1-1试证明两个空间矢量 $\mathbf{r}_1(r_1,\theta_1,\varphi_1)$ 和矢量 $\mathbf{r}_2(r_2,\theta_2,\varphi_2)$ 之间的夹角 $\boldsymbol{\Theta}$ 的余弦为

$$\cos\Theta = \cos\theta_1\cos\theta_2 + \sin\theta_1\sin\theta_2\cos(\varphi_1 - \varphi_2)$$

解: 矢量 $r_1 \{r_1 \sin \theta_1 \cos \varphi_1, r_1 \sin \theta_1 \sin \varphi_1, r_1 \cos \theta_1\}$, 矢量 $r_2 \{r_2 \sin \theta_2 \cos \varphi_2, r_2 \sin \theta_2 \sin \varphi_2, r_2 \cos \theta_2\}$,

$$r_1 \bullet r_2 = r_1 r_2 \cos \Theta, \quad r_1 \bullet r_2 = r_{1x} r_{2x} + r_{1y} r_{2y} + r_{1z} r_{2z}$$

所以

$$\cos\Theta = \frac{r_1 \bullet r_2}{r_1 r_2}$$

$$= \cos\theta_1 \cos\theta_2 + \sin\theta_1 \sin\theta_2 (\cos\varphi_1 \cos\varphi_2 + \sin\varphi_1 \sin\varphi_2)$$

$$= \cos\theta_1 \cos\theta_2 + \sin\theta_1 \sin\theta_2 \cos(\varphi_1 - \varphi_2)$$

1-2 设 $R = \sqrt{(x-x')^2 + (y-y')^2 + (z-z')^2}$ 试证 $\nabla R = \frac{R}{R}$, $\nabla \frac{1}{R} = -\frac{R}{R^3}$ 。 当用算符作用在R或1/R上时,结果为 $\nabla = -\nabla'$ 。

解:

$$\nabla R = \frac{\partial R}{\partial x} \mathbf{e}_{x} + \frac{\partial R}{\partial y} \mathbf{e}_{y} + \frac{\partial R}{\partial z} \mathbf{e}_{z} = \frac{1}{2} \frac{2(x - x')}{R} \mathbf{e}_{x} + \frac{1}{2} \frac{2(y - y')}{R} \mathbf{e}_{y} + \frac{1}{2} \frac{2(z - z')}{R} \mathbf{e}_{z} = \frac{R}{R}$$

$$\nabla \frac{1}{R} = -\frac{1}{2} \frac{2(x - x')}{R^{3}} \mathbf{e}_{x} - \frac{1}{2} \frac{2(y - y')}{R^{3}} \mathbf{e}_{y} - \frac{1}{2} \frac{2(z - z')}{R^{3}} \mathbf{e}_{z} = -\frac{R}{R^{3}}$$

$$\nabla' R = \frac{\partial R}{\partial x'} \mathbf{e}_{x} + \frac{\partial R}{\partial y'} \mathbf{e}_{y} + \frac{\partial R}{\partial z'} \mathbf{e}_{z} = -\frac{1}{2} \frac{2(x - x')}{R} \mathbf{e}_{x} - \frac{1}{2} \frac{2(y - y')}{R} \mathbf{e}_{y} - \frac{1}{2} \frac{2(z - z')}{R} \mathbf{e}_{z} = -\frac{R}{R}$$

$$\nabla' \frac{1}{R} = \frac{1}{2} \frac{2(x - x')}{R^{3}} \mathbf{e}_{x} + \frac{1}{2} \frac{2(y - y')}{R^{3}} \mathbf{e}_{y} + \frac{1}{2} \frac{2(z - z')}{R^{3}} \mathbf{e}_{z} = \frac{R}{R^{3}}$$

结果为 $\nabla = -\nabla'$ 。很明显,这只是对R的函数作用才有此结果,对别的函数作用要另外考虑。

1-4 已知 $\vec{B} = \vec{r}10e^{-2r}\cos\varphi + \vec{Z}10\sin\varphi$ 在 (2, 0, 3)

计算 $\nabla \bullet B$ 和 $\nabla \times B$

解: 柱坐标系下
$$\nabla \bullet a = \frac{1}{\rho} \frac{\partial (\rho a_{\rho})}{\partial \rho} + \frac{1}{\rho} \frac{\partial a_{\varphi}}{\partial \varphi} + \frac{\partial a_{z}}{\partial z}$$

$$\therefore \nabla \bullet B = \frac{10e^{-2r}\cos\varphi}{r} - 20e^{-2r}\cos\varphi\Big|_{(2,0,3)} = -15e^{-4}$$

$$\nabla \times a = e_{\rho} \left(\frac{\partial a_{z}}{\rho \partial \varphi} - \frac{\partial a_{\varphi}}{\partial z} \right) + e_{\varphi} \left(\frac{\partial a_{\rho}}{\partial z} - \frac{\partial a_{z}}{\partial \rho} \right) + \frac{e_{z}}{\rho} \left[\frac{\partial (\rho a_{\varphi})}{\partial \rho} - \frac{\partial a_{\rho}}{\partial \varphi} \right]$$

$$\nabla \times B = \frac{10\cos\varphi}{r}r + \frac{10e^{-2r}\sin\varphi}{r}z\Big|_{(2,0,3)} = 5r$$

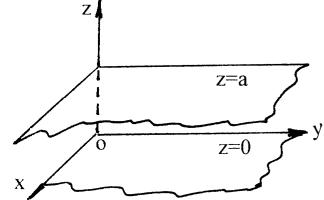
或者 (5, 0, 0) 矢量

2-1两无限大相互平行的理想导体平板,间距为a,其间存在一随时间变化的电场。当取其中一块板为z=0平面时,电场强度为

$$E = A \sin \frac{\pi z}{a} \cos \frac{\pi ct}{a}$$

式中 c 是光速。试求

- a. 磁感应强度B
- b. 导电板上的面电荷密度
- c. 导电板上的面电流密度



解: a. $E = A\sin\frac{\pi c}{a}\cos\frac{\pi ct}{a}e_x$ 满足在z=0,和z=a面上的边界条件,

由Maxwell方程可作运算:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} = \mathbf{e}_y \frac{\partial}{\partial z} \left[A \sin \frac{\pi z}{a} \cos \frac{\pi c t}{a} \right] = \mathbf{e}_y \left[A \frac{\pi}{a} \cos \frac{\pi z}{a} \cos \frac{\pi c t}{a} \right]$$

$$\mathbf{B} = -\frac{A}{c} \mathbf{e}_y \cos \frac{\pi z}{a} \int \left[\cos \frac{\pi c t}{a} \right] d \frac{\pi c t}{a} = -\frac{A}{c} \mathbf{e}_y \cos \frac{\pi z}{a} \sin \frac{\pi c t}{a}$$
式中积分常数因与时间无关,系静磁场,所以令其为零。
b.在 $z = 0$,和 $z = a$ 面上,由边界条件 $\sigma_f = E_n \varepsilon_0 = 0$ c.由边界条件 $\mathbf{k}_f = \mathbf{e}_n \times (\mathbf{H}_2 - \mathbf{H}_1)$, \mathbf{e}_n 由1指向2,在 $z = 0$ 面上。

$$\begin{cases} \boldsymbol{e}_{n} = \boldsymbol{e}_{z}, \boldsymbol{H}_{2} = \frac{\boldsymbol{B}}{\mu_{0}}, \boldsymbol{H}_{1} = 0 \\ \boldsymbol{k}_{f} = -\frac{A}{\mu_{0}c} \cos \frac{\pi z}{a} \sin \frac{\pi ct}{a} \boldsymbol{e}_{z} \times \boldsymbol{e}_{y} \Big|_{z=0} = \boldsymbol{e}_{x} \frac{A}{\mu_{0}c} \sin \frac{\pi ct}{a} \end{cases}$$

在z=a面上,

$$\begin{cases} \mathbf{e}_{n} = \mathbf{e}_{z}, \mathbf{H}_{1} = \frac{\mathbf{B}}{\mu_{0}}, \mathbf{H}_{2} = 0 \\ \mathbf{k}_{f} = \frac{A}{\mu_{0}c} \cos \frac{\pi z}{a} \sin \frac{\pi ct}{a} \mathbf{e}_{z} \times \mathbf{e}_{y} \Big|_{z=a} = \mathbf{e}_{x} \frac{A}{\mu_{0}c} \sin \frac{\pi ct}{a} \end{cases}$$

2-2 设在一载有稳恒电流*i*的长直导线附近,有一矩形闭合回路,边长为*a*和*b*,其中*b*边平行于长导线。当回路在包含长导线的平面内以匀速*v*离长导线而运动时,求回路中的感应电动势。

