# MPS-based quantum impurity solvers DMFT + DMRG

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Chan Group, Princeton University, May 13 2015







### Overview: quantum embedding techniques

#### DMFT, CDMFT, DCA

- CTQMC Gull, Millis, Lichtenstein, Rubtsov, Troyer & Werner, RMP 83, 349 (2011)
- NRG Bulla, Costi & Pruschke, RMP 80, 395 (2008)
- ED Caffarel & Krauth, PRL 72, 1545 (1994) / Granath & Strand, PRB 86, 115111 (2012) / Lu, Höppner,
   Gunnarsson & Haverkort, PRB 90, 085102 (2014)
- Truncated CI Zgid, Gull & Chan, PRB 86, 165128 (2012)
- DMRG García, Hallberg & Rozenberg, PRL 93, 246403 (2004)

#### **DMET**

- Knizia & Chan, PRL 109, 186404 (2012)
- dynamic formulation Booth & Chan, arXiv:1309.2320 (2013)

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Chebyshev and Fourier expansions: cheaper and precise

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FAW, McCulloch, Parcollet & Schollwöck, PRB 90, 115124 (2014a)

> two-site cluster!

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Balzer, FAW, McCulloch, Werner & Eckstein, arXiv:1504.02461 (2015)

#### Outline

- Overview: impurity solvers and quantum embedding
- Compute Green's / spectral functions using MPS: algorithms, cost, error control, computability
- DMFT: bath discretization, bath geometry and long range Hamiltonian
- Benchmark: two-site DCA in different setups

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Let  $|\psi_0\rangle = a^{\dagger}|E_0\rangle$  be a single-particle excitation:

$$A(\omega) = \langle \psi_0 | \delta(\omega - (H - E_0)) | \psi_0 \rangle$$

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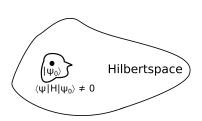
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- $|\psi_0\rangle$  is lowly entangled.
- Need to compute neighborhood of  $|\psi_0\rangle$

$$\{ |\psi\rangle \mid \langle \psi|H|\psi_0\rangle \neq 0 \}$$

 Hope that there is a basis for this neighborhood that is not too strongly entangled!



• How to construct this basis?

• DDMRG: solve  $G_{\eta}(\omega) = \langle \psi_0 | \underbrace{\frac{1}{\omega + i\eta - (H - E_0)} | \psi_0 \rangle}_{= | \mathsf{cv}_{\eta}(\omega) \rangle}$ 

Jeckelmann, PRB 66, 045114 (2002)

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Lanczos: represent H in orthog. Krylov basis  $[1, H, H^2, \dots]|\psi_0\rangle$ 

$$|\psi_{n+1}\rangle = (H - \alpha_n)|\psi_n\rangle - \beta_n|\psi_{n-1}\rangle$$

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• Chebyshev expansion:  $A(\omega) \sim \sum_n \langle \psi_0 | \psi_n \rangle \cos(n \arccos(\omega))$ 

$$|\psi_n\rangle = \cos(n \arccos H')|\psi_0\rangle$$

Weiße, Wellein, Alvermann & Fehske, RMP 78, 275 (2006)

MPS: Holzner, Weichselbaum, McCulloch, Schollwöck & von Delft, PRB 83, 195115 (2011)

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• Fourier expansion:  $A(\omega) \sim \sum_{n} \langle \psi_0 | \psi_n \rangle e^{i\omega t_n}$ 

$$|\psi_n\rangle = e^{-i(H-E_0)t_n}|\psi_0\rangle$$

FAW, Justiniano, McCulloch & Schollwöck, PRB 91, 115144 (2015)

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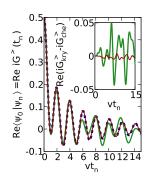
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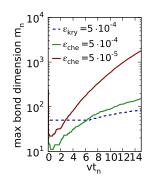
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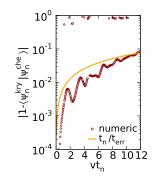
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Truncation error bounded by

$$\sum_{n=n_{\max}+1}^{\infty} \langle \psi_0 | \psi_n \rangle \sim \left\{ \begin{array}{ll} \langle \psi_0 | \psi_{n_{\max}} \rangle & \text{if} \quad \langle \psi_0 | \psi_n \rangle \stackrel{\text{exponentially}}{\to} 0, \\ n_{\max} \langle \psi_0 | \psi_{n_{\max}} \rangle & \text{if} \quad \langle \psi_0 | \psi_n \rangle \stackrel{\text{algebraically}}{\to} 0. \end{array} \right.$$

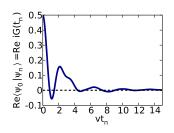
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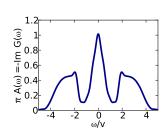
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FAW, Justiniano, McCulloch & Schollwöck, PRB 91, 115144 (2015)

Can we reach high enough values of  $n_{\text{max}}$ ?

