Simplified Hodgkin-Huxley Model

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1 Introduction

Alan Lloyd Hodgkin and Andrew Huxley, in 1963, won the Noble Prize in Physiology and Medicine for their work developing the Hodgkin-Huxley Equations, equations that modeled to incredible accuracy the propagation of action potentials through a squid axon.

2 Premises of the Hodgkin-Huxley Model

The Hodgkin-Huxley model models the cell membrane as a simple circuit.

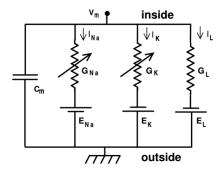


Figure 1: Hodgkin-Huxley Circuit

Two voltage-dependent sodium and potassium channels allow the ions to flow through the membrane. The membrane has a capacitance indicated by c_m . Additionally, the sodium and potassium channels have a conductance parameter G_{Na} and G_K , as well as a leakage conductance as shown. These conductances can be modeled by various functions as will be shown later, and model the buildup in the passage of ions that create the action potential. The circuit is in parallel to mimic the axon.

Using classic current conservation in a circuit, and adding the parameter of an external current I_{ext} , we can derive the current equation for the cell.

$$c_m \frac{dV_m}{dt} + I_{ion} = I_{ext} \tag{1}$$

In our particular model, we will assume that G_L is constant and the capacitance c_m as well. V_m is the potential within the cell mebrane, and I_{ion} is the net ionic current within the cell.

The derivation of each ionic current through the different channels depends on the conductance parameter times the driving potential.

$$I_{ion} = G_{ion}(V_m - E_{ion}) \tag{2}$$

Conductance is modeled in a certain way. The physical meaning of conductance is the probability that the channels that the ions pass through will be considered open, or when the gates that the ions can pass through are all permissive in that particular ion channel. This is dependent on the voltage of the cell membrane.

Hodgkin and Huxley used three different types of gates to model the conductance channels. Using, m, h, and n, G_{Na} and G_{K} are modeled as followed.

$$G_{Na} = \bar{G}_{Na} m^3 h \tag{3}$$

$$G_K = \bar{G}_K n^4 \tag{4}$$

Where m, h, and n follow the differential equation:

$$\frac{dp}{dt} \tag{5}$$

3 Deriving a Reduced-State Model

Evidently, the state variables for the model are V, m, h, and n. Hence, we can write

$$C\frac{dV}{dt} = I_{ext} - \bar{G}_{Na}m^3h(V - E_{Na}) - \bar{G}_Kn^4(V - E_K) - \bar{G}_Lm^3h(V - E_L)$$

$$\frac{dm}{dt} = \frac{m_{\infty} - m}{\tau_m}$$

$$\frac{dh}{dt} = \frac{h_{\infty} - h}{\tau_h}$$

$$\frac{dn}{dt} = \frac{n_{\infty} - n}{\tau_n}$$

While analyzing a four-dimensional nonlinear system is certainly feasible, it is difficult to then visualize our results. Furthermore, the number of combinations of parameters grows exponentially with each new parameter introduced.

Hence, we seek to reduce the number of state variables in our model to gain better insights into our model both visually and analytically.

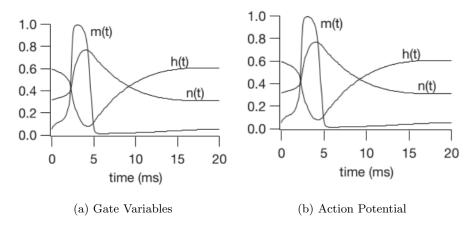


Figure 2: Dynamics of Action Potential on Shared Time-Scale (ms)

Firstly, it has been observed empirically that the time constants, τ_i vary widely between these equations. In fact, reasonable estimates are $\tau_m \approx 0.4ms$, $\tau_h \approx 8ms$, and $\tau_n \approx 5ms$.[2] Recall that τ_m is a measure of the speed with which m converges to its steady state value m_{∞} . Therefore, it may be a reasonable assumption to assume that τ_m is so fast compared to the other parameters that $m=m_{\infty}$. Then, we can either take n or h to be constant which would then give us a two-dimensional nonlinear system. Since h fluctuates over the largest time-scale, we'll take it to be constant. As a result, we have

$$C\frac{dV}{dt} = I_{ext} - \bar{G}_{Na} m_{\infty} (V - E_{Na}) - \bar{G}_K n (V - E_K) - \bar{G}_L (V - E_L)$$
$$\frac{dn}{dt} = \frac{n_{\infty} - n}{\tau_n}$$

The assumption of taking h to be constant can be judged visually in Figure 2, borrowed from Keener and Sneyd[1]. We'll develop this assumption further in a later section.

4 Finding a van der Pol Oscillator

In our previous simplification, taking m to converge infinitely fast to m_{∞} lent itself as a natural assumption due to its magnitude in speed relative to the other convergence speeds. However, taking h to be constant seemed to be a shakier assumption, which was evaluated using Figure 2.

5 Conclusion

References

- [1] James Keener and James Sneyd. Mathematical Physiology. Springer, 2009.
- [2] Eric Young. Nonlinear dynamics and hodgkin-huxley equations. Models of the Neuron, 2014.