解: 在矩形线圈内, **B**的方向与线圈平面的法线方向一致, 所以有:

$$\Phi = \int \mathbf{B} \bullet \, \mathrm{d} \, \mathbf{s} = \int \frac{i\mu_0}{2\pi\rho'} b \, \mathrm{d} \rho' = \frac{i\mu_0}{2\pi} b \int_{\rho}^{\rho+a} \frac{\mathrm{d} \rho'}{\rho'} = \frac{i\mu_0}{2\pi} b \ln \rho' \Big|_{\rho}^{\rho+a}$$

$$= \frac{i\mu_0}{2\pi} b \ln \frac{\rho + a}{\rho}$$

$$\varepsilon = -\frac{\mathrm{d} \Phi}{\mathrm{d} t} = -\frac{i\mu_0 b}{2\pi} \frac{\rho}{\rho + a} \frac{\mathrm{d}}{\mathrm{d} t} (1 + \frac{a}{\rho}) = \frac{i\mu_0 abv}{2\pi\rho(\rho + a)}$$

2-3试利用坡印亭矢量分析稳恒载流直导线中的能量传输 问题,并证明由此导线周围流入导线的功率恒等于该 导线单位时间内的焦耳损耗。

解:设稳恒载流直导线中的电流强度是I,导线表面的磁场强度是 $H = \frac{I}{2\pi a}e_{\varphi}$,一段I长导线周围流入导线的功率:

$$P = -\int (\mathbf{E} \times \mathbf{H}) \bullet ds = -\frac{I}{\pi a^2 \sigma} (\mathbf{e}_z \times \mathbf{e}_\varphi) \bullet \mathbf{e}_\varphi \frac{I}{2\pi a} 2\pi a l = \frac{I^2 l}{\pi a^2 \sigma}$$
$$= I^2 R, (\sharp + R = \frac{l}{\pi a^2 \sigma})$$

这正好是该导线单位时间内的焦耳损耗。

2-4太阳在正午入射地球表面,与入射方向垂直的单位面积上所具有的能量为1.53×10⁶ 耳格/秒·厘米² 称为太阳常数。试求在地球表面上太阳光的电磁场强度。设太阳半径R_s等于7×10¹⁰ 厘米,太阳中心与地面间的距离

 R_{s-e} 是 1.5×10^{13} 厘米。求太阳表面上的电磁场强度。

解: 化为国际单位制。1焦耳=107尔格

$$g_e = E_e H_e = \frac{E_e^2}{c\mu_0} = 1.53 \times 10^6 \, \text{r/m} \, \text{k/m}. \text{m/m}^2$$

$$=\frac{10^4}{10^7}1.53\times10^6$$
焦耳/秒.米²=1.53×10³焦耳/秒.米²

所以,
$$E_e^2 = c\mu_0 g_e = 3 \times 10^8 (\text{米/秒}) \times 4\pi \times 10^{-7} (亨利/米) \times 1.53 \times 10^3 焦耳 / 秒.米^2$$

$$E_e = 759$$
伏/米, $H_e = \frac{E_e}{c\mu_0} = 2.01$ 安/米

$$g_e = E_e H_e = \frac{E_e^2}{c\mu_0} = 1.53 \times 10^6 \text{ 尔格 / 秒.} 厘米^2$$

$$= \frac{10^4}{10^7} 1.53 \times 10^6 \text{ 焦耳 / 秒.} **^2 = 1.53 \times 10^3 \text{ 焦耳 / 秒.} **^2$$

$$\frac{g_s}{g_e} = \frac{4\pi r_{s-e}^2}{4\pi r_s^2} = \frac{\frac{E_s^2}{c\mu_0}}{\frac{E_e^2}{c\mu_0}}, \qquad E_s = 1.63 \times 10^5 \text{ 伏 / } **,$$

$$H_s = 431 \text{ 安 / } **$$

2-5 一平面电磁波
$$\begin{cases} \mathbf{E} = E_0 \cos \omega (t - r/c) \mathbf{e}_x \\ \mathbf{B} = B_0 \cos \omega (t - r/c) \mathbf{e}_y \end{cases}$$
 垂直入于 $z = 0$ 平面。

求作用在此平面上的压强(单位面积的辐射压力)。 设(a)此平面为完全吸收体;(b)此平面为理想导体.

解: 因为
$$\ddot{\Phi} = \varepsilon_0 EE + \frac{1}{\mu_0} BB - \frac{1}{2} (\varepsilon_0 E^2 + \frac{1}{\mu_0} B^2) \ddot{S}$$
,

由电磁场的方向可确定入射波的方向在z方向,即由下半空间入射到 z=0 平面,则平面所受压强为 $-e_z \bullet \ddot{\boldsymbol{\sigma}}$,式中的电磁场是边界面上的总场。

(a)平面为完全吸收体,场在边界上没有变化,进入平面后才转变为热能.

$$P_{1} = -\boldsymbol{e}_{z} \bullet \left\{ \boldsymbol{\varepsilon}_{0} E^{2} \boldsymbol{e}_{x} \boldsymbol{e}_{x} + \frac{B^{2}}{\mu_{0}} \boldsymbol{e}_{y} \boldsymbol{e}_{y} - \frac{1}{2} (\boldsymbol{\varepsilon}_{0} E^{2} + \frac{B^{2}}{\mu_{0}}) (\boldsymbol{e}_{x} \boldsymbol{e}_{x} + \boldsymbol{e}_{y} \boldsymbol{e}_{y} + \boldsymbol{e}_{z} \boldsymbol{e}_{z}) \right\}$$

$$= \frac{1}{2} (\boldsymbol{\varepsilon}_{0} E^{2} + \frac{B^{2}}{\mu_{0}}) \boldsymbol{e}_{z}$$

(b)
$$E_{\mathbb{H}} = 0$$
 $B_{\mathbb{H}} = 2B$

$$P_2 = -e_z \bullet \left\{ -\frac{1}{2} \left(\frac{4B^2}{\mu_0} \right) (e_z e_z) \right\} = \frac{2B^2}{\mu_0} e_z$$

因为
$$\varepsilon_0 E^2 = \frac{B^2}{\mu_0}$$

所以
$$P_2 = 2P_1$$

3-2 对于在任意方向传播的电磁波,试证其群速度为
$$\mathbf{v}_{g} = \nabla_{k} \omega|_{\omega_{0}}$$
 。式中 $\nabla_{k} = \mathbf{e}_{x} \frac{\partial}{\partial k_{x}} + \mathbf{e}_{y} \frac{\partial}{\partial k_{y}} + \mathbf{e}_{z} \frac{\partial}{\partial k_{z}}, \omega_{0}$ 为波包的中心频率。

解:由傅利叶分析知,在k方向传播的频率在 ω_0 $-\frac{\Delta\omega}{2}$ 到 $\omega_0 + \frac{\Delta\omega}{2}$ 之间的所有单色平面波的迭加是

$$\boldsymbol{E} = \frac{1}{\Delta\omega} \int_{\omega_0 - \frac{\Delta\omega}{2}}^{\omega_0 + \frac{\Delta\omega}{2}} \boldsymbol{E}_0(\omega) \exp[j(\boldsymbol{k} \cdot \boldsymbol{r} - \omega t)] d\omega$$

设在所考虑的区间变化不大,且有

$$\mathbf{k}(\omega) = \mathbf{k}(\omega_0) + (\omega - \omega_0)(\frac{\mathrm{d}\mathbf{k}}{\mathrm{d}\omega})_0 + \cdots \approx \mathbf{k}_0 + (\omega - \omega_0)(\frac{\mathrm{d}\mathbf{k}}{\mathrm{d}\omega})_0$$

将其代入上式,得

$$\begin{split} & \boldsymbol{E} = \frac{\boldsymbol{E}_0}{A\omega} \int_{\omega_0 - \frac{A\omega}{2}}^{\omega_0 + \frac{A\omega}{2}} \exp \left[j(\boldsymbol{k}_0 \bullet \boldsymbol{r} + (\omega - \omega_0) (\frac{\mathrm{d}\,\boldsymbol{k}}{\mathrm{d}\,\omega})_0 \bullet \boldsymbol{r} - \omega\,t) \right] \mathrm{d}\omega \\ & = \frac{\boldsymbol{E}_0}{A\omega} \exp \left[j(\boldsymbol{k}_0 \bullet \boldsymbol{r} - \omega_0\,t) \right] \int_{\omega_0 - \frac{A\omega}{2}}^{\omega_0 + \frac{A\omega}{2}} \exp \left\{ j(\omega - \omega_0) \left[(\frac{\mathrm{d}\,\boldsymbol{k}}{\mathrm{d}\,\omega})_0 \bullet \boldsymbol{r} - t \right] \right] \mathrm{d}\omega \\ & = \frac{\boldsymbol{E}_0}{A\omega} \exp \left[j(\boldsymbol{k}_0 \bullet \boldsymbol{r} - \omega_0\,t) \right] \frac{\sin\psi}{\psi} \\ & = \frac{1}{2} \left[\left(\frac{\mathrm{d}\,\boldsymbol{k}}{\mathrm{d}\,\omega} \right)_0 \bullet \boldsymbol{r} - t \right] \frac{\Delta\omega}{2} = \mathrm{const.} = \left[\left(\frac{\mathrm{d}\boldsymbol{k}_x}{\mathrm{d}\,\omega} \right)_0 \boldsymbol{x} + (\frac{\mathrm{d}\boldsymbol{k}_y}{\mathrm{d}\,\omega})_0 \boldsymbol{y} + (\frac{\mathrm{d}\boldsymbol{k}_z}{\mathrm{d}\,\omega})_0 \boldsymbol{z} - t \right] \frac{\Delta\omega}{2} \\ & = \left(\frac{\mathrm{d}\,\boldsymbol{k}}{\mathrm{d}\,\omega} \right)_0 \bullet \frac{\mathrm{d}\,\boldsymbol{r}}{\mathrm{d}\,t} = 1 \\ & = \left(\frac{\partial\omega}{\partial k_x} \right)_0 \left(\frac{\mathrm{d}\boldsymbol{k}_x}{\mathrm{d}\,\omega} \right)_0 + \left(\frac{\partial\omega}{\partial k_y} \right)_0 \left(\frac{\mathrm{d}\boldsymbol{k}_y}{\mathrm{d}\,\omega} \right)_0 + \left(\frac{\partial\omega}{\partial k_z} \right)_0 \left(\frac{\mathrm{d}\boldsymbol{k}_z}{\mathrm{d}\,\omega} \right)_0 = \left(\frac{\mathrm{d}\omega}{\mathrm{d}\,\omega} \right)_0 = 1 \end{split}$$