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Convergence  $\langle \psi_0 | \psi_n \rangle \to 0$  depends on degree of differentiability of  $A(\omega)$ 

- smooth ▷ exponential convergence
- step function  $\triangleright \frac{1}{t_n}$  convergence
- singular ▷ no convergence

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#### What does this mean in practice?

• Real axis  $A(\omega) \triangleright van \ Hove \ kinks$  smoothed out for fermionic interacting systems  $\rightarrow$  fast convergence except at criticality

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- Imag axis  $G(i\omega) 
  ightharpoonup$  metallic phase o algebraic convergence / insulating phase o exponential

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#### Can we reach high enough values of $n_{\text{max}}$ ?

Real axis  $A(\omega)$ 

$$|\psi_n\rangle = e^{-i(H-E_0)t_n}|\psi_0\rangle$$

entanglement explodes, but convergence is often reached earlier

Imag axis  $G(i\omega)$ 

$$|\psi_n\rangle = e^{-(H-E_0)\tau_n}|\psi_0\rangle$$

- no entanglement generated, can compute arbitrarily long times
- metallic phase: very long times have to be computed
- insulating phase: only extremely short times need to be computed

What if entanglement explodes before convergence?

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Windowing / broadening approach: enforce convergence

$$A_{\eta}(\omega) \sim \sum_{n} \langle \psi_0 | \psi_n \rangle e^{i\omega t_n} e^{-(\eta t_n)^2/2}$$

- broadened version  $A_{\eta}(\omega)$  of  $A(\omega)$  conserves sum rules
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- is the usual approach in DDMRG or dynamic DMET

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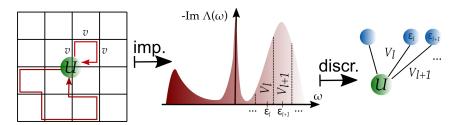
Linear prediction (extrapolation) White & Affleck, PRB 77, 134437 (2008)

- analytically continue to convergence FAW et al., PRB 91, 115144 (2015)
- enhance resolution
- requires exponential convergence

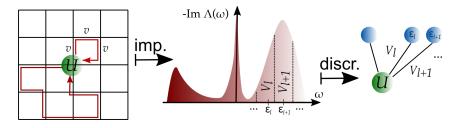
#### Outline

- Overview: impurity solvers and quantum embedding
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- Benchmark: two-site DCA in different setups

### DMFT: optimal bath discretization



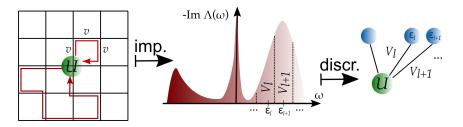
### DMFT: optimal bath discretization



#### Imag axis

- minimize  $\sum_{\omega_n}\left|\Lambda_{\mu
  u}(i\omega_n)-\sum_lrac{V_{\mu l}V_{\nu l}^*}{i\omega_n-arepsilon_l}
  ight|^2$  Caffarel & Krauth, PRL 72, 1545 (1994)
- unstable for many bath sites and/or off-diagonal couplings?
   Liebsch & Ishida, J. Phys.: Condens. Matter 24, 053201 (2012)
   Go & Millis, PRL 114, 016402 (2015)
- extremely fast convergence: 8 bath sites suffice for perfect fit!

### DMFT: optimal bath discretization



#### Real axis

discretization strategy

Bulla, Costi & Pruschke, RMP 80, 395 (2008)

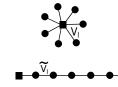
$$|V_{\mu l}|^2 = \int_{\omega_l}^{\omega_{l+1}} d\omega \left(-\mathrm{Im}\Lambda_{\mu\mu}(\omega)\right), \ \varepsilon_l = \frac{1}{|V_{\mu l}|^2} \int_{\omega_l}^{\omega_{l+1}} d\omega \,\omega \left(-\mathrm{Im}\Lambda_{\mu\mu}(\omega)\right)$$

- orthogonal polynomial strategy e.g. Shenvi, Schmidt, Edwards & Tully, PRA 78, 022502 (2008)
- ullet no notion of optimality! no off-diagonal  $\Lambda_{ij}(\omega)$  de Vega & FAW (in progress)
- much slower convergence: many bath sites needed!

