

# PAW Notes

ffr

$$V_{\text{NL}} = \sum_{nm,I} D_{nm}^{(0)} |\beta_n^I\rangle \langle \beta_m^I|$$

$n, m$ : beta func indices

$$\beta_n^I(\mathbf{r}) = \beta_n(\mathbf{r} - \mathbf{R}_I)$$

Electron density:

$$n(\mathbf{r}) = \sum_i \left[ |\psi_i(\mathbf{r})|^2 + \sum_{nm,I} Q_{nm}^I(\mathbf{r}) \langle \psi_i | \beta_n \rangle \langle \beta_m | \psi_i \rangle \right]$$

Orthonormality conditions:

$$\langle \psi_i | S | \psi_j \rangle = \delta_{ij}$$

overlap operator:

$$S = 1 + \sum_{nm,I} q_{nm} |\beta_n^I\rangle \langle \beta_m^I|$$

$$q_{nm} = \int Q_{nm}(\mathbf{r}) \, d\mathbf{r}$$

KS equation:

$$H |\psi_i\rangle = \epsilon_i S |\psi_i\rangle$$

Hamiltonian operator:

$$H = -\frac{\hbar^2}{2m} \nabla^2 + V_{\text{eff}} + \sum_{nm,I} D_{nm}^I |\beta_n^I\rangle \langle \beta_m^I|$$

Screened coefficients:

$$D_{nm}^I = D_{nm}^{(0)} + \int V_{\text{eff}}(\mathbf{r}) Q_{nm}^I(\mathbf{r}) d\mathbf{r}$$

$$E = \sum_n f_n \left\langle \psi_n \left| -\frac{1}{2} \nabla^2 \right| \psi_n \right\rangle + E_H [n + n_Z] + E_{xc} [n]$$

$$|\psi_n\rangle = |\tilde{\psi}_n\rangle + \sum_i \left( |\phi_i\rangle - |\tilde{\phi}_i\rangle \right) \langle \tilde{p}_i | \tilde{\psi}_n \rangle$$

Index  $i$ : atomic site  $\mathbf{R}$ , angular momentum  $L = (l, m)$ , and additional index  $k$  for reference energy  $\epsilon_{kl}$ .

AE partial waves  $\phi_i$  are obtained for a reference atom

PS partial waves  $\tilde{\phi}_i$  are equivalent to the AE partial waves outside a core radius  $r_c^l$  and match continuously onto  $\phi_i$  inside the core radius.

Projector functions  $\tilde{p}_i$  are dual to the partial waves:

$$\langle \tilde{p}_i | \tilde{\phi}_j \rangle = \delta_{ij}$$

$$n(\mathbf{r}) = \tilde{n}(\mathbf{r}) + n^1(\mathbf{r}) - \tilde{n}^1(\mathbf{r})$$

$$\tilde{n}(\mathbf{r}) = \sum_n f_n \langle \tilde{\psi}_n | \mathbf{r} \rangle \langle \mathbf{r} | \tilde{\psi}_n \rangle$$

Onsite electron densities  $n^1$  and  $\tilde{n}^1$  are treated on a radial support grid, that extends up to  $r_{\text{rad}}$  around each ion.

$$n^1(\mathbf{r}) = \sum_{i,j} \rho_{ij} \langle \phi_i | \mathbf{r} \rangle \langle \mathbf{r} | \phi_j \rangle$$

$$\tilde{n}^1(\mathbf{r}) = \sum_{i,j} \rho_{ij} \langle \tilde{\phi}_i | \mathbf{r} \rangle \langle \mathbf{r} | \tilde{\phi}_j \rangle$$

$\rho_{ij}$  are the occupancies of each augmentation channel  $(i, j)$ :

$$\rho_{ij} = \sum_n f_n \langle \tilde{\psi}_n | \tilde{p}_i \rangle \langle \tilde{p}_i | \tilde{\psi}_n \rangle$$

Assuming frozen core approximation.

Introduce four quantities that will be used for the description of the core charge density:

$n_c$ : charge density of frozen core all-electron wave functions in the reference atom.

$\tilde{n}_c$ : partial electronic core density  $\tilde{n}_c$  is equivalent to the frozen core AE charge density outside a certain radius  $r_{pc}$ , which lies inside the argumentation radius. Used in order to calculate nonlinear core corrections

$n_{Zc}$ : point charge density of the nuclei  $n_Z$  plus the frozen core AE charge density  $n_c$ :

$$n_{Zc} = n_Z + n_c$$

$\tilde{n}_{Zc}$ : pseudized core density, equivalent to  $n_{Zc}$  outside the core radius and shall have the same moment as  $n_{Zc}$  inside the core region:

$$\int_{\Omega_r} n_{Zc}(\mathbf{r}) \, d\mathbf{r} = \int_{\Omega_r} \tilde{n}_{Zc}(\mathbf{r}) \, d\mathbf{r}$$