$$\mathbf{v}_{g} = \frac{\mathrm{d}\mathbf{r}}{\mathrm{d}t} = \left(\frac{\partial\omega}{\partial k_{x}}\right)_{0} \mathbf{e}_{x} + \left(\frac{\partial\omega}{\partial k_{y}}\right)_{0} \mathbf{e}_{y} + \left(\frac{\partial\omega}{\partial k_{z}}\right)_{0} \mathbf{e}_{z}$$

3-3 一平面电磁波 $\boldsymbol{E}_{in} = \boldsymbol{E}_0 \exp[\mathrm{j}(\boldsymbol{k} \bullet \boldsymbol{r} - \omega t)]$ 射入一导

体平面,设入射角为 θ ,试求其折射角的大小。(导体的导电率为 σ ,导磁率为 μ)

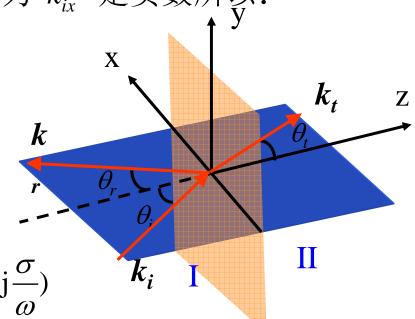
解:取入射平面为xz面,入射角为 θ_i ,折射角为 θ_t 。由边界条件 $k_{ix} = k_{tx} = \beta_{tx} + j\alpha_{tx}$,因为 k_{ix} 是实数所以:

$$\alpha_{tx} = 0$$
 $k_{ix} = \beta_{tx}$

$$k_{iy} = 0 = \beta_{ty} + j\alpha_{ty} ,$$

所以 α_t 垂直于导体表面

$$k_t^2 = \beta_t^2 - \alpha_t^2 + 2j\vec{\alpha}_t \bullet \vec{\beta}_t = \omega^2 \mu(\varepsilon + j\frac{\sigma}{\omega})$$



因为良导体内
$$\frac{\sigma}{\omega\varepsilon}$$
 >>1,所以 $\beta_t \approx \alpha_t \approx \sqrt{\frac{\omega\mu\sigma}{2}}$ $\vec{\alpha}_t \bullet \vec{\beta}_t = \alpha_{tz}\beta_{tz} = \frac{1}{2}\omega\mu\sigma = \frac{1}{2}\omega^2\mu_0\varepsilon_0 \frac{\sigma}{\omega\varepsilon_0}$ >> $\frac{1}{2}\frac{\omega^2}{c^2} = \frac{1}{2}k_i^2$ $\alpha_{tz}\beta_{tz} >> \frac{1}{2}k_i^2 = \frac{1}{2}(k_{ix}^2 + k_{iz}^2)$

所以
$$\alpha_{tz}\beta_{tz} >> k_{ix}^2 = \beta_{tx}^2$$
 所以 $\beta_{tz} >> \beta_{tx}$ 所以 $\beta_{tz} \approx \beta_t$

以上证明了在任意入射角情形下, α_t 垂直于导体表面, β_t 也

接近法线方向,
所以折射定理为
$$\frac{\sin \theta_{t}}{\sin \theta_{i}} = \frac{k_{i}}{\beta_{t}} = \frac{\frac{\omega}{c}}{\sqrt{\frac{\omega \mu \sigma}{2}}} \sqrt{\frac{2\omega}{c^{2} \mu \sigma}}$$

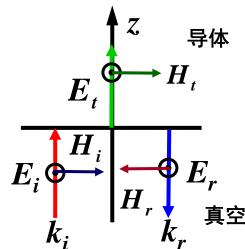
3-5 1GHz x方向极化的平面波沿+Z方向从空气入射到位于x-y平面的金属面(铜质, ε_r =1, μ_r =1, σ =5.8×10 7 s/m)上,电场幅度为12mV/m,求金属中的电场磁场时间表达式。

解: 首先考虑电磁波在空气和金属表面发生折射和透射的情况:

$$\begin{cases} E_i + E_r = E_t \\ H_i - H_r = H_t \end{cases} \quad \not \square \quad \mu_r = 1 \quad \therefore E_i - E_r = \sqrt{\frac{\sigma}{2\omega\varepsilon_0}} (1-j)E_t$$
 结合两式得到
$$\frac{2}{1 + \sqrt{\frac{\sigma}{2\omega\varepsilon_0}} (1-j)}$$
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$$\therefore \sqrt{\frac{\sigma}{2\omega\varepsilon_0}} = \sqrt{\frac{5.8 \times 10^7}{4\pi \times 10^9 \times 8.854 \times 10^{-12}}} \approx 22832$$

$$\therefore E_{t} \approx \frac{2}{22832(1-j)} E_{i} = \frac{1}{16145} E_{i} \times e^{j\frac{\pi}{4}}$$



再利用在导体中的电场表达式:

$$E = E_0 e^{-\alpha z} e^{j(\omega t - \beta z)}$$

在1GHz频率远远小于10¹⁷Hz可以看作良导体

$$\alpha \approx \beta \approx \sqrt{\omega\mu\sigma/2} = \sqrt{\pi f\mu\sigma} = 4.785 \times 10^5$$

$$\therefore E = 0.743e^{-4.785 \times 10^5 z} \exp[j(2\pi \times 10^9 t - 4.785 \times 10^5 z + \frac{\pi}{4})] \mu V / m$$

$$\therefore H = \sqrt{\frac{\sigma}{\omega\mu}} e^{-j\frac{\pi}{4}} \frac{\vec{k}}{k} \times E$$

$$\therefore H = 63.68e^{-4.785 \times 10^5 z} \exp[j(2\pi \times 10^9 t - 4.785 \times 10^5 z)] \quad \mu A / m$$

注意单位

- 4-3 一特性阻抗为50Ω的无耗线,与一200Ω的负载相接,其工作频率为100MHz,求:
 - a. 线上的驻波比。
 - b. 如采用一 λ/4 阻抗变换器进行匹配,则该匹配段的特性阻抗Z和长度*l*为多少?
 - c. 当工作频率变为80MHz时,如仍采用上述匹配段,则 线上的驻波比变为多少?

