analogue  $v(q) \rightarrow V(q)$ .
    """
    return self.get_chi0_dyson_eqs(
        truncation=truncation, **kwargs).customized_dielectric_function(
            direction=direction)

def get_inverse_dielectric_function(self, direction='x', **kwargs
    ) -> InverseDielectricFunction:
    """Calculate the inverse dielectric function  $\epsilon^{-1}(q,\omega) = v(q) \chi(q,\omega)$ .
    """
    return self.get_chi0_dyson_eqs(
        **kwargs).inverse_dielectric_function(direction=direction)

class DielectricFunction(DielectricFunctionCalculator):
    """This class defines dielectric function related physical quantities."""

    def __init__(self, calc, *,
        frequencies=None,
        ecut=50,
        hilbert=True,
        nbands=None, eta=0.2,
        intraband=True, nblocks=1, world=mpi.world, txt=sys.stdout,
        truncation=None,
        qsymmetry=True,
        integrationmode=None, rate=0.0,
        eshift: float | None = None):
        """Creates a DielectricFunction object.

        calc: str
            The ground-state calculation file that the linear response
            calculation is based on.
        frequencies:
            Input parameters for frequency_grid.
            Can be an array of frequencies to evaluate the response function at
            or dictionary of parameters for build-in nonlinear grid
            (see :ref:`frequency grid`).
        ecut: float
            Plane-wave cut-off.
        hilbert: bool
            Use hilbert transform.
        nbands: int
            Number of bands from calculation.
        eta: float
            Broadening parameter.
        intraband: bool
            Include intraband transitions.
        world: comm
            mpi communicator.
        nblocks: int
            Split matrices in nblocks blocks and distribute them G-vectors or

```

```

        frequencies over processes.
    txt: str
        Output file.
    truncation: str or None
        None for no truncation.
        '2D' for standard analytical truncation scheme.
        Non-periodic directions are determined from k-point grid
    eshift: float
        Shift unoccupied bands
    """
    gs, context = get_gs_and_context(calc, txt, world, timer=None)
    wd = get_frequency_descriptor(frequencies, gs=gs, nbands=nbands)

    chi0calc = Chi0Calculator(
        gs, context, nblocks=nblocks,
        wd=wd,
        ecut=ecut, nbands=nbands, eta=eta,
        hilbert=hilbert,
        intraband=intraband,
        qsymmetry=qsymmetry,
        integrationmode=integrationmode,
        rate=rate, eshift=eshift
    )
    super().__init__(chi0calc)
    self.truncation = truncation

def get_frequencies(self) -> np.ndarray:
    """Return frequencies (in eV) that the  $\chi$  is evaluated on."""
    return self.chi0calc.wd.omega_w * Hartree

def get_dynamic_susceptibility(self, *args, xc='ALDA',
                               filename='chiM_w.csv',
                               **kwargs):
    dynsus = self.get_inverse_dielectric_function(
        *args, xc=xc, truncation=self.truncation,
        **kwargs).dynamic_susceptibility()
    if filename:
        dynsus.write(filename)
    return dynsus.unpack()

def get_dielectric_function(self, *args, filename='df.csv', **kwargs):
    """Calculate the dielectric function.

    Generates a file 'df.csv', unless filename is set to None.

    Returns
    -----
    df_NLFC_w: np.ndarray
        Dielectric function without local field corrections.
    df_LFC_w: np.ndarray
        Dielectric function with local field corrections.
    """
    df = self.get_inverse_dielectric_function(
        *args, truncation=self.truncation,
        **kwargs).macroscopic_dielectric_function()
    if filename:
        df.write(filename)
    return df.unpack()

def get_eels_spectrum(self, *args, filename='eels.csv', **kwargs):
    """Calculate the macroscopic EELS spectrum.

    Generates a file 'eels.csv', unless filename is set to None.