### DMFT: geometry of bath

FAW, McCulloch & Schollwöck, PRB 90, 23513 (2014b)

#### star or chain?



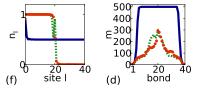




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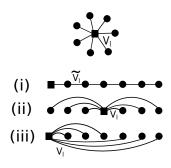
ground state

- star: local, lowly entangled
- chain: delocalized, highly entangled



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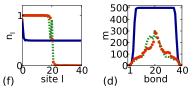
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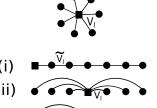
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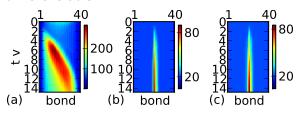


### star or chain?

(iii)



#### time evolution



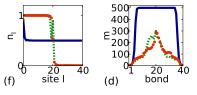
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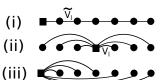
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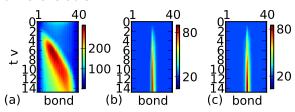


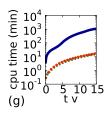
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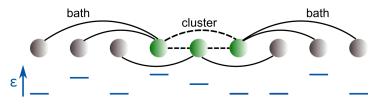




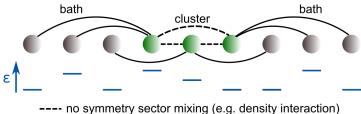
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- ---- no symmetry sector mixing (e.g. density interaction)
   symmetry sector mixing (e.g. hopping)
  - ε potential energy



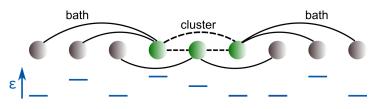
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use perturbation techniques

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    White, PRB 72, 180403 (2005)
    Dolgov & Savostyanov, SIAM J. Sci. Comput. 36, A2248 (2014)
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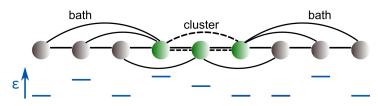
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- • use U=0 solution FAW, McCulloch, Parcollet & Schollwöck, PRB 90, 115124 (2014a)
- use additional hoppings / annealing

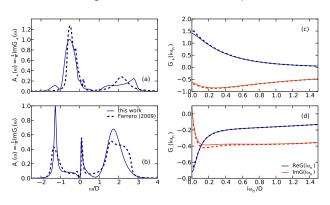
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## Benchmark: two-site cluster DCA in K-space repres.

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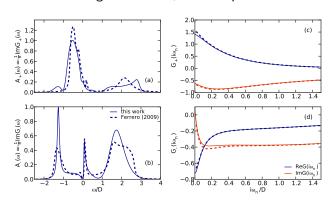
- model: hole-doped (4%) Hubbard model on 2 dim. square lattice
- spectral function: adaptive Chebyshev expansions, linear prediction
- bath discretization: linear discretization,  $L_b/L_c=30-40$
- geometry: chain
- ullet cpu time:  $\sim$  60 min ground state, 300 min spectral function



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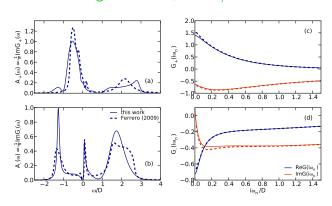
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- spectral function: time evolution
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- geometry: star
- cpu time:  $\sim$  60 min ground state, 40 min spectral function



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- model: hole-doped (4%) Hubbard model on 2 dim. square lattice
- Matsubara Green's function: imag time evolution
- bath discretization: fitting,  $L_b/L_c=8$
- geometry: star
- cpu time:  $\sim 15 \, \text{min}$  ground state,  $15 \, \text{min}/60 \, \text{min}$  Green's function



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- A. Millis (Columbia U.)
- O. Parcollet (CEA Saclay)
- I.P. McCulloch (U. Queensland)
- C. Hubig (LMU Munich)

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# Thank you!

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