解: a.

$$\Gamma_l = \frac{Z_l - Z_0}{Z_l + Z_0} = \frac{200 - 50}{200 + 50} = 0.6, \ \rho = \frac{1 + |\Gamma_l|}{1 - |\Gamma_l|} = \frac{1 + 0.6}{1 - 0.6} = 4$$

b.
$$\lambda = \frac{c}{f} = \frac{3 \times 10^8}{10^8} = 3 \text{cm}, \ l = \frac{\lambda}{4} = 0.75 \text{m}$$
$$Z'_0 = \sqrt{Z_1 Z_{in}} = \sqrt{50 \times 200} = 100 \Omega$$

$$\lambda = \frac{c}{f} = \frac{3 \times 10^8}{10^8} = 3$$
cm, $l = \frac{\lambda}{4} = 0.75$ m

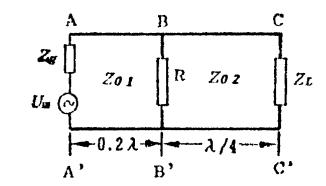
C.
$$f_1 = 80 \times 10^6 \,\text{Hz}, \lambda_1 = 3.75 \,\text{m}, \ l = 0.75 \,\text{m} = 0.2 \lambda_1,$$

$$Z_{in} = Z_0' \frac{Z_l + jZ_0' \tan \frac{2\pi}{\lambda_1} 0.2\lambda_1}{Z_0' + jZ_l \tan \frac{2\pi}{\lambda_1} 0.2\lambda_1} = 54 - j23.7$$

$$\Gamma = \frac{Z_l - Z_0}{Z_l + Z_0} = \frac{54 - j23.7 - 50}{54 - j23.7 + 50}, \ \left|\Gamma\right| = 0.228, \ \rho = \frac{1 + \left|\Gamma\right|}{1 - \left|\Gamma\right|} = 1.59$$

4-5 在如图所示的传输线电路中:

$$U_m = 900V$$
 $Z_g = Z_{01} = 450\Omega$ $Z_{02} = 600\Omega$ $R = 900\Omega$ $Z_L = 400\Omega$



a. 确定各线段上的工作状态。 b. 画出沿线电压、电流和阻抗的振幅分布,并标出它们的最大值和最小值。

解: a.
$$Z_{inB} = \frac{Z_{02}^2}{Z_I} = \frac{600^2}{400} = 900\Omega$$
 图4-5BB'处的总阻抗

$$R_{\text{BB}} = \frac{RZ_{inB}}{R + Z_{inB}} = \frac{900}{2} = 450\Omega$$

因为 $Z_g = Z_{01} = 450\Omega$,所以AB段工作在行波状态。 又 $Z_{02} = 600\Omega$, $Z_L = 400\Omega$,所以BC段工作在行驻波状态。

b. AB段: A点的输入阻抗
$$Z_{inB} = Z_{01} = 450\Omega$$

在AA'处的电流
$$I_A = \frac{U_m}{Z_g + Z_{inA}} = \frac{900}{450 + 450} = 1A = J(s)$$

在AA'处的电压
$$U(s) = 450V = U_l^+$$
 $Z_{in}(s) = \frac{U(s)}{J(s)} = 450\Omega$

BCET:
$$\Gamma_l = \frac{Z_l - Z_{02}}{Z_l + Z_{02}} = \frac{400 - 600}{400 + 600} = -0.2 = 0.2e^{j\pi}, \varphi_l = \pi$$

曲(4.3.12)式
$$|U(s)| = |U_{lC}^{+}| [1 + |\Gamma_{l}| e^{j(\pi-2\beta s)}],$$

在BB'处的电压为450V

$$= |U_{lC}^{+}| \left[1 + 0.2e^{j(\pi - 2\frac{2\pi}{\lambda}\frac{\lambda}{4})}\right] = |U_{lC}^{+}| \left[1 + 0.2\right] = 1.2|U_{lC}^{+}| \left[1 + 0.2\right] = 375V,$$

$$|U(s)| = 375\sqrt{1 + |\Gamma_l|^2 + 2|\Gamma_l|\cos(\pi - 2\beta s)}$$

$$\begin{cases} s = 0, |U_{\min}| = 300V \\ s = \frac{\lambda}{4}, |U_{\max}| = 450V \end{cases} \qquad |J(s)| = \frac{375}{600} \sqrt{1 + |\Gamma_l|^2 - 2|\Gamma_l| \cos(\pi - 2\beta s)}$$

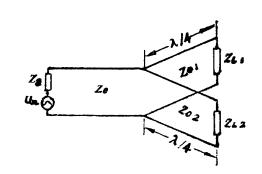
$$\begin{cases} s = 0, |J_{\text{max}}| = 0.75 \text{A} \\ s = \frac{\lambda}{4}, |J_{\text{min}}| = 0.5 \text{A} \end{cases}$$

$$\begin{aligned} |Z_{in}| &= \left| Z_{02} \frac{Z_l + j Z_{02} \tan \beta s}{Z_{02} + j Z_l \tan \beta s} \right| = 600 \left| \frac{2 + j 3 \tan \beta s}{3 + j 2 \tan \beta s} \right| \\ &= 600 \sqrt{\frac{13}{4 + 5 \cos \beta s}} - 1 \end{aligned}$$

所以
$$\begin{cases} s = 0, |Z_{\text{in,min}}| = 400\Omega \\ s = \frac{\lambda}{4}, |Z_{\text{in,max}}| = 900\Omega \end{cases}$$

4-6 一匹配信号源 $U_m = 10V$,通过一特性阻抗为 50Ω 的无耗传输线,以相等的功率馈送给两个分别为和 $Z_{L1} = 64\Omega$, $Z_{L2} = 25\Omega$ 的并联负载,并用 λ 4 变换器来实现负载与主传输线的匹配,如图所示。求:

- $a. \lambda/4$ 变换器的特性阻抗 Z_{01} 、 Z_{02} 。
- b. 在 $\lambda/4$ 匹配段上的驻波比。
- c. 负载 Z_{L1} 、 Z_{L2} 吸收的功率。



解: a.为了负载与主传输线匹配,要求 Z_{11} , Z_{12} 在AA'处的并联输入阻抗为 Z_0 ,为了以相等的功率馈送给两个并联负载,所以要求两个负载在AA'处的输入阻抗都等于2 Z_0 =100 Ω ,所以,

$$Z_{01} = \sqrt{2Z_0Z_{l1}} = \sqrt{100 \times 64} = 80\Omega,$$

 $Z_{02} = \sqrt{2Z_0Z_{l2}} = \sqrt{100 \times 25} = 50\Omega$

b.在匹配的情形下, 主线上的ρ=1, 而两支线上的:

1号线:
$$\Gamma_1 = \frac{Z_{l1} - Z_0}{Z_{l1} + Z_0} = \frac{64 - 80}{64 + 80} = -0.11;$$

$$\rho_1 = \frac{1 + |\Gamma_1|}{1 - |\Gamma_1|} = \frac{1 + 0.11}{1 - 0.11} = 1.25$$

2号线:
$$\Gamma_2 = \frac{Z_{l2} - Z_0}{Z_{l2} + Z_0} = \frac{25 - 50}{25 + 50} = -0.33;$$

$$\rho_2 = \frac{1 + |\Gamma_2|}{1 - |\Gamma_2|} = \frac{1 + 0.33}{1 - 0.33} = 2$$

c.因为是匹配的信号源,

所以
$$Z_g = 50\Omega$$
, $Z_{AA'} = 50\Omega$; $Z_{in} = 50\Omega$,

所以,
$$U_{AA'} = \frac{U_m Z_{AA'}}{Z_g + Z_0} = \frac{10 \times 50}{50 + 50} = 5V$$
, $I_A = \frac{10}{50 + 50} = 0.1A$,
$$P_{AA'} = \frac{1}{2} 5 \times 0.1 = 0.25W$$

所以,
$$P_1 = P_2 = \frac{1}{2}P_{AA'} = 0.125W$$

- 4-7有一特性阻抗为50Ω的无耗传输线,测得第一个电压最小点和最大点距离负载分别为5cm和15cm,振幅分别为5V和10V,求:
 - a. 负载反射系数 Γ_L 和负载阻抗 Z_L 。
 - b. 输入阻抗为 Z_L 的等效传输线的长度l及其终端电阻R。

解: (a) 15cm-5cm=λ/4,所以λ=40cm; 电压驻波比

$$\rho = \frac{10}{5} = 2 \qquad \left| \Gamma_l \right| = \frac{\rho - 1}{\rho + 1} = \frac{1}{3}, \quad \varphi_l = 2\beta d_{\min} + \pi = \pi + \frac{4\pi}{40}5 = \frac{3}{2}\pi$$

$$\Gamma_{l} = \left| \Gamma_{l} \right| e^{j\varphi_{l}} = \frac{1}{3} e^{j\frac{3}{2}\pi} = -\frac{1}{3} j, \quad Z_{l} = Z_{0} \frac{1 + \Gamma_{l}}{1 - \Gamma_{l}} = 50 \frac{1 - \frac{1}{3} j}{1 + \frac{1}{3} j} = (40 - j30)\Omega$$

(b)
$$Z_{in} = Z_l = 40 - j30 = 50 \frac{R + j50 \tan \beta l}{50 + jR \tan \beta l}$$

实部相等得 $200 + 3R \tan \beta l = 5R$

虚部相等得 $4R \tan \beta l - 150 = 250 \tan \beta l$

联立解得 $\begin{cases} R_1 = 100\Omega & l_1 = 5\text{cm} \\ R_2 = 25\Omega & l_2 = 15\text{cm} \end{cases}$

4-8 有一长为8cm、特性阻抗为50Ω 的无耗传输线,测得线上驻波比为2,相邻两电压最小点的距离为2.5cm及第一个电压最小点距负载的距离为1.5cm,利用圆图求负载阻抗 Z_L 和始端的输入阻抗 Z_m 。