    Returns
    -----
    eels0_w: np.ndarray
        Spectrum in the independent-particle random-phase approximation.
    eels_w: np.ndarray

```

```

        Fully screened EELS spectrum.
    """
    eels = self.get_inverse_dielectric_function(
        *args, truncation=self.truncation, **kwargs).eels_spectrum()
    if filename:
        eels.write(filename)
    return eels.unpack()

def get_polarizability(self, q_c: list | np.ndarray = [0, 0, 0],
                      direction='x', filename='polarizability.csv',
                      **kwargs):
    """Calculate the macroscopic polarizability.

    Generate a file 'polarizability.csv', unless filename is set to None.

    Returns:
    -----
    alpha0_w: np.ndarray
        Polarizability calculated without local-field corrections
    alpha_w: np.ndarray
        Polarizability calculated with local-field corrections.
    """
    chi0_dyson_eqs = self.get_chi0_dyson_eqs(
        q_c, truncation=self.truncation, **kwargs)
    if self.truncation:
        # eps: BareDielectricFunction
        method = chi0_dyson_eqs.bare_dielectric_function
    else:
        # eps: CustomizableDielectricFunction
        method = chi0_dyson_eqs.customized_dielectric_function
    eps = method(direction=direction)
    pol = eps.polarizability()
    if filename:
        pol.write(filename)
    return pol.unpack()

def get_macroscopic_dielectric_constant(self, xc='RPA', direction='x'):
    """Calculate the macroscopic dielectric constant.

    The macroscopic dielectric constant is defined as the real part of the
    dielectric function in the static limit.

    Returns:
    -----
    eps0: float
        Dielectric constant without local field corrections.
    eps: float
        Dielectric constant with local field correction. (RPA, ALDA)
    """
    return self.get_inverse_dielectric_function(
        xc=xc, direction=direction).dielectric_constant()

# ----- Serialized dataclasses and IO ----- #

@dataclass
class ScalarResponseFunctionSet:
    """A set of scalar response functions  $\text{rf}_0(\omega)$  and  $\text{rf}(\omega)$ ."""
    wd: FrequencyDescriptor
    rf0_w: np.ndarray
    rf_w: np.ndarray

    @property
    def arrays(self):
        return self.wd.omega_w * Hartree, self.rf0_w, self.rf_w

    def unpack(self):
        # Legacy feature to support old DielectricFunction output format

```

```

# ... to be deprecated ...
return self.rf0_w, self.rf_w

def write(self, filename):
    if mpi.rank == 0:
        write_response_function(filename, *self.arrays)

@property
def static_limit(self):
    """Return the value of the response functions in the static limit."""
    w0 = np.argmin(np.abs(self.wd.omega_w))
    assert abs(self.wd.omega_w[w0]) < 1e-8
    return np.array([self.rf0_w[w0], self.rf_w[w0]])

def write_response_function(filename, omega_w, rf0_w, rf_w):
    with open(filename, 'w') as fd:
        for omega, rf0, rf in zip(omega_w, rf0_w, rf_w):
            if rf0_w.dtype == complex:
                print('%0.6f, %0.6f, %0.6f, %0.6f, %0.6f' %
                      (omega, rf0.real, rf0.imag, rf.real, rf.imag),
                      file=fd)
            else:
                print(f'{omega:.6f}, {rf0:.6f}, {rf:.6f}', file=fd)

def read_response_function(filename):
    """Read a stored response function file"""
    d = np.loadtxt(filename, delimiter=',')
    omega_w = np.array(d[:, 0], float)

    if d.shape[1] == 3:
        # Real response function
        rf0_w = np.array(d[:, 1], float)
        rf_w = np.array(d[:, 2], float)
    elif d.shape[1] == 5:
        rf0_w = np.array(d[:, 1], complex)
        rf0_w.imag = d[:, 2]
        rf_w = np.array(d[:, 3], complex)
        rf_w.imag = d[:, 4]
    else:
        raise ValueError(f'Unexpected array dimension {d.shape}')

    return omega_w, rf0_w, rf_w

```