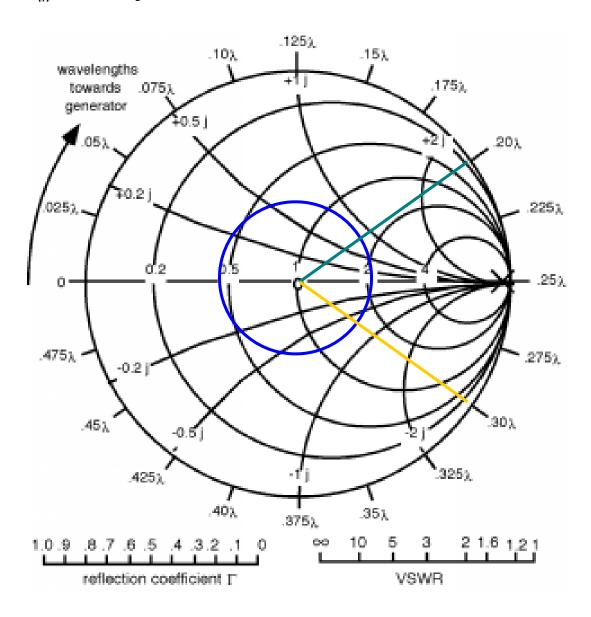
解:相邻两电压最小点的距离为 $\lambda/2=2.5$ cm,所以 $\lambda=5$ cm第一个电压最小点距负载的距离为 $d_{min}=1.5$ cm 所以 $\frac{d_{min}}{\lambda}=\frac{1.5}{5}=0.3$

输入端距负载的距离为8cm,即1.6λ

(1) 因为 $\rho = \tilde{R}_{max} = 2$,作等 ρ 圆,由电压最小点向负载方向 (逆时钟方向)转0.3得归一化负载阻抗 $\tilde{Z}_{l} = 1.57 + j0.69$, 或

$$Z_l = 50\tilde{Z}_l = (78.5 + j34.5)\Omega$$

(2)由负载点向顺时钟方向旋转1.6λ或0.1λ,得归一化输入阻抗 $\tilde{z}_{in} = 1.57 - j0.69$ 或 $Z_{in} = 50\tilde{Z}_{in} = (78.5 - j34.5)Ω$



- 4-9 传输线的特性阻抗 Z_0 = 300Ω ,负载阻抗 Z_L = (450 j150)Ω
- ,工作频率为1GHz,如利用 $\lambda/4$ 阻抗变换器来匹配这传输线。
- a. 求 λ/4 变换器的接入位置和阻抗特性。
- b. 如将 $\lambda/4$ 变换器直接接在负载与主传输线之间,则需在负载处并联一短路分支。求短路分支的长度和 $\lambda/4$ 变换器的特性阻抗。

解: a.
$$\lambda = \frac{c}{f} = 0.3 \text{m}, \Gamma_l = \frac{Z_l - Z_0}{Z_l + Z_0} = \frac{450 - \text{j}150 - 300}{450 - \text{j}150 + 300} = \frac{3 - \text{j}2}{13},$$
$$|\Gamma_l| = 0.277, \rho = \frac{1 + |\Gamma_l|}{1 - |\Gamma_l|} = 1.77$$
$$\varphi_l = \tan^{-1}(\frac{-2}{3}) = -0.588$$

如 $\lambda/4$ 变换器在电压最大点处接入,则该处的 $Z_{in(max)}=Z_0\rho$,

所以其特性阻抗
$$Z_{01} = \sqrt{Z_0 Z_0 \rho} = 300\sqrt{1.77} = 399.12\Omega$$

接入位置在
$$\varphi_l - 2\beta s = \pm 2n\pi$$
 $s_{max} = \frac{\lambda \varphi_l}{4\pi} + \frac{\lambda}{2} = 0.136$ m

如 $\lambda/4$ 变换器在电压最小点处接入,该处的 $Z_{in(\min)}=Z_0/\rho$

其特性阻抗
$$Z'_{01} = \sqrt{Z_0 Z_0 / \rho} = 300 / \sqrt{1.77} = 225.5\Omega$$

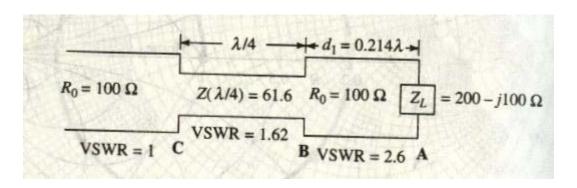
接入位置在
$$\varphi_l - 2\beta s = \pm (2n+1)\pi$$
 处,或 $s_{min} = \frac{\lambda \varphi_l}{4\pi} + \frac{\lambda}{4} = 0.061$ m

归一化负载阻抗
$$\tilde{Z}_l=Z_l/Z_0=1.5-\mathrm{j}0.5$$
,由圆图 $\tilde{Y}_l=0.6+\mathrm{j}0.2$ 。

如将 $\lambda/4$ 变换器直接接在负载与主传输线之间,则需在负载处并联一短路分支,要求它提供 - j0.2, 导纳圆图的短路在0.25处,所以 = 0.25 - 0.032 = 0.218 $l \div 0.0654$ m

这样
$$\tilde{Y}_{l}^{\&} = 0.6, \tilde{Z}_{l}^{\&} = 1.67, Z_{l}^{\&} = 500\Omega, Z_{01}$$
 $= \sqrt{Z_{0}Z_{l}^{\&}} = \sqrt{500 \times 300} = 387.3\Omega$

 $\lambda/4$ **Transformer matching.** A 100 Ω line is terminated in a load impedance $Z_L = 200 - j100\Omega$. Referring to Fig. 3_19, find: (a) d1(min) (b) $Z(\lambda/4)$ (c) VSWR on d1 line (d) VSWR on $\lambda/4$ line



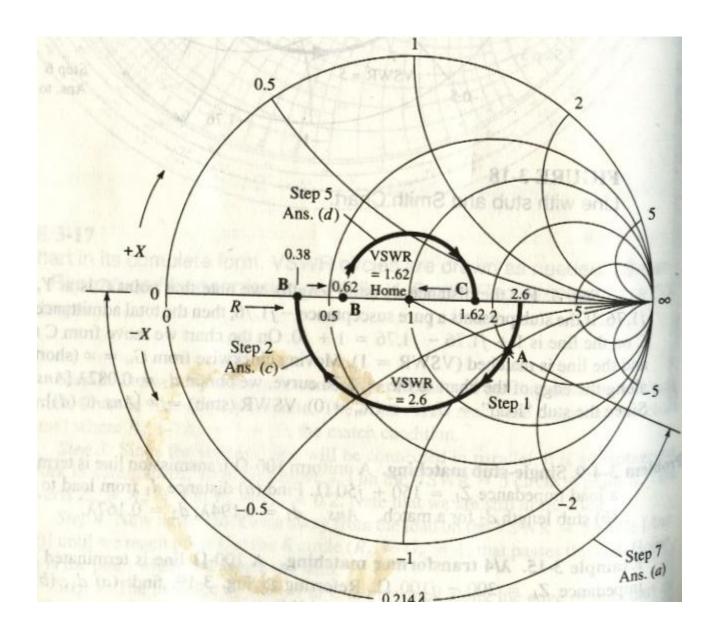
VSWR

Voltage standing wave ratio

$$Z_{nA-B} = \frac{200 - j100}{100} = 2 - j1 \qquad d_1 = 0.214\lambda$$

$$Z(\lambda/4) = \sqrt{Z_B R_0} = \sqrt{38 \times 100} = 61.6\Omega$$

$$Z_{nB-C} = \frac{Z_B}{Z(\lambda/4)} = \frac{38}{61.6} = 0.62$$



5-3

用BJ-32波导作馈线。求:

- a. 当工作波长分别为10cm、7cm和6cm时,波导中可能传输哪些波型?
- b. 波导单模工作的频率范围。
- c. 如果该波导中填充以 $\varepsilon_r = 2.25$ 的理想介质,其单模工作频率的范围如何变 化?

解: 波导的宽边
$$a$$
=7.214cm,窄边b=3.404cm。由截止波长 $\lambda_c = \frac{2}{\sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2}}$ 知

波型 TE₁₀ TE₂₀ TE₀₁ TE₁₁(TM₁₁) TE₂₁(TM₂₁) (cm) 14.43 7.21 6.81 6.16 4.95

- \mathbf{a} . 当 $\lambda < \lambda_c$ 时,波才能传输,所以有
- 工作波长为10cm时波导中可能传输 TE₁₀ 型波。
- 工作波长分别为7cm时波导中可能传输TE₁₀ TE₂₀型波。
- 工作波长分别为6cm时波导中可能传输 TE_{10} TE_{20} TE_{01} $TE_{11}(TM_{11})$ 型波。
- **b**. 波导单模工作的波长范围为7.21(cm)< λ <14.43(cm),由 $f=c/\lambda$,得它相应的单 模工作的频率范围为: 4.16(GHz)>f>2.079(GHz)

• c. 如果该波导中填充以 $\varepsilon_r = 2.25$ 的理想介质,其单模工作波长的范围不变,而截止频率 ε v c

 $f_c = \frac{v}{\lambda_c} = \frac{c}{\lambda_c \sqrt{\varepsilon_r}}$

• 则波型 TE_{10} , $f_c = 1.39 GHz$; TE_{20} 的 $f_c = 2.77 GHz$ 。 所以其单模工作频率的范围为 1.39 GHz < f < 2.77 GHz

5-4

矩形波导的工作频率 f = 5GHz ,传输 TE_{10} 的截止频率 $f_c = 0.8f$, 宽高比为 2,如通过波导的平均功率为1kW,求:

- a. 波导中电场和磁场强度的幅值。
- b. 波导壁上纵向和横向壁电流面密度的幅值。

解: a.
$$f_c = 0.8f = 4$$
GHz, $\lambda_c = \frac{3 \times 10^{10}}{4 \times 10^9} = 7.5$ (cm) = $2a$,

所以
$$a = 3.75$$
(cm), $b = \frac{a}{2} = 1.875$ (cm), $\lambda = \frac{3 \times 10^{10}}{5 \times 10^9} = 6$ (cm) 由(5.1.24)式,在

宽壁中心
$$x = \frac{a}{2}$$
 处, $|E_y| = |E_0| = \frac{\omega\mu a}{\pi} H_0$

• 通过波导的平均功率为
$$1kW$$
时
$$P = \frac{ab}{4\sqrt{\frac{\mu}{\varepsilon}}} \sqrt{1 - (\frac{\lambda}{2a})^2 \left| E_0 \right|^2}$$
 (5.1.28)

所以
$$\left| E_y \right| = \sqrt{\frac{4P\sqrt{\frac{\mu}{\varepsilon}}}{ab\sqrt{1-(\frac{\lambda}{2a})^2}}} = 597.87 \times 10^2 \,\text{V/m}$$
 (5.1.24)

$$H_0 = \frac{\pi}{\omega \mu a} |E_0| = \frac{\pi}{2\pi f a 4\pi \times 10^{-7}} = 1.269 \times 10^2 \text{ A/m} = |H_z|_{\text{max}}$$

$$k_z = \frac{2\pi}{\lambda} \sqrt{1 - (\frac{f_c}{f})^2} = 0.2\pi, |H_x|_{\text{max}} = \frac{k_z a}{\pi} H_0 = 95.2 \text{A/m}$$

b.在波导窄壁上,只有横向电流 $|J_{y}| = |H_{0}| = 126.9 \text{A/m}$

在波导宽壁上,
$$\left|J_{x}\right|_{\text{max}} = \left|H_{0}\right| = 126.9 \text{A/m}, \left|J_{z}\right|_{\text{max}} = \frac{k_{z}a}{\pi} \left|H_{0}\right| = 95.2 \text{A/m}$$
(5.1.27)

空气同轴线内外导体的直径分别为d=32mm,D=75mm,求:

- a. 该同轴线的特性阻抗。
- b. 当它采用 $\varepsilon_r = 2.25$ 的介质环支撑时,如D不变,则应为多少才能保证匹配?
- c. 该同轴线中不产生高次模的最高工作频率。

解: a. 空气同轴线的特性阻抗
$$Z_0 = \frac{60}{\sqrt{\varepsilon_r}} \ln \frac{b}{a} = 60 \ln \frac{75}{32} = 51.1Ω$$

b. 为了保证匹配,就要求介质环支撑段的特性阻抗保持不变。

$$\mathbb{Z}_0' = \frac{60}{\sqrt{2.25}} \ln \frac{75}{d'} = 51.1\Omega ,$$

所以 d' = 20.9mm

c. 在空气同轴线中,不产生高次模的最小工作波长为

$$\lambda_{\min} \ge 1.1\pi(a+b) = 1.1\pi(\frac{32+75}{2}) = 184.9 \text{mm}$$
 $f_{\max} = \frac{c}{\lambda_{\min}} = 1.62 \text{GHz}$

在介质环支撑段
$$\lambda_{\min} \ge 1.1\pi \frac{d'+D}{2} = 1.1\pi (\frac{20.9+75}{2}) = 165.7$$
mm

$$f_{\text{max}} = \frac{c/\sqrt{\varepsilon_r}}{\lambda_{\text{min}}} = \frac{3 \times 10^{10}}{\sqrt{2.25} \times 16.57} = 1.21 \text{GHz}$$
 所以取该频率能在整个同轴 线中都不产生高次模。

- 设计一特性阻抗为 75Ω 的同轴线,要求它的最高工作频率为4.2GHz,求当分别以空气和的介质填充时同轴线的尺寸。
- 解:特性阻抗, $Z_0 = \frac{60}{\sqrt{\varepsilon_r}} \ln \frac{b}{a}$ 最高工作频率为4.2GHz

$$\lambda_{\min} = \frac{c}{\sqrt{\varepsilon_r} f_{\max}} = \frac{7.143}{\sqrt{\varepsilon_r}} (\text{cm}) \quad a + b \le \frac{\lambda_{\min}}{1.1\pi}$$

- 当以空气填充时, $75 = 60 \ln \frac{b}{a}, a+b \le \frac{7.143}{1.1\pi} = 2.07 \text{ (cm)}$
- 可联立解得 $a \le 0.46$ cm, $b \le 1.61$ cm
- 当以的介质填充时, $\ln \frac{b}{a} = \frac{75}{60} \sqrt{2.25} = 1.875, \lambda_{\min} = \frac{7.143}{\sqrt{2.25}} = 4.76 \text{(cm)}$ $a+b \leq \frac{\lambda_{\min}}{1.1\pi} \leq 1.38 \text{(cm)},$
- 可联立解得 $a \le 0.184$ cm, $b \le 1.198$ cm

- 一空气同轴线内外径尺寸分别为d=3cm,D=7cm,当它的终端接的负载200 Ω 时,负载吸收的功率为1W,求:
- a. 为保证只传输TEM波的最高工作频率。
- b. 线上的驻波比和入射功率与反射功率。
- c.为使线上无反射,采用 $\lambda/4$ 线进行匹配,如保持D不变,则 $\lambda/4$ 线的内径d 为多少?匹配后负载吸收的功率为多少?
- 解: 为保证同轴线中只传输TEM模,就必须抑制最低次波导模TE₁₁模。
- A. $\lambda_{\min} \ge \frac{1.1\pi}{2} (D+d) = 17.279 (\text{cm}), f_{\max} = 1.74 \text{GHz}$ 同轴线的特性阻抗 $Z_0 = \frac{60}{\sqrt{\varepsilon_r}} \ln \frac{b}{a} = 60 \ln 2.3 = 50 \Omega$
- B. $|\Gamma| = \left| \frac{Z_l Z_0}{Z_l + Z_0} \right| = \frac{200 50}{200 + 50} = 0.6$ $\rho = \frac{1 + |\Gamma|}{1 |\Gamma|} = 4$
- 线上为行驻波状态,由式(4.3.17)(4.3.18) $P_l(s) = P^+ \left[1 \left| \Gamma_l \right|^2 \right] = 1(\mathbf{W})$

所以
$$P^+ = \frac{1}{\left[1 - \left|\Gamma_l\right|^2\right]} = 1.56(W), P^- = P^+ - P_l(s) = 0.56(W)$$

- C.为使线上无反射,采用 $\lambda/4$ 线进行匹配,
- 它的特性阻抗 $Z_{01} = \sqrt{Z_1 Z_0} = \sqrt{200 \times 50} = 100\Omega$
- $\pm 100 = 60 \ln \frac{7}{d'}, d' = 1.32 \text{(cm)}$

这样负载吸收全部入射功率 $P_l = 1.56(W)$

• 6-4

- 要求在厚度 h = 0.8mm, $t \rightarrow 0$, $\varepsilon_r = 9$ 的基片上制作特性阻抗分别为
- 50Ω 和 100Ω 的微带线, 求它们的导带宽度W。
- 解: 由判断参数 $A = \frac{Z_0}{60} \sqrt{\frac{\varepsilon_r + 1}{2}} + \frac{\varepsilon_r 1}{\varepsilon_r + 1} (0.23 + \frac{0.11}{\varepsilon_r})$

•
$$\begin{cases} Z_0 = 50\Omega, A = 2.06 \\ Z_0 = 100\Omega, A = 3.92 \end{cases} A > 1.52 \text{ MU}$$

$$\frac{W}{h} = \frac{8}{e^A - 2e^{-A}}, \begin{cases} Z_0 = 50\Omega, \frac{W}{h} = 1.05, W = 0.84 \text{mm} \\ Z_0 = 100\Omega, \frac{W}{h} = 0.159, W = 0.127 \text{mm} \end{cases}$$

- 已知微带线的参数为 h = 1mm, W = 1mm, $t \to 0$, $\varepsilon_r = 9$, 求:
- a. 它的最高工作频率。
- b. 当 f = 5GHz 微带中波的导波波长和相速度。
- 解: A. 为了避免TE和TM表面波和准TEM波之间的强耦合,
- 必须使工作频率低于 f_{TM} 和 f_{TE} 又因为 $f_{\text{TM}} < f_{\text{TE}}$,
- 所以只需要计算 $f_{\text{TM}} = \frac{c\sqrt{2}}{4h\sqrt{\varepsilon_r 1}} = 37.5 \text{GHz}$ 再看高次模的截止波长

• 它的最高工作频率为26.5GHz。

• B.
$$\varepsilon_e = \frac{\varepsilon_r + 1}{2} + \frac{\varepsilon_r - 1}{2} \left[(1 + 12\frac{h}{W})^{-\frac{1}{2}} + 0.04(1 - \frac{W}{h})^2 \right] = 6.11,$$

$$Z_0 = \frac{60}{\sqrt{\varepsilon_e}} \ln(8\frac{h}{W} + \frac{W}{4h}) = 51.2\Omega$$

$$f_0 = \frac{0.95}{(\varepsilon_r - 1)^{\frac{1}{4}}} \sqrt{\frac{Z_0}{h}} (\text{GHz}) = 4.03\text{GHz}$$

在此频率以下,色散效应可基本不考虑。因为 $t \to 0$,所以 $\Delta W = 0$, W' = W

$$\varepsilon_{e}' = 3 \times 10^{-6} (1 + \varepsilon_{r})(\varepsilon_{r} - 1) h \left[Z_{0} \frac{W'}{h} \right]^{\frac{1}{2}} (f - f_{0}) + \varepsilon_{e} = 6.2$$

$$\lambda_{0} = \frac{3 \times 10^{10}}{5 \times 10^{9}} = 6 \text{(cm)} \quad \lambda_{g} = \frac{\lambda_{0}}{\sqrt{\varepsilon_{e}'}} = \frac{6}{\sqrt{6.2}} = 2.41 \text{(cm)}$$

$$v_{p} = \frac{c}{\sqrt{\varepsilon_{e}'}} = 1.2 \times 10^{10} \text{(cm/s)}$$

已知由铜导体构成的微带线,W/h=1, t/h=0.02, $\varepsilon_r=9.6$, $\tan\delta=2\times10^{-4}$,工作频率 f=10GHz, 求微带线的介质衰减 α_d 和导体衰减 α_c 。

•
$$\mathbb{A}: \ \mathbb{B} \rightarrow \frac{W}{h} = 1 \ge \frac{1}{2\pi} = 0.159$$

• 所以
$$\frac{\Delta W}{h} = \frac{1.25}{\pi} \frac{t}{h} (1 + \ln \frac{2h}{t}) = 0.0446, \frac{W_e}{h} = \frac{W}{h} + \frac{\Delta W}{h} = 1.0446$$

$$\varepsilon_e = \frac{\varepsilon_r + 1}{2} + \frac{\varepsilon_r - 1}{2} (1 + 10\frac{h}{W})^{-\frac{1}{2}} - \frac{\varepsilon_r - 1}{4.6} \frac{t/h}{\sqrt{W/h}} = 6.56,$$

$$Z_0 = \frac{120\pi/\sqrt{\varepsilon_e}}{\frac{W_e}{h} + 1.393 + 0.667\ln(\frac{W_e}{h} + 1.444)} = 48.16\Omega$$

• 查图6. 4. 5, 得 $\frac{\alpha_c' Z_0 h}{R_s} = 4 \text{dB}$, 铜的 $R_s = 2.6 \times 10^{-7} \sqrt{f} \, (\Omega/\text{cm}^2)$ $\alpha_c' = \frac{4 \times 2.6 \times 10^{-7} \sqrt{10 \times 10^9}}{50h} = 0.002 / h \, (\text{dB/cm})$

$$\lambda_{g} = \frac{\frac{c}{\sqrt{\varepsilon_{e}}}}{f} = \frac{3 \times 10^{10}}{10 \times 10^{9} \times 2.57} = 1.17 \text{ (cm)},$$

$$\alpha'_{d} = 27.3 \frac{\varepsilon_{e} - 1}{\varepsilon_{r} - 1} \frac{\varepsilon_{r}}{\varepsilon_{e}} \frac{\tan \delta}{\lambda_{g}} \text{ (dB/cm)}$$

$$= 27.3 \frac{5.6 \times 9.6 \times 2 \times 10^{-4}}{8.6 \times 6.6 \times 1.17} \text{ (dB/cm)} = 0.0044 \text{ (dB/cm)}$$

• 8-1

有一矩形谐振腔,它沿方向的尺寸分别为a,b,l,试求在(1) a > b > l; (2) a > l > b; (3) a = b = l三种情形下腔的主模和它们的谐振频率。

解: 矩形谐振腔的谐振波长为
$$\lambda_0 = \frac{2}{\sqrt{\left(\frac{m}{a}\right)^2 + \left(\frac{n}{b}\right)^2 + \left(\frac{p}{l}\right)^2}}$$

对于TEmnp模:m,n中只能一个为零,p不能为零。

对于TM_{mmp}模: m, n都不能为零, p可为零。

• (1)
$$TM_{110}$$
 模为主模, $\lambda_0 = \frac{2}{\sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{b}\right)^2}} = \frac{2ab}{\sqrt{a^2 + b^2}}$

• 谐振频率
$$f_0 = \frac{c}{\lambda_0} = \frac{\sqrt{a^2 + b^2}}{2ab\sqrt{\varepsilon_0 \mu_0}}$$

• (2)
$$\text{TE}_{101}$$
 模为主模, $\lambda_0 = \frac{2}{\sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{l}\right)^2}} = \frac{2al}{\sqrt{a^2 + l^2}}$

• 谐振频率
$$f_0 = \frac{c}{\lambda_0} = \frac{\sqrt{a^2 + l^2}}{2al\sqrt{\varepsilon_0 \mu_0}}$$

• (3)
$$\text{TE}_{101} \text{ TE}_{011} \text{ TM}_{110}$$
 模为主模, $\lambda_0 = \frac{2}{\sqrt{\left(\frac{1}{a}\right)^2 + \left(\frac{1}{a}\right)^2}} = \sqrt{2} \ a$
• 它们互为简并

• 谐振频率
$$f_0 = \frac{c}{\lambda_0} = \frac{1}{\sqrt{2a\sqrt{\varepsilon_0\mu_0}}}$$

- 有一矩形谐振腔 (b = a/2), 已知当 f = 3GHz时它谐振于TE₁₀模; 当
- f = 6GHz 时它谐振于 TE_{103} 模,求此谐振腔的尺寸。
- 解: f=3GHz相应于 $\lambda_1 = \frac{3 \times 10^{10}}{2 \times 10^9} = 10$ (cm)

$$f=6\text{GHz相应于} \lambda_2 = \frac{3 \times 10^{10}}{6 \times 10^9} = 5\text{(cm)}$$

$$\lambda_0 = \frac{2}{\sqrt{(\frac{m}{a})^2 + (\frac{n}{b})^2 + (\frac{p}{d})^2}}, \begin{cases} \lambda_{0(\text{TE}_{101})} = \frac{2}{\sqrt{(\frac{1}{a})^2 + (\frac{1}{d})^2}} = 10\text{(cm)} \\ \lambda_{0(\text{TE}_{103})} = \frac{2}{\sqrt{(\frac{1}{a})^2 + (\frac{3}{d})^2}} = 5\text{(cm)} \end{cases}$$

联立解得 a = 6.32(cm), d = 8.15(cm), $b = \frac{a}{2} = 3.16$ (cm)

- 一空气填充的矩形谐振腔尺寸为 3×1.5×4cm3, 求:
- a. 当它工作于TE₁₀₁模时的谐振频率。
- b. 如腔中最大电场强度幅值 $E_m = 10^3 (\text{V/m})$,求腔中储存的总能量。
- c. 若在腔中全填充某种介质后,在同一工作频率上它谐振于**TE**₁₀₂ 模,则该介质的相对介电常数为多少**?**
- \widehat{H} : a. $\lambda_0 = \frac{2}{\sqrt{(\frac{m}{a})^2 + (\frac{n}{b})^2 + (\frac{p}{d})^2}} = \frac{2}{\sqrt{(\frac{1}{a})^2 + (\frac{1}{d})^2}}$ $= 4.8(\text{cm}), f_0 = \frac{3 \times 10^{10}}{4.8} = 6.25 \text{GHz}$
- b.因谐振腔中电场只有 $E_y = E_0 \sin \frac{\pi}{a} x \sin \frac{\pi}{d} z$, 当在 $x = \frac{a}{2}$, $z = \frac{d}{2}$ 处,

$$E_y = E_0 = E_{\text{max}} = 10^3 \text{ (V/m)}$$

$$W = \frac{\mathcal{E}_0}{2} \iiint E_y^2 dV = \frac{\mathcal{E}_0}{2} \int_0^a dx \int_0^b dy \int_0^d dz \left[E_0 \sin \frac{\pi x}{a} \sin \frac{\pi z}{d} \right]^2$$

• 腔中总能量

$$= \frac{\varepsilon_0 E_0^2 abd}{8} = 19.9 \times 10^{-12} (J)$$

• C. $f_0 = \frac{c}{\sqrt{\varepsilon_r \lambda_0}}, \quad \sqrt{\varepsilon_r} = \frac{c}{\lambda_0 f_0} = 4.8 \frac{\sqrt{(\frac{1}{a})^2 + (\frac{2}{d})^2}}{2} = 1.44, \quad \varepsilon_r = 2.08$

• 9-1

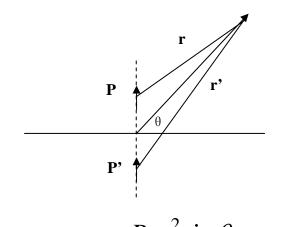
- 试证明一个均匀带电的球壳在作径向振动时将不会辐射。
- 解:一个均匀带电的球壳,所产生的电场是相当于电荷位于球心的静电场,作径向振动时也不改变。
- 因为电流的对球心的对称分布, 矢量势

$$\mathbf{A} = \frac{\mu_0}{4\pi} \int \frac{\mathbf{j}(x', y', z', t - \frac{r}{c})}{r} d\tau' = \frac{\mu_0}{4\pi r} \int \mathbf{j}^{ret} d\tau' = 0, 所以 \mathbf{B} = 0$$

• 9-3

- 一电偶极子的电偶极矩为P,以频率 ω 振荡。它被垂直地放置于离一个无限大理想导体平面距离为 $\lambda/2$ 处,这里 λ 是和频率 ω 相应的波长。设此电偶极子的尺度远较 λ 为小,并指向正z方向。试求其辐射场,辐射功率分布,及辐射总功率。
- 解:由电象法可知,这偶极子的电象为指向正z方向的偶极子,如图取近似

$$\begin{cases} \theta \approx \theta' \approx \theta_0, \\ r = r_0 - \frac{\lambda}{2} \cos \theta_0 \\ r' = r_0 + \frac{\lambda}{2} \cos \theta_0 \end{cases}$$
在分母中 $r \approx r_0 \approx r'$
由单个偶极子的辐射场 $E = \frac{\ddot{P}^{ret} \sin \theta}{4\pi\varepsilon_0 c^2 r} e_{\theta} = -\frac{\mu_0 P \omega^2 \sin \theta \cos \omega (t - \frac{r}{c})}{4\pi r} e_{\theta}$



$$\begin{split} & \boldsymbol{E}_{total} = -\boldsymbol{e}_{\theta} \frac{\mu_{0} P \omega^{2} \sin \theta_{0}}{4\pi r_{0}} \left\{ \cos \left[\omega t - \frac{\omega}{c} (r_{0} - \frac{\lambda}{2} \cos \theta_{0}) \right] \right. \\ & + \cos \left[\omega t - \frac{\omega}{c} (r_{0} + \frac{\lambda}{2} \cos \theta_{0}) \right] \right\} \\ & = -\boldsymbol{e}_{\theta} \frac{\mu_{0} P \omega^{2} \sin \theta_{0}}{4\pi r_{0}} 2 \cos(\omega t - \frac{\omega}{c} r_{0}) \cos(\frac{\omega \lambda}{2c} \cos \theta_{0}) \\ & = -\boldsymbol{e}_{\theta} \frac{\mu_{0} P \omega^{2} \sin \theta_{0}}{2\pi r_{0}} \cos(\omega t - \frac{\omega}{c} r_{0}) \cos(\pi \cos \theta_{0}) \end{split}$$

$$\mathbf{B}_{total} = -\mathbf{e}_{\varphi} \frac{\mu_0 P \omega^2 \sin \theta_0}{2\pi r_0 c} \cos(\omega t - \frac{\omega}{c} r_0) \cos(\pi \cos \theta_0)$$

$$g = E \times H = e_r \frac{\mu_0 P^2 \omega^4 \sin^2 \theta_0}{4\pi^2 r_0^2 c} \cos^2 (\omega t - \frac{\omega}{c} r_0) \cos^2 (\pi \cos \theta_0)$$

$$\frac{\mathrm{d}P}{\mathrm{d}\Omega} = \frac{\mathbf{g} \cdot \mathrm{d}\mathbf{s}}{\frac{\mathrm{d}\mathbf{s}}{r_0^2}} = \frac{\mu_0 P^2 \omega^4 \sin^2 \theta_0}{4\pi^2 c} \cos^2(\pi \cos \theta_0) \cos^2(\omega t - \frac{\omega r_0}{c})$$

$$P = \int g \cdot ds = \frac{\mu_0 P^2 \omega^4}{2\pi c} \cos^2(\omega t - \frac{\omega r_0}{c}) \int_0^{\frac{\pi}{2}} \sin^3\theta_0 \cos^2(\pi \cos\theta_0) d\theta_0$$
$$= (\frac{1}{3} - \frac{1}{4\pi^2}) \frac{\mu_0 P^2 \omega^4}{2\pi c} \cos^2(\omega t - \frac{\omega r_0}{c})$$

- 在一长度为 $2\pi a$ 的圆形导线中,通一电流 $i = I_0 \cos \omega t$ 计算辐射场,辐射功率
- 角分布和辐射总功率。假定λ>>a.。
- 解:

$$i = I_0 \cos \omega t, M = i s = I_0 \cos \omega t \pi a^2 e_z$$

$$\ddot{M}^{ret} = -I_0 \pi a^2 \omega^2 (\cos \theta e_r - \sin \theta e_\theta) \cos \omega (t - \frac{r}{c})$$

$$\boldsymbol{E} = \frac{\mu_0 \, \boldsymbol{r} \times \boldsymbol{\dot{M}}^{ret}}{4\pi c r^2} = \frac{\mu_0 I_0 a^2 \omega^2 \sin \theta \cos \omega (t - \frac{r}{c})}{4cr} \boldsymbol{e}_{\varphi},$$

$$\boldsymbol{B} = \frac{1}{c}\boldsymbol{e}_r \times \boldsymbol{E} = -\frac{\mu_0 I_0 a^2 \omega^2 \sin \theta \cos \omega (t - \frac{r}{c})}{4c^2 r} \boldsymbol{e}_{\theta},$$

$$\mathbf{g} = \frac{1}{\mu_0} \mathbf{E} \times \mathbf{B} = \frac{\mu_0 I_0^2 a^4 \omega^4 \sin^2 \theta \cos^2 \omega (t - \frac{r}{c})}{16c^3 r^2} \mathbf{e}_r,$$

$$\frac{dP}{d\Omega} = \frac{\mu_0 I_0^2 a^4 \omega^4 \sin^2 \theta \cos^2 \omega (t - \frac{r}{c})}{16c^3},$$

$$P = \int dP = \frac{\mu_0 I_0^2 a^4 \omega^4 \cos^2 \omega (t - \frac{r}{c})}{16c^3} \int \sin^3 \theta \, d\theta \, d\phi$$

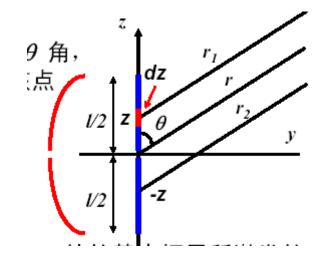
$$= \frac{\mu_0 \pi I_0^2 a^4 \omega^4 \cos^2 \omega (t - \frac{r}{c})}{6c^3}$$

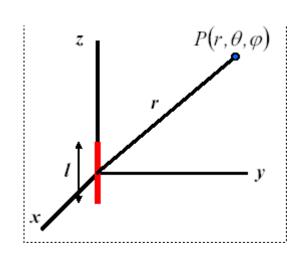
 一长为Nλ/2的细的线天线(λ为波长)。计算辐射场,单位立体角辐射功率和 辐射总功率。

• #:
$$I = I_m e^{j\omega t}$$
 $r = r_0 - \frac{\lambda}{2} \cos \theta_0$

$$A_{z} = \frac{\mu}{4\pi} \int \frac{N\lambda/4}{-N\lambda/4} \frac{I_{m}e^{j\omega(t-\frac{R}{v})}}{R} dl' = \frac{\mu I_{m}e^{j\omega t}}{4\pi} \int \frac{N\lambda/4}{-N\lambda/4} \frac{e^{-jk(r-Z\cos\theta)}}{R} dz$$

$$= \frac{\mu I_m e^{-jkr} \sin(\frac{N\pi}{2} \cos \theta)}{2\pi kr \cos \theta}$$





$$E_{\theta} = j\omega\mu\sin\theta \cdot A_{z}$$

$$H_{\phi} = \frac{1}{\eta}E_{\theta}$$

$$g_r = E_\theta \times H_\phi$$

$$\frac{\mathrm{d}P}{\mathrm{d}\Omega} = \frac{g \bullet ds}{\frac{ds}{r_0^2}}$$

$$P = \int g \bullet ds$$